# S<sub>4</sub> Flavored CP Symmetry for Neutrinos

#### Mohapatra, Nishi, PRD86, 073007(2012)

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#### 2 The model





#### • Harrison, Perkins, Scott ('02) $\rightarrow$ TBM

$$V_{
m MNS} pprox {\cal U}_{
m TB} = egin{pmatrix} \sqrt{2\over 3} & rac{1}{\sqrt{3}} & 0 \ -rac{1}{\sqrt{6}} & rac{1}{\sqrt{3}} & -rac{1}{\sqrt{2}} \ -rac{1}{\sqrt{6}} & rac{1}{\sqrt{3}} & rac{1}{\sqrt{2}} \end{pmatrix} \,.$$

- Very good approximation until recently!
- Daya-Bay, RENO, 2012:  $heta_{13} \approx 9^\circ$
- Assume still good approximation in some limit
- TBM from flavor symmetry?
- 3 families  $\rightarrow$  groups with 3-dim irreps



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# Tribimaximal mixing & flavor symmetries

- $U_l$  diagonalizes  $\overline{M}_l = M_l M_l^{\dagger}$ •  $V_{\rm MNS} = U_l^{\dagger} U_{\nu}$  $U_{\mu}$  diagonalizes  $M_{\mu}$  (Majorana)
- Fix a basis  $\rightarrow$  Flavor basis  $\rightarrow$  diagonal  $\bar{M}_l = \text{diag}(m_e^2, m_\mu^2, m_\tau^2)$
- Symmetries of TBM

 $G_l \simeq U(1) \times U(1)$  of  $\bar{M}_l$ :  $T^{\dagger} \bar{M}_l T = \bar{M}_l$  $T = egin{pmatrix} e^{ilpha_1} & 0 & 0 \ 0 & e^{ilpha_2} & 0 \ 0 & 0 & e^{-i(lpha_1+lpha_2)} \end{pmatrix}$  $G_{\nu} \simeq \mathbb{Z}_2 \times \mathbb{Z}_2$  of  $M_{\nu}$ :  $G_i^{\mathsf{T}} M_{\nu} G_i = M_{\nu}$  $G_2 = -rac{1}{3} egin{pmatrix} 1 & -2 & -2 \ -2 & 1 & -2 \ -2 & -2 & 1 \ \end{pmatrix} \quad G_3 = egin{pmatrix} 1 & 0 & 0 \ 0 & 0 & 1 \ 0 & 1 & 0 \ \end{pmatrix} \quad G_1 = G_2 G_3$ "magic" symmetry  $\mu$ - $\tau$  symmetry automatic (Grimus,Lavoura,Ludl,'09) ٠ specific to TBM

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#### Tribimaximal mixing & finite flavor symmetries

- Restrict to finite groups  $\rightarrow$  finite *T*
- Model:  $G_F \rightarrow G_I$  or  $G_{\nu}$  in each sector
- Flavor group  $G_F$  containing  $G_l, G_{\nu}|_{TR} \rightarrow$  should contain

 $S_4$ .am.'08

$$G_{F} = S_{4}$$

$$G_{I} = \langle T \rangle \simeq \mathbb{Z}_{3} \qquad T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^{2} \end{pmatrix} \qquad \omega \equiv e^{i2\pi/3}$$

$$G_{\nu} = \langle G_{2}, G_{3} \rangle \simeq \mathbb{Z}_{2} \times \mathbb{Z}_{2} \qquad G_{2} = -\frac{1}{3} \begin{pmatrix} 1 & -2 & -2 \\ -2 & 1 & -2 \\ -2 & -2 & 1 \end{pmatrix}$$

$$G_{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

• In another basis of  $G_F = S_4$ 

$$\begin{array}{c} G_{l} = \langle T \rangle \simeq \mathbb{Z}_{3} \\ \hline G_{\nu} = \langle G_{2}, G_{3} \rangle \simeq \mathbb{Z}_{2} \times \mathbb{Z}_{2} \\ \hline G_{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \\ \hline G_{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \end{array}$$

usual **3** representation of  $S_4$ 

• can be generated by T and

nd 
$$S = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}$$



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- Flavor symmetries v.s. residual symmetries of  $\overline{M}_l$  and  $M_{\nu}$
- Some residual symmetries may be accidental
- Different groups  $A_4, S_4, \Delta(27), T_7, \ldots$
- $\theta_{13} \neq 0 \implies \text{TBM is not exact!}$ 
  - Include corrections
  - Different symmetries

• 
$$\mu$$
- $\tau$  symmetry  $\implies \theta_{13} = 0$ 

•  $\mu$ - $\tau$  reflection  $\implies$   $\theta_{23} = 45^{\circ}, \delta_D = \pm 90^{\circ}$  if  $\theta_{13} \neq 0$ 

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} M_{\nu}^{*} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} = M_{\nu}$$

 $\mu$ - $\tau$  interchange w/ c.c. Harrison & Scott, '02 Grimus & Lavoura, '03

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# Incorporating **CP** in $S_4$

$$S_{4} \text{ generated by } S = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \qquad T^{3} = S^{4} = 1, \quad ST^{2}S = T$$

$$\tilde{S}_{4} \text{ generated by } \tilde{S} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \cdot CP \qquad T = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\tilde{S}_{4} \simeq S_{4} \text{ once we factor } CP^{2} = -1 \text{ for fermions}$$

$$\text{erecipe: extract } \tilde{S}_{4} \text{ subgroup from } S_{4} \otimes \langle CP \rangle$$

$$\text{e irreps of } \tilde{S}_{4} : 3, 1, 1_{\omega}, 1_{\omega^{2}} \text{ from}$$

$$\text{e irreps of } S_{4} : 3, 3', 1, 1', 2$$

# The model

3 families of	leptons $L_i, I_i$	<ul><li>3 Higgs doublets</li><li>4 Higgs triplets</li></ul>	$egin{array}{c c} \phi_i \ \Delta_0, \Delta_i \end{array} \sim$	(3,2)
Fransforming	under $\tilde{S}_4$ as			
$L_i \sim 3$	$L_i(x) \stackrel{\tilde{S}}{\longrightarrow} S_{ij}CL_j^*$	$f^*(\hat{x}), \ L_i(x) \stackrel{T}{\longrightarrow} T_{ij}$	$_{j}L_{j}(x); \hat{x}$	$\dot{x} = (x_0, -\mathbf{x})$
$\phi_i \sim 3$	$\phi_i(\mathbf{x}) \stackrel{\tilde{\mathbf{S}}}{\longrightarrow} \mathbf{S}_{ij}\phi_j^*(\mathbf{x})$	$\hat{x}$ ), $\phi_i(x) \stackrel{T}{\longrightarrow} T_{ij}\phi$	$\phi_j(x); \sim$	$\Delta_i \sim 3$
$l_1 \sim 1$	$l_1(x) \stackrel{\tilde{s}}{\longrightarrow} Cl_1^*(\hat{x})$	$, I_1(x) \xrightarrow{T} I_1(x);$	$\sim$	$\Delta_0 \sim \textbf{1}$
$I_2 \sim 1_\omega$	$l_2(x) \stackrel{\tilde{S}}{\longrightarrow} Cl_2^*(\hat{x})$	$\mathbf{k}), \ \mathbf{l}_2(\mathbf{x}) \xrightarrow{T} \omega  \mathbf{l}_2(\mathbf{x})$	<b>x</b> );	$\omega = e^{i2\pi/3}$
$I_3 \sim 1_{\omega^2}$	$l_3(x) \stackrel{\tilde{\mathbf{S}}}{\longrightarrow} Cl_3^*(x)$	$\hat{x}$ ), $l_3(x) \xrightarrow{T} \omega^2 l_3$	3( <i>X</i> );	문 ( 문 ) 문

#### The model: charged leptons

• Yukawa Lagrangian for charged leptons

$$\begin{split} -\mathcal{L}'_{Y} &= y_{1}(\bar{L}_{1}\phi_{1} + \bar{L}_{2}\phi_{2} + \bar{L}_{3}\phi_{3})l_{1} + y_{2}(\bar{L}_{1}\phi_{1} + \omega^{2}\bar{L}_{2}\phi_{2} + \omega\bar{L}_{3}\phi_{3})l_{2} \\ &+ y_{3}(\bar{L}_{1}\phi_{1} + \omega\bar{L}_{2}\phi_{2} + \omega^{2}\bar{L}_{3}\phi_{3})l_{3} + h.c., \end{split}$$

 $y_i$  real due to  $\tilde{S}$ 

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• doublet vevs 
$$\langle \phi_i \rangle = \frac{v}{\sqrt{3}}(1,1,1)$$
  $\tilde{S}_4 \rightarrow \tilde{S}_3$   
• mass matrix  $M_l = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \omega^2 & \omega \\ 1 & \omega & \omega^2 \end{pmatrix} \operatorname{diag}(m_e, m_\mu, m_\tau)$   
 $U_\omega^* \implies U_l^\dagger = U_\omega$ 



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# The model: neutrinos

#### • Type II seesaw

$$\begin{split} -\mathcal{L}^{\nu} &= \frac{1}{2} \textit{f}_0 \overline{\textit{L}_i^c} \epsilon \Delta_0 \textit{L}_i + \textit{f}_1 \left( \overline{\textit{L}_2^c} \epsilon \Delta_1 \textit{L}_3 + \overline{\textit{L}_3^c} \epsilon \Delta_2 \textit{L}_1 + \overline{\textit{L}_1^c} \epsilon \Delta_3 \textit{L}_2 \right) + \textit{h.c.}, \\ & \textit{f}_0, \textit{f}_1 \text{ real due to } \tilde{\textit{S}} \end{split}$$

• Triplet vevs 
$$\langle \Delta_0^{(0)} \rangle = u_0, \ \langle \Delta_i^{(0)} \rangle = u_i$$
  
• Mass matrix  $M_{\nu} = \begin{pmatrix} a & f & e \\ f & a & d \\ e & d & a \end{pmatrix}$   $a, d, e, f$  real  
• In flavor basis  $U_{\omega}^{\dagger} M_{\nu} U_{\omega}^* = \begin{pmatrix} x & z & z^* \\ z & -2z^* & y \\ z^* & y & -2z \end{pmatrix}$ 



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• In flavor basis  $U_{\omega}^{\dagger} M_{\nu} U_{\omega}^* = \begin{pmatrix} x & z & z^* \\ z & -2z^* & y \\ z^* & y & -2z \end{pmatrix}$   $\mu$ - $\tau$  reflection!  
accidental



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- Decompose  $V_{\text{MNS}} = U_l^{\dagger} U_{\nu} = U_{\text{TB}} \operatorname{diag}(1, 1, i) U_{\epsilon}$   $U_{\nu} = U U_{\epsilon}$
- $U_{\epsilon}$  diagonalizes  $M'_{\nu} = U^{\mathsf{T}} M_{\nu} U = \begin{pmatrix} a+d & b & 0 \\ b & a & c \\ 0 & c & a-d \end{pmatrix}$  $U = U^*_{\omega} U_{\mathsf{TB}} \operatorname{diag}(1,1,i)$   $b = \frac{e+f}{\sqrt{2}}, c = \frac{e-f}{\sqrt{2}}$
- TBM : *b* = *c* = 0
  - $m_1 = |d| a, m_2 = a, m_3 = a + |d|$
- $c \neq 0$  controls  $\theta_{13} \neq 0$
- $\bullet \ \text{maximal} \ \ \text{Dirac CP phase} \ \leftrightarrow \text{one maximal} \ \ \text{Majorana phase}$
- approximate sum-rule  $m_3 2m_2 m_1 \approx 0$
- $a, b, c, d \rightarrow \theta_{12}, \theta_{13}, m_1, m_2, m_3$



no phases

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 $m_1 = |d| - a$ ,  $m_2 = a$ ,  $m_3 = a + |d|$  only normal hierarchy

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 $m_1 = |d| - a$ ,  $m_2 = a$ ,  $m_3 = a + |d|$  only normal hierarchy

- $c \neq 0$  controls  $\theta_{13} \neq 0$
- maximal Dirac CP phase  $\leftrightarrow$  one maximal Majorana phase
- approximate sum-rule  $m_3 2m_2 m_1 \approx 0$  no phases

•  $a, b, c, d \rightarrow \theta_{12}, \theta_{13}, m_1, m_2, m_3$  correlations! 
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#### $(m_1, m_2, m_3) \approx (13, 16, 52) \,\mathrm{meV}$

- pts compatible within 1- $\sigma$
- "large masses", small mee
- similar to Ishimori & Ma ('12) but more restrictive



### Relation with more general approaches

- Holthausen, Lindner, Schmidt, arXiv:1211.6953
  - General theory: **CP**  $\leftrightarrow$  outer automorphism of  $G_F$
  - Clarification of geometric/calculable CP phases
  - Analysis of several relevant cases  $A_4, \Delta(27), T_7, \ldots$
  - Consequence to our work:

Ŝ₄ is the only consistent CP extension of  $A_4$  (if 3, 1' are present)

- Feruglio, Hagedorn, Ziegler, arXiv:1211.5560
  - General consequences of including CP
  - Focus on residual CP symmetry in neutrino sector  $G_{\nu}$

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- Dirac and Majorana phases depending on one angle
- Specific analysis for S<sub>4</sub> (not a complete model)



# Conclusions

- A consistent way of incorporating CP into flavor groups
- The only way of defining CP to A<sub>4</sub> (Holthausen, et.al., '12)
- $\tilde{S}_4$  example, similar to  $A_4$  but more constraining
- Can be further explored for flavor model-building
- Accidental  $\mu$ - $\tau$  reflection leading to maximal  $\delta_D$ ,  $\theta_{23}$
- NH and approximate sum-rule (10-15%)
- One maximal Majorana phase



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