

“Higgses Masses in some Supersymmetric
3-3-1 Models”

Marcos C. Rodriguez

*Universidade Federal do Rio Grande
Instituto de Matemática, Estatística e Física
Av. Itália km 8, Campus Carreiros
96201-900, Rio Grande, RS, Brazil*

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models $SU(3)_C \otimes SU(3)_L \otimes U(1)_N$ 3-3-1

$$\frac{Q}{e} = \frac{1}{2}(\lambda_3 - \vartheta\lambda_8) + N I$$

∅ two different models (1992-1994)

motivations to study this kind of model

1. The family number must be three
2. It explains why $\sin^2 \theta_W < \frac{1}{4}$ is observed at the Z-pole
3. It is the simplest model that includes bileptons of both types: scalar and vector ones
4. The model has several sources of CP violation

First Model **MSUSY331** $\vartheta = \sqrt{3}$

- T. V. Duong and E. Ma, Phys. Lett. **B316**, 307 (1993) (Scalar Sector)
- H. N. Long and P. B. Pal, Mod. Phys. Lett. **A13**, 2355 (1998) (Proton Decay)
- J. C. Montero, V. Pleitez and M. C. Rodriguez, Phys. Rev. **D65**, 035006 (2002)
- M. Capdequi-Peyranère and M.C. Rodriguez, Phys. Rev. **D65**, 035001 (2002)

Second Model **SUSY331RN** $\vartheta = 1/\sqrt{3}$

- J. C. Montero, V. Pleitez and M. C. Rodriguez, Phys. Rev. **D70**, 075004 (2004)
- D. T. Huong, M. C. Rodriguez and H. N. Long, [arXiv:hep-ph/0508045]

More References on these models will be presented during this seminar

MSUSY331 Particle Content

$$L_{aL} \sim \begin{pmatrix} \nu_a \\ l_a \\ l_a^c \end{pmatrix}_L \sim (\mathbf{1}, \mathbf{3}, 0) \quad a = e, \mu, \tau$$

Higgs

$$\eta = \begin{pmatrix} \eta^0 \\ \eta_1^- \\ \eta_2^+ \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, 0)$$

$$\rho = \begin{pmatrix} \rho^+ \\ \rho^0 \\ \rho^{++} \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, +1)$$

$$\chi = \begin{pmatrix} \chi^- \\ \chi^{--} \\ \chi^0 \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, -1)$$

$$S \sim (\mathbf{1}, \mathbf{6}^*, 0)$$

Supersymmetric Partners:

$$\tilde{\eta} = \begin{pmatrix} \tilde{\eta}^0 \\ \tilde{\eta}_1^- \\ \tilde{\eta}_2^+ \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, 0)$$

$$\tilde{\rho} = \begin{pmatrix} \tilde{\rho}^+ \\ \tilde{\rho}^0 \\ \tilde{\rho}^{++} \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, +1)$$

$$\tilde{\chi} = \begin{pmatrix} \tilde{\chi}^- \\ \tilde{\chi}^{--} \\ \tilde{\chi}^0 \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, -1)$$

$$\tilde{S} \sim (\mathbf{1}, \mathbf{6}^*, 0)$$

Cancel Chiral Anomaly

$$\tilde{\eta}' = \begin{pmatrix} \tilde{\eta}'^0 \\ \tilde{\eta}'_1^+ \\ \tilde{\eta}'_2^- \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}^*, 0)$$

$$\tilde{\rho}' = \begin{pmatrix} \tilde{\rho}'^- \\ \tilde{\rho}'^0 \\ \tilde{\rho}'^{--} \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}^*, -1)$$

$$\tilde{\chi}' = \begin{pmatrix} \tilde{\chi}'^+ \\ \tilde{\chi}'^{++} \\ \tilde{\chi}'^0 \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}^*, +1)$$

$$\tilde{S}' \sim (\mathbf{1}, \mathbf{6}^*, 0)$$

Lagrangian

$$\mathcal{L}_{331S} = \mathcal{L}_{SUSY} + \mathcal{L}_{\text{soft}}$$

$$\begin{aligned} \mathcal{L}_{SUSY} = & \mathcal{L}_{\text{Lepton}} + \mathcal{L}_{\text{Quarks}} + \mathcal{L}_{\text{Gauge}} \\ & + \mathcal{L}_{\text{Scalar}} \end{aligned}$$

symmetry breaking

$$\begin{array}{l} \text{MSUSY}_{331} \xrightarrow{\mathcal{L}_{\text{soft}}} \text{SU}(3)_C \otimes \text{SU}(3)_L \otimes \text{U}(1)_N \\ \xrightarrow{\langle \chi \rangle \langle \chi' \rangle} \text{SU}(3)_C \otimes \text{SU}(2)_L \otimes \text{U}(1)_Y \\ \xrightarrow{\langle \rho, \eta, \rho', \eta' \rangle} \text{SU}(3)_C \otimes \text{U}(1)_Q \end{array}$$

$$\begin{aligned} \mathcal{L}_{\text{Fermion}} &= \mathcal{L}_{ffV} + \mathcal{L}_{\tilde{f}\tilde{f}V} + \mathcal{L}_{f\tilde{f}\tilde{V}} + \mathcal{L}_{\tilde{f}\tilde{f}VV} \\ &+ \mathcal{L}_{\text{cin}}^{\text{Fermion}} + \mathcal{L}_F^{\text{Fermion}} + \mathcal{L}_D^{\text{Fermion}} \end{aligned}$$

$$\mathcal{L}^{\text{gauge}} = \mathcal{L}_{dc} + \mathcal{L}_D^{\text{gauge}}$$

$$\begin{aligned} \mathcal{L}^{\text{scalar}} &= \mathcal{L}_F^{\text{scalar}} + \mathcal{L}_D^{\text{scalar}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Higgsinos}} \\ &+ \mathcal{L}_{H\tilde{H}\tilde{V}} \end{aligned}$$

W is the superpotential

$$W = \frac{W_2}{2} + \frac{W_3}{3}$$

$$W_2 = \mathcal{L}_F^{W^2} + \mathcal{L}_{\tilde{H}L} + \mathcal{L}_{HMT},$$

$$\begin{aligned} W_3 &= \mathcal{L}_F^{W^3} + \mathcal{L}_{uH} + \mathcal{L}_{\tilde{l}\tilde{H}} + \mathcal{L}_{l\tilde{H}H} + \mathcal{L}_{\tilde{l}\tilde{H}\tilde{H}} + \mathcal{L}_{H\tilde{H}\tilde{H}} \\ &+ \mathcal{L}_{qqH} + \mathcal{L}_{q\tilde{q}\tilde{H}} + \mathcal{L}_{lq\tilde{q}} + \mathcal{L}_{qq\tilde{q}} + \mathcal{L}_{qq\tilde{l}} \end{aligned}$$

J. C. Montero, V. Pleitez and M. C. Rodriguez, Phys. Rev. **D65**, 095008 (2002).

$$W_2 = \mu_{0a} \hat{L}_a \hat{\eta}' + \mu_\eta \hat{\eta} \hat{\eta}' + \mu_\rho \hat{\rho} \hat{\rho}' + \mu_\chi \hat{\chi} \hat{\chi}' + \mu_S \hat{S} \hat{S}'$$

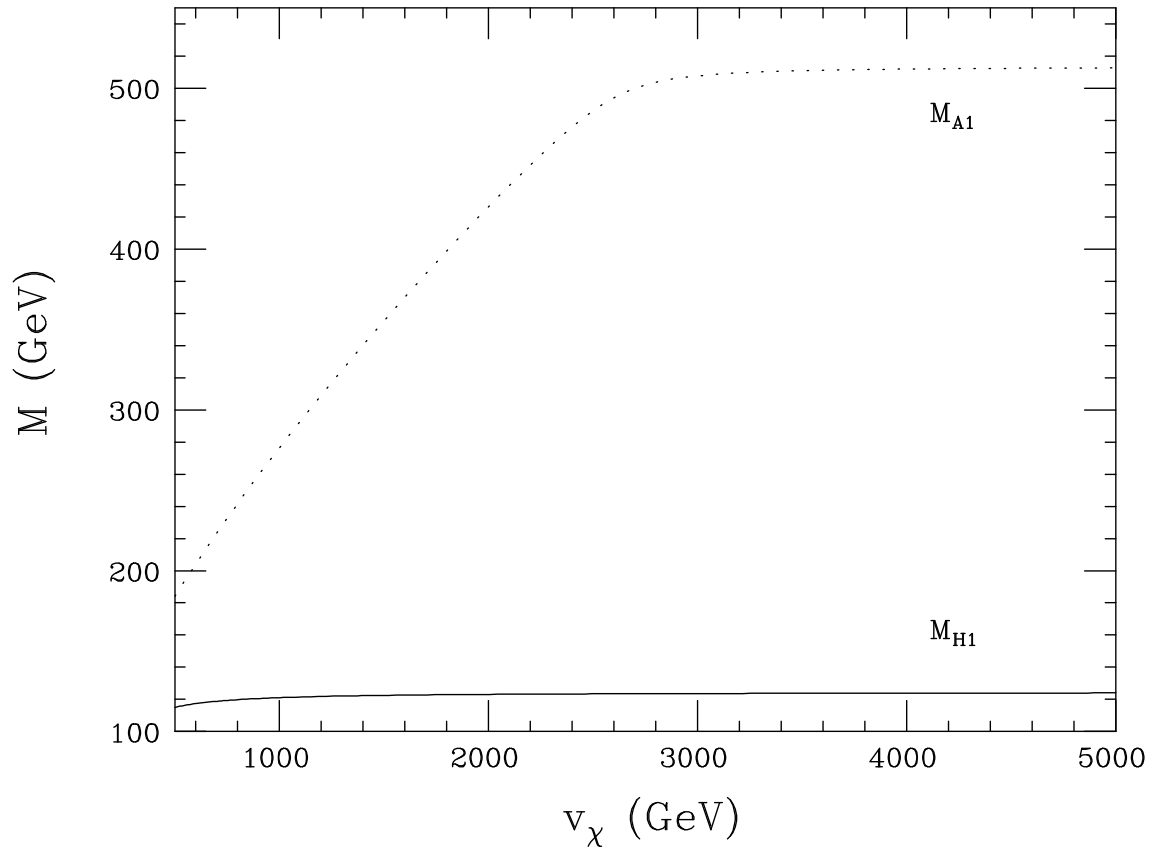
$$\begin{aligned} & y \quad \hat{L} \hat{L} \hat{S} + \lambda_{1abc} \epsilon \hat{L}_a \hat{L}_b \hat{L}_c + \lambda_{2ab} \epsilon \hat{L}_a \hat{L}_b \hat{\eta} + \lambda_{3a} \hat{L}_a \hat{\rho} \hat{\chi} \\ & + f_1 \epsilon \hat{\eta} \hat{\rho} \hat{\chi} + f'_1 \epsilon \hat{\eta}' \hat{\rho}' \hat{\chi}' + \lambda'_{\alpha ai} \hat{Q}_\alpha \hat{L}_a \hat{d}_i^c \\ & + \kappa_{1i} \hat{Q}_1 \hat{\eta}' \hat{u}_i^c + \kappa_{2i} \hat{Q}_1 \hat{\rho}' \hat{d}_i^c + \kappa_3 \hat{Q}_1 \hat{\chi}' \hat{J}^c \\ & + \kappa_{4\alpha i} \hat{Q}_\alpha \hat{\eta} \hat{d}_i^c + \kappa_{5\alpha i} \hat{Q}_\alpha \hat{\rho} \hat{u}_i^c + \kappa_{6\alpha\beta} \hat{Q}_\alpha \hat{\chi} \hat{J}_\beta^c \end{aligned}$$

$\mu_0, \lambda_2, \lambda_3$ interactions give masses to charged leptons and neutrinos complex generate Leptogenesis $\rho^0 \rightarrow \tilde{\chi}^0 l$

M. C. Rodriguez, *Journ. Mod. Phys.* **2**, 1193 (2011)

1. The proton is stable at tree level (No $\widehat{u}\widehat{u}\widehat{d}$) $p \rightarrow e^+\pi^0$;
2. Allow Double Beta Decay without Neutrinos $dd \rightarrow uuee$;
3. New contributions to the muon decay;
4. $D^+ \rightarrow \overline{K^0}l_i^+\nu_i$; $K^+ \rightarrow \pi^+\nu\bar{\nu}$;
5. New contributions to the Neutrals $K\bar{K}$ and also $B\bar{B}$ Systems (SUPER CKM matrix);
6. Hadronic B Meson Decay.

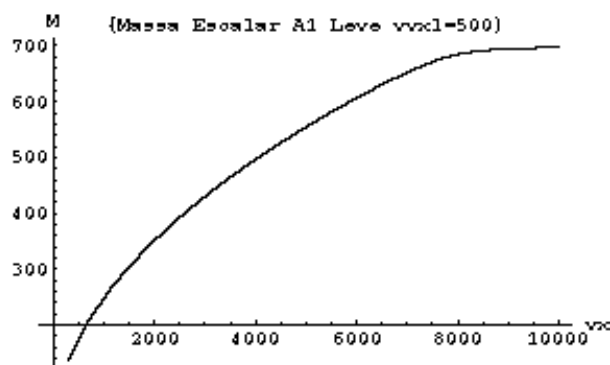
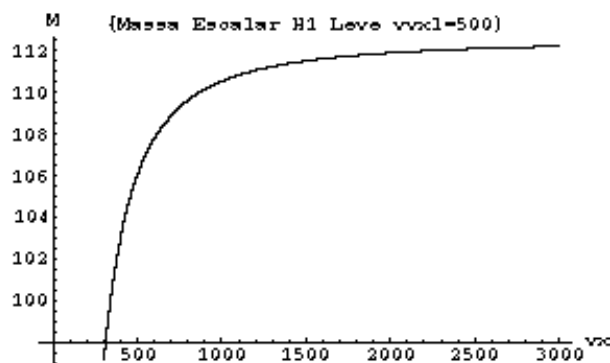
scalar sector with anti-sextet



$$M_{H_1}^2 \leq (124.5 + \epsilon)^2 \text{ GeV}^2$$

124.5 GeV tree value ($\epsilon = 0$).

M. C. Rodriguez, Int. J. Mod. Phys. **A21**,
4303 (2006)



$$M_{H_1}^2 \leq (114 + \epsilon)^2 \text{ GeV}^2$$

$$M_{A_1} > M_Z$$

SUSY331RN Particle Content

$$L_{aL} = \begin{pmatrix} \nu_a \\ l_a \\ \nu_a^c \end{pmatrix}_L \sim (\mathbf{1}, \mathbf{3}, -1/3)$$

$$a = 1, 2, 3$$

$$l_{aL}^c \sim (\mathbf{1}, \mathbf{1}, \mathbf{1})$$

$$\eta = \begin{pmatrix} \eta_1^0 \\ \eta^- \\ \eta_2^0 \end{pmatrix}$$
$$\chi = \begin{pmatrix} \chi_1^0 \\ \chi^- \\ \chi_2^0 \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, -1/3)$$
$$\rho = \begin{pmatrix} \rho_1^+ \\ \rho^0 \\ \rho_2^+ \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, 2/3)$$

Supersymmetric Partners:

$$\begin{aligned}\tilde{\eta} &= \begin{pmatrix} \tilde{\eta}_1^0 \\ \tilde{\eta}^- \\ \tilde{\eta}_2^0 \end{pmatrix} \\ \tilde{\chi} &= \begin{pmatrix} \tilde{\chi}_1^0 \\ \tilde{\chi}^- \\ \tilde{\chi}_2^0 \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, -1/3) \\ \tilde{\rho} &= \begin{pmatrix} \tilde{\rho}_1^+ \\ \tilde{\rho}^0 \\ \tilde{\rho}_2^+ \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, 2/3)\end{aligned}$$

Cancel Chiral Anomaly

$$\begin{aligned}\tilde{\eta}' &= \begin{pmatrix} \tilde{\eta}'^0 \\ \tilde{\eta}'^+ \\ \tilde{\eta}'^0 \end{pmatrix} \\ \tilde{\chi}' &= \begin{pmatrix} \tilde{\chi}'^0 \\ \tilde{\chi}'^+ \\ \tilde{\chi}'^0 \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}^*, 1/3) \\ \tilde{\rho}' &= \begin{pmatrix} \tilde{\rho}'^- \\ \tilde{\rho}'^0 \\ \tilde{\rho}'^- \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}^*, -2/3)\end{aligned}$$

Lagrangian

$$\mathcal{L}_{SUSY331RN} = \mathcal{L}_{SUSY} + \mathcal{L}_{soft}$$

$$\begin{aligned} \mathcal{L}_{SUSY} = & \mathcal{L}_{Lepton} + \mathcal{L}_{Quark} + \mathcal{L}_{Gauge} \\ & + \mathcal{L}_{Scalar} \end{aligned}$$

\mathcal{L}_{Gauge} equal as MSUSYS331 P8

symmetry breaking

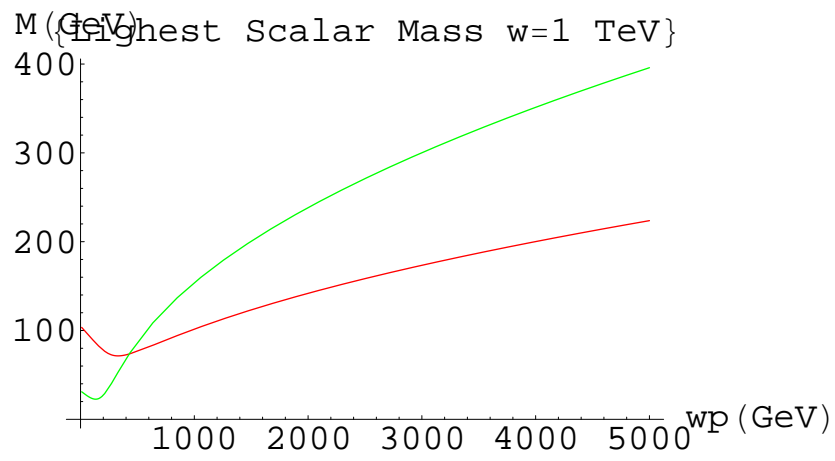
$$\begin{array}{l} \text{SUSY331RN} \xrightarrow{\mathcal{L}_{soft}} \text{SU}(3)_C \otimes \text{SU}(3)_L \otimes \text{U}(1)_N \\ \begin{array}{l} \langle \chi \rangle \langle \chi' \rangle \\ \xrightarrow{\quad} \end{array} \text{SU}(3)_C \otimes \text{SU}(2)_L \otimes \text{U}(1)_Y \\ \begin{array}{l} \langle \rho, \eta, \rho', \eta' \rangle \\ \xrightarrow{\quad} \end{array} \text{SU}(3)_C \otimes \text{U}(1)_Q \end{array}$$

P. V. Dong, D. T. Huong, M. C. Rodriguez and H. N. Long, Eur. Phys. J. **C48**, 229 (2006)

$$\begin{aligned}
W_{2RC} &= \mu_0 \hat{L} \hat{\eta}' + \mu_1 \hat{L} \hat{\chi}' + \mu_\eta \hat{\eta} \hat{\eta}' + \mu_\chi \hat{\chi} \hat{\chi}' \\
&+ \mu_2 \hat{\eta} \hat{\chi}' + \mu_3 \hat{\chi} \hat{\eta}' + \mu_\rho \hat{\rho} \hat{\rho}' \\
W_{3RC} &= \lambda_{1ab} \hat{L}_a \hat{\rho}' \hat{l}_b^c + \lambda_{2a} \hat{L}_a \hat{\chi} \hat{\rho} + \lambda_{3a} \hat{L}_a \hat{\eta} \hat{\rho} \\
&+ \lambda_{4ab} \epsilon \hat{L}_a \hat{L}_b \hat{\rho} + f_1 \epsilon \hat{\rho} \hat{\chi} \hat{\eta} + f_1' \epsilon \hat{\rho}' \hat{\chi}' \hat{\eta}' \\
&+ \kappa_{1i} \hat{Q}_3 \hat{\eta}' \hat{u}_i^c + \kappa_1' \hat{Q}_3 \hat{\eta}' \hat{T}^c + \kappa_{2i} \hat{Q}_3 \hat{\chi}' \hat{u}_i^c \\
&+ \kappa_2' \hat{Q}_3 \hat{\chi}' \hat{T}^c + \kappa_{3\alpha i} \hat{Q}_\alpha \hat{\eta} \hat{d}_i^c + \kappa_{3\alpha\beta}' \hat{Q}_\alpha \hat{\eta} \hat{D}_\beta^c \\
&+ \kappa_{4\alpha i} \hat{Q}_\alpha \hat{\rho} \hat{u}_i^c + \kappa_{4\alpha}' \hat{Q}_\alpha \hat{\rho} \hat{T}^c + \kappa_{5i} \hat{Q}_3 \hat{\rho}' \hat{d}_i^c \\
&+ \kappa_{5\beta}' \hat{Q}_3 \hat{\rho}' \hat{D}_\beta^c + \kappa_{6\alpha i} \hat{Q}_\alpha \hat{\chi} \hat{d}_i^c + \kappa_{6\alpha\beta}' \hat{Q}_\alpha \hat{\chi} \hat{d}_\beta^c \\
&+ \lambda'_{\alpha ai} \hat{Q}_\alpha \hat{L}_a \hat{d}_i^c + \xi_{2\alpha\alpha\beta} \hat{Q}_\alpha \hat{L}_a \hat{d}_\beta^c
\end{aligned}$$

Same nice phenomenological aspect as in MSUSY331
 Pags(9,10)

P.V.Dong, D.T.Huong, M.C.Rodriguez and H.N.Long, *Journ. of Modern Physics* **2**, 792, (2011)



$$M_{A_1} \approx 5\text{GeV}^2$$

A_1 does not couple to the Z^0