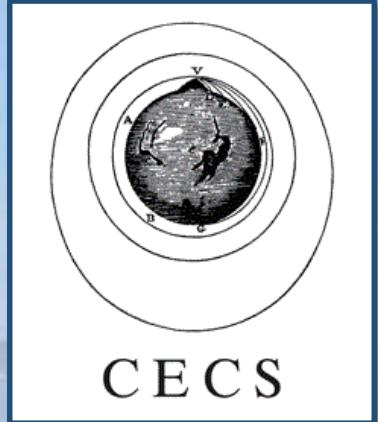




PONTIFICIA UNIVERSIDAD
CATHOLICA
DE VALPARAISO



Gravity around a spherically symmetric Spinorial Source

Adolfo Toloza Díaz (PUCV-CECs)
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Outline

- Motivation
- Einstein-Cartan-Holts Theory
- Stationary & Spherically Symmetric Solutions
- Conclusions and further work

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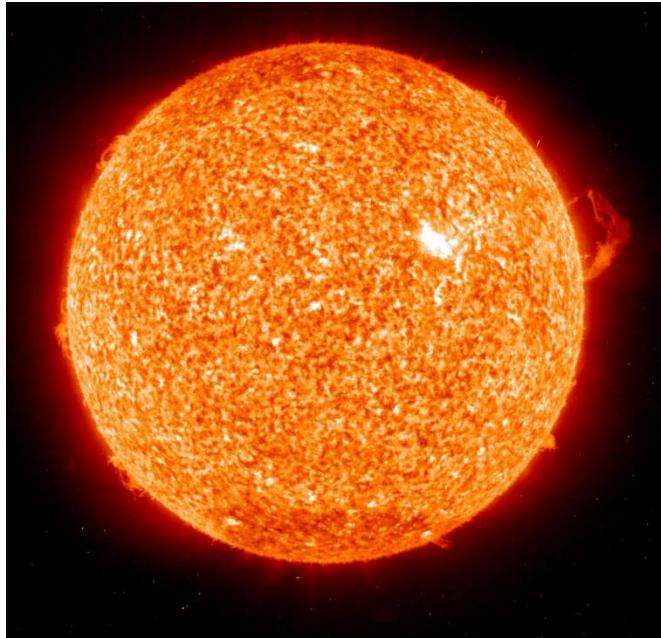
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Motivation

- Cartan formalism (1922)
- Poincaré Gauge Theory (Kibble-Sciama 1961-1962)



Élie Cartan (1869-1951)

¿What is Matter made of?

First Order Formalism

	Translations	Rotations
Field	$e^a = e^a_{\mu} dx^{\mu}$	$\omega^{ab} = \omega^{ab}_{\mu} dx^{\mu}$
Structure equations	$T^a = de^a + \omega^a_b \wedge e^b$	$R^{ab} = d\omega^{ab} + \omega^a_c \wedge \omega^{cb}$
Bianchi Identities	$D T^a = R^a_b e^b$	$D R^{ab} = 0$

Einstein-Cartan-Holts Theory

Action Principle ($\kappa = 8\pi G$)

$$I[e, \omega, \psi] = \frac{1}{4\kappa} \int (\epsilon_{abcd} R^{ab} \wedge e^c \wedge e^d + 2\gamma R_{ab} \wedge e^a \wedge e^b) + \int L_M[\psi]$$

Equations of Motion

$$\delta e : -\frac{1}{2} \epsilon_{abcd} R^{ab} \wedge e^c + \gamma R_{da} \wedge e^a = \kappa \tau_d$$

$$\delta \omega : \epsilon_{abcd} T^c \wedge e^d + 2\gamma T_{[a} \wedge e_{b]} = \kappa \sigma_{ab}$$

$$\delta \psi : \frac{\delta L_M}{\delta \psi} = 0$$

where

$$\delta L_M = -\delta e^a \wedge \tau_a - \frac{1}{2} \delta \omega^{ab} \wedge \sigma_{ab} + \frac{\delta L_M}{\delta \psi} \delta \psi$$

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(non-Grassmann) Dirac Field

Consider

$$\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix} \quad , \quad \bar{\psi} = -i\psi^\dagger \Gamma^0$$

$$\Gamma = \Gamma_a e^a \quad , \quad \{\Gamma_a, \Gamma_b\} = 2\eta_{ab} \mathbb{I} \quad , \quad \Gamma_{ab} = \frac{1}{2} [\Gamma_a, \Gamma_b]$$

Quiral Representation

$$\Gamma_0 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \quad , \quad \Gamma_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} \quad , \quad \Gamma^5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

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Covariant Derivative ($\Gamma_{ab} = \frac{1}{2}[\Gamma_a, \Gamma_b]$)

$$D\psi = d\psi + \frac{1}{4}\omega^{ab}\Gamma_{ab}\psi$$

$$D\bar{\psi} = d\bar{\psi} - \frac{1}{4}\omega^{ab}\bar{\psi}\Gamma_{ab}$$

$$D\Gamma^a = 0$$

define

(vector current) $J^a = i\bar{\psi}\Gamma^a\psi$

(quiral current) $J^{5a} = i\bar{\psi}\Gamma^5\Gamma^a\psi$

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Matter Lagrangian

Impose conditions:

- Lorentz & diffeomorphism invariance
- First order equations
- Minimal Coupling
- Real Action

$$I_M[\psi] = \int \left(\frac{1}{2} \bar{\psi} \star \Gamma \wedge D\psi + \frac{1}{2} D\bar{\psi} \wedge \star \Gamma \psi - \frac{1}{2} m\bar{\psi}\psi e^4 \right) \\ + in \left(\frac{1}{2} \bar{\psi} \star \Gamma \wedge D\psi - \frac{1}{2} D\bar{\psi} \wedge \star \Gamma \psi \right)$$

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 $e^4 = |e|d^4x$: Volume 4-form

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Additional term

$$\begin{aligned}in \left(\frac{1}{2} \bar{\psi} \star \Gamma \wedge D\psi - \frac{1}{2} D\bar{\psi} \wedge \star \Gamma\psi \right) &= -\frac{in}{2} d(\bar{\psi} \star \Gamma\psi) + \frac{in}{2} \bar{\psi} D \star \Gamma\psi \\&= \frac{n}{2} T^a \bar{\psi} \star (\Gamma e_a) \psi + b.t. \\&= \frac{n}{2} i(\bar{\psi} \Gamma^a \psi) T^b {}_{ab} |e| d^4x + b.t.\end{aligned}$$

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where

$$\alpha = 1 + in$$

Equations of Motion

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$$(\text{Dirac}) \quad \bar{\not{D}}\psi - m\psi - \frac{3i}{4}A_a\Gamma^5\Gamma^a\psi + \frac{3ni}{4}V_a\Gamma^a\psi = 0$$

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Splitting Torsion

$$T_{abc} = \epsilon_{abcd} A^d + \eta_{a[b} V_{c]} + M_{abc}$$

$$\eta^{ab} M_{abc} = \epsilon^{abcd} M_{abc} = 0$$

From Cartan Equation

$$A_a = \frac{\kappa}{2(1+\gamma^2)} (-J_a^5 + n\gamma J_a) \quad , \quad M^a{}_{bc} = 0$$

$$V_a = -\frac{\kappa}{2(1+\gamma^2)} (n J_a + \gamma J_a^5)$$

$$T^a{}_{bc} = \frac{\kappa}{2(1+\gamma^2)} (\epsilon^a{}_{bcd} (-J^{5d} + n\gamma J^d) - n\delta^a_{[b} J_{c]} - \gamma\delta^a_{[b} J_{c]}^5)$$

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Stationary & Spherically Symmetric Solutions

Metric:

$$ds^2 = -f(r)dt^2 + h(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

Torsion:

$$T^a_{bc} = T^a_{bc}(r)$$

$$A_2 = A_3 = V_2 = V_3 = 0$$

first constraints

$$\bar{\psi}\Sigma_2\psi = \bar{\psi}\Sigma_3\psi = 0$$

$$n\bar{\psi}\Gamma_{02}\psi = n\bar{\psi}\Gamma_{03}\psi = 0$$

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Stationary & Spherically Symmetric Solutions

Metric:

$$ds^2 = -f(r)dt^2 + h(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

Torsion:

$$T^a_{bc} = T^a_{bc}(r)$$

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Stationary & Spherically Symmetric Solutions

Vielbein: $(g_{\mu\nu} = \eta_{ab} e^a{}_\mu e^b{}_\nu, \eta_{ab} = \text{diag}(-1,1,1,1))$

$$e^0 = f^{1/2} dt, \quad e^1 = h^{1/2} dr, \quad e^2 = r d\theta, \quad e^3 = \sin \theta d\phi^2$$

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$$V_0 = \frac{\kappa}{2(1+\gamma^2)} (-n \psi^\dagger \psi + \gamma \psi^\dagger \Gamma^5 \psi)$$

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Stationary & Spherically Symmetric Solutions

Spin Connection: from first structure equation $T^a = de^a + \omega^a{}_b \wedge e^b$

$$\omega^{01} = \frac{1}{2} \left(\frac{f'}{fh^{1/2}} + V_1 \right) e^0 + \frac{1}{2} V_0 e^1$$

$$\omega^{02} = \frac{1}{2} V_0 e^2 + \frac{1}{2} A_1 e^3$$

$$\omega^{03} = \frac{1}{2} V_0 e^3 - \frac{1}{2} A_1 e^2$$

$$\omega^{12} = - \left(\frac{1}{rh^{1/2}} + \frac{1}{2} V_1 \right) e^2 - \frac{1}{2} A_0 e^3$$

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Einstein Equations

$$\frac{h'}{rh^2} + \frac{1}{r^2} \left(1 - \frac{1}{h}\right) = \kappa \left(\frac{i}{2h^{1/2}} \psi^\dagger \Gamma_{01} \overleftrightarrow{\partial}_r \psi + \frac{i}{2r} \psi^\dagger \Gamma_{02} \overleftrightarrow{\partial}_\theta \psi + \frac{i}{2r \sin \theta} \psi^\dagger \Gamma_{03} \overleftrightarrow{\partial}_\phi \psi + \Delta + m \bar{\psi} \psi \right)$$

$$-\frac{f'}{rfh} + \frac{1}{r^2} \left(1 - \frac{1}{h}\right) = \kappa \left(\frac{i}{2f^{1/2}} \psi^\dagger \overleftrightarrow{\partial}_t \psi + \frac{i}{2r} \psi^\dagger \Gamma_{02} \overleftrightarrow{\partial}_\theta \psi + \frac{i}{2r \sin \theta} \psi^\dagger \Gamma_{03} \overleftrightarrow{\partial}_\phi \psi + \Delta + m \bar{\psi} \psi \right)$$

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where

$$\Delta = -\frac{3n}{8} V_0 \psi^\dagger \psi - \frac{3n}{8} V_1 \psi^\dagger \Gamma_{01} \psi - \frac{3}{8} A_0 \psi^\dagger \Gamma^5 \psi + \frac{3}{8} A_1 \psi^\dagger \Sigma_1$$

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Dirac Equation

$$\begin{aligned} & \frac{1}{f^{1/2}} \partial_t \psi + \frac{1}{h^{1/2}} \Gamma_{01} \partial_r \psi + \frac{1}{r} \Gamma_{02} \partial_\theta \psi + \frac{1}{r \sin \theta} \Gamma_{03} \partial_\phi \psi + \frac{1}{h^{1/2}} \left(\frac{f'}{4f} + \frac{1}{r} \right) \Gamma_{01} \psi \\ & + \frac{\cot \theta}{2r} \Gamma_{02} \psi - m \Gamma_0 \psi + \frac{3i}{4} (nV_0 + nV_1 \Gamma_{01} + A_0 \Gamma^5 - A_1 \Sigma_1) \psi = 0 \end{aligned}$$

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additional constraints

plus additional constraints

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Case $n \neq 0$

Solution of constraints

$$\psi_L = \Phi(r)e^{i\lambda(x)} \begin{pmatrix} 1 \\ \pm 1 \end{pmatrix}, \quad \psi_R = \Phi(r)e^{-i\lambda(x)} \begin{pmatrix} 1 \\ \mp 1 \end{pmatrix}$$

here

$$\psi^\dagger \Gamma^5 \psi = \psi^\dagger \Sigma_1 \psi = 0 \quad \text{but} \quad \psi^\dagger \psi = 4\Phi^2 \quad \& \quad \psi^\dagger \Gamma_{01} \psi = \mp 4\Phi^2$$

so

$$j_a = 4\Phi^2(1, \mp 1, 0, 0) \quad \& \quad j_a^5 = (0, \vec{0})$$

$$V_a = -\frac{2n\kappa}{(1+\gamma^2)} \Phi^2(1, \pm 1, 0, 0)$$

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Final Equations for $n \neq 0$

Einstein Equations

$$\begin{aligned}\frac{h'}{rh^2} + \frac{1}{r^2} \left(1 - \frac{1}{h}\right) &= 0 \\ -\frac{f'}{rfh} + \frac{1}{r^2} \left(1 - \frac{1}{h}\right) &= 0 \\ -\frac{1}{2(fh)^{1/2}} \left(\frac{f'}{(fh)^{1/2}}\right)' + \frac{1}{2rf} \left(\frac{f}{h}\right)' &= 0\end{aligned}$$

Dirac Equation

$$\begin{aligned}\Phi' + \left(\frac{f'}{4f} + \frac{1}{r}\right)\Phi &= 0 \\ \Phi \left(\partial_t \lambda \mp \left(\frac{f}{h}\right)^{\frac{1}{2}} \partial_r \lambda\right) &= 0 \\ \Phi (\partial_\theta \lambda \pm mr \cos 2\lambda) &= 0 \\ \Phi \left(\partial_\phi \lambda \mp \frac{1}{2} \cos \theta - rm \sin 2\lambda \sin \theta\right) &= 0\end{aligned}$$

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Solutions for $n \neq 0$

Metric

$$f = \frac{1}{h} = 1 - \frac{2GM}{r}$$

Spinor

$$\Phi = \frac{\Phi_0}{rf^{1/4}}$$

Torsion

$$V_a = -\frac{2n\kappa}{(1 + \gamma^2)} \frac{\Phi_0^2}{r^2 f^{1/2}} (1, \pm 1, 0, 0)$$

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Spinor Density

non trivial spinor density

$$j^\mu = \rho(r)(1, \mp 1, 0, 0)$$

$$\rho(r) = \frac{\Phi_0^2}{r^2 |1 - R_H/r|}$$

$$Q(R) \equiv 4\pi \int_0^R \rho(r) r^2 dr$$

$$Q(R_H - \epsilon) \rightarrow -4\pi R_H \Phi_0^2 \text{Log} \left(\frac{\epsilon}{R_H} \right)$$

$$Q(R \rightarrow \infty) \rightarrow 4\pi R \Phi_0$$

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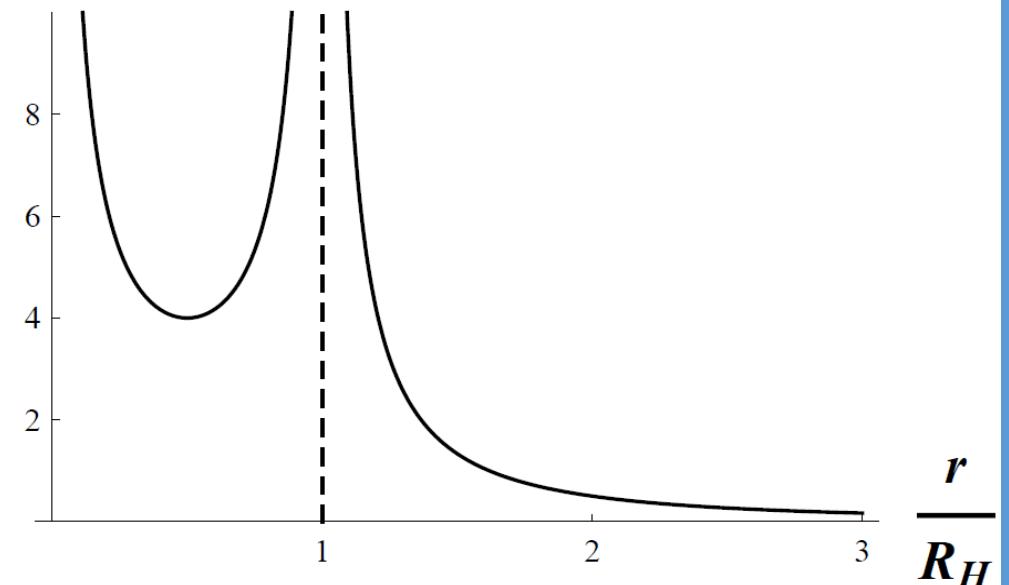
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Spinor Density

$$\rho(r)(R_H/\Phi_0)^2$$



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non trivial spinor density

$$j^\mu = \rho(r)(1, \mp 1, 0, 0)$$

$$\rho(r) = \frac{\Phi_0^2}{r^2 |1 - R_H/r|}$$

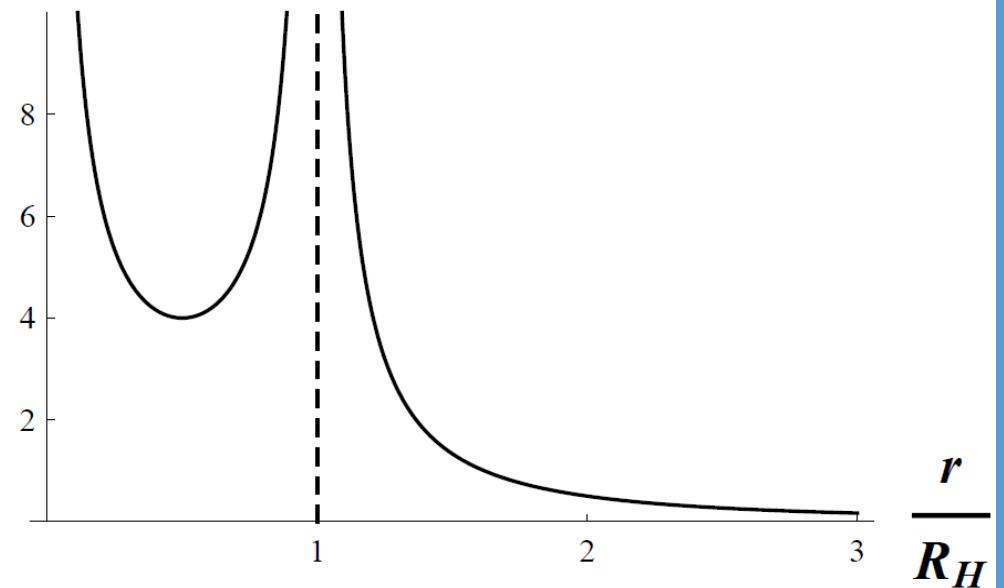
$$Q(R) \equiv 4\pi \int_0^R \rho(r) r^2 dr$$

$$Q(R_H - \epsilon) \rightarrow -4\pi R_H \Phi_0^2 \log\left(\frac{\epsilon}{R_H}\right)$$

$$Q(R \rightarrow \infty) \rightarrow 4\pi R \Phi_0$$

Spinor Density

$$\rho(r)(R_H/\Phi_0)^2$$



Torsion

Effect of Immirzi parameter

$$V_a \sim \frac{1}{(1 + \gamma^2)} n \rho(r)$$

$$A_a \sim \frac{\gamma}{(1 + \gamma^2)} n \rho(r)$$

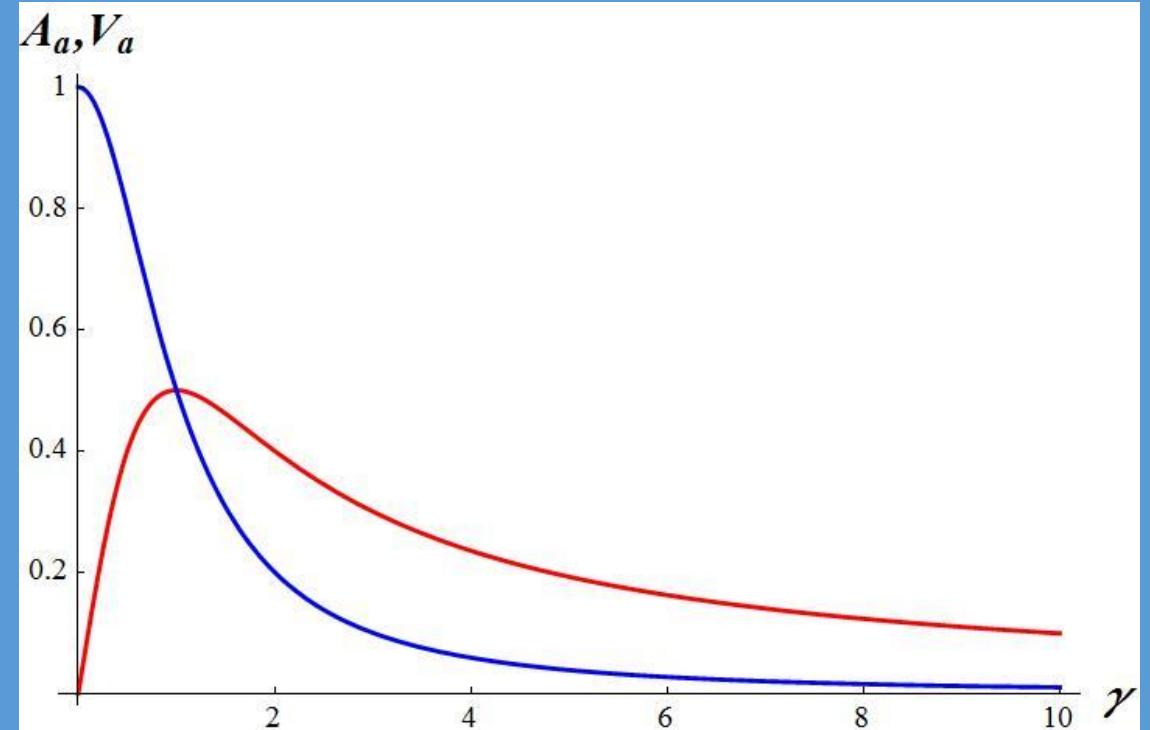
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A and V vectors



Conclusions and further work

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- For $n \neq 0$ the spinorail source has no effect on the metric.
- There is a non-trivial spinor configuration
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- Study the $n = 0$ case
- Reconsider Grassmann Dirac Spinor
- Include rotation &/or coupling with Maxwell field

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The background image shows a wide-angle view of a tropical coastline. On the left, a sandy beach runs along the shore, dotted with small buildings and trees. The water is a vibrant turquoise color, transitioning to a darker blue further out. In the distance, a range of green mountains rises against a sky filled with scattered white clouds. A small, isolated island is visible on the right side of the frame.

Thank You