Worldsheet dilatation operator for the *AdS* superstring

Israel Ramírez

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- ▶ How to make computation in the worldsheet

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- OPEs

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▶ A method to extract the divergent part of an operator

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we find

$$\left\langle \frac{\delta S}{\delta \phi(z)} \mathcal{O}(y) \right\rangle = \left\langle \frac{\delta \mathcal{O}(y)}{\delta \phi(z)} \right\rangle.$$

A simple example: the exact propagator

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A simple example: the exact propagator For a free field $S[\phi] = \frac{1}{2} \int d^d x \partial^\mu \phi(x) \partial_\mu \phi(x)$, using $\mathcal{O} = \phi$, $-\Box_z \left\langle \phi(z) \phi(y) \right\rangle = \delta^d(z-y)$

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Here I will show how to compute amplitudes for a simple sigma model.

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The algebra for this model is just

$$[T_a, T_b] = f_{ab}^c T_c$$

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Step 3: Apply fluctuations on a classical background $g = g_0 e^X$

$$J = J_0 + \nabla X + \frac{1}{2} [\nabla X, X] + \cdots$$



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$$\left\langle \frac{\delta S_{kin}}{\delta X^{a}(z)} X^{b}(y) \right\rangle + \int d^{2}w \frac{\delta^{2} S_{int}}{\delta X^{c}(w) \delta X^{a}(z)} \left\langle X^{c}(w) X^{b}(y) \right\rangle$$
$$= \delta^{b}_{a} \delta^{2}(z - y)$$

$$\Box_{z}G^{ab}(z,y) = -\frac{1}{T}\eta^{ab}\delta^{2}(z-y) + \frac{f_{ce}^{a}}{2}\left(\partial_{z}G^{cb}(z,y)\bar{J}_{0}^{e} + \bar{\partial}_{z}G^{cb}(z,y)J_{0}^{e}\right)$$

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Working a little bit more

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$$G^{ab}(z,k) = \frac{\eta^{ab}}{T} \frac{1}{|k|^2} + \frac{\Box}{|k|^2} G^{ab} + i \left(\frac{\partial}{k} + \frac{\bar{\partial}}{\bar{k}}\right) G^{ab}$$
$$- \frac{f_{ce}^a}{2} \left(\bar{J}_0^e \left[\frac{i}{\bar{k}} + \frac{\partial}{|k|^2}\right] + J_0^e \left[\frac{i}{k} + \frac{\bar{\partial}}{|k|^2}\right]\right) G^{cb}$$

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: Solve G(z, k) iteratively in powers of 1/k

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$$G^{ab}(z,k) = \frac{\eta^{ab}}{T} \frac{1}{|k|^2} - \frac{\eta^{bd}}{T} \frac{i}{2|k|^2} f_{de}^a \left(\frac{J_0^e}{k} + \frac{\overline{J}_0^e}{\overline{k}} \right) + \cdots$$

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Step 8: Go back to coordinates space using Fourier transform, deal with divergences using dimensional regularization

$$\lim_{z \to y} \int \frac{d^d k}{(2\pi)^d} \frac{e^{ik(z-y)}}{|k|^2} = -\frac{1}{2\pi\epsilon}$$

where $d = 2 - \epsilon$.

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$$\langle XX \rangle \equiv \lim_{y \to z} \langle X(z)X(y) \rangle$$

$$= \lim_{z \to y} \int \frac{d^d k}{(2\pi)^d} e^{ik(z-y)} \frac{A}{|k|^2} + \dots = IA$$

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For a term with a derivative

$$\begin{split} \left\langle X \bar{\partial} X \right\rangle &\equiv \lim_{y \to z} \bar{\partial}_y \left\langle X(z) X(y) \right\rangle \\ &= \lim_{z \to y} \int \frac{d^d k}{(2\pi)^d} e^{ik(z-y)} \frac{-iC}{|k|^2} + \dots = -ilC \end{split}$$

For our simple case

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$$\begin{split} \left\langle X^a X^b \right\rangle = & I \eta^{ab} \\ \left\langle X^a \partial X^b \right\rangle = & -\frac{1}{2} \eta^{bd} f^a_{de} J^e_0 \\ \left\langle X^a \bar{\partial} X^b \right\rangle = & -\frac{1}{2} \eta^{bd} f^a_{de} \bar{J}^e_0 \end{split}$$

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Step 10: For a multiplication of general operators \mathcal{O} , \mathcal{O}' use

$$\left\langle \mathcal{O}\mathcal{O}'\right\rangle = \left\langle \mathcal{O}\right\rangle \mathcal{O}' + \mathcal{O}\left\langle \mathcal{O}'\right\rangle + \left\langle \mathcal{O}, \mathcal{O}'\right\rangle$$

where

$$\langle \mathcal{O} \rangle = \frac{1}{2} \int d^2 z d^2 y \frac{\delta^2 \mathcal{O}}{\delta X^a(z) \delta X^b(y)} \left\langle X^a(x) X^b(y) \right\rangle$$
$$\langle \mathcal{O}, \mathcal{O}' \rangle = \int d^2 z d^2 y \frac{\delta \mathcal{O}}{\delta X^a(z)} \frac{\delta \mathcal{O}'}{\delta X^b(y)} \left\langle X^a(x) X^b(y) \right\rangle$$

What is a Pure Spinor?

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Pure Spinor model \rightarrow a Super String model on a coset space.

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$$\lambda^{A}\gamma^{\mu}_{AB}\lambda^{B}=0$$

Pure Spinor model \rightarrow a Super String model on a coset space. Equivalent to Green-Schwartz model.

The Pure Spinors action is

$$S = \frac{R^2}{2\pi} \int d^2z \operatorname{STr} \left[\frac{1}{2} J_2 \bar{J}_2 + \frac{1}{4} J_1 \bar{J}_3 + \frac{3}{4} \bar{J}_1 J_3 + \omega \bar{\nabla} \lambda + \hat{\omega} \nabla \hat{\lambda} - N \hat{N} \right]$$

$$N = \{\omega, \lambda\}$$

$$\nabla = \partial + [J_0,]$$

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$$\begin{split} S = & \frac{R^2}{2\pi} \int d^2z \text{STr} \left[\frac{1}{2} J_2 \bar{J}_2 + \frac{1}{4} J_1 \bar{J}_3 + \frac{3}{4} \bar{J}_1 J_3 + \omega \bar{\nabla} \lambda + \hat{\omega} \nabla \hat{\lambda} - N \hat{N} \right] \\ N = & \{\omega, \lambda\} \\ \nabla = & \partial + [J_0,] \end{split}$$

The $AdS_5 \times S^5$ is described by the group

$$\frac{PSU(2,2|4)}{SO(1,4)\times SO(5)}$$

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$$J = J_0 + J_1 + J_2 + J_3 = J_0 + K$$

= $J_0^i T_i + J_1^{\alpha} T_{\alpha} + J_2^m T_m + J_3^{\hat{\alpha}} T_{\hat{\alpha}}$

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And since we have

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we can choose

$$X = X_1 + X_2 + X_3$$



How does everything look like now?

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$$G^{AB} = I^{AB} + F_C^A G^{CB}$$

How does everything look like now? Matrix multiplication!

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Some results

$$\langle K \rangle = 0$$

$$\langle j \rangle = \left\langle g \left(J_2 + \frac{3}{2} J_3 + \frac{1}{2} J_1 + 2N \right) g^{-1} \right\rangle = 0$$

$$\langle b \rangle = \left\langle (\lambda \hat{\lambda})^{-1} \mathsf{STr} \left(\hat{\lambda} [J_2, J_3] + \{\omega, \hat{\lambda}\} [\lambda, J_1] \right) - \mathsf{STr} \left(\omega J_1 \right) \right\rangle = 0$$

Dilatation operator

$$\mathcal{O} = \langle J, J \rangle \frac{\delta}{\delta J} \frac{\delta}{\delta J} + \langle J, X \rangle \frac{\delta}{\delta J} \frac{\delta}{\delta X} + \langle X, X \rangle \frac{\delta}{\delta X} \frac{\delta}{\delta X} + \langle J \rangle \frac{\delta}{\delta J}$$

Conclusions

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- ▶ It is fast and can be done by a computer
- ► Can be extended to a 3-point function straightforward

Future Work

- Compute 3 and 4-point functions
- Do it with a computer program
- Look for signs of integrability in the string
- Find physical states of the string

Obrigado