

Homework Problems  
for the String Theory Lectures

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# 1 Problems for Lecture 1

## Problem 1

Consider a  $p$ -brane action of the form

$$S_\sigma = -\frac{T_p}{2} \int d^{p+1}\sigma \sqrt{-h} h^{\alpha\beta} G_{\alpha\beta}(X) + \Lambda_p \int d^{p+1}\sigma \sqrt{-h}, \quad (1)$$

where the induced metric is given by

$$G_{\alpha\beta} = g_{\mu\nu}(X) \partial_\alpha X^\mu \partial_\beta X^\nu \quad \alpha, \beta = 0, \dots, p \quad (2)$$

where  $\partial_\alpha$  denotes the partial derivative with respect to  $\sigma^\alpha$ , and  $g_{\mu\nu}$  is the spacetime metric tensor in  $D$  dimensions. Also,  $h_{\alpha\beta}(\sigma)$  denotes a world-volume metric tensor in  $p+1$  dimensions (with Lorentzian signature),  $h^{\alpha\beta}$  denotes its inverse, and  $h$  denotes its determinant.

- Derive the equation of motion obtained by varying  $h^{\alpha\beta}$ .
- Show that this equation is solved by  $h_{\alpha\beta} = G_{\alpha\beta}$  provided that  $\Lambda_p = c_p T_p$  for a certain choice of  $c_p$ . Verify that for this value of  $c_1$  the  $p=1$  (or string) action is invariant under the Weyl transformation  $h_{\alpha\beta}(\sigma) \rightarrow \exp(\phi(\sigma)) h_{\alpha\beta}(\sigma)$ .
- Use this solution to establish classical equivalence to the “volume action”

$$S_p = -T_p \int \sqrt{-\det G_{\alpha\beta}} d^{p+1}\sigma, \quad (3)$$

## Problem 2

Consider an on-shell open-string state of the bosonic string theory in  $D$ -dimensional Minkowski spacetime of the form

$$|\phi\rangle = (A\alpha_{-1} \cdot \alpha_{-1} + B\alpha_0 \cdot \alpha_{-2} + C(\alpha_0 \cdot \alpha_{-1})^2) |0; k\rangle,$$

where  $A$ ,  $B$  and  $C$  are constants.

- Determine the eigenvalue of  $\alpha_0 \cdot \alpha_0$  and relations among the coefficients  $A$ ,  $B$  and  $C$  so that  $|\phi\rangle$  satisfies the physical-state conditions  $(L_0 - 1)|\phi\rangle = L_1|\phi\rangle = L_2|\phi\rangle = 0$  for arbitrary  $D$ . Recall that  $[\alpha_m^\mu, \alpha_n^\nu] = m\eta^{\mu\nu} \delta_{m+n,0}$  and

$$L_0 = \frac{1}{2}\alpha_0^2 + \alpha_{-1} \cdot \alpha_1 + \alpha_{-2} \cdot \alpha_2 + \dots, \quad L_1 = \alpha_0 \cdot \alpha_1 + \alpha_{-1} \cdot \alpha_2 + \dots, \quad L_2 = \frac{1}{2}\alpha_1 \cdot \alpha_1 + \alpha_0 \cdot \alpha_2 + \dots$$

- Compute the norm of  $|\phi\rangle$ . What conclusion can you draw from the result?

## 2 Problems for Lecture 2

### Problem 3

Show that there are the same number of physical degrees of freedom in the NS and R sectors at the first massive level of the ten-dimensional open superstring after GSO projection.

Suggestion: The counting is most easily carried out in light-cone gauge. First, show that at the first massive level, the NS states have  $N = 3/2$  and the R states have  $N = 1$ .

### Problem 4

The conclusion that T-duality interchanges

$$X(\tau, \sigma) = X_L(\tau + \sigma) + X_R(\tau - \sigma)$$

and

$$\tilde{X}(\tau, \sigma) = X_L(\tau + \sigma) - X_R(\tau - \sigma)$$

can be understood from a world-sheet viewpoint. Consider the following world-sheet action:

$$\int \left( \frac{1}{2} V^\alpha V_\alpha - \epsilon^{\alpha\beta} X \partial_\beta V_\alpha \right) d^2\sigma, \quad (4)$$

where an overall constant coefficient is omitted, because the considerations that follow are classical.

a) Show that varying  $X$ , which acts as a Lagrange multiplier, gives an equation of motion that is solved by  $V_\alpha = \partial_\alpha \tilde{X}$ . This can then be used to obtain an action that only depends on  $\tilde{X}$ . Alternatively, varying  $V_\alpha$  in the original action gives an equation of motion that can be used to obtain an action that only depends on  $X$ . Show that the resulting actions for  $X$  and  $\tilde{X}$  have the same form. Thus, the (nonlocal) transformation from  $X$  to  $\tilde{X}$  is a symmetry.

b) Use the equations you found in part a) to deduce a relationship between the derivatives of  $\tilde{X}$  and  $X$ . Verify that this is exactly the relationship expected for T-duality.

### 3 Problems for Lecture 3

#### Problem 5

Let  $A$  be a positive definite  $m \times m$  symmetric matrix and define

$$f(A) = \sum_{\{M\}} \exp(-\pi M^T A M). \quad (5)$$

Here  $M$  represents a vector made of  $m$  integers  $M_1, M_2, \dots, M_m$  each of which is summed from  $-\infty$  to  $+\infty$ . Derive the Poisson resummation formula:

$$f(A) = \frac{1}{\sqrt{\det A}} f(A^{-1}). \quad (6)$$

Suggestion: it is helpful to add dependence on  $m$  variables  $x^i$  and define

$$f(A, x) = \sum_{\{M\}} \exp(-\pi(M+x)^T A(M+x)). \quad (7)$$

This function is periodic, with period 1, in each of the  $x^i$ . Therefore, it must have a Fourier series expansion of the form

$$f(A, x) = \sum_{\{N\}} C_N(A) \exp(2\pi i N^T x). \quad (8)$$

Evaluate the Fourier coefficients  $C_N(A)$  explicitly. (Hint: the infinite sum of integrals over an interval of length one can be converted to a single integral over an infinite interval.) Finally, the desired formula is obtained by setting  $x = 0$ .

The Poisson resummation formula is crucial for establishing the modular invariance of one-loop closed-string amplitudes. Such amplitudes are given as an integral over conformal equivalence classes of a torus, which is the one-loop world sheet. These classes are parametrized by a modular parameter  $\tau$ . The integrand (including the measure) should be invariant under  $SL(2, \mathbb{Z})$  transformations of the form  $\tau \rightarrow (a\tau + b)/(c\tau + d)$  in order that the integral not depend on the choice of a fundamental domain. The Poisson resummation formula is required for proving the symmetry for the S transformation  $\tau \rightarrow -1/\tau$ . The entire group is generated by the S transformation and the T transformation:  $\tau \rightarrow \tau + 1$ , which is usually trivial.

## Problem 6

The type IIB theory has an infinite spectrum of half-BPS  $(p, q)$  strings in ten dimensions, which can be regarded as bound states of  $p$  fundamental strings and  $q$  D-strings.  $p$  and  $q$  are integer charges that measure the strengths with which the strings couple to the two-form gauge fields  $B_2$  and  $C_2$ . These strings are stable whenever  $p$  and  $q$  are coprime. Their tensions (in the Einstein frame) are given by

$$T_{p,q} = |p + q\tau|T_F, \quad (9)$$

where  $T_F$ , the tension of the fundamental  $(1, 0)$  string, is

$$T_F = \frac{1}{2\pi l_S^2} = \frac{1}{2\pi\sqrt{\tau_2}l_B^2}. \quad (10)$$

Here  $l_S$  denotes the fundamental string length scale and  $l_B$  denotes the ten-dimensional Planck length in the type IIB superstring theory.  $\tau$  is the vacuum value of the complex scalar field  $C_0 + i \exp(-\Phi)$  in the type IIB theory, and  $\tau_2$  denotes the imaginary part of  $\tau$ .

Verify that the  $\tau$  dependence of these equations is consistent with the fact that the string charges  $p$  and  $q$  transform linearly under  $SL(2, \mathbb{Z})$  transformations  $\tau \rightarrow (a\tau + b)/(c\tau + d)$ . This is required because the 2-form gauge fields also transform as a doublet. Remark: It is important that the Planck length  $l_B$ , and not the string length scale  $l_S$ , is  $SL(2, \mathbb{Z})$  invariant.