MHD turbulence in the solar corona and solar wind

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- Relation to magnetic reconnection.
- Clues for dissipation mechanisms.
- Differential energization: $T^e_{\parallel} \gg T^e_{\perp}, \quad T^i_{\perp} \gg T^i_{\parallel}$
- **Direct approach**: Test particles trajectories are followed in the turbulent fields obtained from a direct numerical solution of the MHD equations.

Dmitruk, Matthaeus, Seenu, ApJ 617, 667 (2004) Dmitruk, Matthaeus, Seenu, Brown, ApJ 597, L81 (2003) Ambrosiano, Matthaeus, Goldstein, Plante, JGR 93, 14383 (1988)

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Disadvantages

- Extremely computationally demanding if we want to fully resolve from turbulent (MHD) to particle scales
- Lack of self-consistency in the MHD-kinetic physics interplay at the "dissipative scales"

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Run the code for two $t_0 = L/v_0$ (eddy turnover times), fully turbulent state developed, take a **snapshot** for pushing the test particles.



MHD spatial structure: parallel current density J_z



Cross-sections: magnetic field over current density



xz plane

xy plane

Particles

Equations of motion for charged particles:

$$\frac{d\mathbf{u}}{dt} = \frac{q}{m}(\frac{1}{c}\mathbf{u} \times \mathbf{B} + \mathbf{E}), \qquad \frac{d\mathbf{x}}{dt} = \mathbf{u}$$

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with electric field

$$\mathbf{E} = -\frac{1}{c}\mathbf{v} \times \mathbf{B} + \frac{v_0 L}{R_m c} \mathbf{J} = \text{inductive} + \text{resistive}$$

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Particle gyrofrequency $w_g = qB/(mc) \rightarrow \text{(small)}$ lengthscale $r_g = u_\perp/w_g = u_\perp mc/(qB)$. If $u_\perp = v_0, B = B_0 \rightarrow \text{gyroradius} r_0 = v_0 mc/(qB_0)$.

Dissipation length $l_d \approx \rho_{ii}$ (ion skin depth)

Supported by solar wind observations, linear Vlasov theory, kinetic physics reconnection studies.

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Particle (nominal) gyroradius vs turbulent dissipative scale

$$\frac{r_0}{l_d} = Z \frac{m}{m_p} \frac{\delta B}{B_0}$$

electrons $\rightarrow r_0^e/l_d = 6 \times 10^{-5}$

protons $\rightarrow r_0^p/l_d = 0.1$

Electrons (rms displacement)



Run stops when $<\Delta z^2 >^{1/2} \approx L$

Electrons (rms velocity)





Electrons (velocity distribution)



In $t = 10^4 \tau_e = 0.1 t_0$ with $\tau_e = 2\pi m_e c / (eB_0)$

Electrons (velocity scatter plot)



Electrons (gyroradii)



 $r_g^e = u_\perp m_e c / (eB_0)$

Electrons (trajectories)



$$\frac{d\mathbf{u}_{\parallel}}{dt} = -\frac{e}{m_e} \mathbf{E}_{\parallel} = -\frac{e}{m_e c} \frac{v_0 L}{R_m} \mathbf{J}_{\parallel}, \qquad \qquad \frac{dz}{dt} = u_{\parallel}$$

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which can be rewritten using $J_0 = \delta B/L$ and ρ_{ii} as

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Using $\bar{J}_{\parallel}/J_0 = 4.5$ we get $t_{\parallel} \approx 0.09t_0$ and $\Delta u_{\parallel} \approx 20v_0$ for the average velocity



Protons (velocity distribution function)



At $t = 90\tau_p = 1.8t_0$ with $\tau_p = 2\pi m_p c/(eB_0)$

Protons (velocity scatter plot)



Protons: gyroradii at $t = 2\tau_p = 0.04t_0$



Protons: gyroradii at $t = 20\tau_p = 0.4t_0$



Protons: gyroradii at $t = 90\tau_p = 1.8t_0$



Protons: trajectories



Protons: trajectories top view



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Using the numerical values $u_{\perp} \approx 60v_0$

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- Another issue: influence of the Hall effect \rightarrow , negligible for electrons, it affects more (but not much) the behavior of ions (see Dmitruk & Matthaeus, Jour. Geophys. Res. 111, A12110, 2006)

Hall MHD turbulence

Electric field

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Spectrum of the electric field



Distribution functions



Effect on reconnection \rightarrow changes the reconnection rate



Dmitruk and Matthaeus, Phys. Plasmas 2006; also Reduced Hall MHD in Gomez, Dmitruk, Mahajan, Phys. Plasmas 2009; Martin, Dmitruk, Gomez, Phys. Plasmas 2010



Electrons (momentum distribution function)



At $t = 10^4 \tau_e = 0.1 t_0$, for Hall and non-Hall (light lines)

Protons (momentum distribution function)



At $t = 90\tau_p = 1.8t_0$, for Hall and non-Hall (light lines)

Results with time dependent fields

Electrons (velocity distribution function)



At $t = 10^4 \tau_e = 0.1 t_0$



At $t = 200\tau_p = 4t_0$