# Expansion Method, its general properties and some applications

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arXiv:1308.4832: colaboration with L. Andrianopoli, M. Trigiante (Politecnico di Torino) and F. Nadal (Universidad de Valencia)



#### **Contents**

- I) Introduction and motivations
- II) S-expansion method
- **III) Some applications**
- IV) Results and conclusions





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- 2) A G invariant polynomial of order n is a multilineal map

$$\langle ... \rangle : \mathcal{G}^n \to \mathbb{C}$$

which satisfy the invariant condition

$$\forall \mathbf{T}_A \in \mathcal{G}, g \in G, \langle \mathbf{T}_1, ..., \mathbf{T}_n \rangle = \langle g \mathbf{T}_1 g^{-1}, ..., g \mathbf{T}_n g^{-1} \rangle$$





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3) A gauge connection one form  $\mathbf{A} = A^A \mathbf{T}_A$  and its associated strength field  $\mathbf{F} = d\mathbf{A} + \mathbf{A} \wedge \mathbf{A}$ 





#### Examples:

Yang-Mills Theories:

$$\mathcal{L}_{YM}^{(d)}(\mathbf{A}) = \langle \mathbf{F} \wedge * \mathbf{F} \rangle,$$
  
=  $F^A \wedge * F^B \langle \mathbf{T}_A, \mathbf{T}_B \rangle$ 

or a Chern-Simons Theories in d=3:

$$\mathcal{L}_{CS}^{(3)}(\mathbf{A}) = \left\langle \mathbf{A} \wedge \left( d\mathbf{A} + \frac{2}{3} \mathbf{A}^{\wedge 2} \right) \right\rangle$$
$$= A^{A} \wedge \left( dA^{B} + \frac{1}{3} C_{MN}^{B} A^{M} \wedge A^{N} \right) \left\langle \mathbf{T}_{A}, \mathbf{T}_{B} \right\rangle$$





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Non trivial relations between **groups** and **Lie algebras** 



Non trivial relations between **physical theories** wich are invariants under that symmetries





Non trivial relations between **groups** and **Lie algebras** 

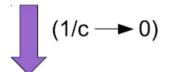


Non trivial relations between **physical theories** wich are invariants under that symmetries

First example: I.E. Segal (1951)

#### **Physics Theories**

Special Theory of Relativity



Newtonian Mechanics

#### **Symmetry Group**

Poincaré's Group



Galileo's group





During the second half of the 20th century appeared in the literature:

Mechanism	Example	
1. Contraction	(Anti)de-Sitter →	Poincaré
<ol><li>Deformation</li></ol>	Galileo	Poincaré
3. Extension	Poincaré	Super-Poincaré





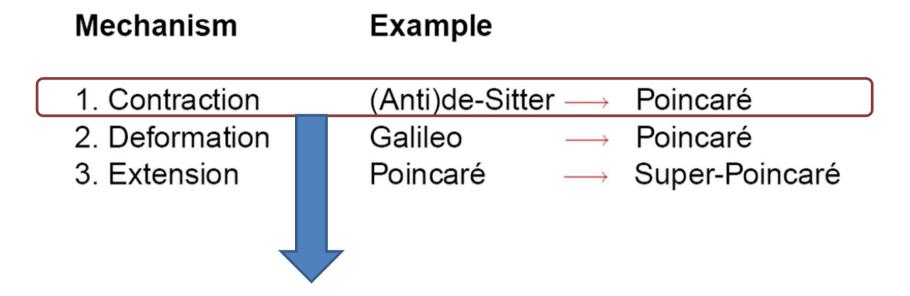
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During the second half of the 20th century appeared in the literature:



We focus on its generalizations (particularly one called **S-expansion**)









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Ex: *ISO*(3,1) from *SO*(3,2)

$$[M_{ab}, M_{cd}] = g_{ac}M_{bd} - g_{bc}M_{ad} - g_{ad}M_{bc} + g_{bd}M_{ac}$$
$$a, c = 1, ...5$$

Defining  $M_{5\mu} = RP_{\mu}$  with  $\mu, \nu = 1, ..., 4$  it leads to

$$\begin{split} [P_{\mu},P_{\nu}] &= \frac{1}{R^2} \left[ M_{5\mu}, M_{5\nu} \right] = \frac{1}{R^2} g_{55} M_{\mu\nu} \\ [M_{\mu\nu},P_{\rho}] &= \frac{1}{R} \left[ M_{\mu\nu}, M_{5\rho} \right] = \frac{1}{R} \left( -g_{\mu\rho} M_{\nu5} + g_{\nu\rho} M_{\mu5} \right) \\ &= \frac{1}{R} \left( -g_{\mu\rho} R P_{\nu} + g_{\nu\rho} R P_{\mu} \right) = -g_{\mu\rho} P_{\nu} + g_{\nu\rho} P_{\mu} \end{split}$$





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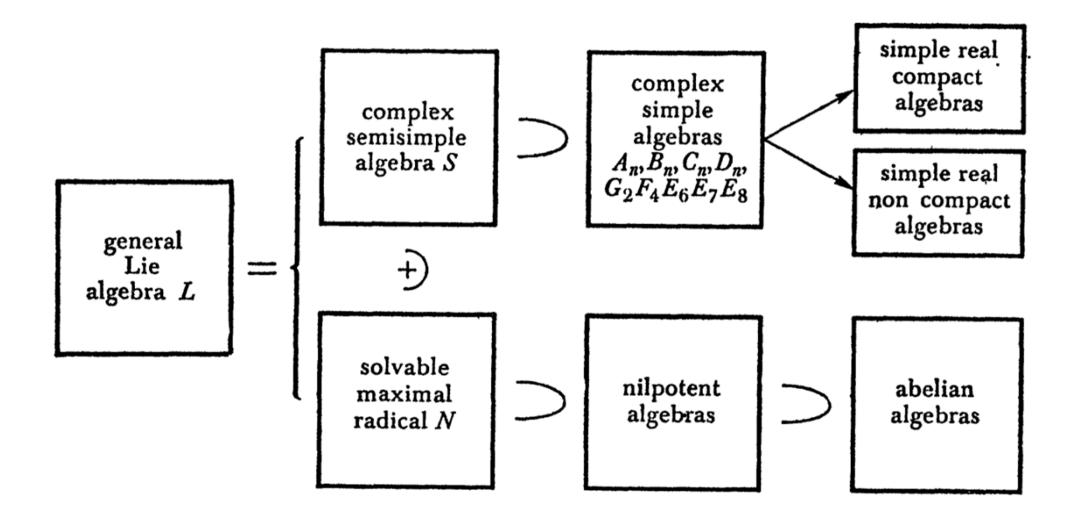
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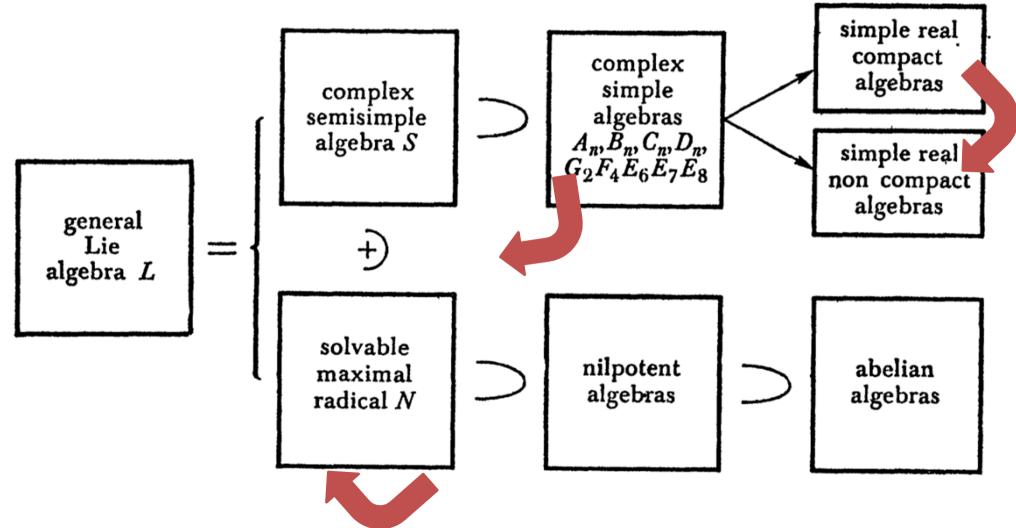
# Classification of Lie algebras







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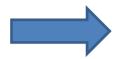


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This motivates the following question:

Given two algebras

A and B

Can these algebras be related by some contraction or expansion method?





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Here we review the main characteristics of the method. For further details, see the original article:

[1] F. Izaurieta, E. Rodríguez, P. Salgado, J.Math.Phys.47:123512,2006 arXiv: 0606215







#### Ingredients:

- 1) A Lie algebra  $\mathcal{G}$  with basis  $\{\mathbf{T}_A\}_{A=1}^{\dim \mathcal{G}}$
- 2) A finite abelian semigroup  $S = \{\lambda_{\alpha}\}_{\alpha=1}^{n}$  (there is a compostion law  $\lambda_{\alpha}\lambda_{\beta}$  which is closed and associative)





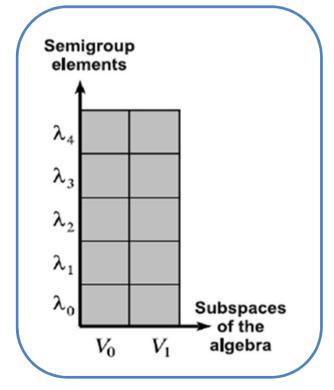
#### Steps:

I) Construction of  $\mathcal{G}_S = S \otimes \mathcal{G}$  with basis  $\{\mathbf{T}_{(A,\alpha)} = \lambda_{\alpha} \otimes \mathbf{T}_A\}$  and define the induced Lie product:

$$\left[\mathbf{T}_{(A,\alpha)},\mathbf{T}_{(B,\beta)}\right] = \lambda_{\alpha}\lambda_{\beta} \otimes \left[\mathbf{T}_{A},\mathbf{T}_{B}\right]$$

First result:  $\mathcal{G}_S = S \otimes \mathcal{G}$  is a Lie algebra, called the **expanded algebra** 





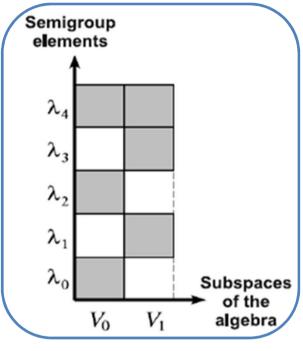






II) Supose that  $\mathcal{G} = \bigoplus_{p \in I} V_p$ ,  $S = \bigcup_{p \in I} S_p$  satisfying respectively:

$$[V_p,V_q] \subset \bigoplus_{r \in i(p,q)} V_r \text{ and } S_p \times S_q \subset \bigcup_{r \in i(p,q)} V_r$$



Perform the following construction  $\mathcal{G}_{S,R} = \bigoplus_{p \in I} S_p \otimes V_p$ 

Second result:  $\mathcal{G}_{S,R}$  is a subalgebra called the **resonant subalgebra** 

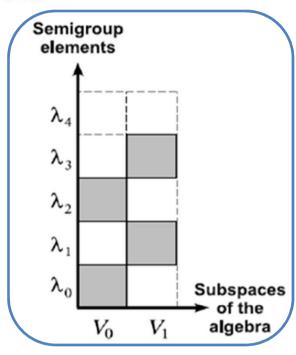






III) Supose that there is an element  $0_S$  wich satisfies  $0_S \lambda_{\alpha} = 0_S \ \forall \lambda_{\alpha} \in S$ 

Third result: the sector  $0_S \otimes \mathcal{G}$  can be removed from the expanded algebra to obtain the so called **reduced algebra**.



Observation: Steps II and III are independent, but can also be applied simultaneously.

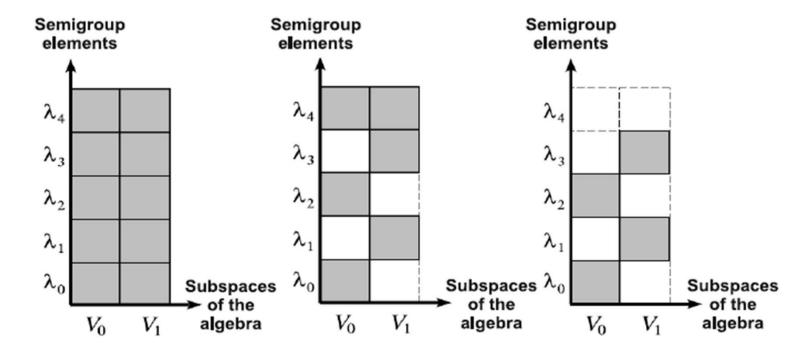




#### S-expansion

In general the S-expansion method consist in a serie of steps:

- Obtain the S-expanded algebra
- Find the (resonant)subalgebras
- Perform a reduction.









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- ⇒ and the dAIPV-expansion method







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$$S_E^{(N)} = \{\lambda_0, \lambda_1, ..., \lambda_N, \lambda_{N+1}\}$$
 with

$$\lambda_{\alpha}\lambda_{\beta} = \lambda_{\alpha+\beta}$$
, if  $\alpha + \beta < N+1$   
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Example: Inönü-Wigner contraction as an  $S_E^{(1)}$ -expansion

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Consider the semigroup 
$$S_E^{(1)} = \{\lambda_0, \lambda_1, \lambda_2\}$$
 with

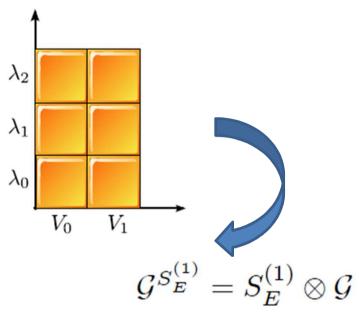
$$\lambda_{\alpha}\lambda_{\beta} = \lambda_{\alpha+\beta}$$
, if  $\alpha + \beta < 2$   
 $\lambda_{\alpha}\lambda_{\beta} = \lambda_{2}$ , if  $\alpha + \beta \ge 2$   $S_{E}^{(1)} = S_{0}$ 

$$S_0, \lambda_1, \lambda_2$$
 with  $S_0 = \{\lambda_0, \lambda_2\}$   $S_1 = \{\lambda_1, \lambda_2\}$   $S_1 = \{\lambda_1, \lambda_2\}$   $S_0 \times S_0 \subset S_0$   $S_0 \times S_1 \subset S_1$   $S_1 \times S_1 \subset S_0$ 





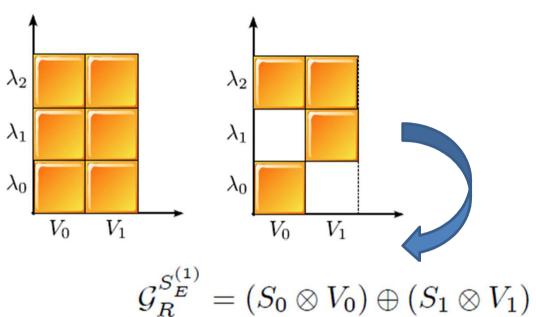








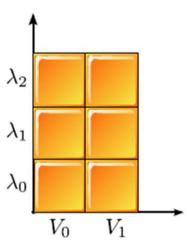


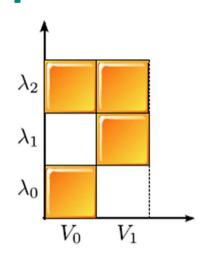


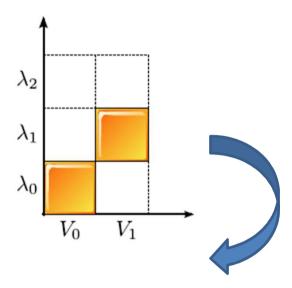








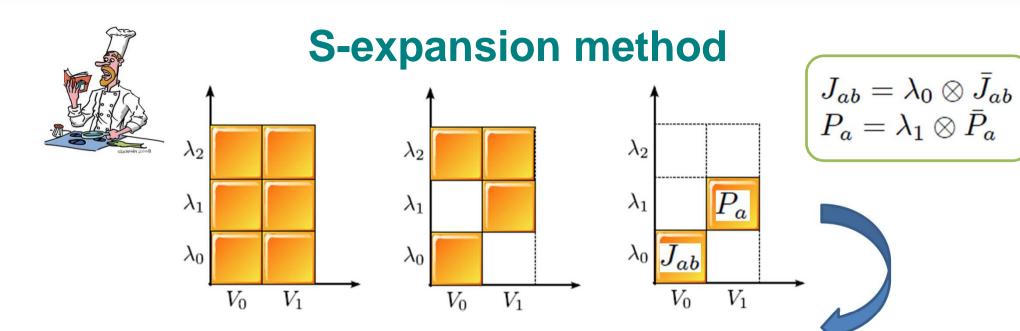




$$\mathcal{G}_{R,\mathrm{red}}^{S_E^{(1)}} = \left[ \left( S_0 \otimes V_0 \right) \oplus \left( S_1 \otimes V_1 \right) \right]_{\mathrm{red}}$$





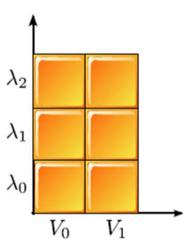


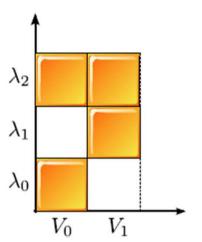
$$\mathcal{G}_{R,\mathrm{red}}^{S_E^{(1)}} = \left[ (S_0 \otimes V_0) \oplus (S_1 \otimes V_1) \right]_{\mathrm{red}} = \left\langle \left\{ \lambda_0 \bar{J}_{ab}, \lambda_1 \bar{P}_a \right\} \right\rangle = \left\langle \left\{ J_{ab}, P_a \right\} \right\rangle$$

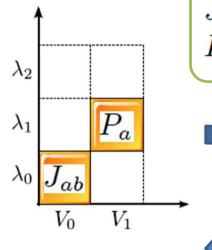




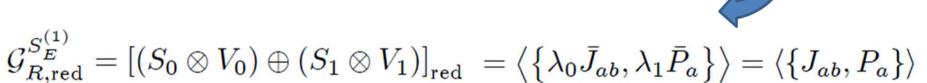








$$J_{ab} = \lambda_0 \otimes \bar{J}_{ab}$$
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ight] \sim \lambda_1ar{P}_c \sim P_c \ [P_a,P_b] &= \lambda_1\lambda_1\otimes \left[ar{P}_a,ar{P}_b
ight] \sim \lambda_2ar{J}_{ab} \sim 0 \end{aligned}$$

IW contraction.







Example: The  $\mathcal{B}$  algebra

Starting from  $\mathcal{G}=SO(4,2)=V_0\oplus V_1$ 

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Is easy to check that  $S = S_0 \cup S_1$  satisfy

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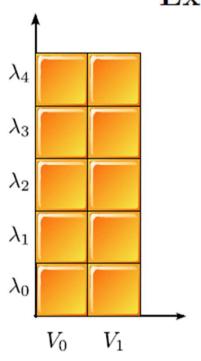
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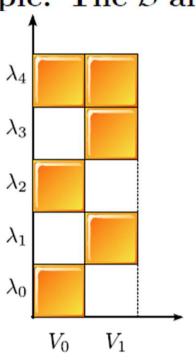


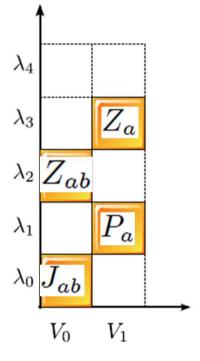




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$$egin{aligned} J_{ab} &= \lambda_0 \otimes ar{J}_{ab} \ Z_{ab} &= \lambda_2 \otimes ar{J}_{ab} \ P_a &= \lambda_1 \otimes ar{P}_a \ Z_a &= \lambda_3 \otimes ar{P}_a \end{aligned}$$



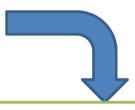




Example: The  $\mathcal{B}$  algebra

$$\mathcal{B} = \{J_{ab}, Z_{ab}, P_a, Z_a\}$$

$$\left[\mathbf{T}_{(A,\alpha)},\mathbf{T}_{(B,\beta)}\right]=\lambda_{\alpha}\lambda_{\beta}\otimes\left[\mathbf{T}_{A},\mathbf{T}_{B}\right]$$



$$[P_a, P_b] = Z_{ab} , \qquad [J_{ab}, P_c] = P_a \eta_{bc} - P_b \eta_{ac} ,$$

$$[J_{ab}, J_{cd}] = -J_{ac} \eta_{bd} + J_{bc} \eta_{ad} - J_{bd} \eta_{ac} + J_{ad} \eta_{bc} ,$$

$$[Z_a, Z_b] = [Z_{ab}, Z_c] = [Z_{ab}, Z_{cd}] = [P_a, Z_b] = 0 ,$$

$$[Z_{ab}, P_c] = [J_{ab}, Z_c] = Z_a \eta_{bc} - Z_b \eta_{ac} ,$$

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- 4) CS for AdS in D=3 is equivalent to GR
- 5) CS for Lorentz in D=3 defines exotic gravity.





- 1) Exist in odd dimensions
- 2) Its fundamental field is a G-valued gauge connection A with associate curvature F
- 3) Given an **invariant tensor** *g*, the dynamics of A is governed by the CS action whose Lagrangian is constructed from:

$$\langle F \wedge \cdots \wedge F \rangle_g = dL_{\rm CS}(A)$$

- 4) CS for AdS in D=3 is equivalent to GR
- 5) CS for Lorentz in D=3 defines exotic gravity.

For further details about CS therories see [2].

[2] J. Zanelli, *arXiv: 0502193* 





#### 1) CS action for the M algebra:

In ref. [3] the M algebra is obtained from the Osp(32/1) algebra using the dAIPV-expansion method.

- [3] J. A. de Azcárraga, J. M. Izquierdo, M. Picón, O. Varela, Nucl. Phys. B **662** (2003) 185 *arXiv:* 0212347
- [4] F. Izaurieta, E. Rodríguez, P. Salgado, Eur. Phys. J. C **54**, 675-684, 2008 *arXiv:* 0606225





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A CS action for the M algebra is constructed.

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- the standard procedure to obtain an invariant tensor of range r is to use the symmetrized (super)trace for the product of r generators in some matrix representation of the algebra
- there are limitations in the case of non-semisimple Lie algebras.
- on [arXiv: 0606225] it is shown that the supertrace is not a good choice





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- then it is not possible to reproduce general relativity neither to include fermionic fields or fields associated with central charges.
- the S-expansion procedure provide an invariant tensor, different from the supertrace, which leads to a theory with a richer structure.
- the S-expansion method does not solve the problem of classifying all invariant tensors for non-semisimple algebras, it at least gives invariant tensors different from the supertrace that are useful for the construction of the mentioned gauge theories of gravity.





#### 2) Dual formulation:

A dual formulation of S-expansion method is given in [5]. permits to perform the expansion at the level of the Lagrangian.

[5] F. Izaurieta, A. Pérez, E. Rodríguez, P. Salgado, Jour.Math. Phys. 50 7 073511 (2009) arXiv: 0903.4712





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As an example, this permit to obtain 3D CS AdS gravity from exotic gravity (because SO(2,2) can be regarded as an expansion of SO(2,1))

[5] F. Izaurieta, A. Pérez, E. Rodríguez, P. Salgado, Jour.Math. Phys. 50 7 073511 (2009) arXiv: 0903.4712





#### 3) Standard General Relativity from CS theories in d = 5

In ref. [6] it was observed that

$$L_{\text{AdS}}^{(5)} = \kappa \varepsilon_{abcde} \left( \frac{1}{5\ell^5} e^a e^b e^c e^d e^e + \frac{2}{3\ell^3} R^{ab} e^c e^d e^e + \frac{1}{\ell^2} R^{ab} R^{cd} e^e \right).$$

does not leave E-H term alone in any limit of the ℓ parameter!

[6] J. D. Edelstein, M. Hassaine, R. Troncoso, J. Zanelli, Phys. Lett. B **640** (2006) 278. *arXiv:* 0605174





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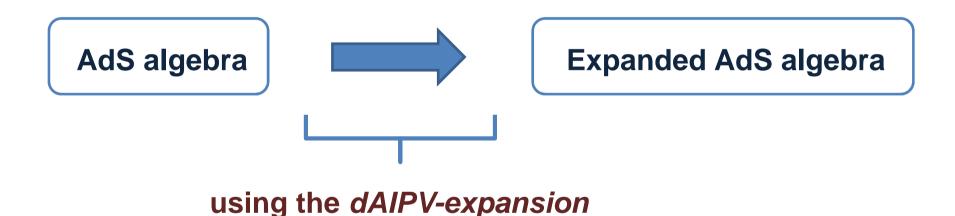
It was proposed to construct a CS action for an expansion of AdS symmetry in D=5 and to study if there is a connection with GR.

[6] J. D. Edelstein, M. Hassaine, R. Troncoso, J. Zanelli, Phys. Lett. B 640 (2006) 278. arXiv: 0605174





3) Standard General Relativity from CS theories in d = 5







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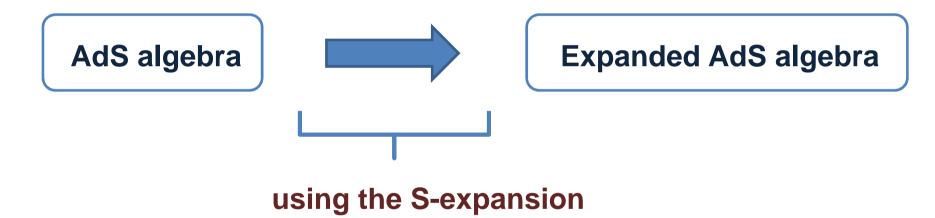
Therefore using the *dAIPV-expansion* method does not permit to connect a CS theory with 5D GR.





### 3) Standard General Relativity from CS theories in d = 5

The same idea was followed in ref. [7], but using S-expansion



[7] F. Izaurieta, P. Minning, A. Pérez, E. Rodríguez and P. Salgado, Phys. Lett. B **678** (2009) 213 *arXiv:* 0905.2187





### 3) Standard General Relativity from CS theories in d = 5

#### Result:

A CS action constructed in terms of the connection

$$\mathbf{A} \to \widetilde{\mathbf{A}} = e^a P_a + \frac{1}{2} \omega^{ab} J_{ab} + h^a Z_a + \frac{1}{2} \kappa^{ab} Z_{ab} .$$

valuated in the algebra  $\mathcal{B}$  introduced before is given by

$$L_{\text{CS}}^{(5)} = \alpha_1 \ell^2 \varepsilon_{abcde} R^{ab} R^{cd} e^e + \alpha_3 \varepsilon_{abcde} \left( \frac{2}{3} R^{ab} e^c e^d e^e + 2\ell^2 k^{ab} R^{cd} T^e + \ell^2 R^{ab} R^{cd} h^e \right)$$





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### 4) Cosmological and Black hole solutions

In refs. [8,9] there were found a *black hole* and a *cosmological* solution for the Chern-Simons theory constructed in terms of the *B* algebra.

- [8] C. A. C. Quinzacara and P. Salgado, Phys. Rev. D **85**, 124026 (2012)
- [9] F. Gomez, P. Minning, and P. Salgado, Phys. Rev. D **84**, 063506 (2011)





### **Contents**

- I) Introduction and motivations
- II) S-expansion method
- **III) Some applications**
- III) Results and conclusions





- Many physical applications have been appeared by using contractions and expansion methods,





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$$S_E^{(N)} = \{\lambda_0, \lambda_1, ..., \lambda_N, \lambda_{N+1}\}$$





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Expansions with **other semigroups** can generate algebras that cannot be reached nor by any *contraction* neither by a *dAIPV expansion*.





An example is given in [10].

[10] R. Caroca, I. Kondrashuk, N. Merino and F. Nadal, J. Phys. A: Math. Theor. 46 (2013) 225201 (24pp), *arXiv: 1104.3541* 

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J. Phys. A: Math. Theor. 46 (2013) 225201 (24pp)

doi:10.1088/1751-8113/46/22/225201

## Bianchi spaces and their three-dimensional isometries as S-expansions of two-dimensional isometries

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<sup>&</sup>lt;sup>4</sup> Instituto de Física Corpuscular (IFIC), Edificio Institutos de Investigación., c/ Catedrático José Beltrán, 2. E-46980 Paterna, Spain

<sup>&</sup>lt;sup>5</sup> Dipartimento di Scienza Applicata e Tecnologia (DISAT), Politecnico di Torino, Corso Duca degli Abruzzi, 24, I-10129 Torino, Italy

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Principal Idea: can be related 2 and 3-dimensional isometries?

Considering the two set of algebras:  $[X_1, X_2] = 0$  and

 $[X_1, X_2] = X_1$ 

	Group	Algebra
-	type I	$[X_1,X_2]=[X_1,X_3]=[X_2,X_3]=0$
-	type II	$[X_1, X_2] = [X_1, X_3] = 0,  [X_2, X_3] = X_1$
-	type III	$[X_1, X_2] = [X_2, X_3] = 0,  [X_1, X_3] = X_1$
	type IV	$ X_1, X_2  = 0,   X_1, X_3  = X_1,   X_2, X_3  = X_1 + X_2$
-	$\mathrm{type}\;\mathrm{V}$	$[X_1,X_2]=0, [X_1,X_3]=X_1, [X_2,X_3]=X_2$
	type VI	$[X_1,X_2]=0,  [X_1,X_3]=X_1,  [X_2,X_3]=hX_2,$ where $h \neq 0,1$
	type $VII_1$	$[X_1, X_2] = 0,  [X_1, X_3] = X_2,  [X_2, X_3] = -X_1$
	type $VII_2$	$[X_1, X_2] = 0,  [X_1, X_3] = X_2,  [X_2, X_3] = -X_1 + hX_2,$ where $h \neq 0 \ (0 < h < 2).$
	type VIII	$[X_1,X_2]=X_1, [X_1,X_3]=2X_2, [X_2,X_3]=X_3$
	type IX	$[X_1,X_2]=X_3, [X_2,X_3]=X_1, [X_3,X_1]=X_2$





Therefore, to answer the general question (given in *Introduction*):

Given two algebras

A and B

Can these algebras be related by some contraction or expansion method?

we need to consider the complete family of abelian semigrops.





We need to consider the *history about finite semigroup*:

order	Q = # semigroups
1	1
2	4
3	18
4	126
5	1,160
6	15,973
7	836,021
8	1,843,120,128
9	52,989,400,714,478

[Forsythe '54]

[Motzkin, Selfridge '55]

[Plemmons '66]

[Jurgensen, Wick '76]

[Satoh, Yama, Tokizawa '94]

[Distler, Kelsey, Mitchell '09]





We have developed a procedure to answer that question in [11].

[11] L. Andrianopoli, N. Merino, F. Nadal, M. Trigiante, *General properties of the S-expansion method*, J. Phys. A: Math. Theor. 46 (2013) 365204 (33pp) arXiv:1308.4832

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#### General properties of the expansion methods of Lie algebras

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1) First by establishing general properties of the expansion method (preservation of some properties of the Lie algebra)

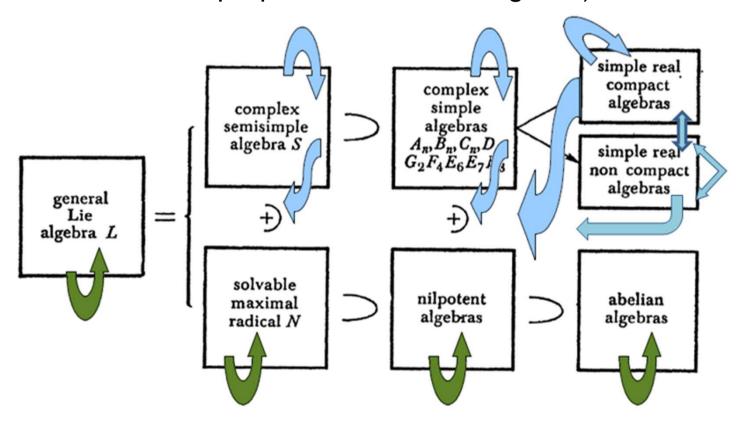
Properties preserved by the action of the S-expansion process

Original $\mathcal G$	Expanded $G_S$	Resonant $\mathcal{G}_{S,R}$	Reduced $\mathcal{G}_{S,R}^{\mathrm{red}}$
Abelian	Abelian	Abelian	Abelian
Solvable	Solvable	Solvable	Solvable
Nilpotent	Nilpotent	Nilpotent	Nilpotent
Compact	Arbitrary	Arbitrary	Arbitrary
Semisimple	Arbitrary	Arbitrary	Arbitrary
G = S	$\mathcal{G}_S = N_{\exp} \uplus S_{\exp}$	$\mathcal{G}_{S,R} = N_{\exp,R} \uplus S_{\exp,R}$	$\mathcal{G}_{S,R}^{\text{red}} = N_{\exp,R}^{\text{red}} \uplus S_{\exp,R}^{\text{red}}$
Arbitrary	Arbitrary	Arbitrary	Arbitrary
$\mathcal{G} = N \uplus S$	$\mathcal{G}_S = N_{\exp} \uplus S_{\exp}$	$\mathcal{G}_{S,R} = N_{\exp,R} \uplus S_{\exp,R}$	$\mathcal{G}_{S,R}^{\mathrm{red}} = N_{\mathrm{exp},R}^{\mathrm{red}} \uplus S_{\mathrm{exp},R}^{\mathrm{red}}$





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2) Then, by implementing computer programs to perform expansion with any semigroup (up to order 6).

order	Q = # semigroups	
1	1	
2	4	
3	18	
4	126	[Forsythe '54]
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#### Those programs allow us:

- i) to study all resonant decomposition of a semigroup,
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- iv) To study further conditions for some specific problem





### **Final remarks**

#### As mentioned in

Nesterenko M. 2012 S-expansions of three dimensional Lie algebras arXiv:1212.1820

it would be interesting to know whether S-expansions fit the classification of solvable Lie algebras of a fixed dimension by means of S-expansions of simple (semisimple) Lie algebras of the same dimension.

The theoretical results proposed in this work could be useful in solving that problem.





### **Final remarks**

The S-expansion method have already been extended to the case of other mathematical structures (higher order Lie algebras, Loop algebras)

Caroca R, Merino N and Salgado P 2009 S-Expansion of higher-order Lie algebras J. Math. Phys. 50 013503

Caroca R, Merino N, Pérez A and Salgado P 2009 Generating higher order Lie algebras by expanding Maurer-Cartan forms J. Math. Phys. 50 123527

Caroca R, Merino N, Salgado P and Valdivia O 2011 *Generating infinite-dimensional algebras from loop algebras by expanding Maurer–Cartan forms* J. Math. Phys. 52 043519





### **Final remarks**

We think that this mathematical tool could be useful other gauge theories such as *Yang-Mills theory, Wess-Zumino models, higher spin theories, gauge/gravity duality,* etc.





## Thanks you for your attention





## Implementing computer programs

### **History about finite semigroup programs:**

#### HANDBOOK OF FINITE SEMIGROUP PROGRAMS

John A Hildebrant

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COMM SORTI					IGF	RO	U.	P												







## Implementing computer programs

### History about finite semigroup programs:

For example, for n=2 the program com.f gives the following list:





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### Checking with computer programs

### Isomorphisms and consistency of the procedure

Algebra	Semigroup used
Type I	many semigroups
Type II	$S_{(4)}^{10}, S_{(4)}^{12}$
Type III	$S_{(4)}^{13}, S_{(4)}^{28}, S_{(4)}^{42}, S_{(4)}^{43}, S_{(4)}^{44}, S_{(4)}^{45} \text{ and } S_{(4)}^{64}$
Type V	$S_{(4)}^{42}$

	isomorphic to		isomorphism
$S_{N1}$	$\iff$	$S_{(4)}^{44}$	$(\lambda_4 \lambda_1 \lambda_2 \lambda_3)$
$S_{N2}$	$\iff$	$S_{(4)}^{12}$	$(\lambda_4 \lambda_3 \lambda_2 \lambda_1)$
$S_{N3}$	$\iff$	$S_{(4)}^{42}$	$(\lambda_4 \lambda_1 \lambda_3 \lambda_2)$
$S_E^{(2)}$	$\iff$	$S_{(4)}^{43}$	$(\lambda_4 \lambda_3 \lambda_2 \lambda_1)$
$S_K^{(3)}$	$\iff$	$S_{(4)}^{45}$	$(\lambda_4 \lambda_2 \lambda_1 \lambda_3)$



