

Axion-Higgs Unification

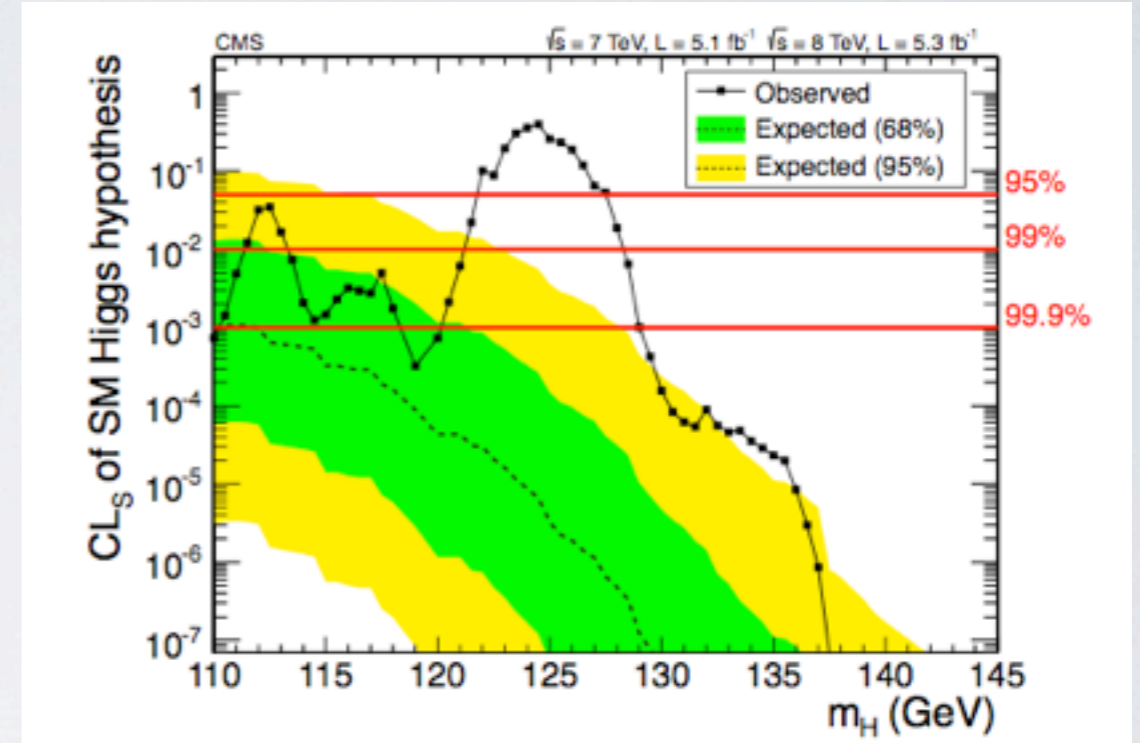
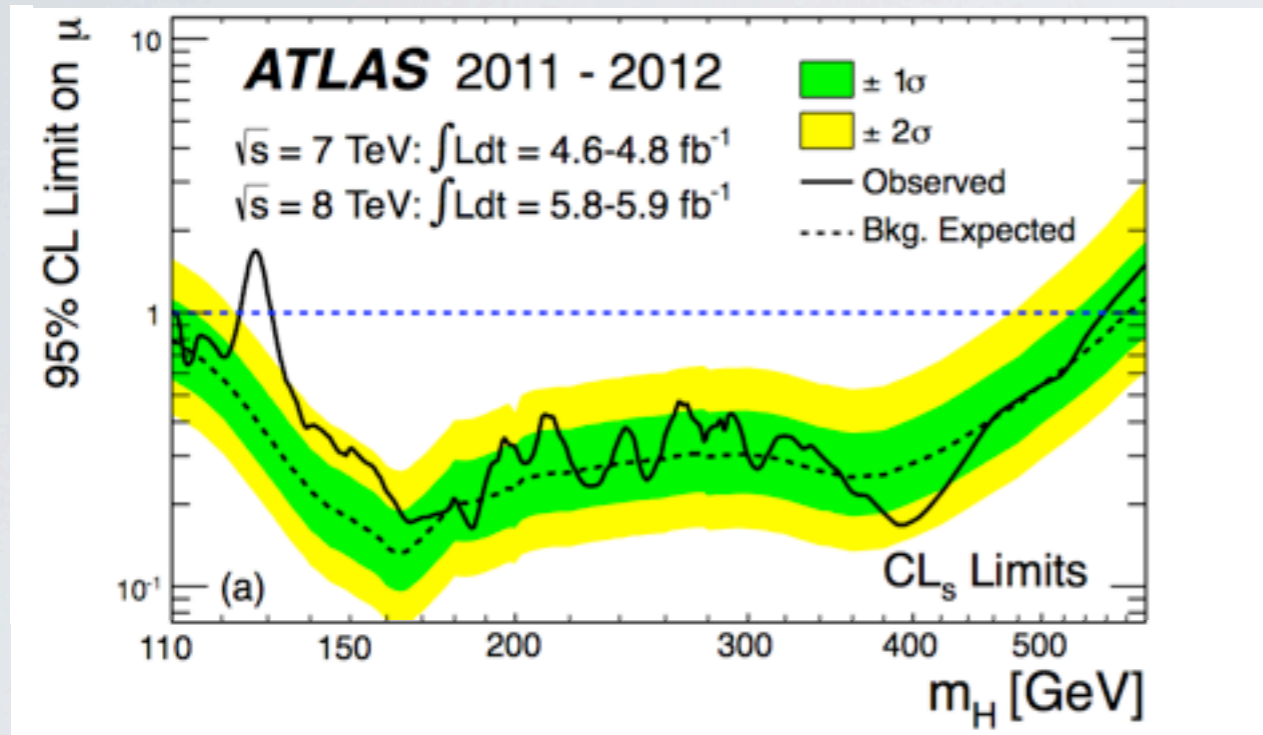
Michele Redi



1208.6013 with A. Strumia

Sao Paulo, October 2013

July 31, 2012 Phys. Lett. B716



$$m_h \approx 125 \text{ GeV}$$

What is the scale of new physics?

FLAVOR HAS FOUND NOTHING

$$\Lambda > 10^5 \text{ TeV}$$

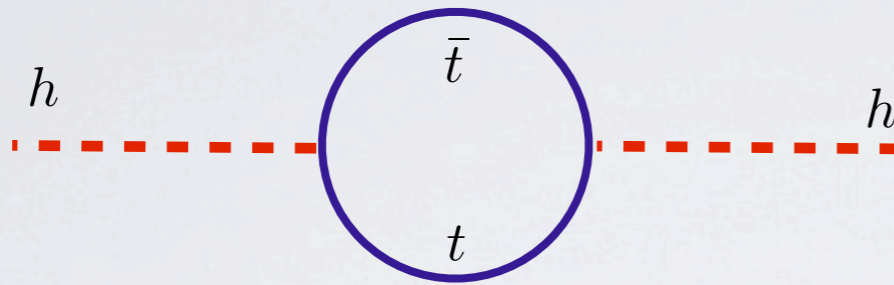
LEP HAS FOUND NOTHING

$$\Lambda > 5 - 10 \text{ TeV}$$

LHC HAS FOUND THE HIGGS
+ NOTHING

$$\Lambda > \text{few} \times \text{TeV}$$

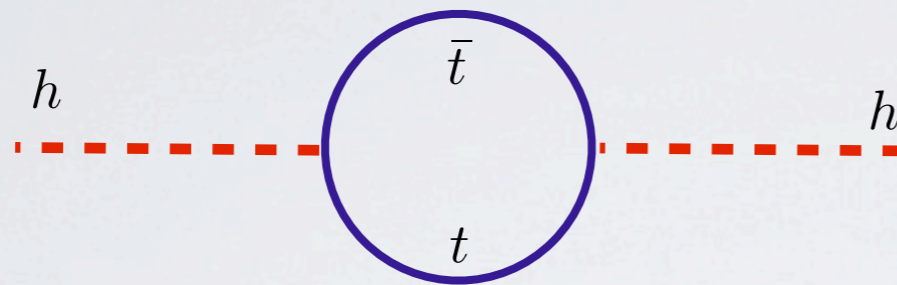
Hierarchy Problem:



$$\delta m_h^2 = -\frac{3\lambda_t^2}{8\pi^2}\Lambda_t^2$$

$$\Lambda < 1 \text{ TeV}$$

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Suspicion:

Maybe naturalness was not a good guide

$\Lambda \gg \text{TeV?}$

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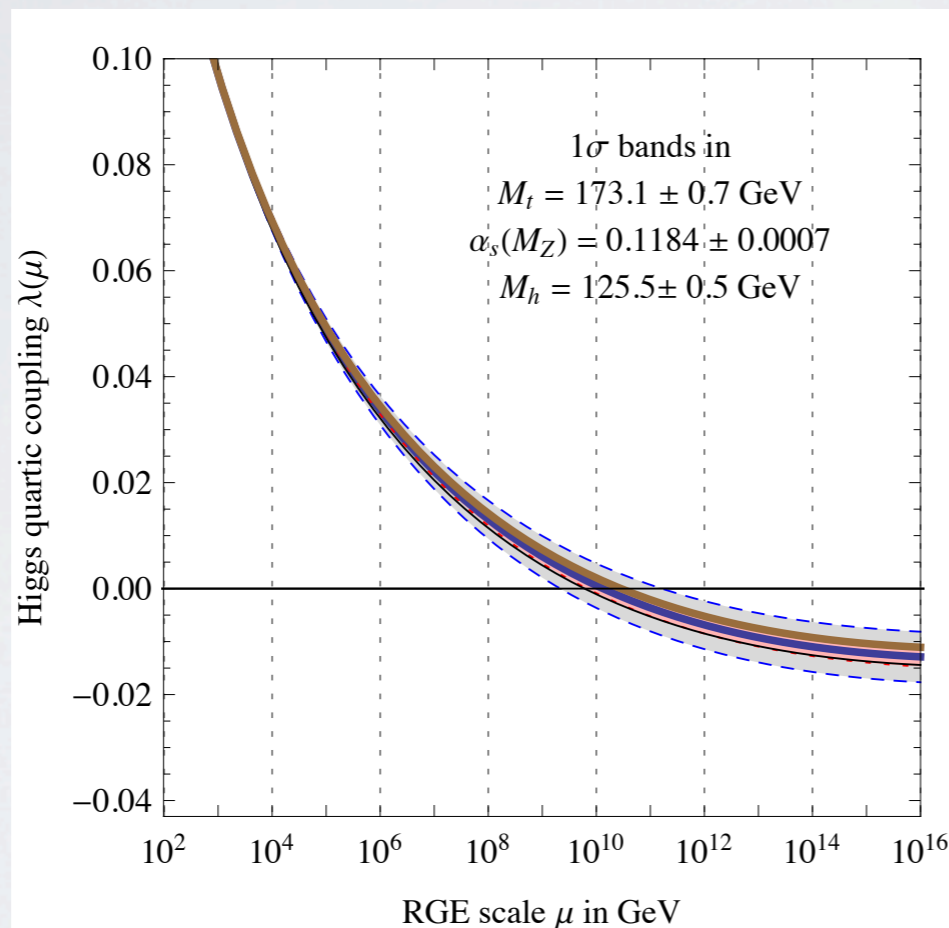
- Explains why we have not seen anything

$$\Lambda \gg \text{TeV?}$$

- Explains why we have not seen anything
- Higgs VEV could be tuned anthropically (as c.c.)

HINTS ?

- Running:



$$V(h) = m^2 h^2 / 2 + \lambda h^4 / 4$$

De Grassi et al.'12

Quartic almost zero at high scale for 125 GeV Higgs

- Strong CP problem:

$$\frac{\theta}{32\pi^2} \int d^4x \epsilon^{\mu\nu\rho\sigma} \text{Tr}[G_{\mu\nu} G_{\rho\sigma}] \qquad \theta < 10^{-10}$$

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Axions are most elegant solution

$$\theta \rightarrow \frac{a(x)}{f}$$

Axions are Goldstone bosons of a symmetry anomalous under QCD

$$m_a \sim \frac{m_\pi f_\pi}{f}$$

$$f > 10^9 \text{ GeV} \longrightarrow m_a < 10^{-3} \text{ eV}$$

Axions can be dark matter

$$\frac{\rho_a}{\rho_{\text{DM}}} \approx \theta_i^2 \left(\frac{f}{2 - 3 \times 10^{11} \text{ GeV}} \right) \xrightarrow{\theta_i \sim 1} f \approx 10^{11} \text{ GeV}$$

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- Neutrino masses

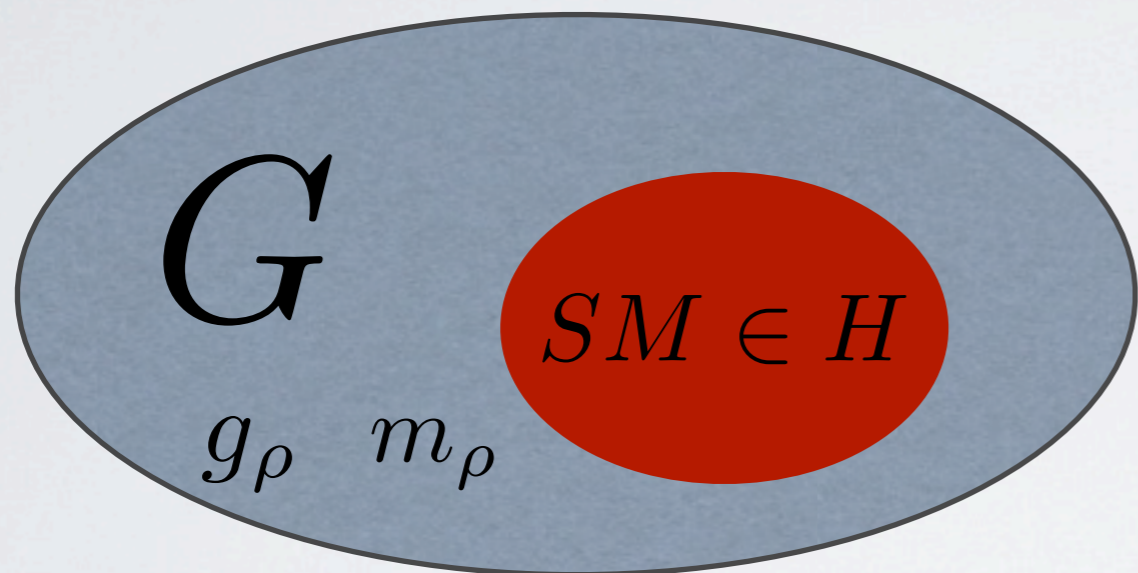
$$\frac{1}{\Lambda} (LH)^2 \qquad m_\nu \propto \frac{v^2}{\Lambda}$$

- Unification

$$\Lambda < M_{GUT}$$

GOLDSTONE HIGGS

The Higgs could be a remnant of strong dynamics.
Most compelling Higgs as Nambu-Goldstone boson:



$$GB = \frac{G}{H}$$

$$\delta m_h^2 \sim N_c \frac{y_t^2}{8\pi^2} m_\rho^2$$

Ex:

$$\frac{SO(5)}{SU(2)_L \otimes SU(2)_R} \xrightarrow{f = \frac{m_\rho}{g_\rho}} GB = (2, 2)$$

Agashe, Contino,
Pomarol, '04

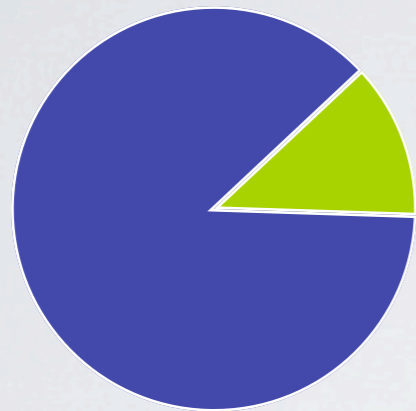
Deviation from SM:

$$\mathcal{O}\left(\frac{v^2}{f^2}\right)$$

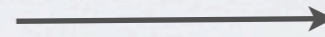
Deviation from SM:

$$\mathcal{O} \left(\frac{v^2}{f^2} \right)$$

Higgs is an angle,



$$0 < h < 2\pi f$$



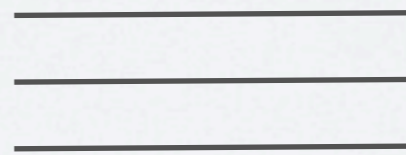
$$\text{TUNING} \propto \frac{f^2}{v^2}$$

Small Tuning

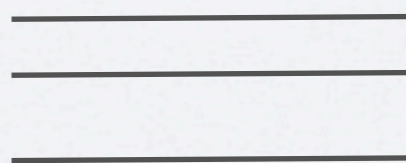
$$f < \text{TeV}$$

Natural spectrum:

$$f = 0.5 - 1 \text{ TeV}$$



$$m_\rho \sim 3 \text{ TeV}$$



$$m_h = 125 \text{ GeV}$$

$$m_W = 80 \text{ GeV}$$

$$0$$

Partial compositeness:

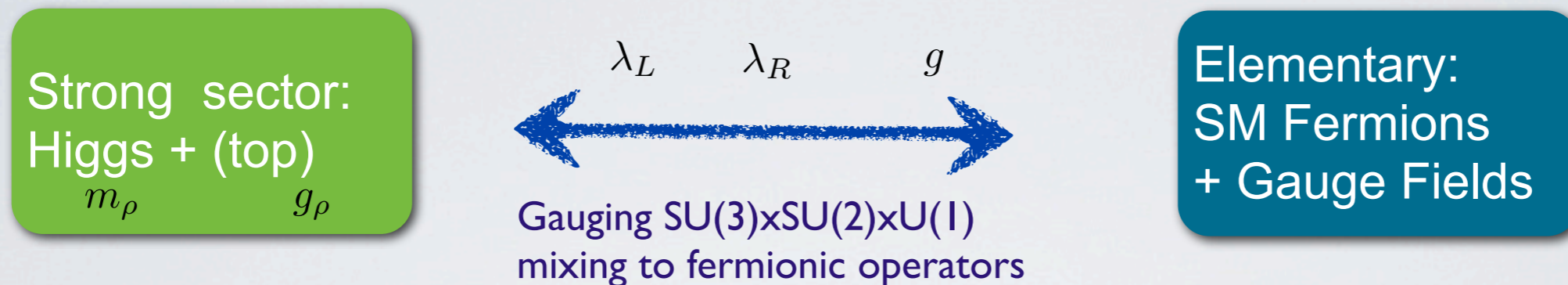
D. B. Kaplan '92
Contino-Pomarol, '04

Strong sector:
Higgs + (top)
 m_ρ g_ρ

Elementary:
SM Fermions
+ Gauge Fields

Partial compositeness:

D. B. Kaplan '92
Contino-Pomarol, '04



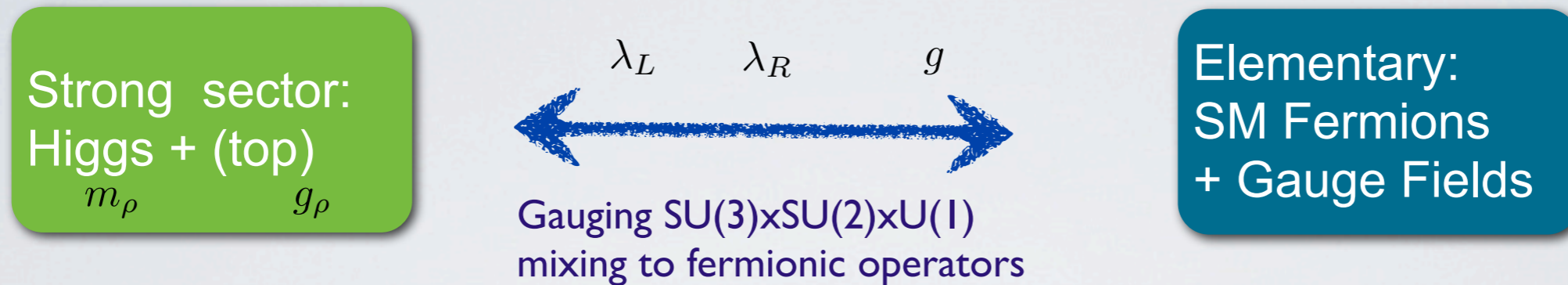
They talk through linear couplings:

$$\mathcal{L}_{gauge} = g A_\mu J^\mu$$

$$\mathcal{L}_{mixing} = \lambda_L \bar{f}_L O_R + \lambda_R \bar{f}_R O_R \quad \xrightarrow{\epsilon \sim \frac{\lambda}{Y}} \quad y_{SM} = \epsilon_L \cdot Y \cdot \epsilon_R$$

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- + phenomenology greatly ameliorated
- big dynamical assumptions

Many challenges:

- flavor

$$m_\rho > 20 \text{ TeV}$$

- precision tests

$$m_\rho > 3 \text{ TeV}$$

- direct exclusion

$$m_f > 0.7 \text{ TeV}$$

$$m_\rho > 2 \text{ TeV}$$

- smart model building

$$f \geq 10^{11} \text{ GeV}$$

- 125 GeV Higgs can be explained
- Axions + Higgs can be unified

HIGGS MASS

Higgs potential is generated by the couplings that break the global symmetry. Minimally gauge and Yukawa couplings.

$$V(h) = \sum_i a_i \sin^{2i} \left(\frac{h}{f} \right)$$

Electro-weak scale:

$$v \ll f$$



a_i must be tuned

Higgs mass is then “predicted”.

Coleman-Weinberg effective potential:

$$V(h)_{gauge} = \frac{9}{2} \int \frac{d^4 p}{(2\pi)^4} \ln \left[1 + \frac{1}{4} \frac{\Pi_1(p^2)}{\Pi_0(p^2)} \sin^2 \frac{h}{f} \right]$$

$$\langle J_\mu^a(q) J_\mu^a(-q) \rangle = (\eta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2}) \Pi^a(q^2)$$

$$\Pi_0 = \Pi^a$$

$$\langle J_\mu^{\hat{a}}(q) J_\mu^{\hat{a}}(-q) \rangle = (\eta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2}) \Pi^{\hat{a}}(q^2)$$

$$\Pi_1 = 2(\Pi^{\hat{a}} - \Pi^a)$$

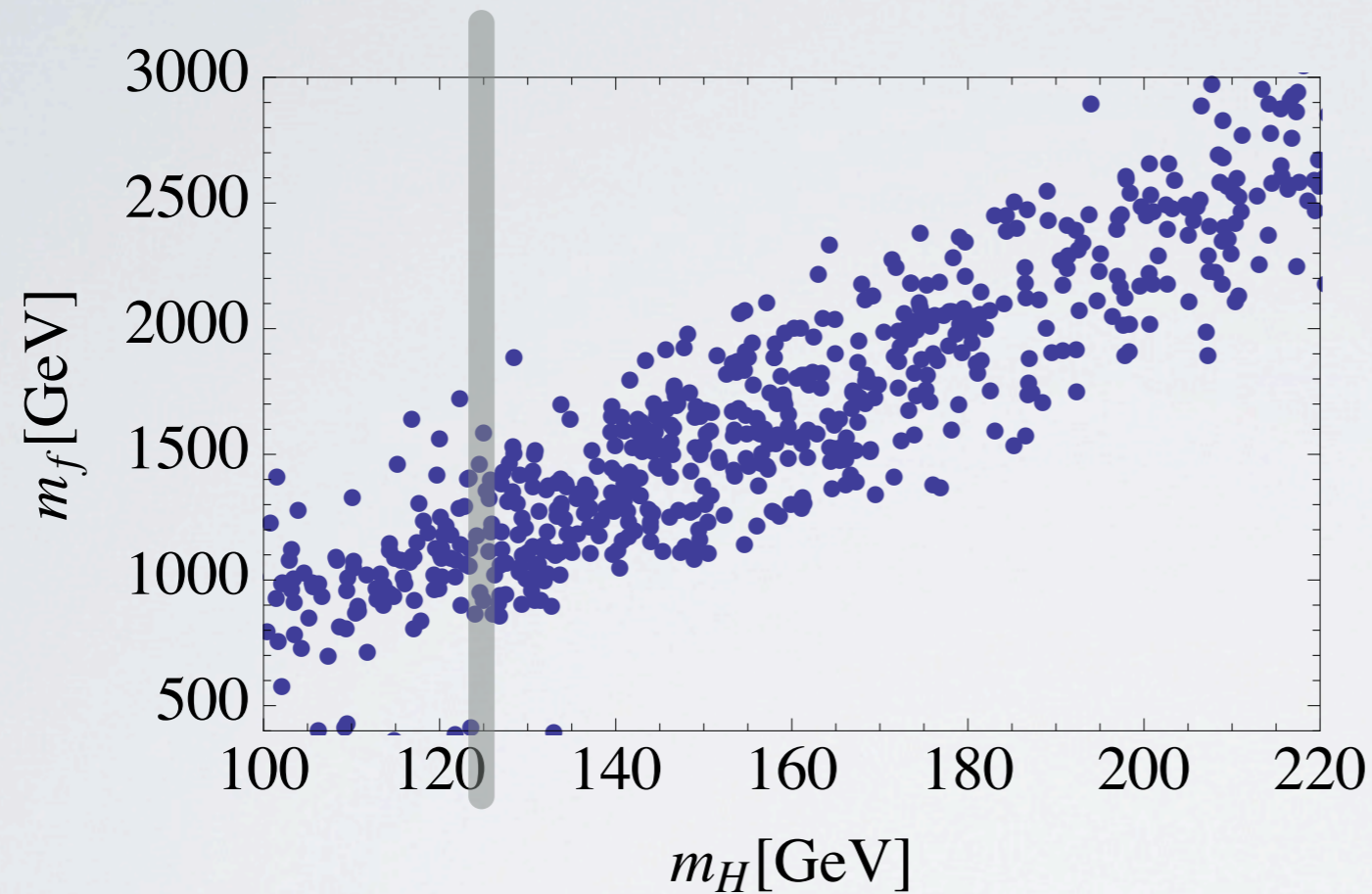
Potential is finite with a single SO(5) multiplet:

$$V(h)_{gauge} \approx \frac{9}{4} \frac{g^2}{(4\pi)^2} \frac{m_\rho^4}{g_\rho^2} \log \left[\frac{m_a^2}{m_\rho^2} \right] \sin^2 \frac{h}{f}$$

Divergences effectively cut-off by resonances

CHM5:

De Curtis, MR, '11
MR, Tesi, '12



$$f = 800 \text{ GeV}$$

$$m_h \sim \sqrt{\frac{N_c}{2}} \frac{y_t}{\pi} \frac{m_f}{f} v$$

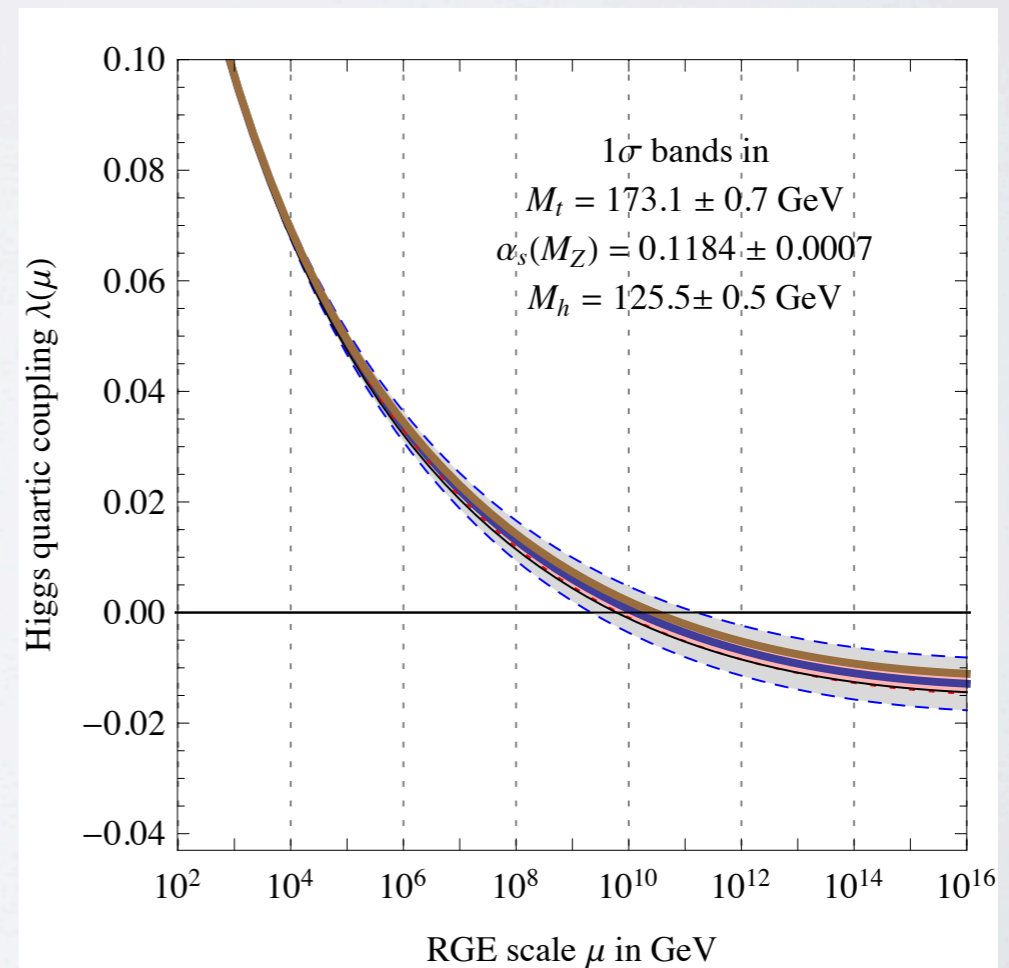
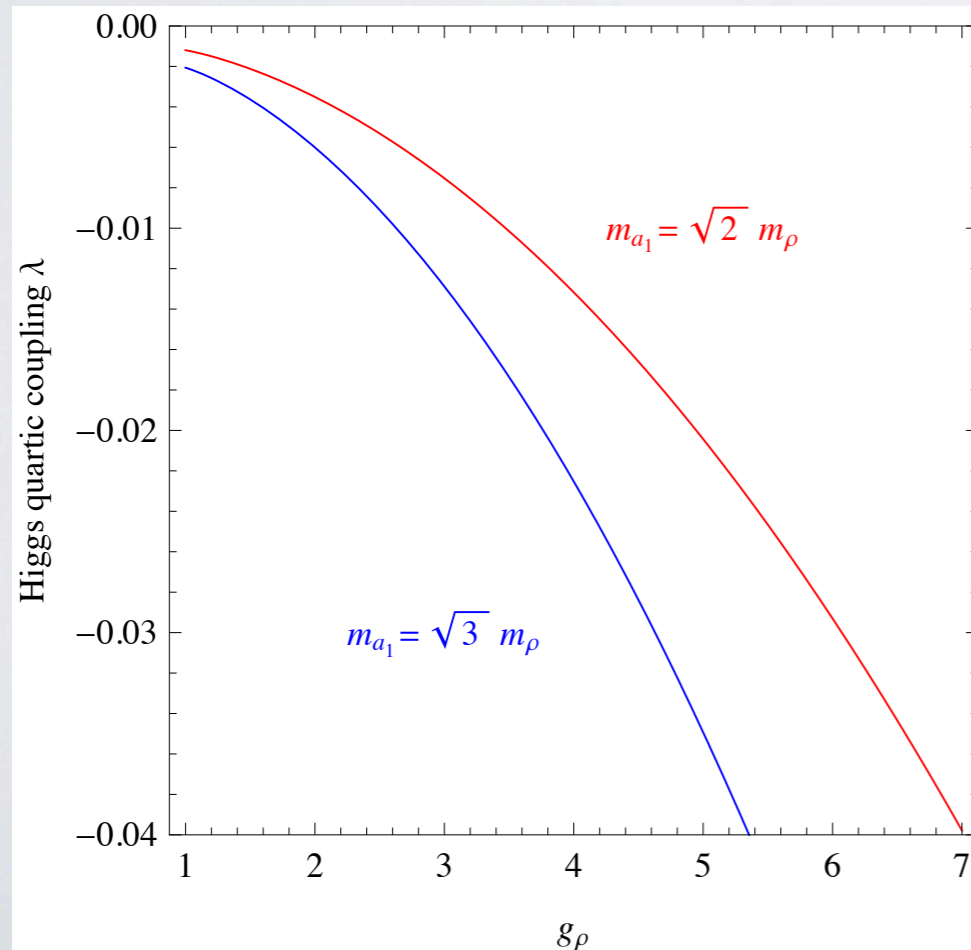
Partners around $\sim 1 \text{ TeV}$

Quartic boundary condition:

$$V(h, m_\rho)_{gauge} = \frac{9}{2} \int_{m_\rho}^{\infty} \frac{d^4 p}{(2\pi)^4} \ln \left[1 + \frac{1}{4} \frac{\Pi_1(p^2)}{\Pi_0(p^2)} \sin^2 \frac{h}{f} \right]$$

$$= \int_{m_\rho}^{\infty} \frac{d^4 p}{(2\pi)^4} \left[\frac{9}{8} \frac{\Pi_1(p^2)}{\Pi_0(p^2)} \sin^2 \frac{h}{f} - \frac{9}{64} \frac{\Pi_1(p^2)^2}{\Pi_0(p^2)^2} \sin^4 \frac{h}{f} + \dots \right]$$

$$\lambda(m_\rho) = -3 a_1 g_{SM}^2 \frac{g_\rho^2}{(4\pi)^2} + O \left(\frac{g_{SM}^4}{(4\pi)^2} \right)$$



Quartic contributions:

$$\cancel{\frac{g_\rho^4}{(4\pi)^2}}$$

$$\frac{g_{SM}^2 g_\rho^2}{(4\pi)^2}$$

$$\frac{g_{SM}^4}{(4\pi)^2}$$

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- **Leading order cancellation:**

$$\lambda(\Lambda) \sim g_{SM}^2 \frac{g_\rho^2}{(4\pi)^2} \sim \text{few } 10^{-2}$$

125 GeV Higgs requires weak coupling (large n)

- **Subleading order cancellation:**

$$\lambda(\Lambda) \sim \frac{g_{SM}^4}{(4\pi)^2} \sim 10^{-3}$$

AXION-HIGGS

Basic idea:

Axion and Higgs originate from the same dynamics.

f is fixed by dark matter and the electro-weak scale is tuned.

$$\frac{G}{H} \xrightarrow{f \approx 10^{11} \text{ GeV}} \text{Higgs} + \text{singlet}$$

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$$\frac{G}{H} \xrightarrow{f \approx 10^{11} \text{ GeV}} \text{Higgs} + \text{singlet}$$

- QCD anomaly from new fermions (KSVZ)
- QCD anomaly from SM fermions (DFSZ)

HIGGS + KSVZ AXION

Kim-Shifman-Vainstein-Zakharov:

Add new colored fermions + complex scalar

$$\Psi_Q \rightarrow e^{i\alpha_Q \gamma_5} \Psi_Q, \quad \sigma \rightarrow e^{-2i\alpha_Q} \sigma$$

$$L = L_{\text{SM}} + \bar{\Psi}_Q \partial \Psi_Q + |\partial_\mu \sigma|^2 + (\lambda \sigma \bar{\Psi}_Q \Psi_Q + \text{h.c.}) - V(\sigma)$$

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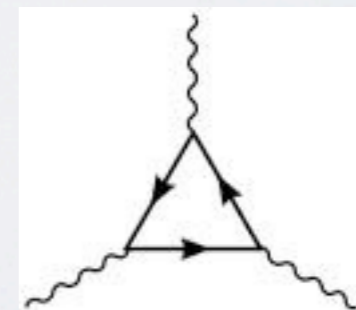
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Spontaneous PQ symmetry breaking

$$f \approx \langle \sigma \rangle$$

$$a = \sqrt{2} \text{Im}[\sigma]$$

PQ symmetry anomalous under QCD



$$\frac{G}{H} = \frac{SU(6)_L \times SU(6)_R}{SU(6)_{L+R}}$$

Under $SU(5)_{SM}$

$$\mathbf{35} = \mathbf{24} \oplus \mathbf{5} \oplus \bar{\mathbf{5}} \oplus \mathbf{1}$$

One Higgs doublet.

Two massless singlets are axion candidates.

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Two massless singlets are axion candidates.

Under SM 33 charged scalars acquire mass.

$$m \approx \frac{g_{SM}}{4\pi} \Lambda$$

UV realization: $SU(n)$ gauge theory with 6 flavors

Fermions	$U(1)_Y$	$SU(2)_L$	$SU(3)_c$	$SU(n)_{TC}$
D	$\frac{1}{3}$	1	$\bar{3}$	n
L	$-\frac{1}{2}$	2	1	n
N	0	1	1	n
\bar{D}	$-\frac{1}{3}$	1	3	\bar{n}
\bar{L}	$\frac{1}{2}$	$\bar{2}$	1	\bar{n}
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$$(\bar{5} + 1, n) \oplus (5 + 1, \bar{n})$$

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$$(\bar{5} + 1, n) \oplus (5 + 1, \bar{n})$$

$$\langle D\bar{D} \rangle = \langle L\bar{L} \rangle = \langle N\bar{N} \rangle \approx \Lambda^3$$

$$H \sim (L\bar{N}) - (\bar{L}N)^*$$

Singlets:

$$D\bar{D}$$

$$L\bar{L}$$

$$N\bar{N}$$

$$D\bar{D} + L\bar{L} + N\bar{N} \xrightarrow[U(1) \times SU(n)_{TC}^2 \text{ anomaly}]{\quad} \frac{g_{TC}}{4\pi} \Lambda$$

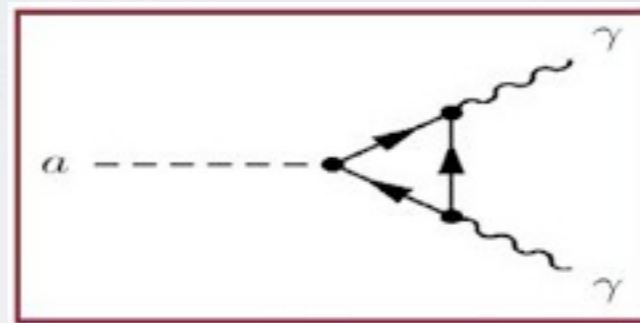
Axions couple to photon and gluons through anomalies

$$\frac{a}{32\pi^2 f} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma}$$

$$\frac{aN}{32\pi^2 f} \epsilon^{\mu\nu\rho\sigma} \text{Tr}[G_{\mu\nu} G_{\rho\sigma}]$$

$$E = \sum Q_{PQ} Q_{em}^2$$

$$N = \sum Q_{PQ} T_{SU(3)}^2$$



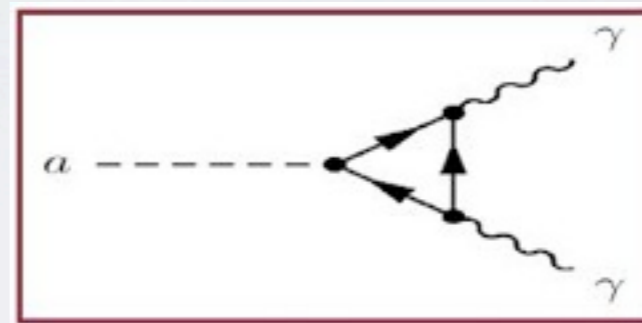
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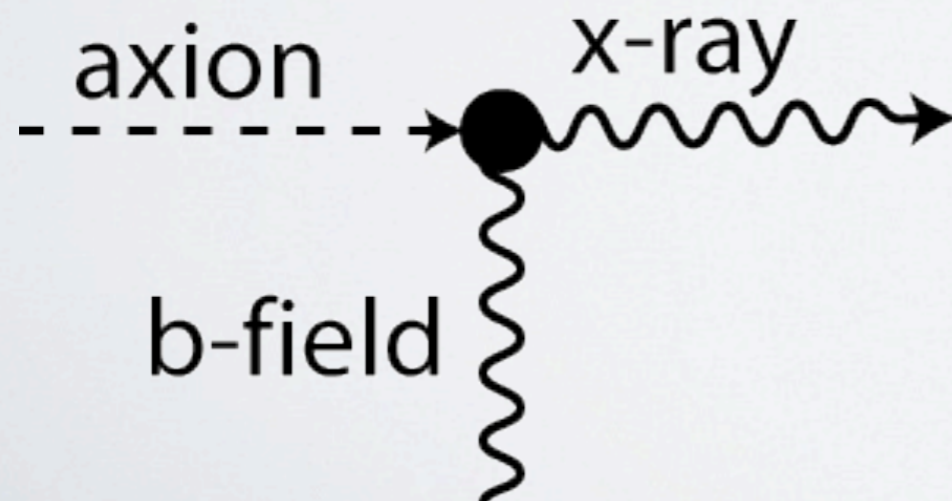
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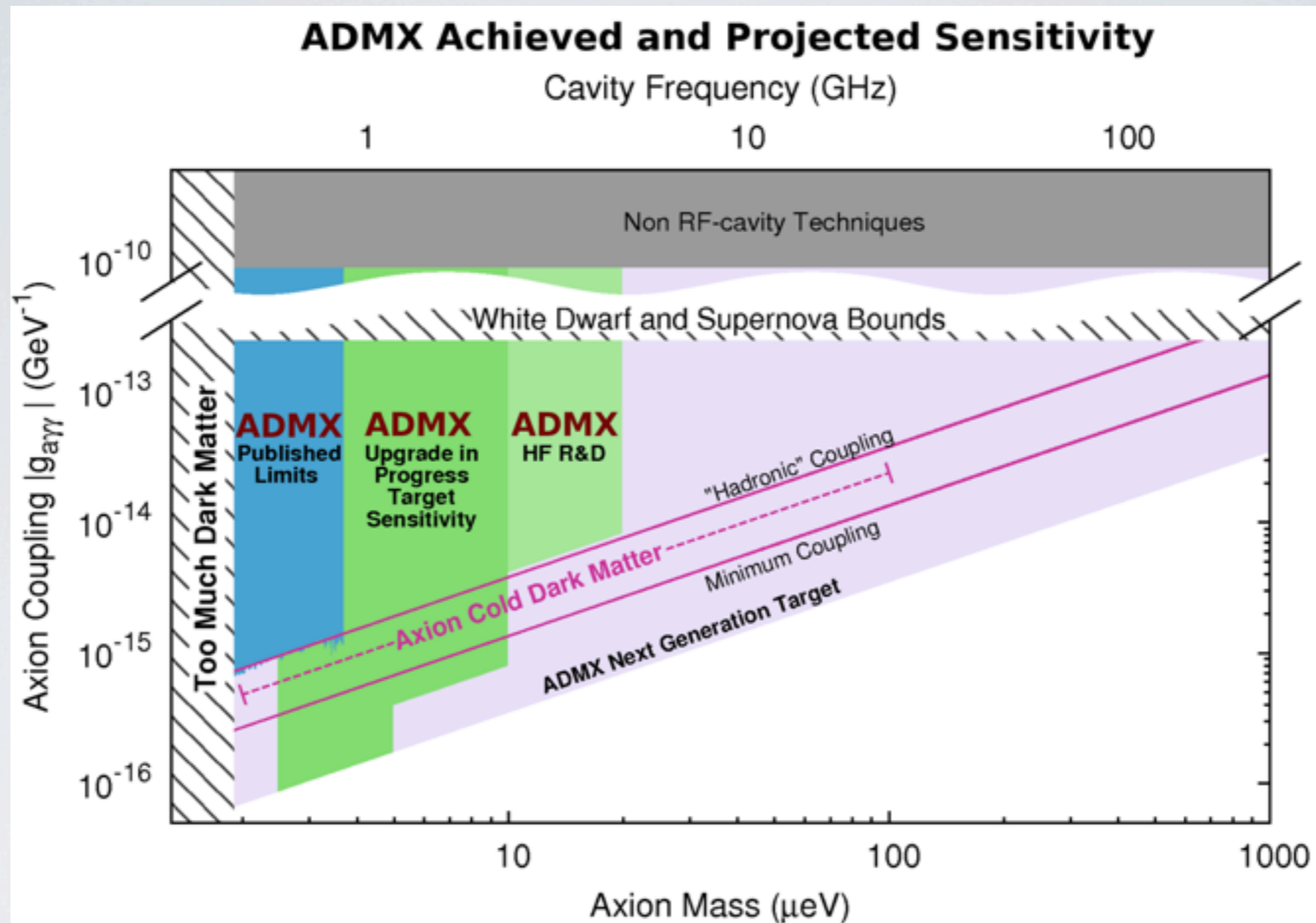


Experiments measure conversion of axion to photons



$$g_{a\gamma\gamma} = \frac{2(E/N - 1.92)}{10^{16} \text{ GeV}} \frac{m_a}{\mu\text{eV}}$$

$$m_a \sim \frac{N f_\pi m_\pi}{2f} \frac{\sqrt{m_u m_d}}{m_u + m_d}$$



$$\frac{E}{N} < 1.92 + 3.5 \sqrt{\frac{0.3 \text{ GeV}/\text{cm}^3}{\rho_{DM}}} \quad (m_a = 1.9 - 3.55 \times 10^{-6} \text{ eV})$$

Flavor:

$$\frac{1}{\Lambda_1^2} (qu)(L\bar{N})$$

$$\frac{1}{\Lambda_2^2} (\bar{q}\bar{u})(\bar{L}N)$$

Flavor: $\frac{1}{\Lambda_1^2} (qu)(L\bar{N}) \quad \frac{1}{\Lambda_2^2} (\bar{q}\bar{u})(\bar{L}N)$

a) If UV interactions respect singlets symmetry

$$\frac{4D - 3L - 6N}{\sqrt{102}}, \quad \frac{L - 2N}{\sqrt{3}} \quad \frac{E}{N} = -\frac{5}{6}$$

b) If all Yukawas allowed

$$\frac{D - 3L + 3N}{\sqrt{30}} \quad \frac{E}{N} = -\frac{16}{3}$$

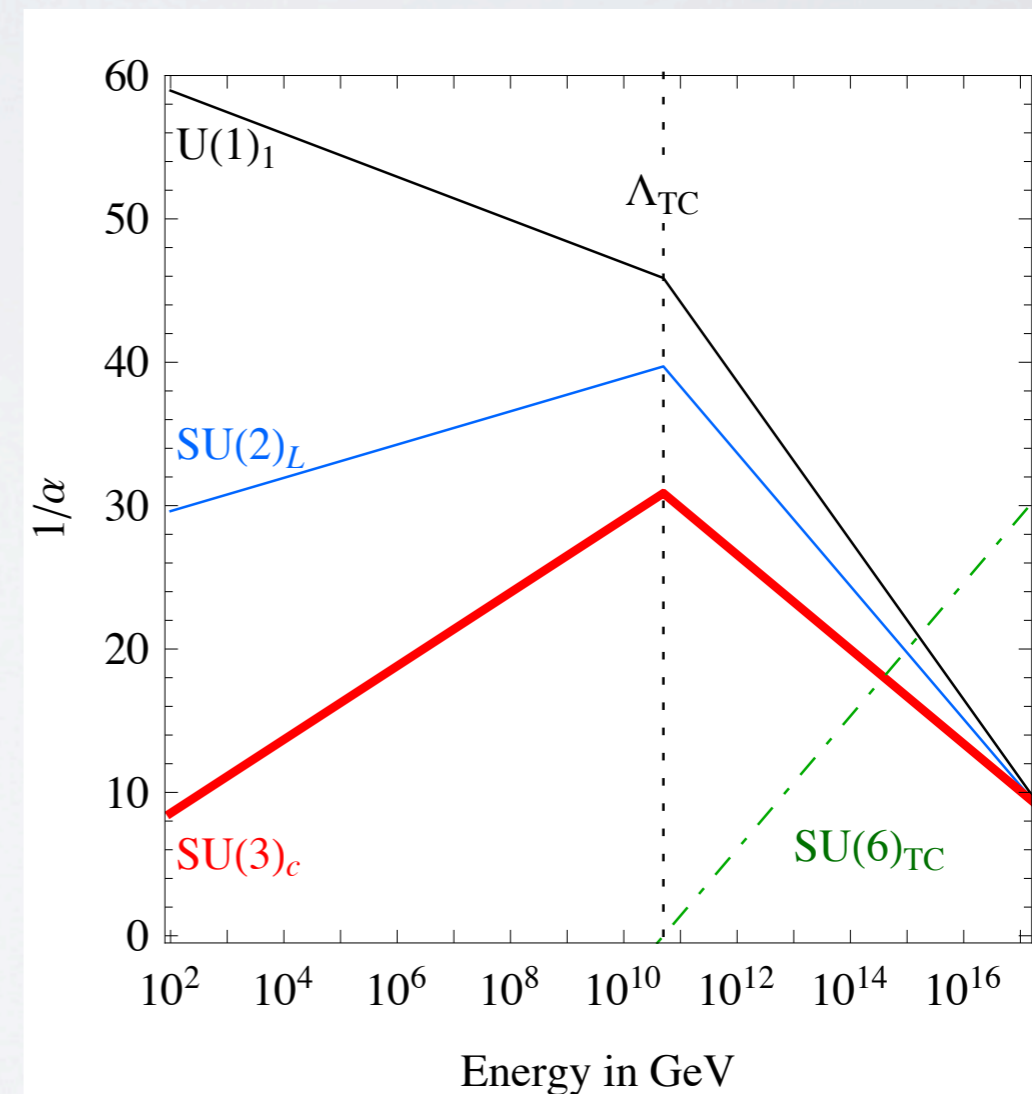
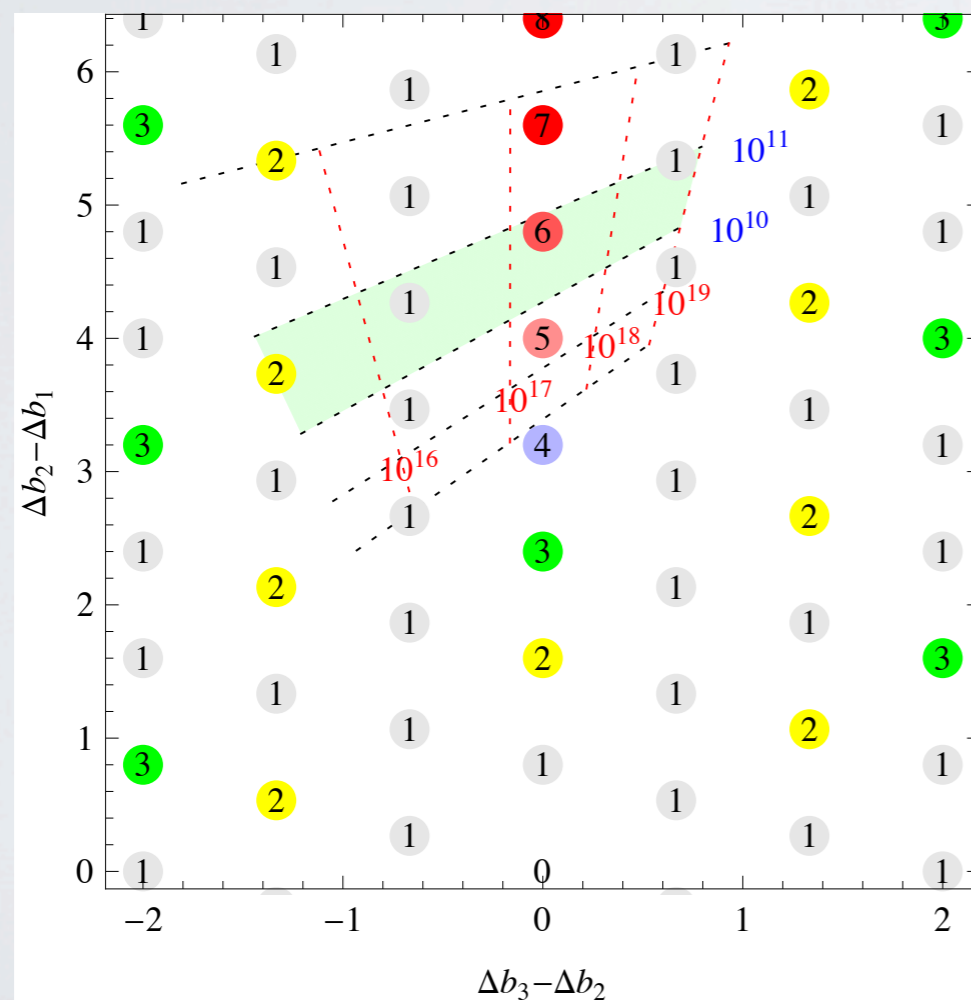
c) SU(5) is gauged

$$\frac{D + L - 5N}{\sqrt{30}} \quad \frac{E}{N} = \frac{8}{3}$$

Incomplete SU(5) multiplets can improve unification (“unificaxion”)

Giudice, Rattazzi, Strumia '12

Ex: D, L, Q, U, N



HIGGS + DFSZ AXION

Dine-Fischler-Srednicki-Zhitnitsky:

Two Higgs doublets and complex singlet

$$\sigma \rightarrow e^{4i\alpha} \sigma, \quad q_{L,R} \rightarrow e^{i\alpha} q_{L,R} \quad H_u \rightarrow e^{-2i\alpha} H_u, \quad H_d \rightarrow e^{-2i\alpha} H_d$$

$$f = \sqrt{v_u^2 + v_d^2 + |\sigma|^2}$$

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Ex:

$$\frac{G}{H} = \frac{SU(6)}{SO(6)}$$

$$SO(6) \supset SO(4) \otimes U(1)_{PQ}$$

$$\mathbf{20}' = (\mathbf{2}, \mathbf{2})_{\pm 2} \oplus (\mathbf{1}, \mathbf{1})_{\pm 4} \oplus (\mathbf{1}, \mathbf{1})_0 \oplus (\mathbf{3}, \mathbf{3})_0$$

UV realization: $SO(n)$ gauge theory with 6 flavors

Fermions	$U(1)_Y$	$SU(2)_L$	$SU(3)_c$	$SO(n)_{TC}$	$U(1)_{PQ}$
L	$-\frac{1}{2}$	2	1	n	0
\bar{L}	$\frac{1}{2}$	$\bar{2}$	1	n	0
N	0	1	1	n	2
\bar{N}	0	1	1	n	-2

$$\langle L\bar{L} \rangle = \langle N\bar{N} \rangle = \Lambda^3$$

$$H_1 \sim LN$$

$$H_2 \sim \bar{L}\bar{N}$$

UV realization: $SO(n)$ gauge theory with 6 flavors

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$$H_1 \sim LN$$

$$\langle L\bar{L} \rangle = \langle N\bar{N} \rangle = \Lambda^3$$

$$H_2 \sim \bar{L}\bar{N}$$

Yukawas must respect PQ

$$\frac{1}{\Lambda_t^2} (q_L t_R^c)^\dagger (L N) + \frac{1}{\Lambda_b^2} (q_L b_R^c)^\dagger (\bar{L} \bar{N}) + \text{h.c.}$$

Anomalies:

$$E_{TC} = 0$$

$$\frac{E}{N} = \frac{8}{3}$$

PARTIAL COMPOSITENESS

$$\frac{G}{H} = \frac{SO(6)}{SO(5)} \simeq \frac{SU(4)}{Sp(4)}$$

Gripaios, Pomarol, Riva, Serra '09

Redi, Tesi '12

Galloway et. al. '10

5 GBs:

$$5 = (2, 2) + 1$$

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5 GBs:

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Gauging of SM gauge symmetry preserves

$$SU(2)_L \times U(1)_Y \times U(1)_{PQ}$$

Under $U(1)_{PQ}$ singlet shifts.

Sp(n) theories with 4 flavors

Fermions	$U(1)_Y$	$SU(2)_L$	$SU(3)_c$	$Sp(n)_{TC}$	$U(1)_{PQ}$
D	0	2	1	n	+1
S	$+\frac{1}{2}$	1	1	n	-1
\bar{S}	$-\frac{1}{2}$	1	1	n	-1

Sp(n) theories with 4 flavors

Fermions	$U(1)_Y$	$SU(2)_L$	$SU(3)_c$	$Sp(n)_{TC}$	$U(1)_{PQ}$
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S	$+\frac{1}{2}$	1	1	n	-1
\bar{S}	$-\frac{1}{2}$	1	1	n	-1

Difficult to generate QCD anomaly

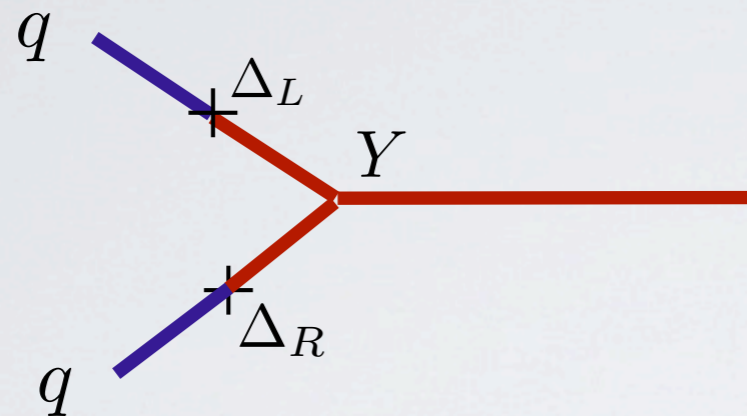
$$(qu)(DS)$$

$$(qu)(DS)(S\bar{S})$$

We can be build models with partial compositeness

$$m\psi\Psi + M\Psi\Psi + g_{TC}\Psi\Psi H$$

Partial compositeness



$$y_{SM} = \epsilon_L \cdot Y \cdot \epsilon_R$$

$$\psi_{SM} \Psi_{comp}$$

$$\Psi_{comp} \in G$$

We can choose

$$10 = (2, 2) + (3, 1) + (1, 3)$$

$$q_L \subset (2, 2)$$

$$u_R, d_R, e_R \subset (1, 3)$$

Mixing induces a PQ charge on the SM fermions.

$$\delta q_L = 0 \quad \delta u_R = -\frac{1}{\sqrt{2}}u_R \quad \delta d_R = -\frac{1}{\sqrt{2}}d_R \quad \delta e_R = -\frac{1}{\sqrt{2}}e_R$$

PQ symmetry is anomalous due to u_R, d_R, e_R rotations

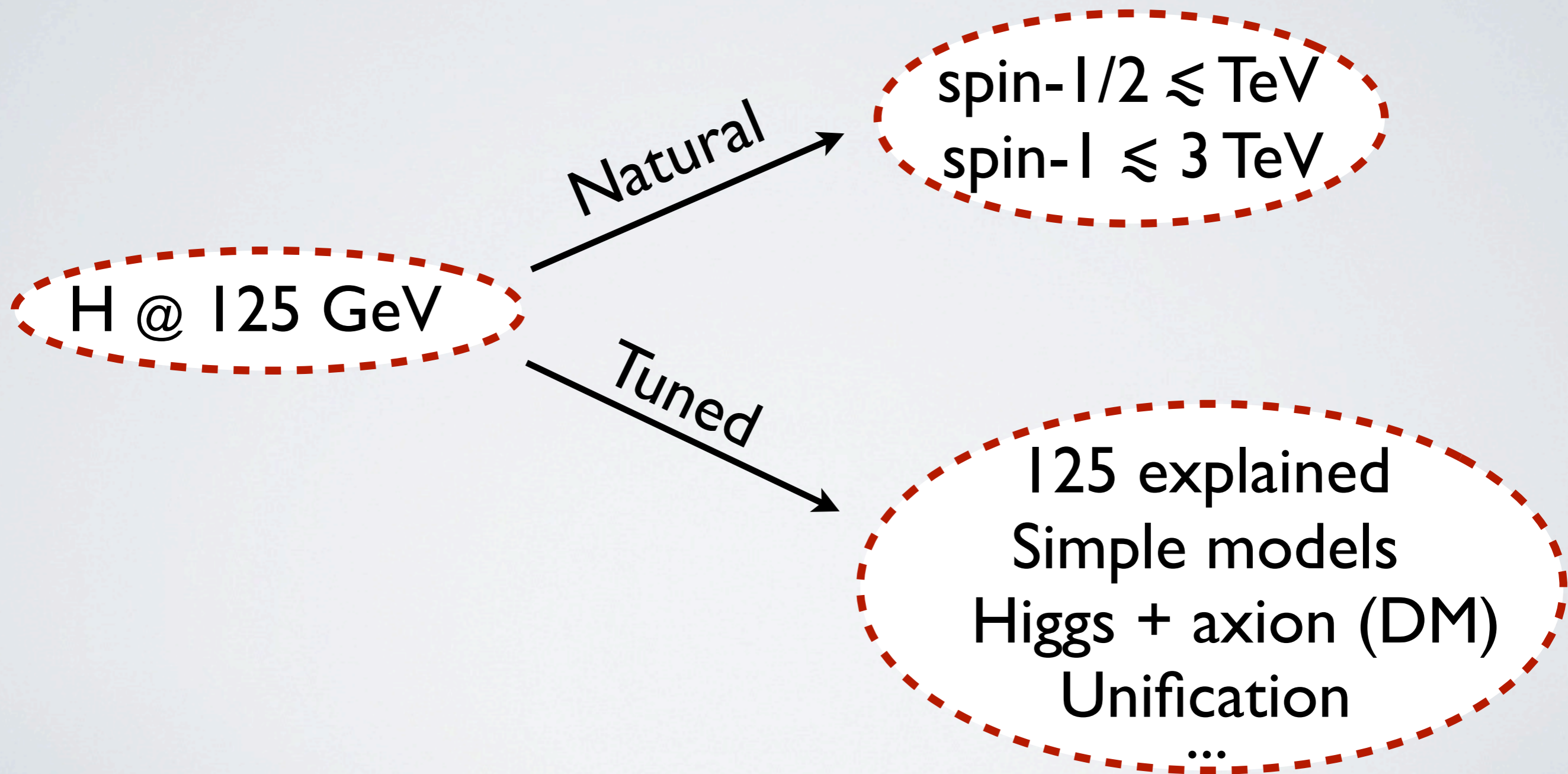
$$N = 2N_F$$

$$E = 2 \left[\left(\frac{4}{9} + \frac{1}{9} \right) 3 + 1 \right] N_F + E_{TC}$$

$$\frac{E}{N} = \frac{8}{3} + \frac{E_{TC}}{6}$$

$$E_{TC} \sim n$$

GOLDSTONE HIGGS



OBRIGADO !

Fermions can couple to $6=(2,2)+2 \times 1$

$$q_L \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} b_L \\ -ib_L \\ t_L \\ it_L \\ 0 \\ 0 \end{pmatrix} \quad t_R \rightarrow \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ i \cos \theta t_R \\ \sin \theta t_R \end{pmatrix}$$

Fermions can couple to $6=(2,2)+2 \times 1$

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For $\theta = \frac{\pi}{4}$ singlet becomes exact GB

PQ symmetry is anomalous due to tR rotations

$$E = 2 \left[\left(\frac{4}{9} + \frac{1}{9} \right) 3 + 1 \right] N_F + E_{\text{TC}}$$

$$\frac{E}{N} = \frac{8}{3} + \frac{E_{\text{TC}}}{6} \quad E_{\text{TC}} \sim n$$

Neutrino masses can be generated by see-saw mechanism

$$\frac{1}{\Lambda_\nu^2} (l\nu_R^c)^\dagger (LN) + m(\nu_R^c)^2 + h.c. \longrightarrow \lambda l H_1 \nu_R^c + m(\nu_R^c)^2 + h.c.$$

$$m_\nu \sim \frac{\lambda^2 v^2}{m}$$

If no right-handed neutrinos

$$\frac{1}{\Lambda_\nu^5} (l\bar{L})^2 N^2 \longrightarrow \frac{1}{\Lambda_\nu^3} (lH_u)^2 \sigma^2 + \dots$$

Leptogenesis?