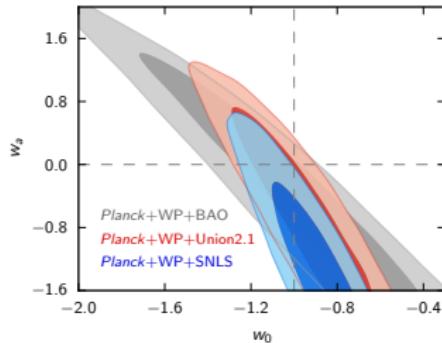
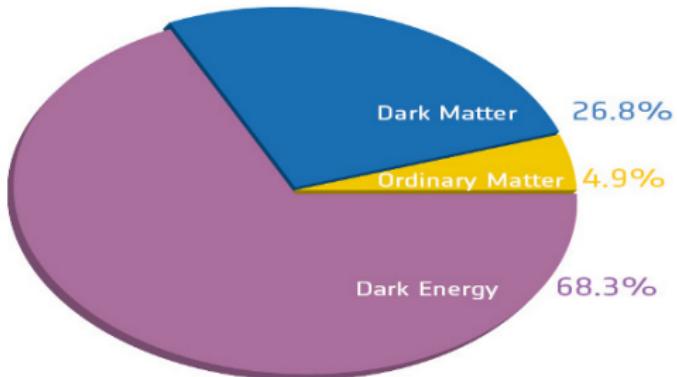


Modified gravity from the laboratory to the stars

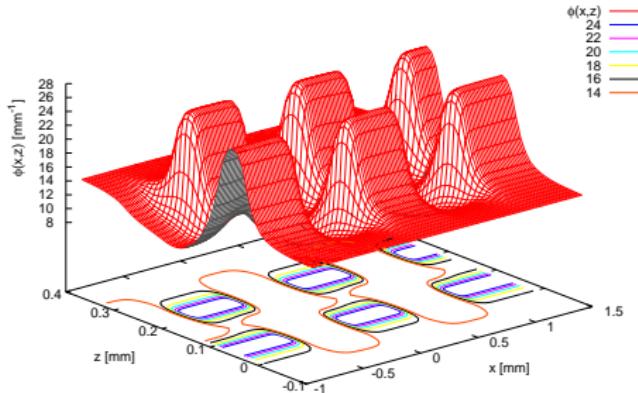
A photograph of a city skyline at night, featuring numerous lit buildings. A bright, branching lightning bolt strikes one of the taller skyscrapers on the right side of the frame.

Amol Upadhye
IEU, Ewha University
February 20, 2014

Cosmic acceleration



Time-variation



Couplings to known particles

Outline

- ① Introduction: Coupled dark energy from modified gravity
- ② Fifth forces from matter couplings
 - “Quantum-stable” chameleon dark energy
 - Laboratory searches
 - Modified gravity in stars
- ③ How dark is dark energy? Searches for photon couplings
 - Scalar-photon oscillation
 - Afterglow experiments
 - Dark energy from the Sun

Coupled dark energy from modified gravity



A phenomenological toolbox:

Modified gravity	Effective scalar	New physics
4-D modified action: $R \rightarrow f(R)$	Conformal trans.: \Rightarrow chameleon	matter coupling, effective $m(\rho)$

Coupled dark energy from modified gravity



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4-D modified action: $\phi \rightarrow -\phi$ symmetry	Conformal trans.: \Rightarrow symmetron	matter coupling, uncoupled phase

Coupled dark energy from modified gravity



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DGP, etc.: non-compact extra dimension	Decoupling limit (weak gravity) ⇒ Galileon	matter coupling, non-canonical kinetic term

Coupled dark energy from modified gravity



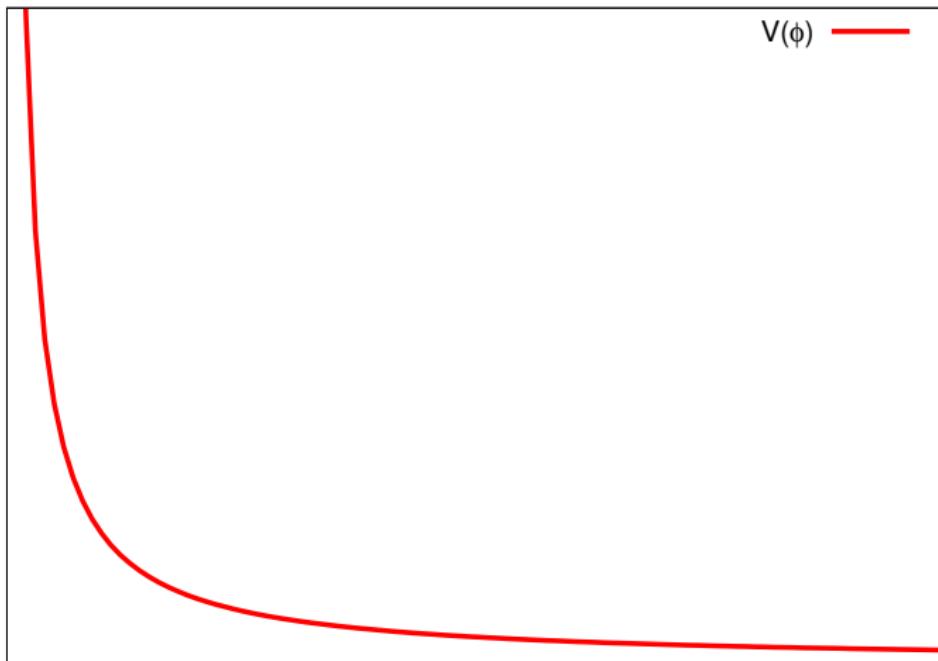
A phenomenological toolbox:

Modified gravity	Effective scalar	New physics
4-D modified action: $R \rightarrow f(R)$	Conformal trans.: ⇒ chameleon	matter coupling, effective $m(\rho)$
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DGP, etc.: non-compact extra dimension	Decoupling limit (weak gravity) ⇒ Galileon	matter coupling, non-canonical kinetic term
Kaluza-Klein, etc.: compact extra dim.	Small extra dim. ⇒ radion	matter coupling, photon coupling

At low energies, dark energy can have a matter coupling, whose fifth force must be screened locally. Dark energy can also have a photon coupling, allowing the production of dark energy particles.

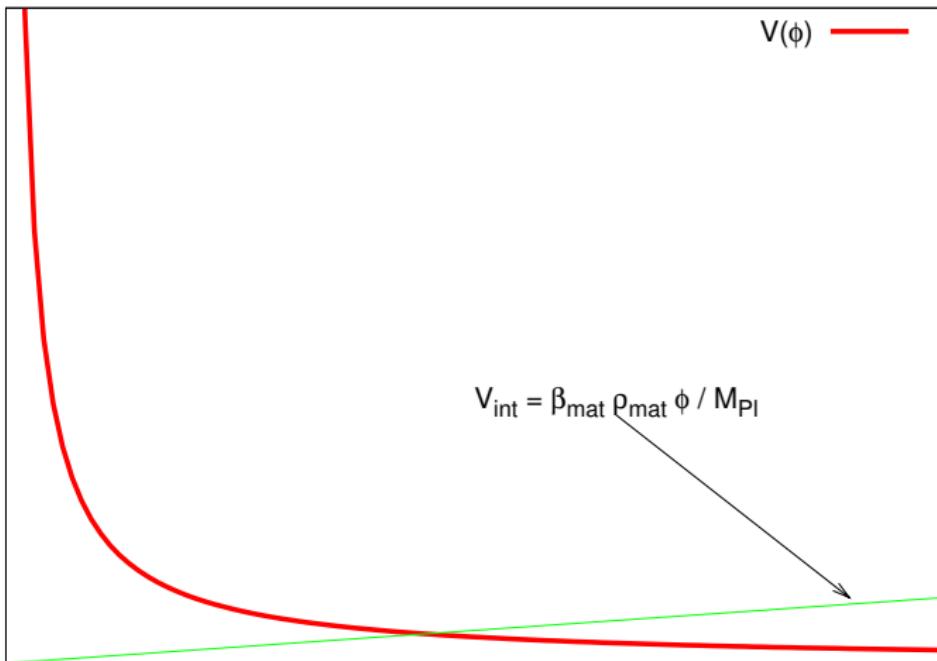
Chameleon mechanism

effective potential: $V_{\text{eff}}(\phi, \rho) = V(\phi) + \beta\rho\phi/M_{\text{Pl}}$



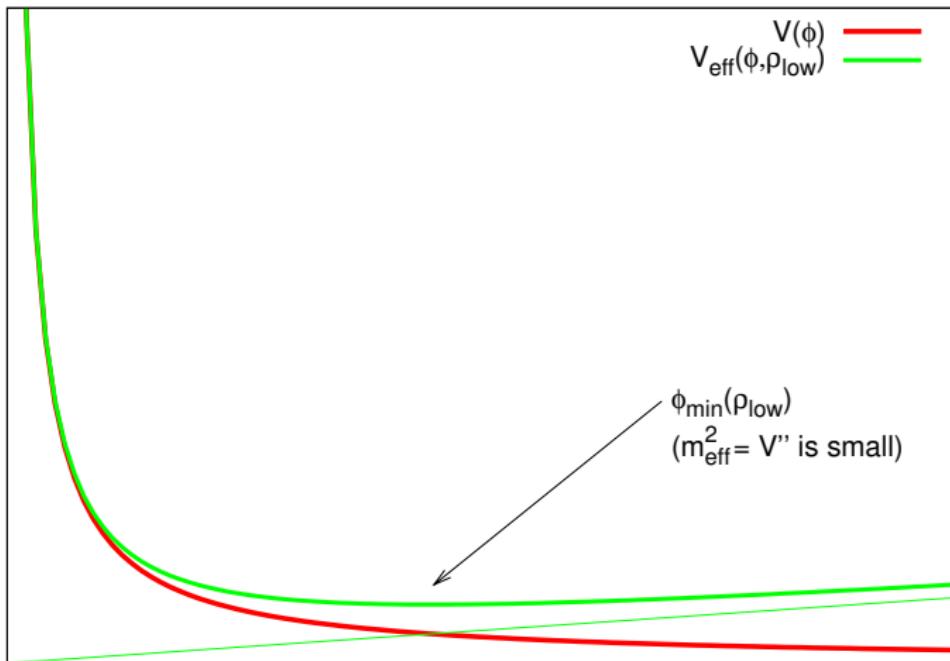
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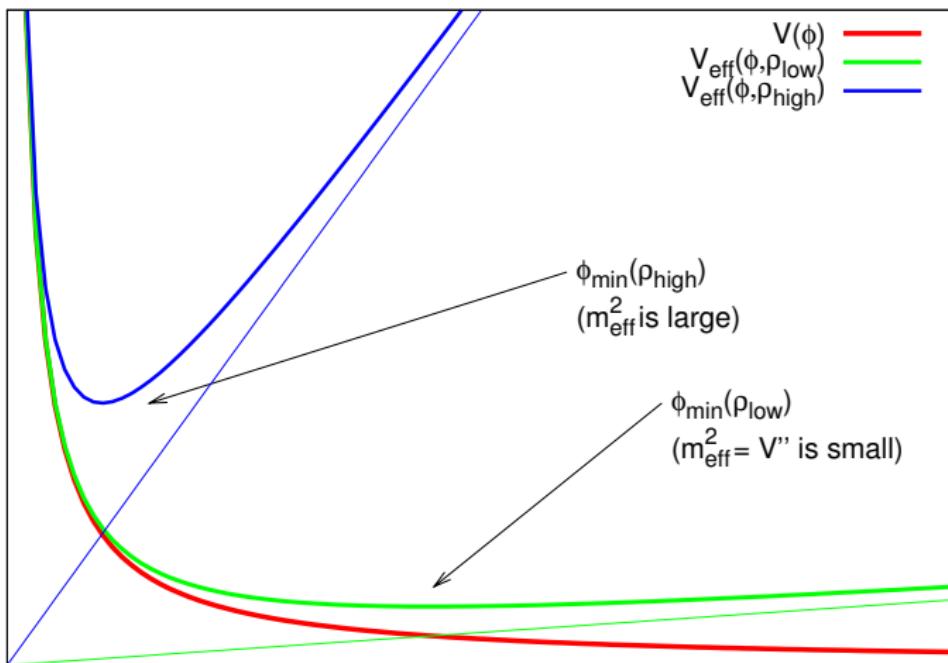
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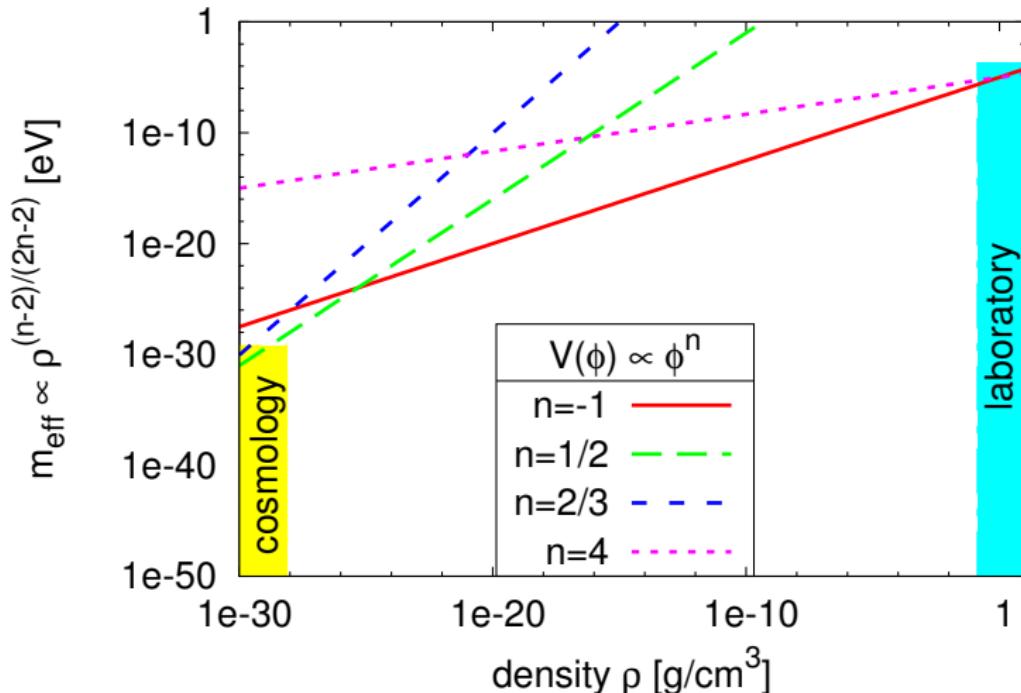
Chameleon mechanism

effective potential: $V_{\text{eff}}(\phi, \rho) = V(\phi) + \beta\rho\phi/M_{\text{Pl}}$



At which scale should we probe each model?

$$V(\phi) \propto \phi^n + \text{const.} \Rightarrow m_{\text{eff}} \propto \rho^{\frac{n-2}{2n-2}} \quad (\text{use lab for } n \lesssim -\frac{1}{2}, n > 2)$$

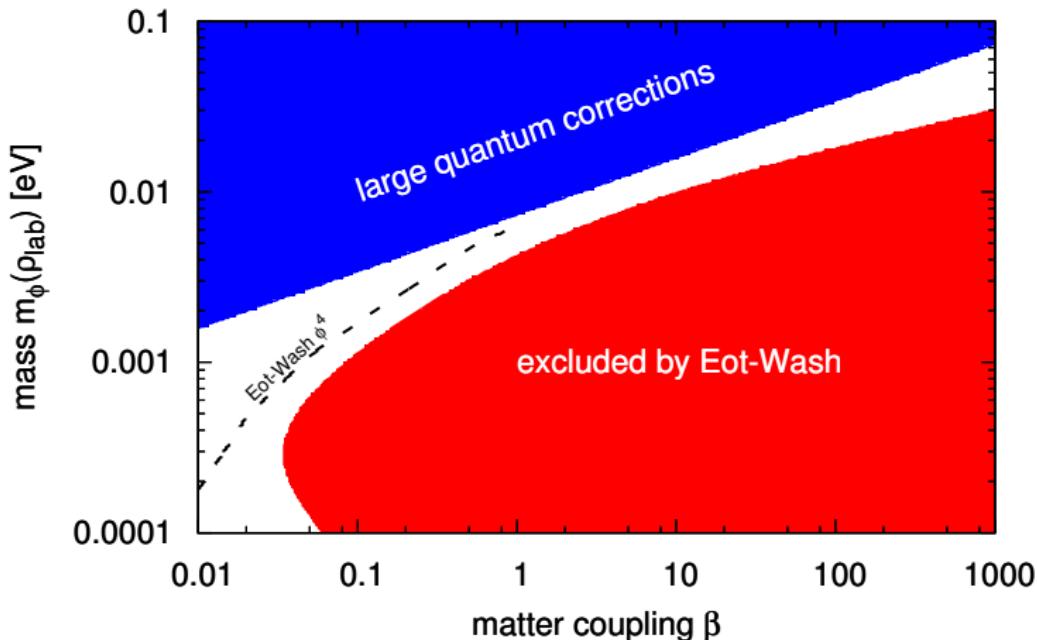


AU, PRD 86:102003(2012)[arXiv:1209.0211].

Laboratory benchmark: “quantum-stable” chameleons

$$\Delta V_{\text{1-loop}}(\phi) = \frac{m_{\text{eff}}(\phi)^4}{64\pi^2} \log\left(\frac{m_{\text{eff}}(\phi)^2}{\mu^2}\right) < V_{\text{tree}}$$

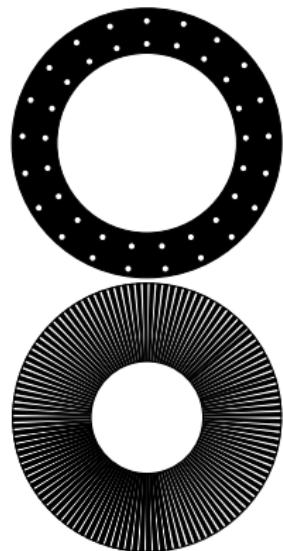
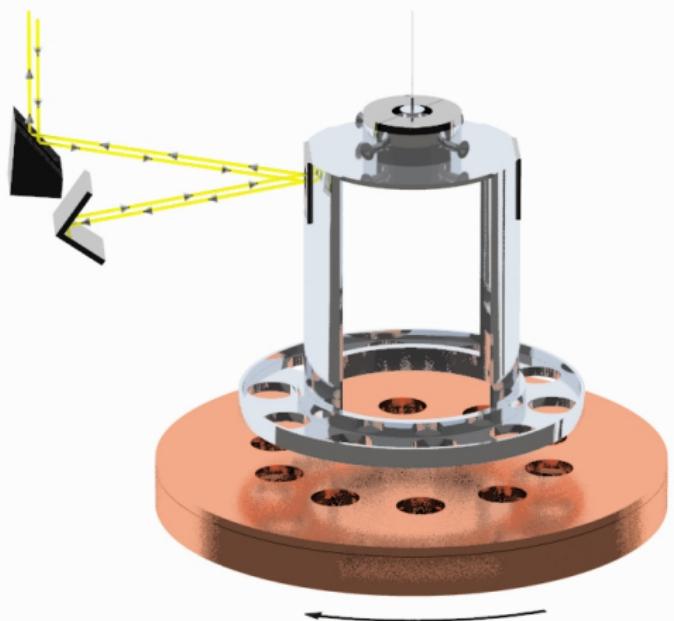
$$\Rightarrow m_{\text{eff}} \leq \left(\frac{48\pi^2 \beta^2 \rho^2}{M_{\text{Pl}}^2} \right)^{1/6} = 0.0073 \left(\frac{\beta \rho}{10 \text{g/cm}^3} \right)^{1/6} \text{eV}$$



AU, Hu, Khoury, PRL 109:041301(2012)[arXiv:1204.3906].

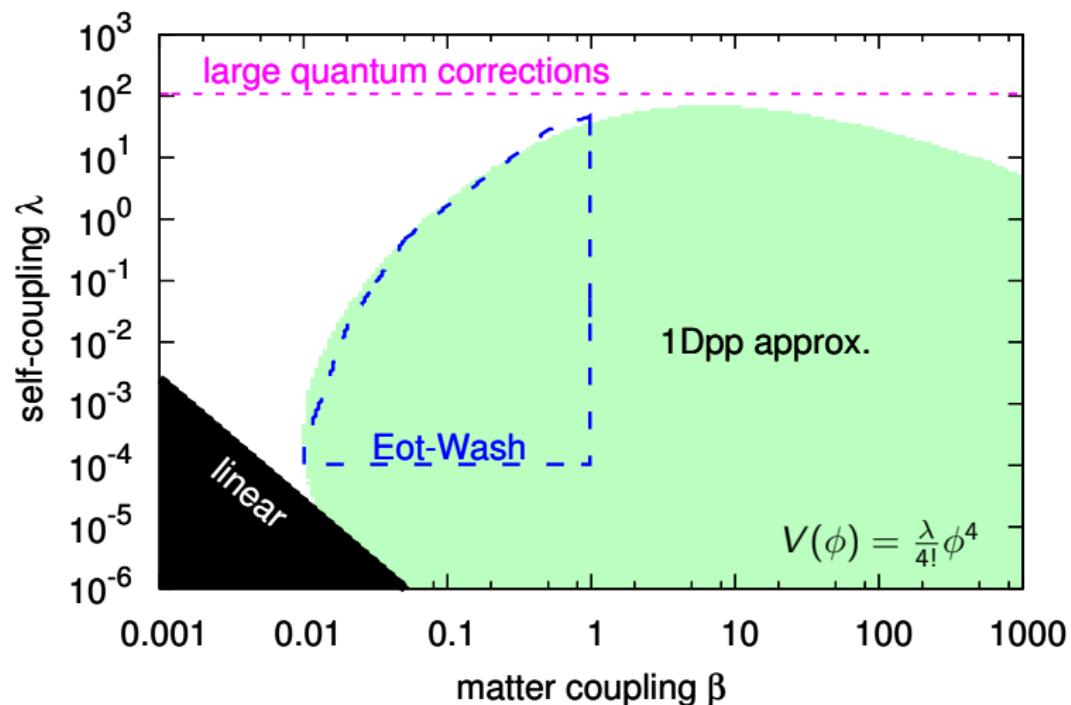
Fifth-force tests using a torsion pendulum

Eöt-Wash Experiment



<http://www.npl.washington.edu/eotwash>

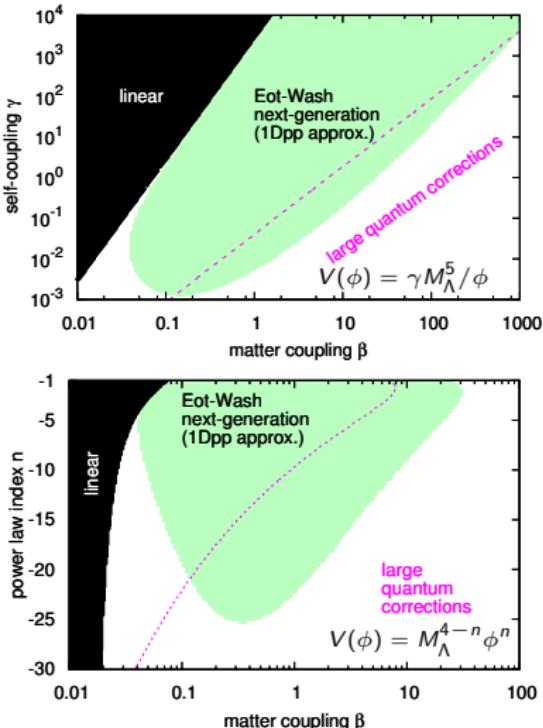
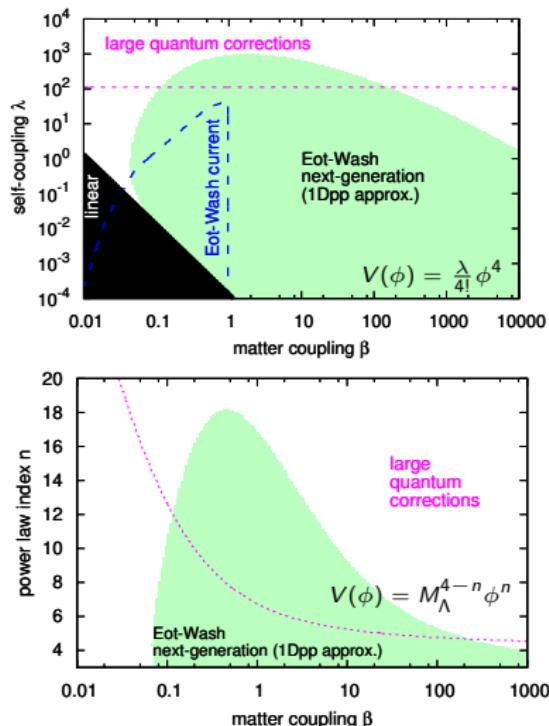
Eöt-Wash constraints on chameleons



Eöt-Wash: Adelberger, Heckel, Hoedl, Hoyle, Kapner, AU. PRL **98** 131104 (2007)

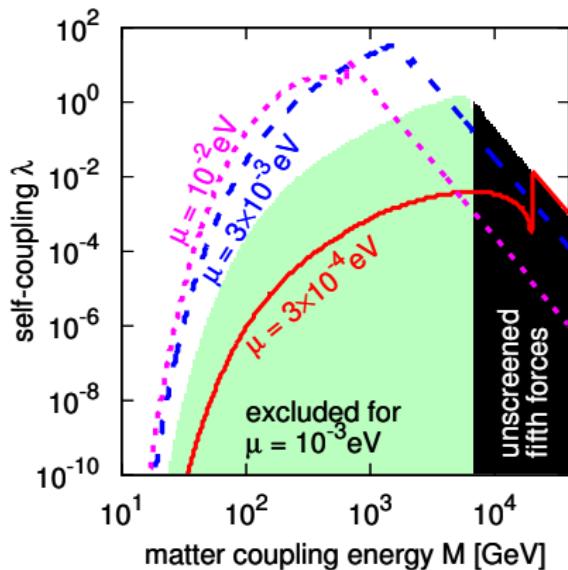
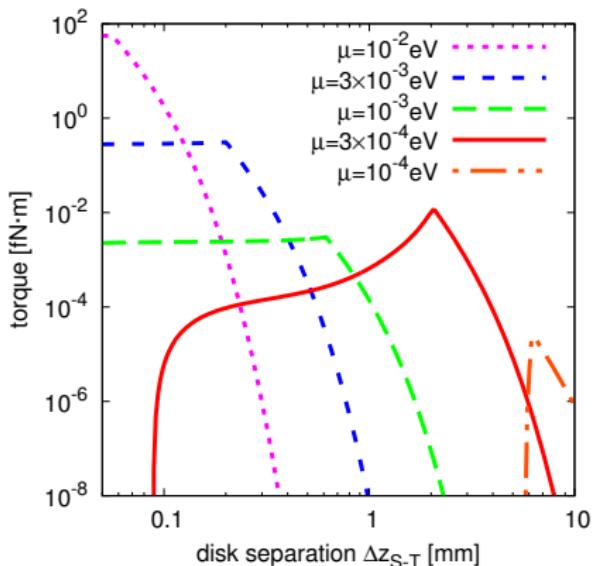
1Dpp: AU, PRD **86** 102003 (2012) [arXiv:1209.0211]

Next-generation Eöt-Wash: chameleon forecasts



AU, PRD 86:102003(2012)[arXiv:1209.0211].

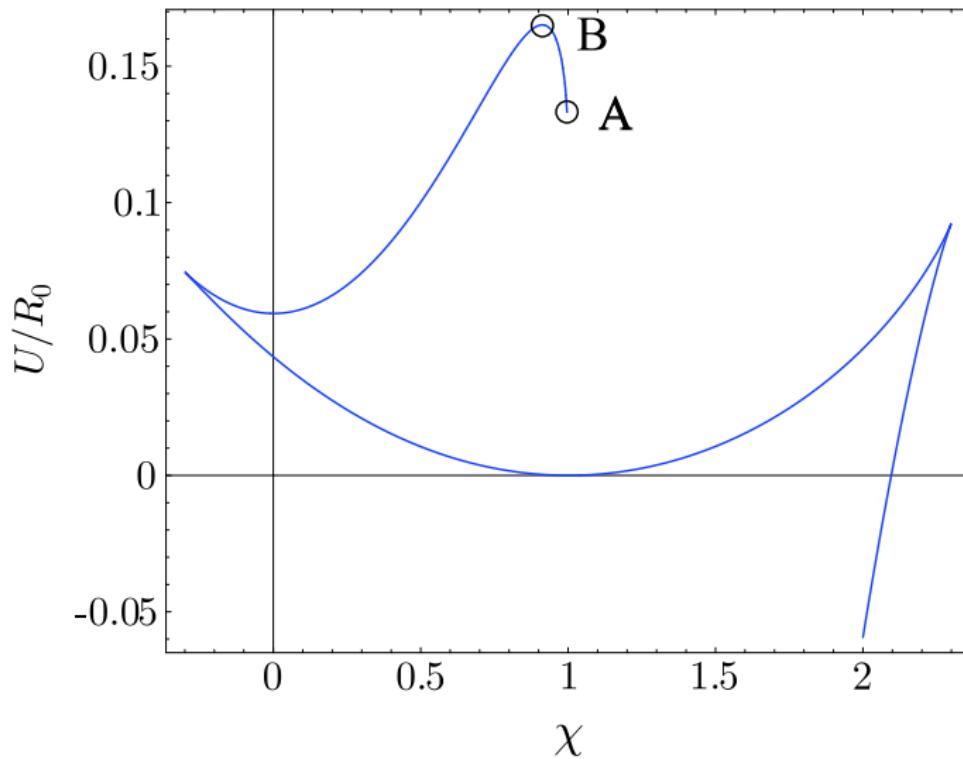
Estimated (1Dpp) Eöt-Wash constraints on symmetrons



Symmetron effective potential: $V_{\text{eff}} = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{\lambda}{4!} \phi^4$
 Eöt-Wash probes $\lambda \sim 1$, $\mu \sim 10^{-3} \text{ eV}$ (dark energy),
 $M \sim 1 \text{ TeV}$ (beyond the Standard Model)

AU, PRL 110:031301(2013)[arXiv:1210.7804].

Does screening work at all in relativistic stars?



Kobayashi and Maeda, PRD 78:064019(2008)[arXiv:0807.2503]

Chameleon screening in stars: equations of motion

metric: $ds^2 = -N(r)dt^2 + \frac{dr^2}{B(r)} + r^2(d\theta^2 + \sin^2\theta d\varphi^2)$

hydrostatic equilibrium: $P'(r) = -\frac{N'}{2N}(\rho + P)$

equation of state: $\rho(r) = \text{constant } (1\text{g/cm}^3)$

modified Einstein eq. (trace, tt , rr), $f_R = \frac{df}{dR}$, $\phi = -\frac{M_{\text{Pl}}}{2\beta_m} \log f_R$:

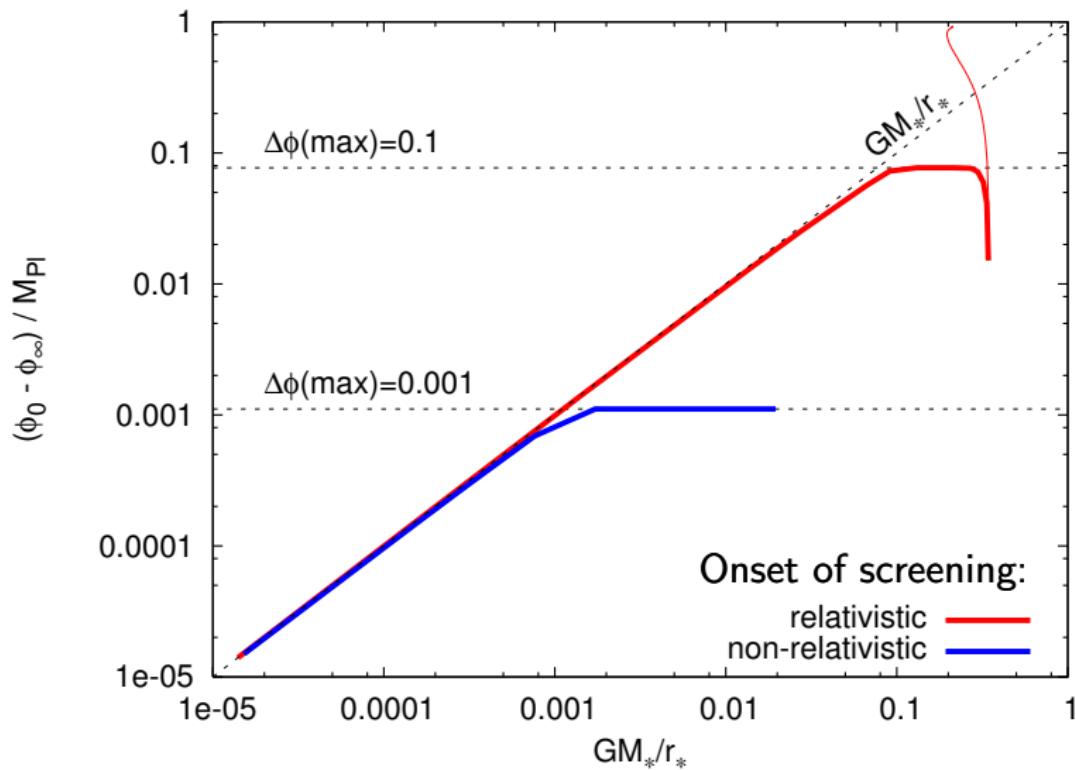
$$\left[f_R'' + \left(\frac{2}{r} + \frac{N'}{2N} + \frac{B'}{2B} \right) f_R' \right] B = \frac{dV}{df_R} - \frac{8\pi G}{3}(\rho - 3P)$$

$$\frac{(-1 + B + rB')f_R}{r^2} + \left[f_R'' + \left(\frac{2}{r} + \frac{B'}{2B} \right) f_R' \right] B = -8\pi G\rho + \frac{f - Rf_R}{2}$$

$$\frac{(-1 + B + rBN'/N)f_R}{r^2} + \left(\frac{2}{r} + \frac{N'}{2N} \right) f_R' B = 8\pi GP + \frac{f - Rf_R}{2}$$

AU and Hu, PRD 80:064002(2009)[arXiv:0905.4055].

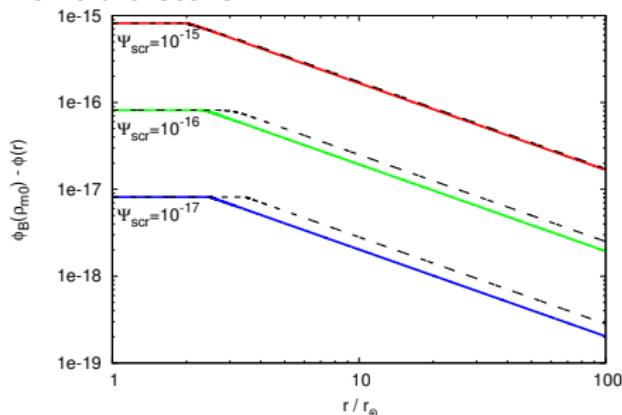
Chameleon screening in relativistic stars



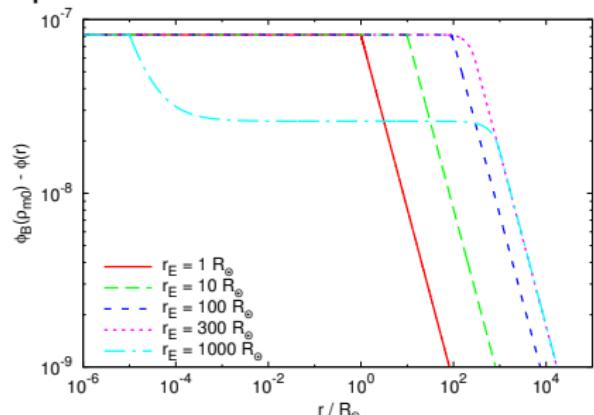
Monopole radiation in modified gravity

Screening implies that the scalar charge associated with an object is not conserved. An oscillating or expanding star can emit scalar monopole radiation.

Variable stars

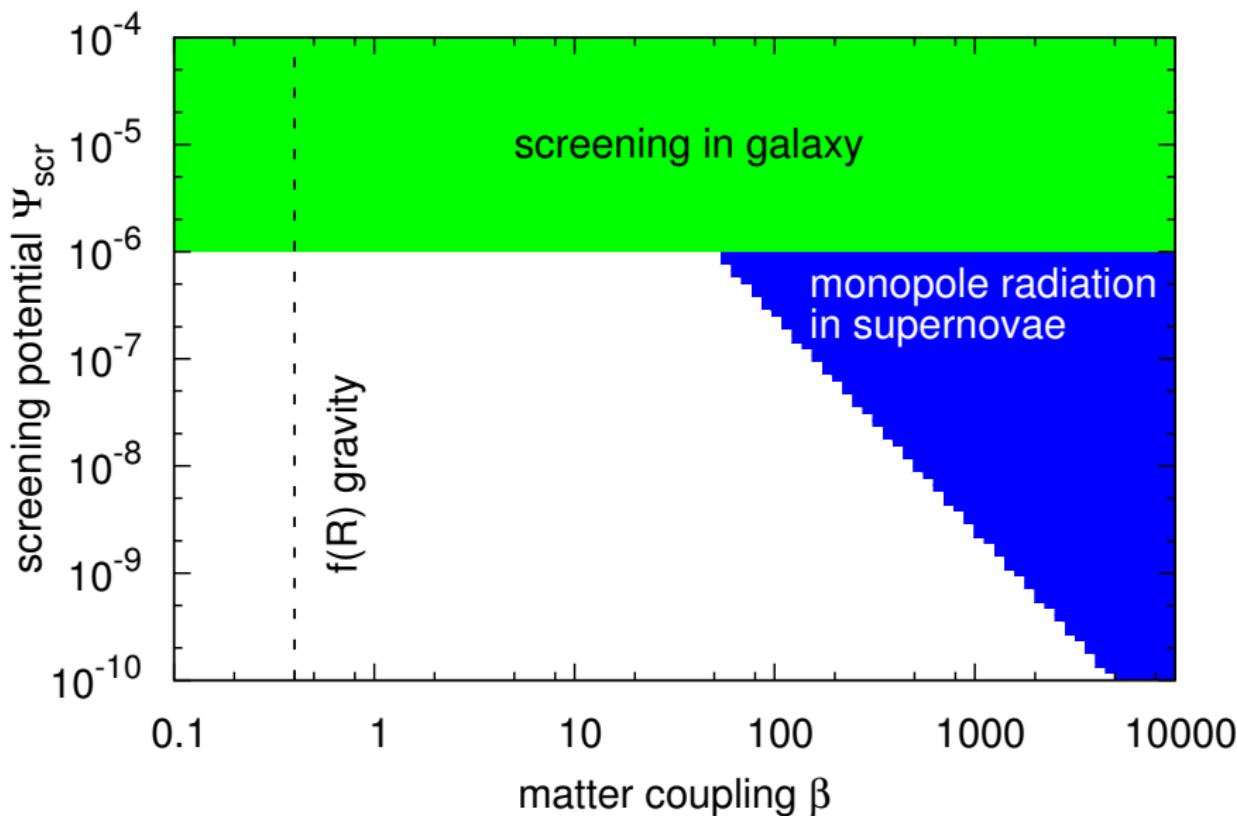


Supernovae



AU and Steffen, arXiv:1306.6113

Estimated constraints from monopole radiation

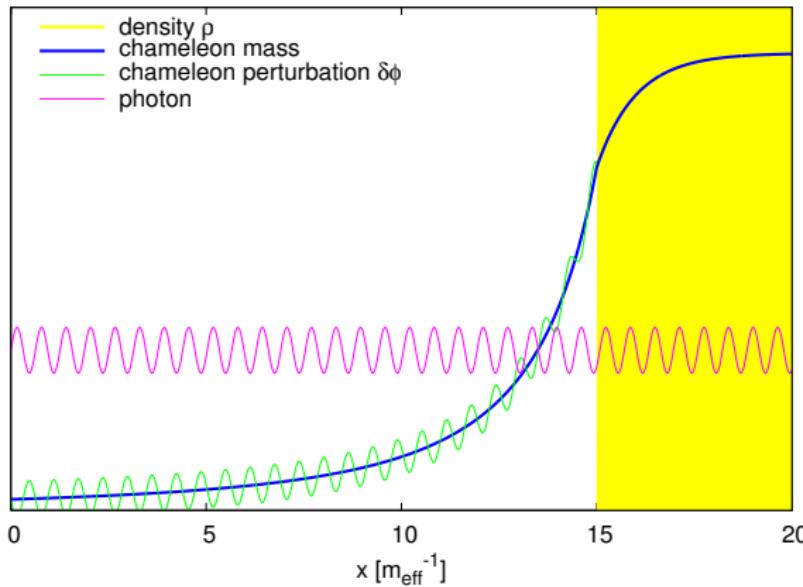


AU and Steffen, arXiv:1306.6113

How dark is dark energy? Searches for photon couplings

Oscillation: Photon coupling term $\frac{\beta\gamma}{4M_{Pl}} F_{\mu\nu} F^{\mu\nu} \phi \Rightarrow$ dark energy particles produced from photons in magnetic field

Containment: Dark energy particles reflect from matter. Windows perform quantum measurements.

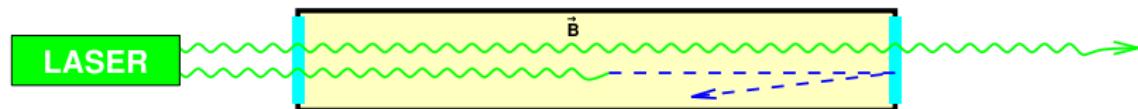


Afterglow experiments

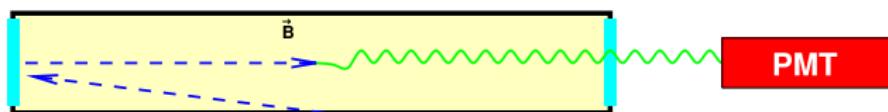
An **afterglow experiment** has two phases:

- (a) Production phase: photons streamed through \vec{B}_0 region; some oscillate into chameleons

a)



b)



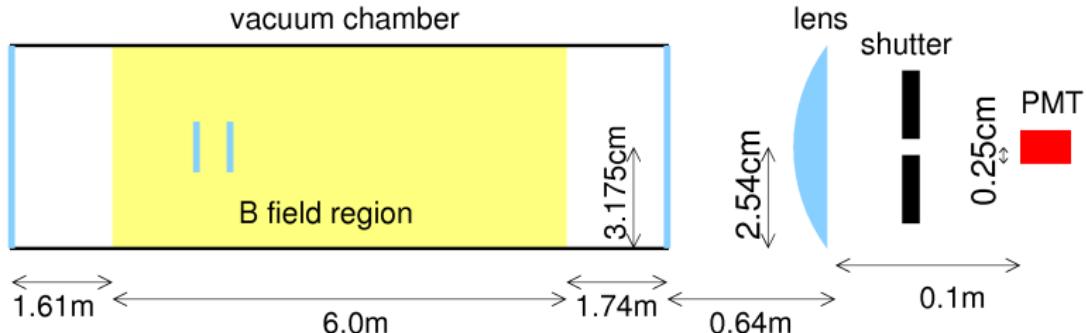
- (b) Afterglow phase: chameleons slowly oscillate back into photons, escaping chamber

Systematics: • adiabatic evolution • emission from vacuum materials

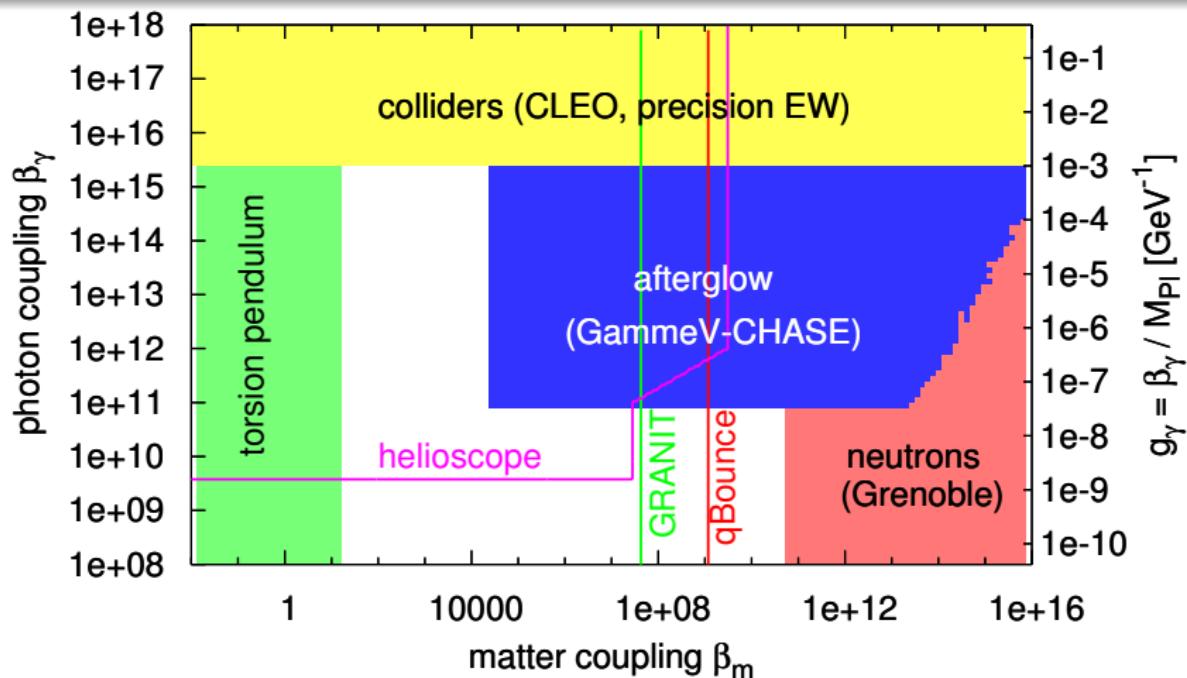
• diffuse reflection • scattering from atoms • effects of chamber geometry

Thorough review: *AU, Steffen, Chou, PRD 86:035006(2012)[arXiv:1204.5476]*..

CHASE (CHameleon Afterglow SEArch)



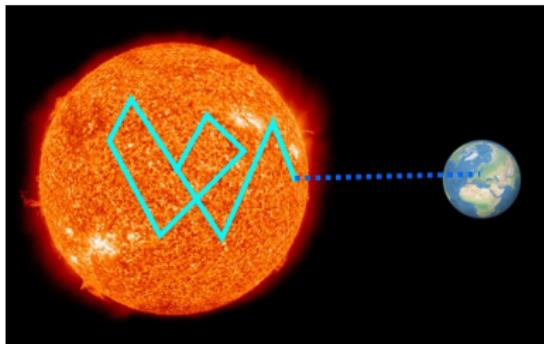
CHASE constraints on $V(\phi) = M_\Lambda^4(1 + M_\Lambda/\phi)$



Theory: AU, Steffen, Chou, PRD **86**:035006(2012)[arXiv:1204.5476].,
 AU, Steffen, Weltman, PRD **81**:015013(2010)[arXiv:0911.3906].

Experiment: Steffen, AU, Baumbaugh, Chou, Mazur, Tomlin, Weltman,
 Wester, PRL **105**:261803(2010)[arXiv:1010.0988].

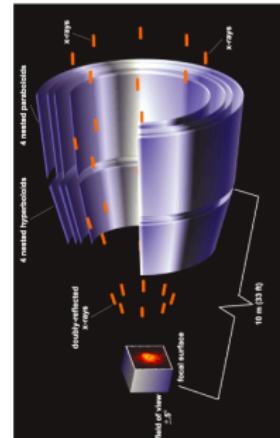
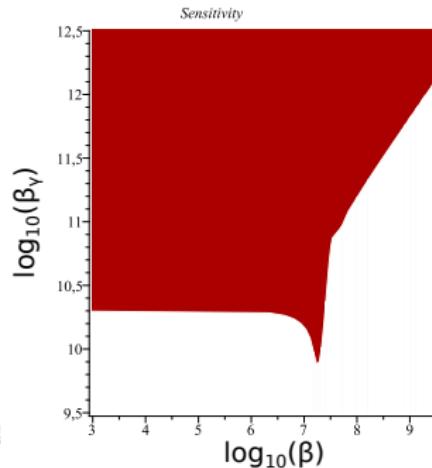
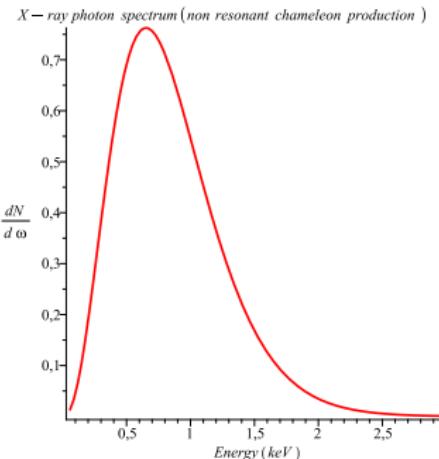
Chameleons from the Sun



- \sim keV photons oscillate into chameleons inside Sun
- chameleon particles reach Earth
- helioscope magnet regenerates photons for detection



Helioscope forecasts



Solar chameleon spectrum peaked at 600 eV.

Forecast constraints.

P. Brax, A. Lindner, K. Zioutas, PRD **85** 043014 (2012)

Increase collecting area using an X-ray mirror.

O. K. Baker, A. Lindner,
AU, K. Zioutas (2012)

Conclusions

- ① The physics responsible for the cosmic acceleration can be tested in the laboratory and stars as well as in the universe.
- ② Modifications to gravity lead to couplings to Standard Model particles.
- ③ Laboratory torsion pendulum experiments will soon rule out “quantum-stable” chameleon models and an interesting class of symmetron models.
- ④ $f(R)$ gravity is consistent with relativistic and variable stars.
- ⑤ Afterglow experiments have placed bounds on the coupling between dark energy and photons.
- ⑥ Dark energy candidate particles from the Sun could be detected by a helioscope such as CAST at CERN.