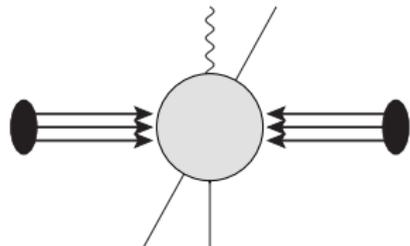


Part 4: Multiple interactions and hadronisation.

- a) Multiparton interactions
- b) Hadronisation

Back to the big picture: Some questions...

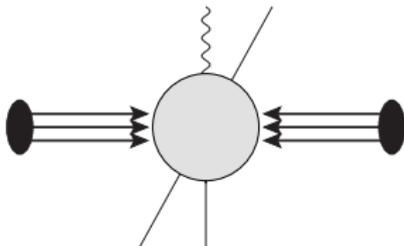
Detector event



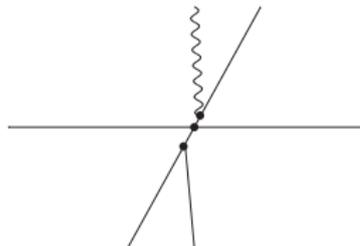
Say an event contains one boson and *three or four* jets. Where do these particles come from?

Back to the big picture: Some questions...

Detector event



Perturbative scattering



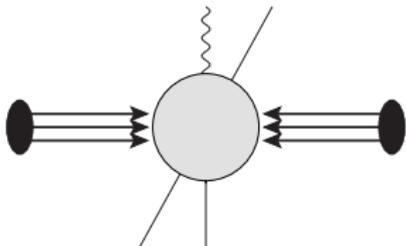
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By now, we know quite well how to get these jets by dressing a complicated hard scattering. But when does this apply?

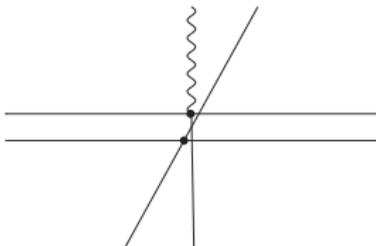
What if two jets merge? What if the boson and a jet are collinear? What if the jets have a low transverse momentum? What if pairs are back-to-back?

Back to the big picture: Some questions...

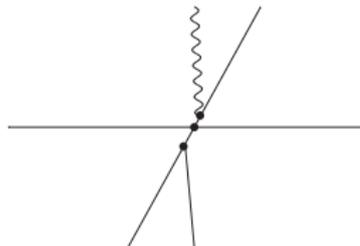
Detector event



Multiple scattering



Perturbative scattering



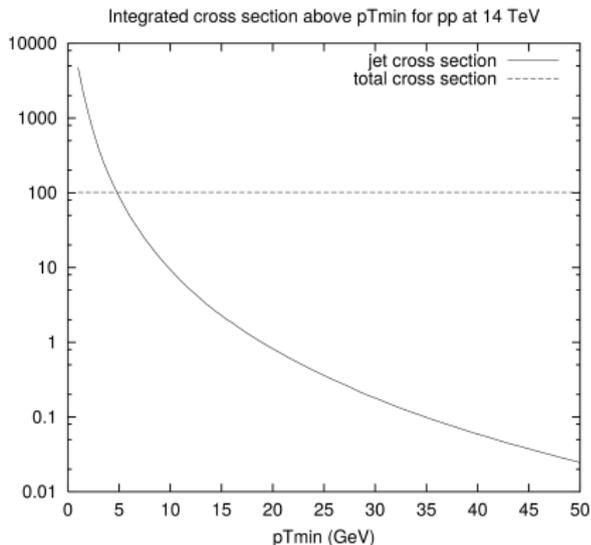
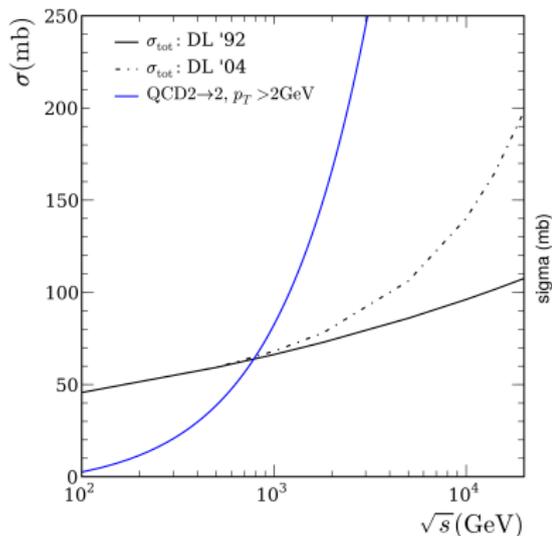
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By now, we know quite well how to get these jets by dressing a complicated hard scattering. But when does this apply?

What if two jets merge? What if the boson and a jet are collinear? What if the jets have a low transverse momentum? What if pairs are back-to-back?

When colliding composite objects, many constituent scatterings "compete" for the collision energy – **and multiple scattering can look like single complicated scatterings!**

The dijet process



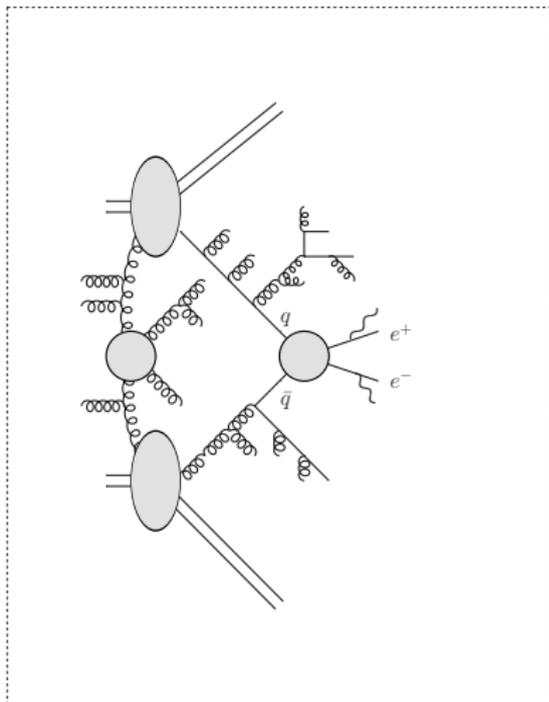
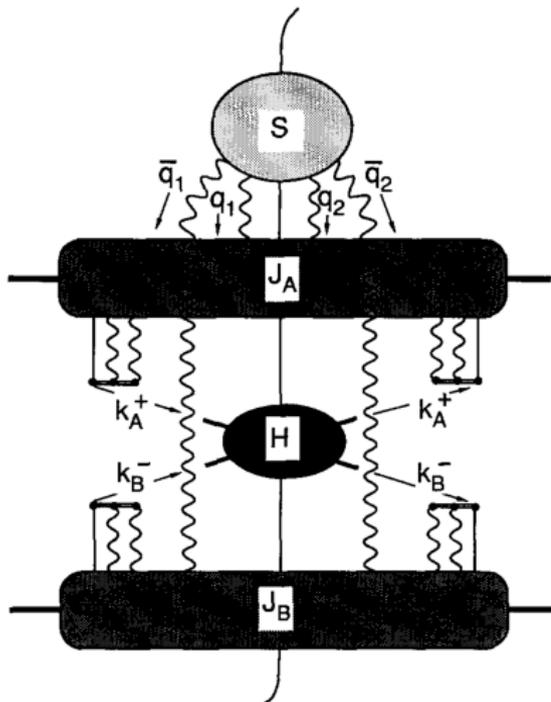
Perturbative cross section

$$\sigma(pp \rightarrow jj + X) = \int_{p_{\perp \text{min}}}^{\frac{E_{\text{cm}}}{2}} dx_1 dx_2 f_1(x_1) f_2(x_2) \frac{d\hat{\sigma}}{dp_{\perp}} dp_{\perp} > \sigma(pp \rightarrow \text{anything}) \text{ for } \frac{p_{\perp \text{min}}}{E_{\text{cm}}} \rightarrow 0$$

as $f(x)$ not small (enough) for low $x \approx \frac{p_{\perp \text{min}}}{E_{\text{cm}}}$ to suppress $\frac{p_{\perp \text{min}}}{E_{\text{cm}}} \rightarrow 0$ divergence!

Back to factorisation

Still consistent with perturbative QCD: PDFs are the *inclusive* probability to find parton at x , with all other interactions above $x \approx \frac{p_{\perp, \min}}{E_{cm}}$ integrated out!



Multiple scatterings and the inelastic cross section

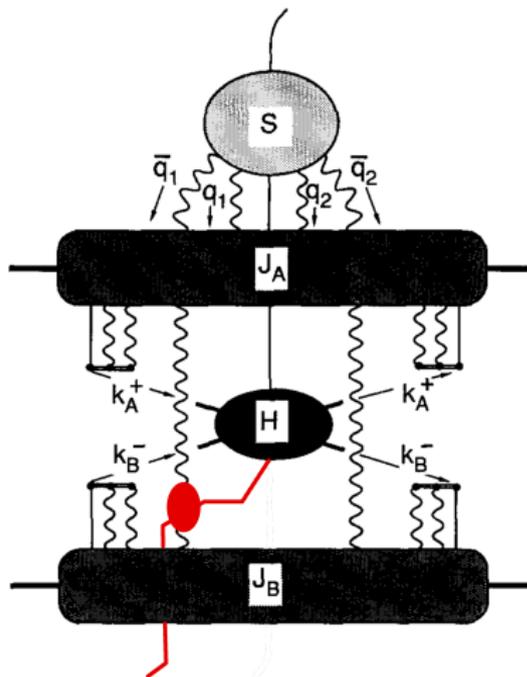
Exclusive observables (i.e. not integrating everything out) "see" these additional interactions!

⇒ An average $\langle n(p_{\perp min}) \rangle$ scatterings accompany one scattering above $p_{\perp min}$, so that

$$\sigma^{inc}(p_{\perp min}, E_{cm}) = \langle n(p_{\perp min}) \rangle \cdot \sigma^{inel}(p_{\perp min}, E_{cm})$$

where

$$\sigma^{inel} < \sigma(pp \rightarrow anything)$$



Multiple interactions

Multiple interactions between the composite protons are supported by 30 years of evidence:

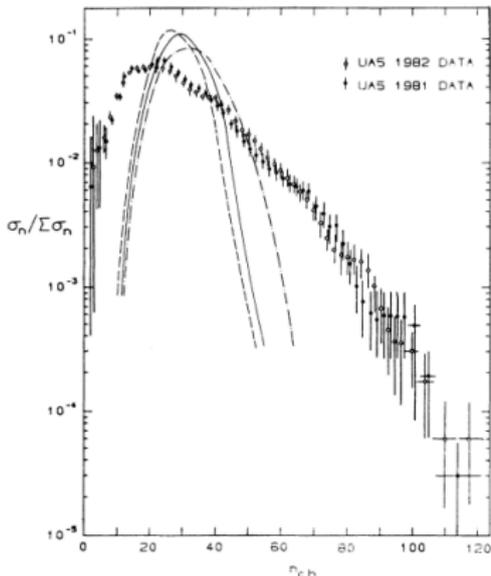


FIG. 3. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs simple models: dashed low p_T only, full including hard scatterings, dash-dotted also including initial- and final-state radiation.

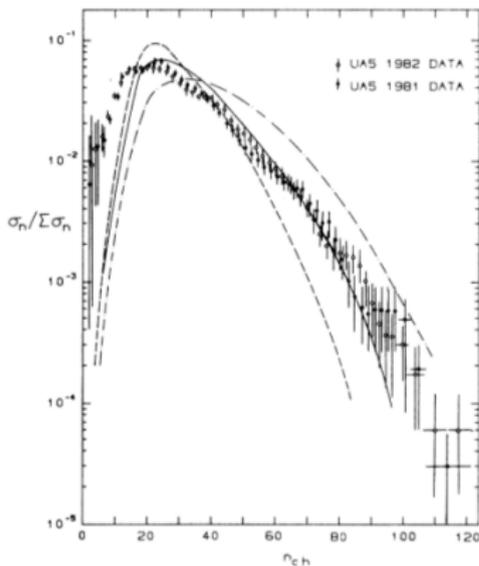


FIG. 5. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs impact-parameter-independent multiple-interaction model: dashed line, $p_{Tmin} = 2.0$ GeV; solid line, $p_{Tmin} = 1.6$ GeV; dashed-dotted line, $p_{Tmin} = 1.2$ GeV.

How do we get there?

Question: Can't we just overlay many scatterings to approximate the result? Just like we do for Pile-Up?

Answer: No!

For large p_{\perp} , model must preserve the perturbative hard scattering cross section, otherwise factorisation of inclusive cross section violated!

How do we get there?

Question: Can't we just overlay many scatterings to approximate the result? Just like we do for Pile-Up?

Answer: No!

For large p_{\perp} , model must preserve the perturbative hard scattering cross section, otherwise factorisation of inclusive cross section violated!

Solution: Subtract what you add!

For every additional scattering, we need "virtual corrections"

This should sound familiar from PS unitarity / multi-jet merging.

”Virtual corrections” for second scattering



”Virtual corrections” for second scattering

$$\left[\begin{array}{c} \text{Diagram 1} \\ \mathcal{O}(S_H) \end{array} + \int \begin{array}{c} \text{Diagram 2} \\ \mathcal{O}(S_H S_{2 \rightarrow 2}) \end{array} \right]$$

The image shows a mathematical expression enclosed in large square brackets. On the left side of the brackets is a blue Feynman diagram consisting of two incoming lines meeting at a vertex, with a wavy line extending from that vertex. Below this diagram is the label $\mathcal{O}(S_H)$. To the right of the first term is a plus sign followed by an integral symbol \int . Inside the integral is a green Feynman diagram of a bubble loop. The bubble is a horizontal oval with two vertices on each side, and four external lines (two on each side) extending outwards. The bubble itself is labeled with 'MP' in the center. Below this diagram is the label $\mathcal{O}(S_H S_{2 \rightarrow 2})$. The entire expression is enclosed in large square brackets.

"Virtual corrections" for second scattering


$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MFP} + \int \text{MFP} \mathcal{O}(S_H S_{2 \rightarrow 2}) \right]$$

The diagram shows a vertex where two blue lines meet. To the right of this vertex is a wavy line. This is followed by a large square bracket containing three terms: the first is $\mathcal{O}(S_H)$, the second is $-\mathcal{O}(S_H)$ multiplied by an integral over a grey oval labeled 'MFP' with four green lines, and the third is an integral over another 'MFP' oval multiplied by $\mathcal{O}(S_H S_{2 \rightarrow 2})$.

"Virtual corrections" for second scattering


$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MFP} + \int \text{MFP} \mathcal{O}(S_H S_{2 \rightarrow 2}) \right]$$

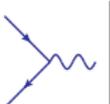
Remember PS unitarity:
First terms in a "no-scattering" factor.

"Virtual corrections" for second scattering



$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MPI} + \int \text{MPI} \mathcal{O}(S_H S_{2 \rightarrow 2}) \right]$$

Remember PS unitarity:
First terms in a "no-scattering" factor.



$$\left[\Pi^{\text{MPI}}(E_{cm}, p_{\perp \min}) \mathcal{O}(S_H) + \int \text{MPI} \Pi^{\text{MPI}}(E_{cm}, p_{\perp 1}) \mathcal{O}(S_H S_{2 \rightarrow 2}) \right]$$

"Virtual corrections" for multiple scatterings


$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MFP} + \int \text{MFP} \mathcal{O}(S_H S_{2 \rightarrow 2}) \right]$$

The diagram shows a vertex with two incoming blue lines and one outgoing wavy line. This vertex is multiplied by a bracketed expression representing virtual corrections. The expression is: $\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MFP} + \int \text{MFP} \mathcal{O}(S_H S_{2 \rightarrow 2}) \right]$. The first term is $\mathcal{O}(S_H)$. The second term is $\mathcal{O}(S_H)$ multiplied by an integral over a self-energy loop labeled "MFP" (Multi-Particle Function), which consists of two green lines and a wavy line. The third term is an integral over a self-energy loop labeled "MFP" multiplied by $\mathcal{O}(S_H S_{2 \rightarrow 2})$.

"Virtual corrections" for multiple scatterings



$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MPP} + \int \text{MPP} \otimes \left\{ \mathcal{O}(S_H S_{2 \rightarrow 2}) \right. \right. \\
 \left. \left. + \int \text{MPP} \mathcal{O}(S_H S_{2 \rightarrow 2} S'_{2 \rightarrow 2}) \right\} \right]$$

The diagram illustrates the structure of virtual corrections for multiple scatterings. It features a vertex on the left with two incoming blue lines and one outgoing wavy line. This vertex is enclosed in a large dashed box. To the right of the vertex, the expression is enclosed in large square brackets. The expression consists of several terms: a leading $\mathcal{O}(S_H)$ term, followed by a subtraction of $\mathcal{O}(S_H)$ multiplied by an integral of a self-energy-like diagram (labeled 'MPP' in a grey oval) with two green external lines. This is followed by a plus sign and another integral of a similar 'MPP' diagram. A tensor product symbol \otimes is then used with a curly bracket containing $\mathcal{O}(S_H S_{2 \rightarrow 2})$. Finally, a plus sign is followed by an integral of a 'MPP' diagram multiplied by $\mathcal{O}(S_H S_{2 \rightarrow 2} S'_{2 \rightarrow 2})$. The entire expression is enclosed in large square brackets.

”Virtual corrections” for multiple scatterings



$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MPI} + \int \text{MPI} \right] \otimes \left\{ \mathcal{O}(S_H S_{2 \rightarrow 2}) - \mathcal{O}(S_H S_{2 \rightarrow 2}) \int \text{MPI} + \int \text{MPI} \mathcal{O}(S_H S_{2 \rightarrow 2} S'_{2 \rightarrow 2}) \right\}$$

The diagram illustrates the structure of virtual corrections for multiple scatterings. It consists of two main parts enclosed in large square brackets. The first part is a sum of three terms: a tree-level operator $\mathcal{O}(S_H)$, a loop correction term $-\mathcal{O}(S_H) \int \text{MPI}$ (where MPI is a grey oval with two green lines), and a self-energy correction term $+\int \text{MPI}$. The second part, enclosed in curly braces and multiplied by the first part, represents higher-order corrections: $\mathcal{O}(S_H S_{2 \rightarrow 2}) - \mathcal{O}(S_H S_{2 \rightarrow 2}) \int \text{MPI} + \int \text{MPI} \mathcal{O}(S_H S_{2 \rightarrow 2} S'_{2 \rightarrow 2})$.

"Virtual corrections" for multiple scatterings

Remember PS unitarity:
First terms in a "no-scattering" factor.

$$\left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MPI} + \int \text{MPI} \right] \otimes \left\{ \mathcal{O}(S_H S_{2 \to 2}) - \mathcal{O}(S_H S_{2 \to 2}) \int \text{MPI} + \int \text{MPI} \mathcal{O}(S_H S_{2 \to 2} S'_{2 \to 2}) \right\}$$

"Virtual corrections" for multiple scatterings

Remember PS unitarity:
First terms in a "no-scattering" factor.

$$\begin{aligned}
 & \left[\mathcal{O}(S_H) - \mathcal{O}(S_H) \int \text{MPI} + \int \text{MPI} \right] \\
 & \otimes \left\{ \mathcal{O}(S_H S_{2 \rightarrow 2}) - \mathcal{O}(S_H S_{2 \rightarrow 2}) \int \text{MPI} + \int \text{MPI} \mathcal{O}(S_H S_{2 \rightarrow 2} S'_{2 \rightarrow 2}) \right\} \\
 & \left[\Pi^{\text{MPI}}(E_{cm}, p_{\perp min}) \mathcal{O}(S_H) + \int \text{MPI} \right] \\
 & \otimes \left\{ \Pi^{\text{MPI}}(p_{\perp 1}, p_{\perp min}) \mathcal{O}(S_H S_{2 \rightarrow 2}) + \int \text{MPI} \Pi^{\text{MPI}}(p_{\perp 1}, p_{\perp 2}) \mathcal{O}(S_H S_{2 \rightarrow 2} S'_{2 \rightarrow 2}) \right\}
 \end{aligned}$$

If you have a hammer...

...everything looks like a parton shower. Assume

$\delta p_{\perp} \langle n(p_{\perp}) \rangle \equiv$ Probability for scattering with $p_{\perp} \in [p_{\perp min}, p_{\perp min} + \delta p_{\perp}]$.

Then the probability of no scattering is

$$1 - \delta p_{\perp} \langle n(p_{\perp}) \rangle$$

or, if δp_{\perp} is divided into m parts, and the scattering probabilities are independent

$$[1 - \delta p_{\perp} / m \langle n(p_{\perp}) \rangle]^m \xrightarrow{m \rightarrow \infty} \exp \left(- \int_{p_{\perp min}}^{p_{\perp min} + \delta p_{\perp}} dp_{\perp} \langle n(p_{\perp}) \rangle \right) \equiv \Pi^{\text{MPI}}(p_{\perp min} + \delta p_{\perp}, p_{\perp min})$$

We can define a no-additional scattering probability which contains "all-order virtual corrections" – just like a parton shower Sudakov factor.

\implies Can recycle the PS algorithm to produce additional scatterings.

Exercise: The inclusive scattering probability

The probability of a hardest scattering at $p_{\perp 1}$ is

$$\Pi^{\text{MPI}}(p_{\perp 0}, p_{\perp 1}) \langle n(p_{\perp 1}) \rangle$$

The probability of having a second hardest scattering at $p_{\perp 2} < p_{\perp 1}$ is

$$\int_{p_{\perp 2}}^{p_{\perp 0}} dp_{\perp 1} \Pi^{\text{MPI}}(p_{\perp 0}, p_{\perp 1}) \langle n(p_{\perp 1}) \rangle \Pi^{\text{MPI}}(p_{\perp 1}, p_{\perp 2}) \langle n(p_{\perp 2}) \rangle$$

Show that the probability of having any partonic scattering (provided the protons scatter) is given by σ^{inc} , i.e. by the perturbative result!

Hints: Look at n scatterings with $p_{\perp n} < p_{\perp n-1} < \dots < p_{\perp 0}$, use the properties of the exponential functions, you can find a short form for nested integrals by finding a differential equation for the sum of all nested integrals, and thinking about solutions to linear differential equations.

Parameters of Multiparton Interaction models

⇒ Perturbative MPI *model* keeping the inc. cross section. Unknowns:

- MPI probability $\langle n(p_{\perp}) \rangle = \frac{\sigma^{inc}}{\sigma^{inel}}$ with σ^{inel} taken from data (tuned)
- Most MPI very soft, but σ^{inc} is still divergent for $p_{\perp min} \rightarrow 0$, i.e. needs regulator ⇒ Extra parameter $p_{\perp 0}$
- Regulator should be larger if E_{cm} becomes larger (to not violate the total cross section) ⇒ Parameters for energy scaling of $p_{\perp 0}$
- ...and some technical parameters specific to implementation.

Current MPI models are much more complicated and differ significantly:

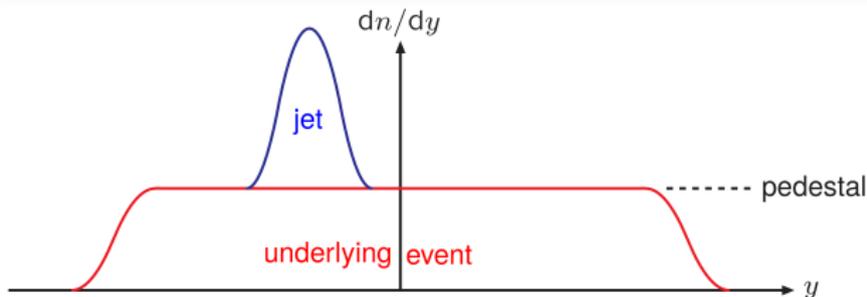
HERWIG: Pick interactions prior to running according to Poissonian,

SHERPA: MPI after hard process evolution in a p_{\perp} -ordered sequence.

PYTHIA: MPI + ISR + FSR combined in one single p_{\perp} sequence.

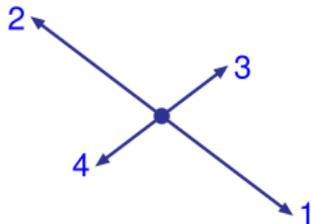
Note: For a complete picture of the *total* cross section, MPI supplemented by non-perturbative dynamics (diffractive physics)

What MPI does for (to?) you



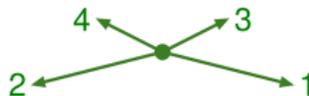
Multiple interactions model the "underlying event" that is present in any hard scattering event. All hadron collider measurement can be sensitive to MPI. However, MPI can be assessed because of its typical kinematics:

Double Parton Scattering



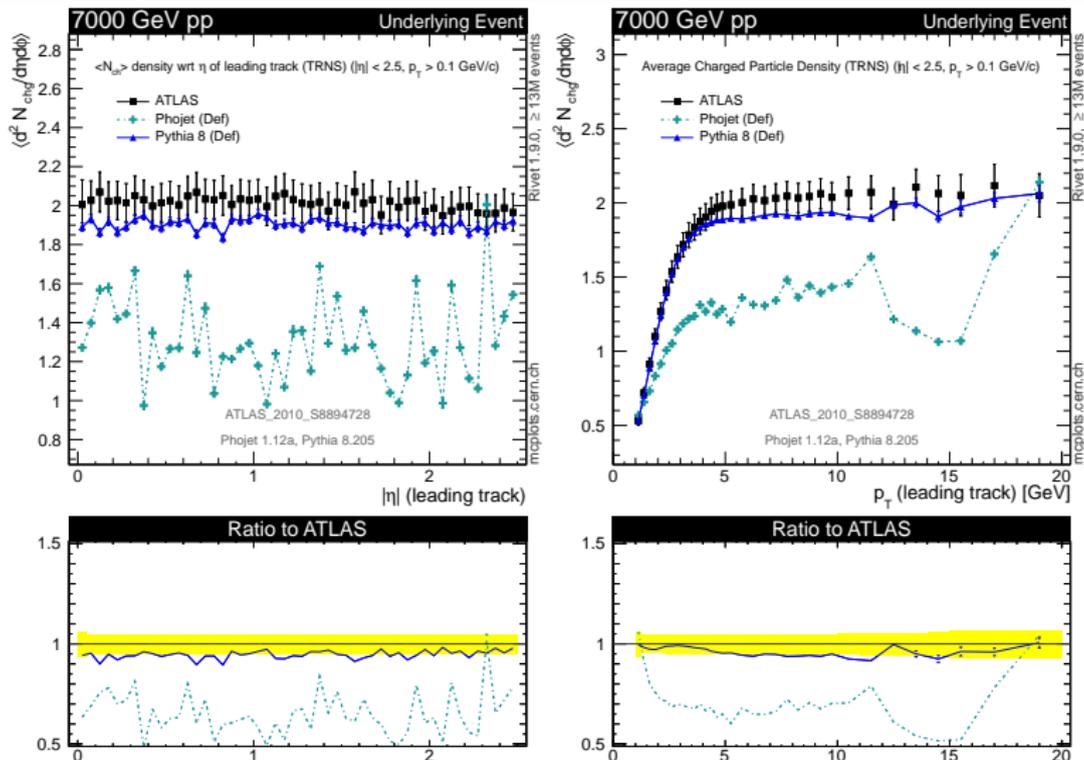
$$|\mathbf{p}_{\perp 1} + \mathbf{p}_{\perp 2}| \approx 0$$
$$|\mathbf{p}_{\perp 3} + \mathbf{p}_{\perp 4}| \approx 0$$

Double BremsStrahlung



$$|\mathbf{p}_{\perp 1} + \mathbf{p}_{\perp 2}| \gg 0$$
$$|\mathbf{p}_{\perp 3} + \mathbf{p}_{\perp 4}| \gg 0$$

MPI "perpendicular" to the hard scattering



Activity uniform in rapidity. More particles for harder scatterings \Rightarrow Trigger bias, and harder collisions more "central" \Rightarrow MPI have impact parameter dependence. PHOJET not MPI, but based on Pomeron picture.

Hadronisation

However, our result still contains coloured partons.

⇒ Need to convert to hadrons! Two prescriptions in use:

Cluster

Form hadrons by decaying "preconfined" colourless clusters of partons.

Gluons split non-perturbatively to $q\bar{q}$

Many-parameter energy-momentum structure.

Few-parameter flavour chemistry.

Used in HERWIG, SHERPA

String

Colour flux tubes (strings, junctions) between partons break to form hadrons.

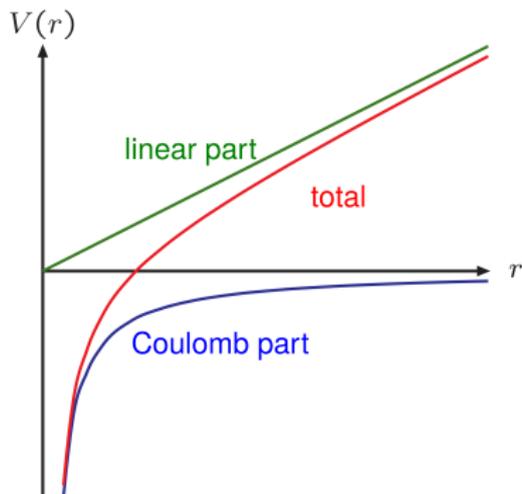
Gluons are kink on string.

Few-parameter energy-momentum structure.

Many-parameter flavour chemistry.

Used in PYTHIA, EPOS (?)

The interquark potential

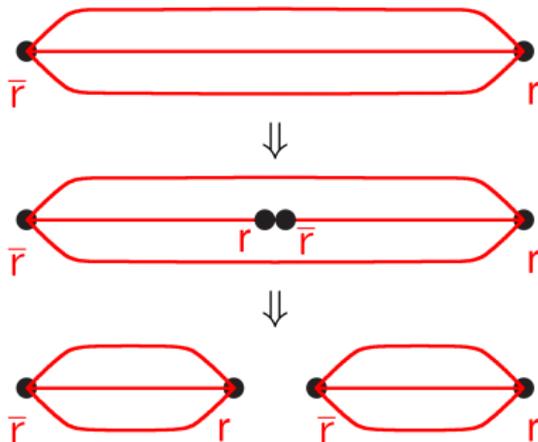
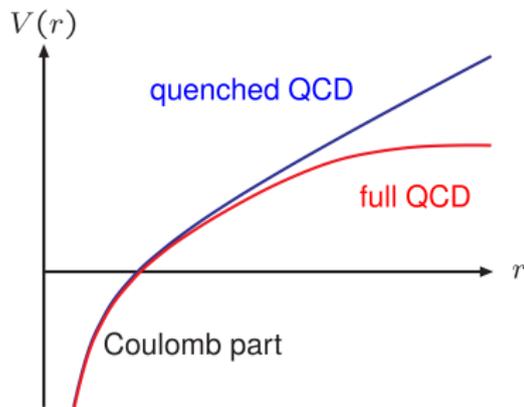


Potential between two quarks assumed linear

⇒ Constant force per unit length (just like a string / flux tube)

⇒ Confining force.

String model



Mesons have yo-yo modes while strings break before yo-yo point.

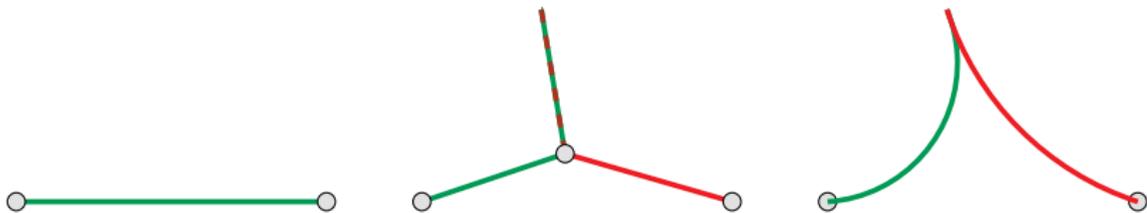
⇒ Linear potential flattens off.

⇒ Breaking gives back-to-back particle production in string CM frame.

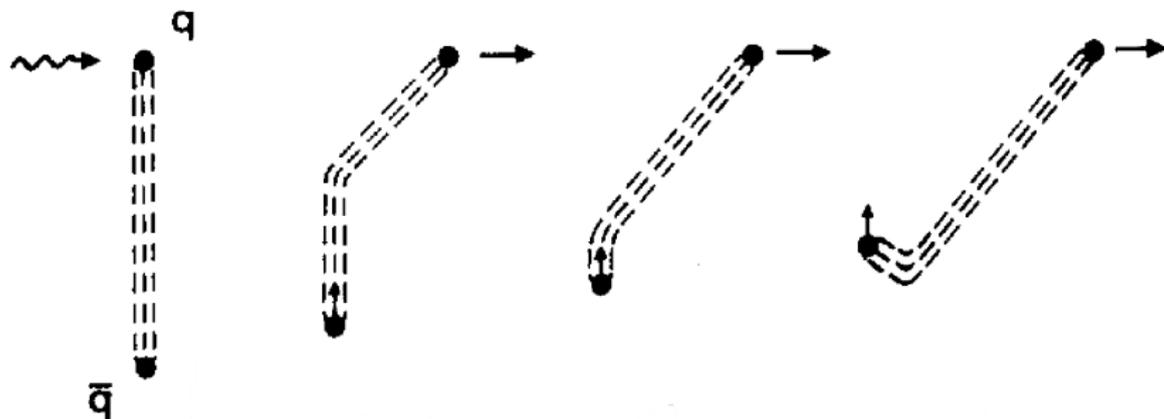
What happens to the gluons?

But we don't only have quarks! Three possibilities when adding gluons:

- Singlett: Gluon does not change colour field. Very unlikely.
- Junction: Gluon is new type string, attached to old string in a junction. Needs new parameters.
- Kink: Exists on massless relativistic string. No extra parameters.



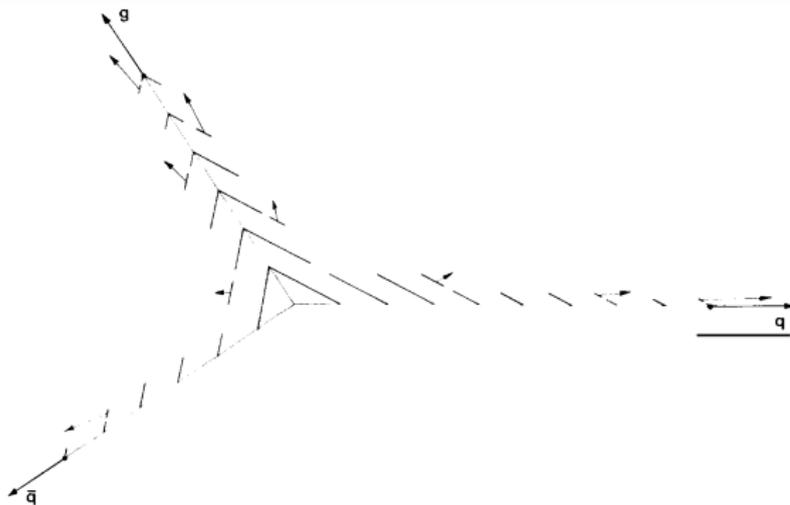
The gluon as kink on a string



What is a kink?

- Large, instantaneous momentum transfer at initial time
→ Stretches string in one direction.
- Kink is connected to two string segments
→ Loses energy twice as fast as "endpoint quarks", like gluon ($C_A/C_F = 9/4$ ($N = 3$), 2 ($N \rightarrow \infty$))
- By causality, string segments fragment as before. String + Kink system fragments as any other string would.

String effect



The addition of gluons leads to the string effect:

Gluon kink drags string along, while decoupling the quark pair
⇒ Hadron production along qg and $g\bar{q}$ strings.

⇒ 4 hadron production regions, two of them from the "kink" end.

⇒ 3 jets in event frame, almost no hadrons opposite of kink.

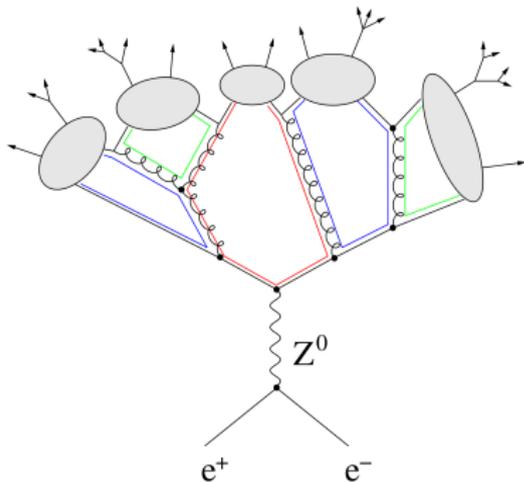
⇒ Dynamical coherence effect!

Coherence and the cluster model

- Note: Color coherence prevents gluon production at comparable angles!
⇒ Approximates string effect at perturbative level.
⇒ Can get away with simpler non-perturbative model?

Cluster model:

- ◇ Use perturbative calculation that preserves coherence.
- ◇ Convert gluons to quarks non-perturbatively.
- ◇ Collect quarks into colour singlet *preconfined* "clusters".
- ◇ Clusters decay isotropically into two hadrons.
- ◇ Heavy clusters need to be split in string-like fashion.



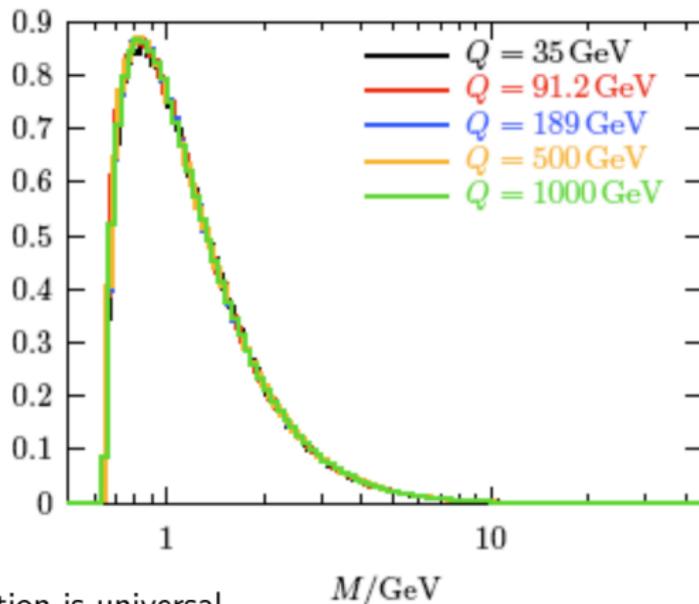
Cluster mass distribution

The main motivation of the cluster model is colour preconfinement.

Following the flow of colour in the $N_c \rightarrow \infty$ limit, we find:

1. Colour-singlet parton pairs are close in phase space.
2. Mass of singlet clusters almost independent of hard scattering scale.

Primary Light Clusters

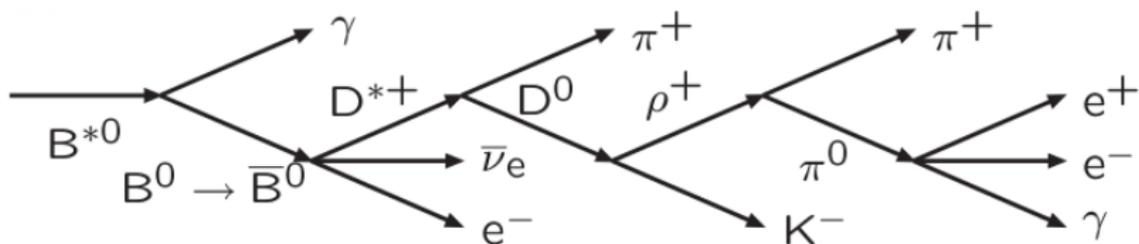


\Rightarrow Hadronisation is universal

\Rightarrow Fix parameters at LEP, then "predict" at another collider.

Hadron decays

But we're still not there yet! Fragmentation can produce excited hadrons, which will then decay. For example



Most particles are produced in this part.

- ⇒ Process has to be modelled for the correct jet structure by
- ...Hadronic matrix elements for some (important) decays.
- ...PDG decay tables for others. If tables are incomplete, be creative.

Summary of Part 4: Soft physics and hadronisation

- Prediction incomplete before assessing/including non-pert. effects.
- Soft physics:
 - $2 \text{ partons} \rightarrow 2 \text{ partons}$ cross section naively exceeds total cross section.
 - Factorisation hints this may be due to additional scatterings.
 - Multiparton interaction models attempt to describe hadron-collider data by "resolving" these additional scatterings.
 - Models fulfill some consistency conditions, e.g. should yield the inclusive cross section when additional scattering are integrated out.
 - Models come with a handful of parameters.
- Hadronisation:
 - A complete event generation needs to convert partons to hadrons.
 - Two generic models exist: The string model and the cluster model.
 - In Lund string model, the constant color field between quarks is split into smaller pieces by string breaking. Includes coherence effects non-perturbatively.
 - In cluster model, quarks form preconfined color singlet clusters which decay isotropically. The perturbative inputs must include coherence.
 - Hadronisation produces excited hadrons which have to be decayed.

Summary of Event Generator Lectures

- To make the most of (collider) data, we want an accurate representation of our physics model
⇒ Event generation.
- Event generation is split up into handy bits:
 - Hard cross section** (perturbative)
 - Parton shower resummation** (perturbative)
 - Multiparton interactions** (non-perturbative, not factorisable)
 - Hadronisation** (non-perturbative, factorisable)
- Parton shower resummation relies on probabilities to produce all-order results, but only for soft / collinear emissions.
- Many improvements in place for accurate jet modelling – by combining with many precise fixed order calculations.
- No first-principle results for non-perturbative components, but still important feature. Modelling requires physics insight!

References

Factorisation and soft gluons:

Collins, Soper, Sterman (Nucl. Phys. B 308 (1988) 833)

The book: Collins, Perturbative Quantum Chromodynamics

Multiparton interactions:

The original article: Sjostrand, van Zijl (PRD D36 (1987) 2019)

A good introduction to the HERWIG model: M. Bähr, Underlying Event Simulation in the Herwig++ Event Generator, Dissertation, ITP Karlsruhe (see https://www.itp.kit.edu/prep/phd/PSFiles/Diss_Baehr.pdf)

Hadronisation:

The Lund string model: Andersson, Gustafson, Ingelman, Sjostrand (Phys.Rept. 97 (1983) 31)

Cluster model: Webber (Nucl. Phys. B 238 (1984) 492)