**Neutrino Astrophysics** 

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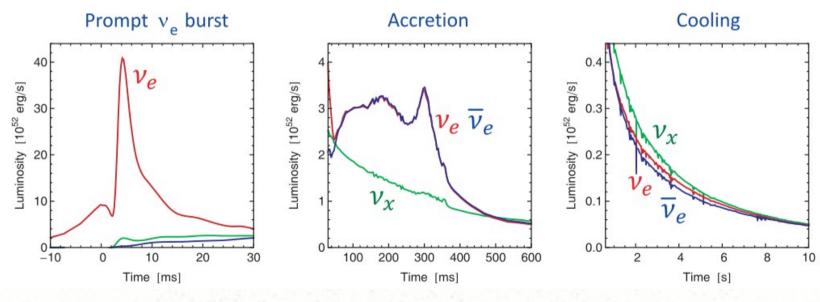
INSS, 08/2015

#### • Outline of the talk

- Supernova Neutrino Detection
- Diffuse Supernova Neutrinos
- High-Energy Neutrinos

#### <u>Supernova Neutrino Emission</u>

- strong emission of neutrinos from colapsing stars

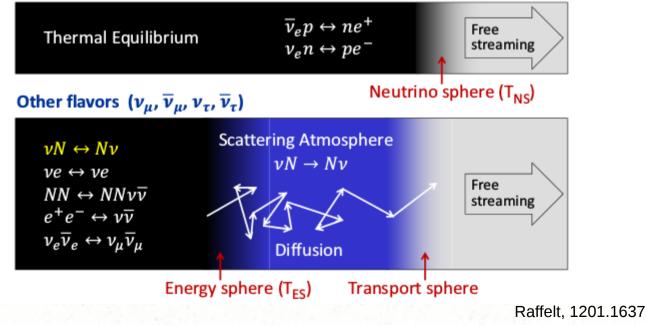


Raffelt, 1201.1637

#### <u>Supernova Neutrino Emission</u>

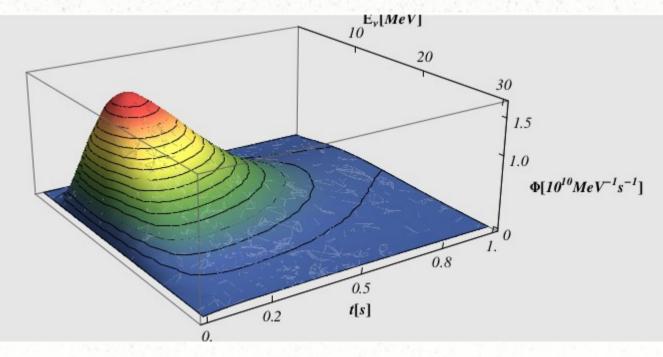
- close to thermal emission, with  $\langle E_{\nu_e} \rangle < \langle E_{\overline{\nu}_e} \rangle < \langle E_{\nu_r} \rangle$ 

#### Electron flavor ( $\nu_e$ and $\overline{\nu}_e$ )



#### <u>Supernova Neutrino Emission</u>

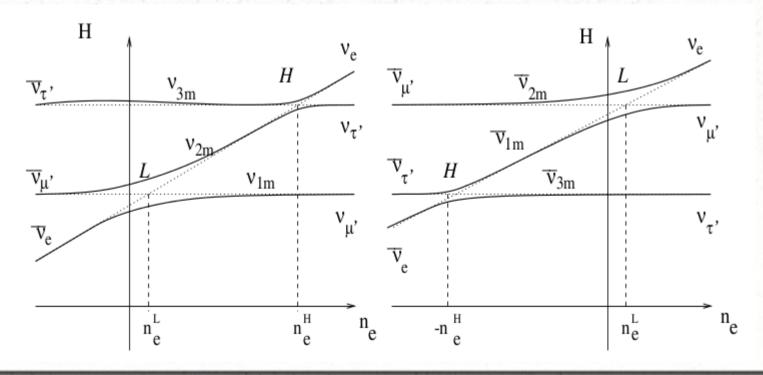
- importance features on time and energy structure!



from Vissani's talk, Paris 2011.

#### <u>Supernova Neutrino Conversion</u>

- like solar neutrinos, supernova neutrinos are produced in a dense medium, in flavour eigenstates that relates to mass eigenstates through proper diagonalization of the Hamiltonian.



normal inverted

#### <u>Supernova Neutrino Conversion</u>

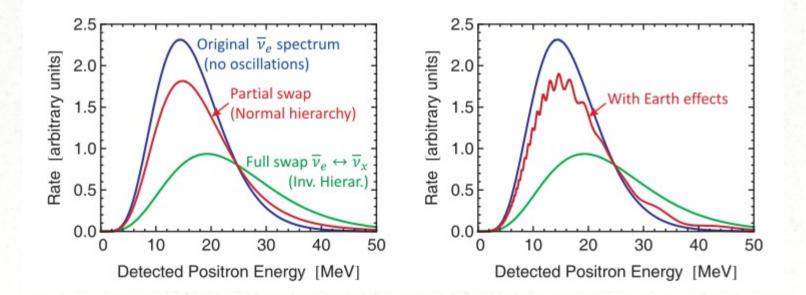
- survival probability depends on hierarchy

Scenario	Hierarchy	p	$ar{p}$	Earth effects
A	Normal	0	$\cos^2 \theta_{12}$	$\bar{\nu}_e$
В	Inverted	$\sin^2 \theta_{12}$	0	$ u_e$

- large fraction of electronic anti-neutrinos flux with normal hierarchy.

#### <u>Supernova Neutrino Conversion</u>

Detection rate:



But things are more complicated:

- very dense medium, flavour freeze, neutrino-neutrino interaction

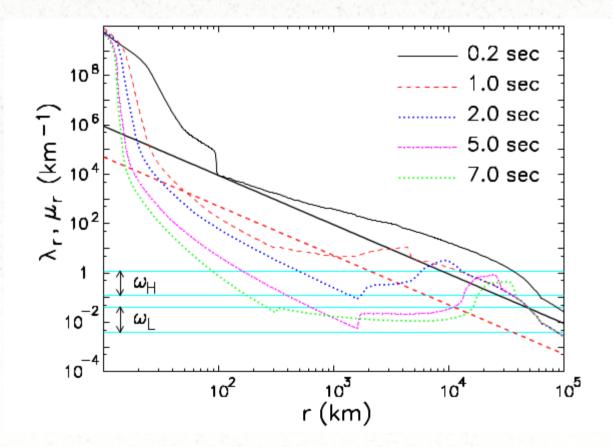
$$\varrho_{\mathbf{p},\mathbf{x}} = \begin{pmatrix} \varrho_{ee} & \varrho_{e\mu} & \varrho_{e\tau} \\ \varrho_{\mu e} & \varrho_{\mu\mu} & \varrho_{\mu\tau} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \end{pmatrix}$$

$$\mathbf{v}_{\mathbf{p}} \cdot \nabla_{\mathbf{x}} \varrho_{\mathbf{p},\mathbf{x}} = -i[\Omega_{\mathbf{p}}^{\text{vac}}, \varrho_{\mathbf{p},\mathbf{x}}] - i[\Omega_{\mathbf{p},\mathbf{x}}^{\text{ref}}, \varrho_{\mathbf{p},\mathbf{x}}]$$
$$\mathbf{v}_{\mathbf{p}} \cdot \nabla_{\mathbf{x}} \overline{\varrho}_{\mathbf{p},\mathbf{x}} = +i[\Omega_{\mathbf{p}}^{\text{vac}}, \overline{\varrho}_{\mathbf{p},\mathbf{x}}] - i[\Omega_{\mathbf{p},\mathbf{x}}^{\text{ref}}, \overline{\varrho}_{\mathbf{p},\mathbf{x}}]$$

$$\Omega_{\mathbf{p}}^{\text{vac}} = U \frac{M^2}{2p} U^{\dagger}$$
$$M^2 = \text{diag}\left(m_1^2, m_2^2, m_3^2\right) = \left(-\frac{\delta m^2}{2}, +\frac{\delta m^2}{2}, \pm \Delta m^2\right)$$

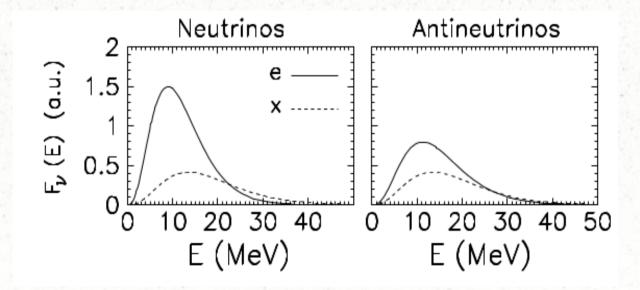
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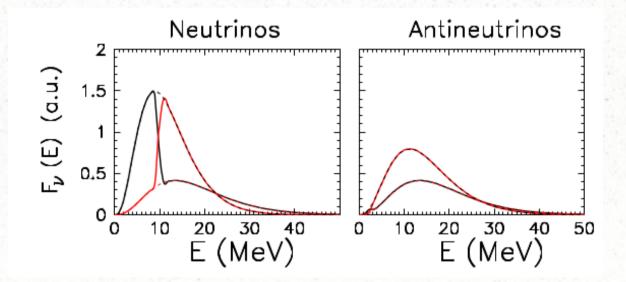
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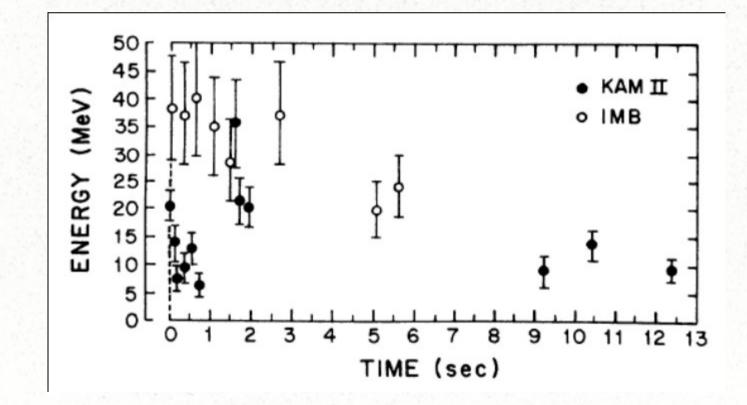


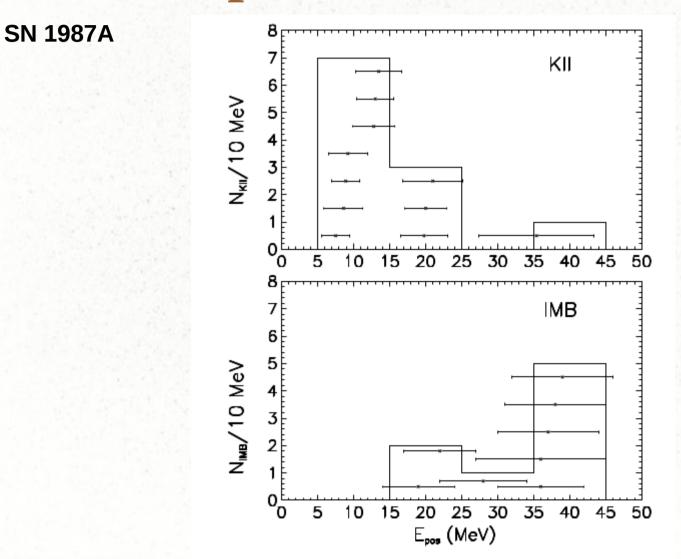
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SN 1987A: a succesfull case for Supernova Neutrino detection





Next Galactic Supernova will provide great statistics, > 1000 articles in the first month, dedicated conferences in nice places and a lot of learning opportunities.

Fortunately, it will happen in our lifetime.

- Why waiting for a gallactic SN, if somewhere else in the Universe a SN is exploding every second?

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 $1/r^{2}$ 

worse then that, we have to account redshift and space expansion!

- But what if we integrate the signal of all neutrinos from SN explosions?

Related question: Olbers' paradox.

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Why is the night dark?

$$L(t) = \int_0^\infty \frac{L_s(t)}{r^2} 4\pi r^2 dr \to \infty$$

We should have a bright sky!

Related question: Olbers' paradox.

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- Universe is expanding, so photons redshift.

- Back to DSNB:

$$\Phi_{(\nu_{\alpha})} = \frac{c}{H_0} \int_{M_0}^{M_{max}} \int_0^{z_{max}} dz \frac{\dot{\rho}_{SN}(z, M) F_{\nu_{\alpha}}(E(1+z), M)}{\sqrt{\Omega_M (1+z)^3 + \Omega_\Lambda}}$$

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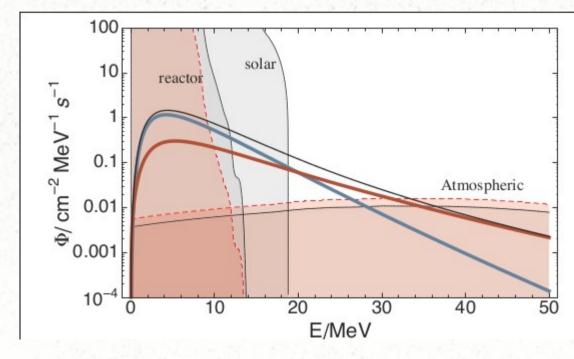
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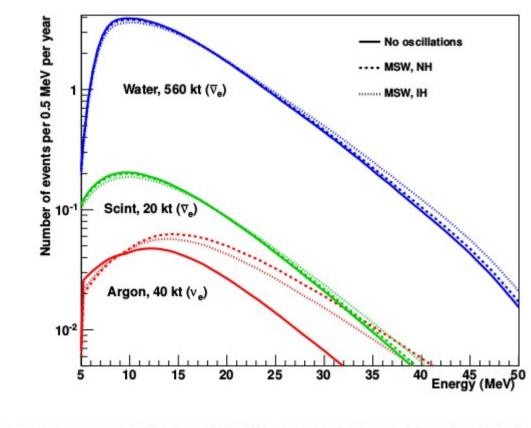
- universe expansion <

- Small window between solar and atmospheric neutrinos to search for DSNB.

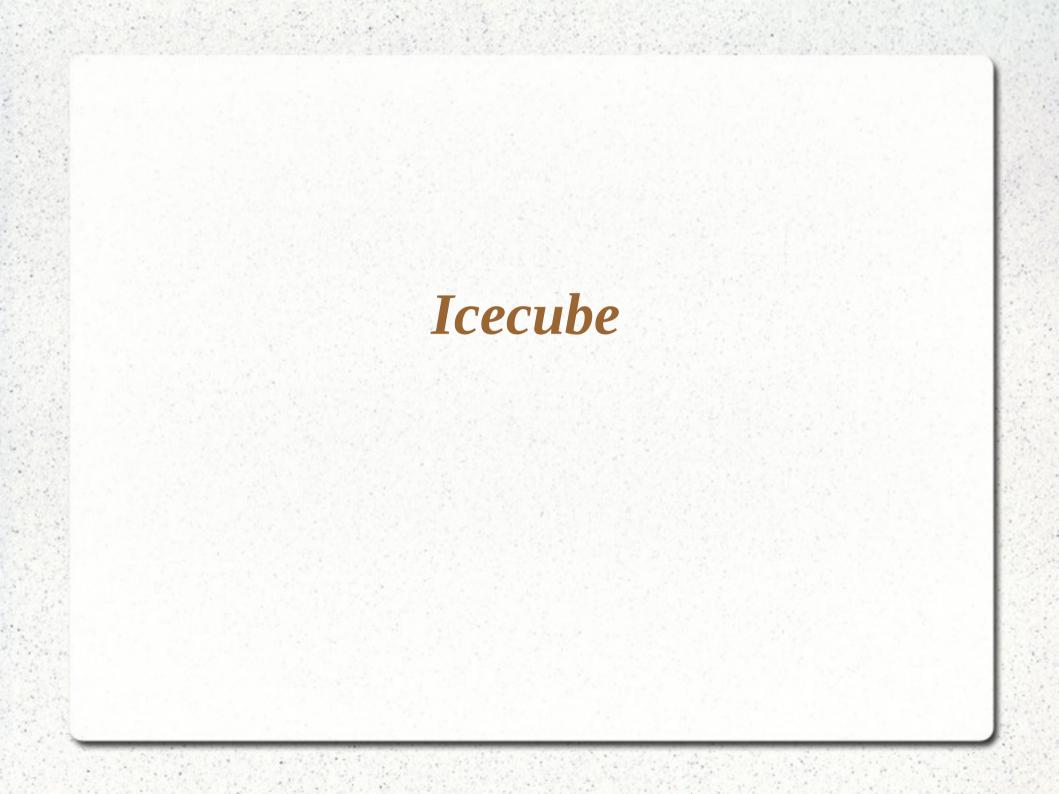


Lunardini et al., ArXiV:1012.1274

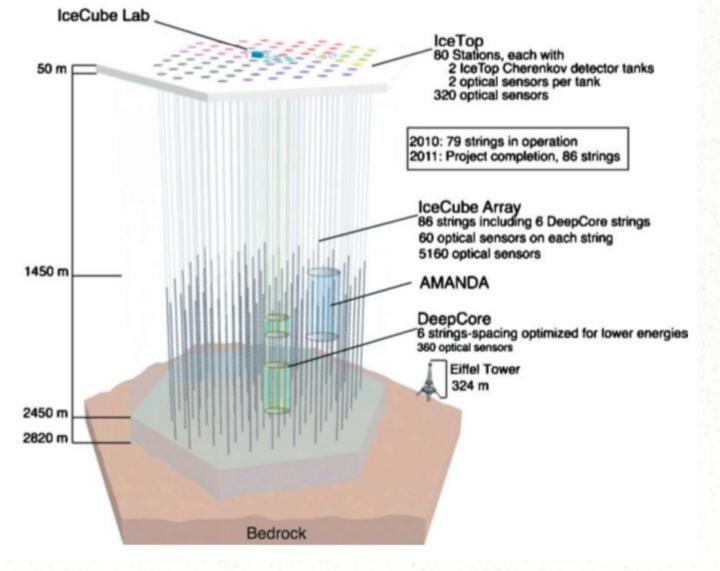
- but with a chalenging event rate...



From ArXiV:1508.00785



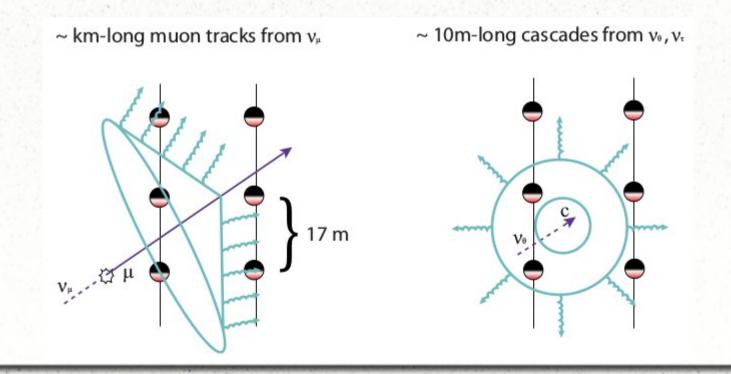




- detect cherenkov light from charged particles.

- incoming neutrino produces the associated charged lepton through charged current.

- Muon neutrinos leave a track + a hadronic shower
- Electron and (low energy) tau neutrinos produce only a shower
- neutral currents produces only the hadronic shower.

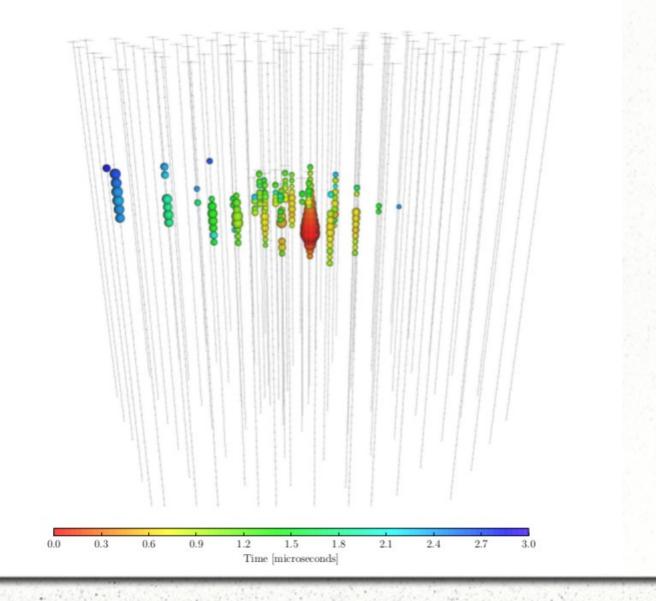


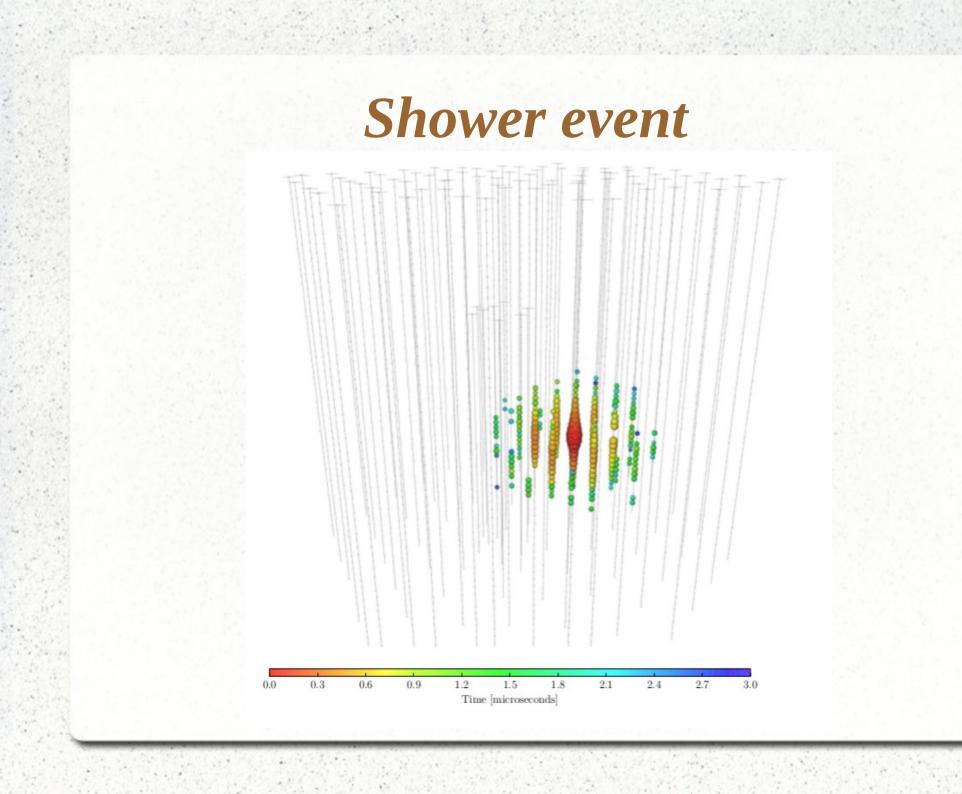
- events with muon tracks has an excelent angular resolution  $< 1^{\circ}$ .
  - But events are not contained, so energy leaks out of the detector, making it difficult to find original neutrino energy.
  - Huge atmospheric background, difficult (but not impossible) to distinguish non-atmospheric neutrinos.
  - At such high energies, neutrino absortion on Earth is relevant, so more sensible to downgoing neutrinos: south hemisphere

- Shower events can do a better job in reconstructing neutrino energy, but with poor direccionality:  $< 30^{\circ}$ 

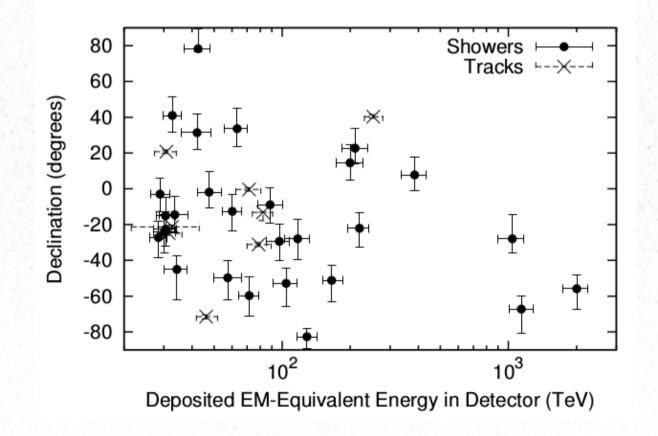
• Sensible to electronic and tauonic neutrinos, not present in atmospheric neutrino flux for high energies  $\rightarrow$  cleaner background

## Track event

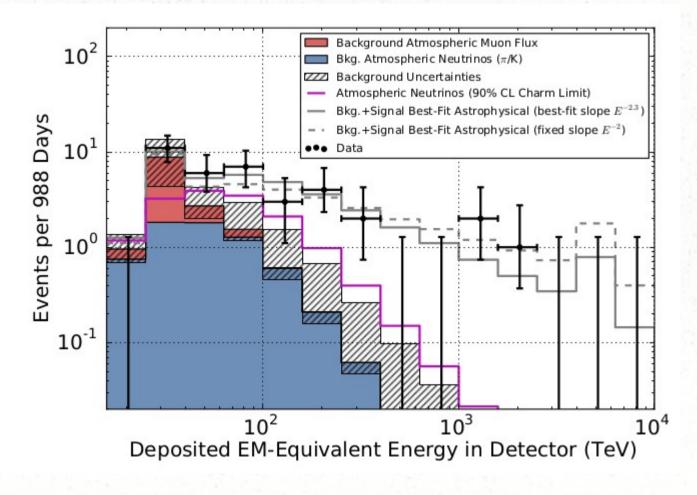


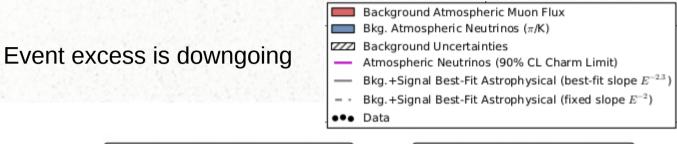


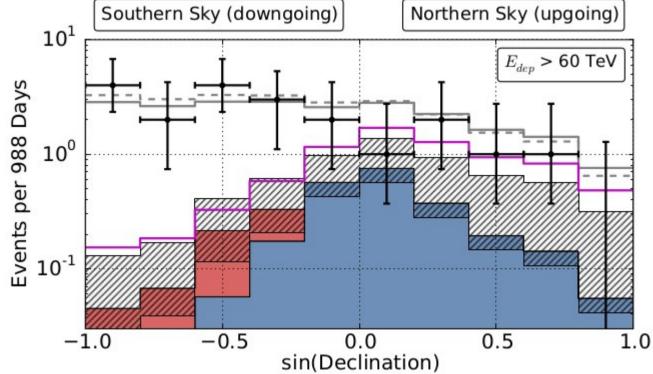
- Selection Criteria: only events that originated inside the detector
- 37 events from 30 to 2000 TeV
- only atmospheric excluded at 5.7 sigma



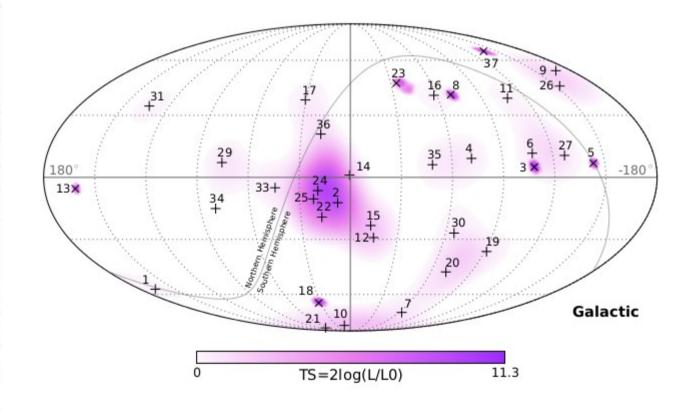
Event excess in high-energy.







No directionality



Due to large mixing angles, flavour distribution in initial neutrino fluxes changes with neutrino evolution.

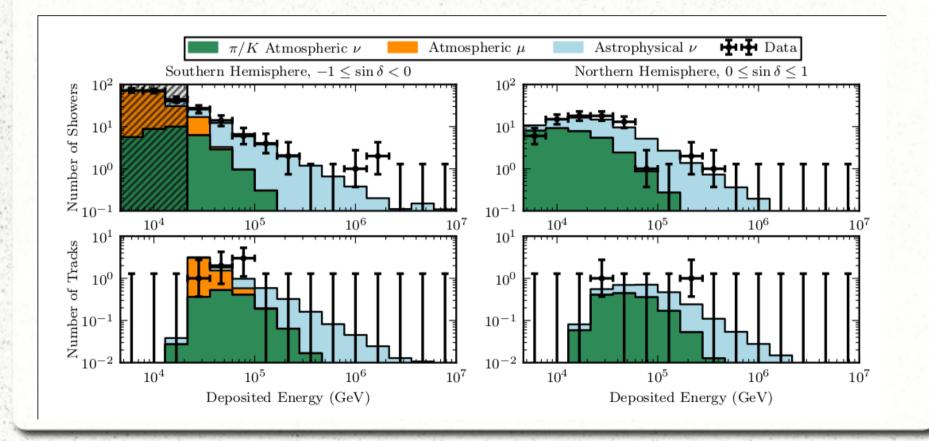
$$(\nu_e:\nu_\mu:\nu_\tau) = (2:1:0) \xrightarrow{\text{neutrino}} (0.93:1.05:1.02)$$

$$(\nu_e : \nu_\mu : \nu_\tau) = (1 - 0 : 0 - 1 : 0)$$
neutrino
evolution
$$\{ (0.6 : 1.3 : 1.1) \\ (1.6 : 0.6 : 0.8) \}$$

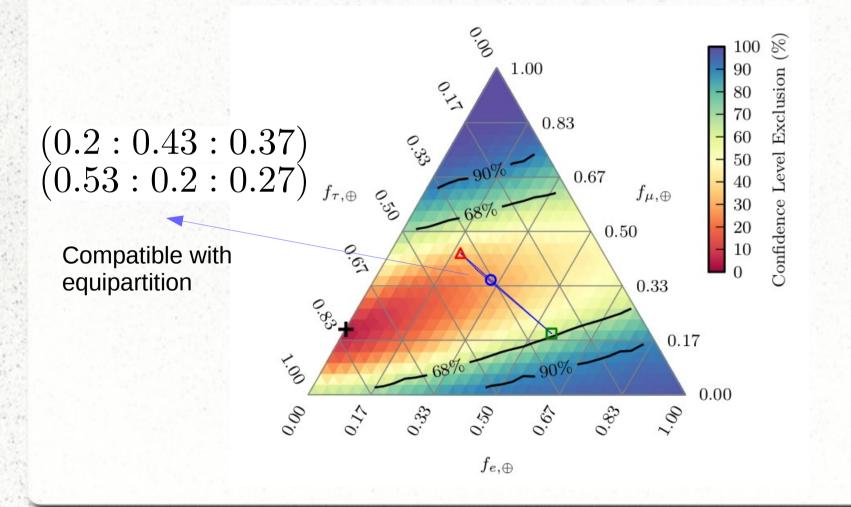
It is reasonable to expect something close to (1:1:1) on Earth.

Physical quantities to be considered:

- event topology
- deposited energy and arrival direction



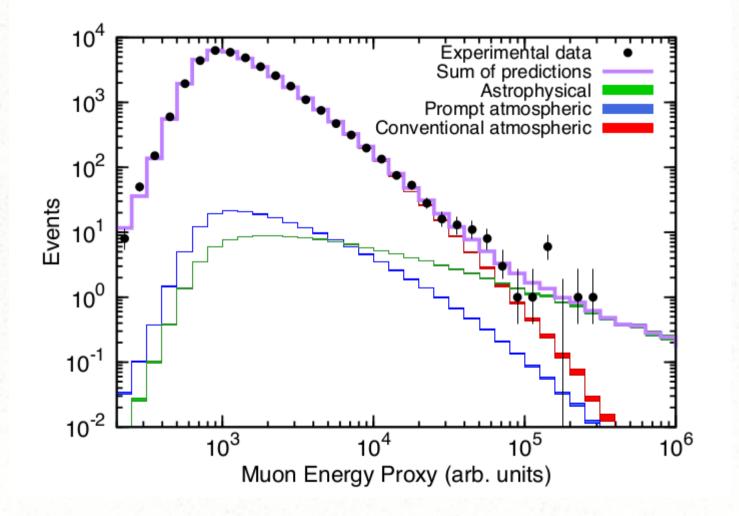
Results (arxiv:1502.03376)

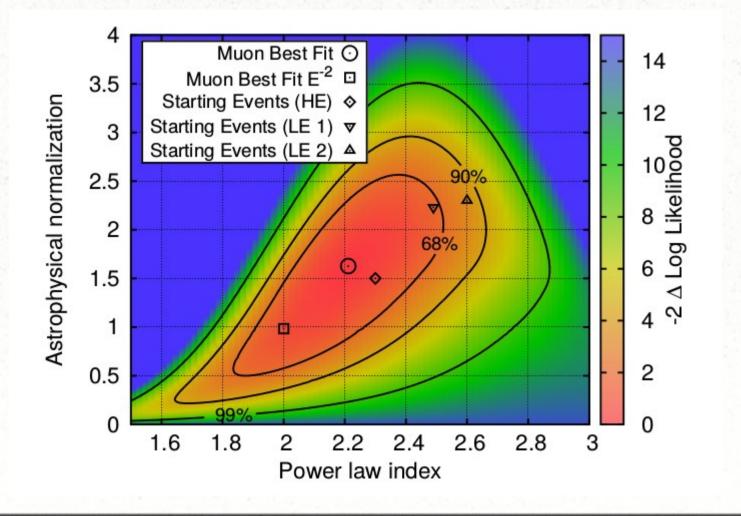


- Recent analysis relax the requirement that the event begins inside the detector  $\rightarrow$  *sensible only to v*<sub>"</sub>.

- Selection of muons that could not have survived all the way to the detector  $\rightarrow$  must have been created through a neutrino interactions halfway.

- Zenith angle > 85 degrees.
- Energy loss before entering the detector.





- "While this work represents the first strong evidence for an astrophysical numu flux in Northern Hemisphere, the sources producing these neutrinos remain unknown."

- starting events:

$$\phi = (9.5 \pm 3) \times 10^{-19} \,\mathrm{GeV \, cm^{-2} \, sr^{-1} s^{-1}}$$

- "northern hemisphere" events:

$$\phi = (9.9^{+3.9}_{-3.4}) \times 10^{-19} \,\text{GeV}\,\text{cm}^{-2}\,\text{sr}^{-1}\text{s}^{-1}$$

- complete analysis is sensible to different flavours, indicating a (1:1:1) flavour distribution

- compatible with some models, incompatible with others.

### **Final Conclusions**

- other lectures: rich phenomenology involving neutrinos theory, production, evolution and detection.

- these lectures:

• which science can we do if we have no control on neutrino production

#### **Neutrino Astrophysics**

• Which science can we do if we can't have a direct neutrino detection

Neutrino Cosmology

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Thank you!