

III. Heterotic model building

heterosis

SAVE  POPULARITY 

noun | het·er·o·sis | \,he-tə-'rō-sis\

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Definition of HETROSIS

: the marked vigor or capacity for growth often exhibited by crossbred animals or plants —called also *hybrid vigor*

— het·er·ot·ic  | \-rā-tik\ | adjective

The heterotic string

world-sheet degrees of freedom

(fermionic formulation $SO(32)$)

right R

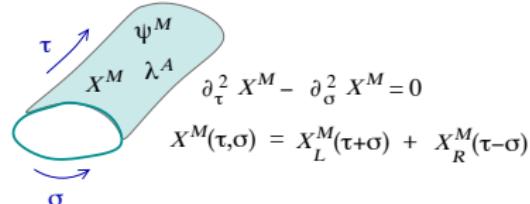
X_R^M , ψ^M 2d fermions

$M = 0, \dots, 9$, $M = 2, \dots, 9$ light-cone

left L

X_L^M , λ^A 2d fermions

$A = 1, \dots, 32$



$$\psi^M(\tau, \sigma + 2\pi) = \mp \psi^M(\tau, \sigma) ; \quad \lambda^A(\tau, \sigma + 2\pi) = \mp \lambda^A(\tau, \sigma) ; \quad \text{--- Neveu-Schwarz(NS), + Ramond}$$

massless states $|R\rangle \otimes |L\rangle$

NS \otimes NS

e.g. $b_{-\frac{1}{2}}^M |0\rangle \otimes \lambda_{-\frac{1}{2}}^A \lambda_{-\frac{1}{2}}^B |0\rangle$

$$M_R^2 = \tilde{N}_X + \tilde{N}_\psi - \frac{1}{2}$$

$$M_L^2 = N_X + N_\lambda - 1$$

496 gauge vectors of $SO(32)$ in $10d$ (gauginos in Ramond \otimes NS)

Full massless spectrum

$R \otimes L$ in light cone

$$[(\mathbf{8}_v \oplus \mathbf{8}_s, \mathbf{1})]_R \otimes [(\mathbf{8}_v, \mathbf{1}) \oplus (\mathbf{1}, \mathbf{496})]_L \\ = (\mathbf{1} \oplus \mathbf{35} \oplus \mathbf{28} + \mathbf{8}_c + \mathbf{56}_s, \mathbf{1}) \oplus (\mathbf{8}_v \oplus \mathbf{8}_s, \mathbf{496})$$

massless fields

$$\{\varphi, G_{MN}, B_{MN}, \Psi, \Psi_M\} \oplus \{A_a^M, \chi_a\} \quad a = 1, \dots, \dim G_{\text{het}}$$

10d, $\mathcal{N}=1$ supergravity \oplus super Yang-Mills G_{het}

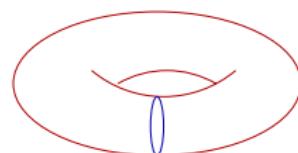
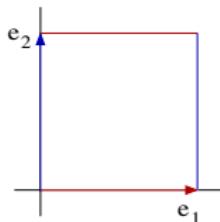
$$G_{\text{het}} = E_8 \times E_8, SO(32)$$

Gauge and gravitational anomalies are cancelled
by the Green-Schwarz mechanism

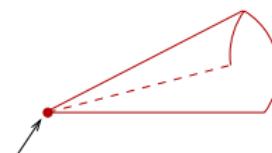
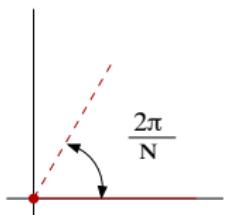
Orbifolds

$\mathcal{O} = \mathcal{M}/\Gamma$; Γ = discrete group of isometries of \mathcal{M}

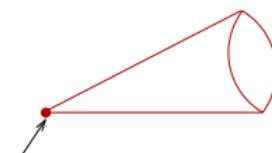
$$\mathbf{T}^2 = \mathbf{R}^2 / \Lambda$$



$$\text{cono} \quad \mathbf{R}^2 / \mathbf{Z}_N$$



punto fijo

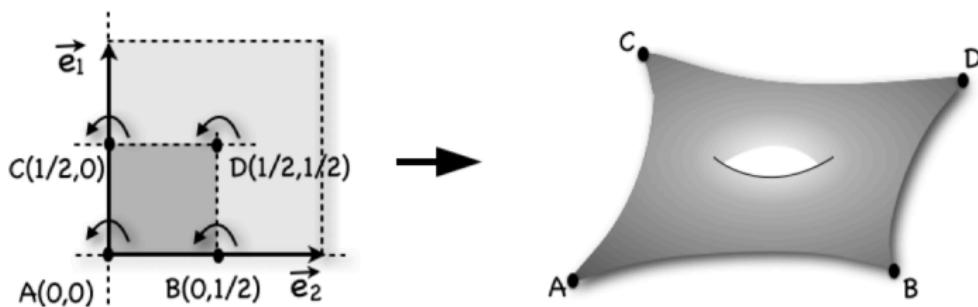


singularidad

T^2/\mathbb{Z}_2 orbifold

$T^2 = \mathbb{R}^2/\Lambda$, $\Lambda = SO(4)$ root lattice

$\mathbb{Z}_2 = \{\mathbf{1}, \theta\}$, $\theta = \text{rotation by } \pi$



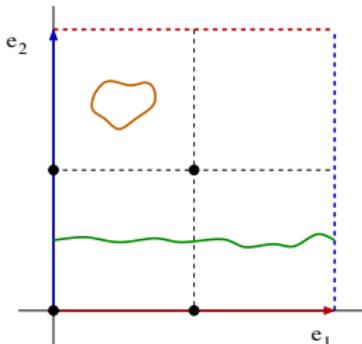
• : fixed points \longrightarrow singularities

Closed strings on orbifolds \mathcal{M}/Γ

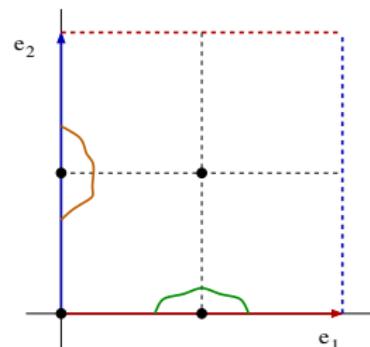
Dixon, Harvey, Vafa, Witten

*

untwisted sector



twisted sector



$$\vec{X}(\tau, \sigma + 2\pi) = \vec{X}(\tau, \sigma) + n_i \vec{e}_i$$

$$\vec{X}(\tau, \sigma + 2\pi) = g \vec{X}(\tau, \sigma) + n_i \vec{e}_i$$

$g \in \Gamma$ admits fixed points

- * Orbifold projection: physical states are invariant under Γ

Both conditions are required by modular invariance

Toroidal orbifolds T^6/\mathbb{Z}_N

$$T^6 = \mathbb{R}^6/\Lambda, \quad \mathbb{Z}_N = \{\mathbf{1}, \theta, \dots, \theta^{N-1}\}, \quad \theta \in SO(6)$$

crystallographic action: $W \in \Lambda, \theta W \in \Lambda$

$$\theta^N = \mathbf{1} \Rightarrow \theta \text{ has eigenvalues } e^{\pm 2\pi i v_i}, \quad v_i = \frac{k_i}{N}, \quad k_i \in \mathbb{Z}, \quad i = 1, 2, 3$$

complex internal coordinates $Z^i = \frac{1}{\sqrt{2}} (X^{2i+2} + iX^{2i+3}), \quad \theta Z^i = e^{2\pi i v_i} Z^i$

$$\theta = \exp 2\pi i(v_1 J_{12} + v_2 J_{34} + v_3 J_{56}), \quad J_{2i-1, 2i} : SO(6) \text{ Cartan generator}$$

action on spinor representation

$$\theta |\pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2}\rangle = e^{i\pi(\pm v_1 \pm v_2 \pm v_3)} |\pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2}\rangle$$

supersymmetry $\Rightarrow \pm v_1 \pm v_2 \pm v_3 = \text{even}$

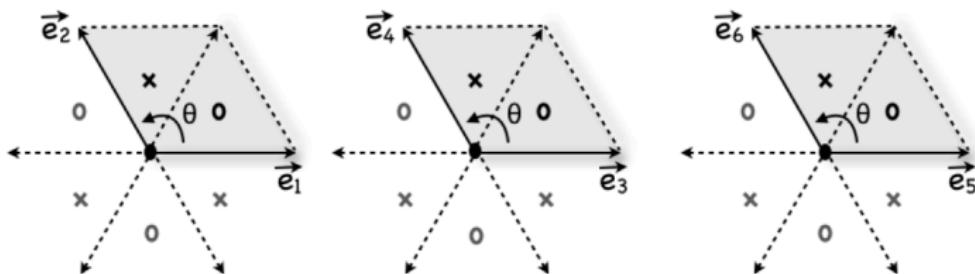
$$\theta^N = \mathbf{1} \text{ acting on fermions} \Rightarrow N(v_1 + v_2 + v_3) = \text{even}$$

Example: T^6/\mathbb{Z}_3 orbifold

$T^6 = \mathbb{R}^6/\Lambda$, Λ = product of three $SU(3)$ root lattices

$\mathbb{Z}_3 = \{\mathbf{1}, \theta, \theta^2\}$, θ = rotation by $\frac{2\pi}{3}$ in each sub-lattice

$$(\nu_1, \nu_2, \nu_3) = (\frac{1}{3}, \frac{1}{3}, -\frac{2}{3})$$



•, ○, ✕ : fixed points

altogether $3 \times 3 \times 3 = 27$ fixed points

Action on 2d fermions

right movers

$$\Psi^i = \frac{1}{\sqrt{2}} (\psi^{2i+2} + i\psi^{2i+3}), \quad \text{θ-twisted sector}$$
$$\Psi^i(\tau, \sigma + 2\pi) = \mp e^{2\pi i v_i} \Psi^i(\tau, \sigma + 2\pi)$$

left movers give gauge degrees of freedom

$$\lambda_{\pm}^A = \frac{1}{\sqrt{2}} (\lambda^{2A-1} \pm i\lambda^{2A}), \quad A = 1, \dots, 16$$

for $E_8 \times E_8$ divide in two groups: $\lambda_{\pm}^A, \lambda'^A, A = 1, \dots, 8$

Modular invariance requires that λ^A transforms under θ . The action can be realized by rotation γ with eigenvalues $e^{\pm 2\pi i V_A}$, $V_A = \frac{K_A}{N}$, $K_A \in \mathbb{Z}$.

$$\gamma^N = 1 \Rightarrow N(V_1 + \dots + V_{16}) = \text{even}$$

Gauge shift vector: $V = (V_1, \dots, V_8) \times (V'_1, \dots, V'_8)$, $E_8 \times E_8$

Modular invariance (level-matching) further requires $N(V^2 - v^2) = \text{even}$

Standard Embedding: $V = (v_1, v_2, v_3, 0, \dots, 0) \times (0, \dots, 0)$

Compactification on Calabi-Yau (CY) manifolds

Recall fields in 10d

$$\{\varphi, G_{MN}, B_{MN}, \Psi, \Psi_M\} \oplus \{A_a^M, \chi_a\}, \quad a = 1, \dots, \dim G_{\text{het}} \quad G_{\text{het}} = E_8 \times E_8, SO(32)$$

Compactification $\mathcal{M}_{10} = \mathcal{M}_4 \times K_6$

Supersymmetry in 4d $\Rightarrow D_m \eta = 0 \Rightarrow R_{mn} = 0$

K_6 Ricci-flat \Rightarrow holonomy $SU(3)$ K_6 is CY

Furthermore, K_6 is Kähler (complex with special property of the metric)

$$x^m \longrightarrow z^i, \bar{z}^{\bar{i}} \quad k\text{-forms: } \omega_{m_1 \dots m_k} \quad (p, q)\text{-forms: } \omega_{i_1 \dots i_p \bar{j}_1 \dots \bar{j}_q}$$

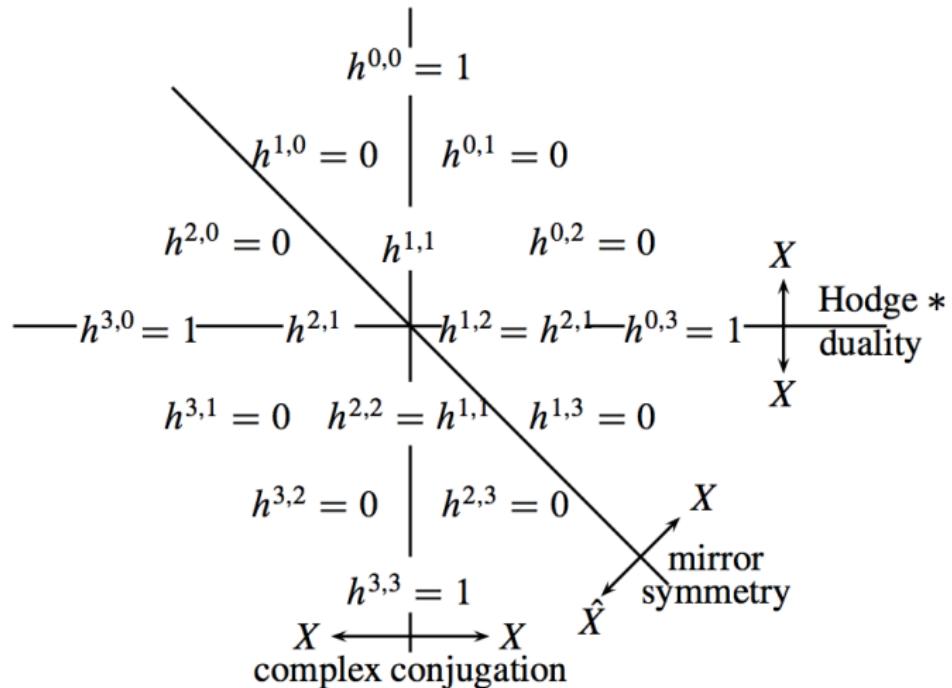
Betti numbers $b_k = \# \text{ closed (mod exact) } k\text{-forms} = \# \text{ harmonic } k\text{-forms}$

Hodge numbers

$h^{p,q} = \# \text{ closed (mod exact) } (p, q)\text{-forms} = \# \text{ harmonic } (p, q)\text{-forms}$

$$b_k = \sum_{p=0}^k h^{p,k-p}$$

Hodge diamond of a CY X



$$\chi = 2(h^{1,1} - h^{1,2}) \quad \text{Euler characteristic of } X$$

Hodge plot

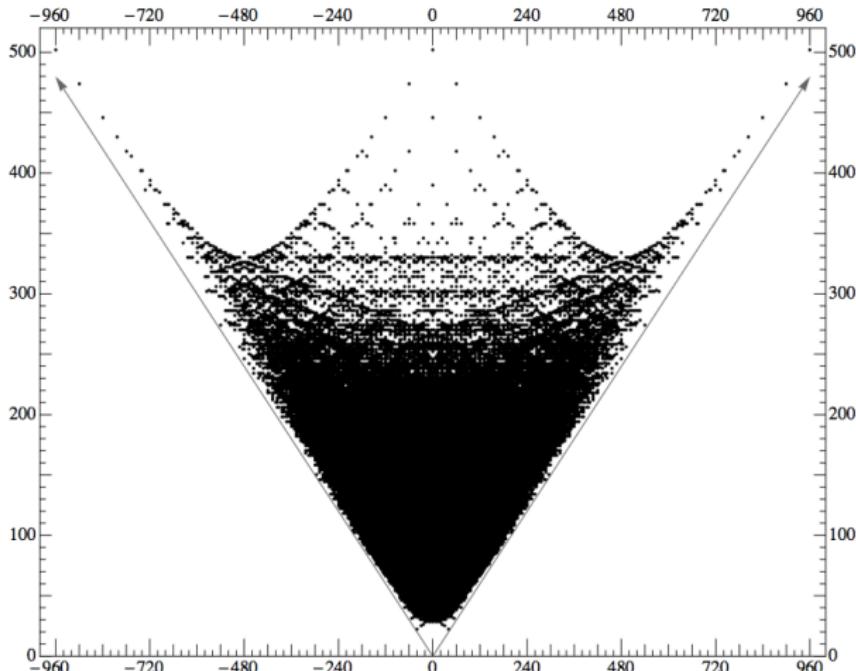


Figure 1: The Hodge plot for the list of reflexive 4-polytopes. The Euler number $\chi = 2(h^{1,1} - h^{1,2})$ is plotted against the height $y = h^{1,1} + h^{1,2}$. The oblique axes correspond to $h^{1,1} = 0$ and $h^{1,2} = 0$.

Taken from arXiv:1207.4792, based on the Kreuzer-Skarke list