

Beyond Ultracold Atoms: Halo Nuclei and Hadronic Molecules

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Agenda



- 1. EFT for Ultracold Atoms I: Effective Field Theories & Universality
- 2. EFT for Ultracold Atoms II: Cold Atoms & the Unitary Limit
- 3. EFT for Ultracold Atoms III: Weak Coupling at Finite Density
- 4. EFT for Ultracold Atoms IV: Few-Body Systems in the Unitary Limit
- 5. Beyond Ultracold Atoms: Halo Nuclei and Hadronic Molecules

Literature

- G.P. Lepage, TASI Lectures 1989, arXiv:hep-ph/0506330
- D.B. Kaplan, arXiv:nucl-th/9506035
- E. Braaten, HWH, Phys. Rep. 428 (2006) 259 [arXiv:cond-mat/0410417]

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- Separation of scales:
 - $1/k = \lambda \gg R$
- Limited resolution at low energy:
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 - \longrightarrow include long-range physics explicitly
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 - \longrightarrow capture in low-energy constants using renormalization
 - \longrightarrow include long-range physics explicitly
- Systematic, model independent
- Very low energies: only short range interactions
- Exploit cluster substructures \implies Halo EFT
- Universal properties





Halo Nuclei



• Low separation energy of valence nucleons: $B_{valence} \ll B_{core}, E_{ex}$

 \longrightarrow close to "nucleon drip line" \longrightarrow scale separation \longrightarrow EFT



C.-B. Moon, Wikimedia Commons

• EFT for halo nuclei

(Bertulani, HWH, van Kolck, 2002; Bedaque, HWH, van Kolck, 2003; ...)

Scales and Antisymmetrization



- Scales: $R_{halo} \gg R_{core} \sim \ell$
- Antisymmetrization with respect to neutrons in core?
- Core neutrons not active dof in halo EFT



Physics: exchange of core nucleon and halo nucleon only contributes to observables if there is spatial overlap between wave functions of core and halo nucleon

 \implies small for $R_{core} \ll R_{halo}$

• Effects subsumed in low-energy constants, included perturbatively in expansion in R_{core}/R_{halo}

Two-Neutron Halo Nuclei

Effective Lagrangian (schematically)

$$\mathcal{L}_{d} = \psi^{\dagger} \left(i\partial_{t} + \frac{\vec{\nabla}^{2}}{2m} \right) \psi + \frac{g_{2}}{4} d^{\dagger}d - \frac{g_{2}}{4} (d^{\dagger}\psi^{2} + (\psi^{\dagger})^{2}d) - \frac{g_{3}}{36} d^{\dagger}d\psi^{\dagger}\psi + \dots$$

- 2- and 3-body interaction at leading order: g_2 , g_3 enhanced!
- 2-body amplitude:
- 3-body amplitude:





• Structure of 2-neutron halo nuclei \rightarrow energies, matter form factors, radii

technische

 $+ \dots$

Renormalization

- \checkmark Observables are independent of regulator/cutoff Λ
- \Rightarrow Running coupling $H(\Lambda) \propto \Lambda^2 g_3(\Lambda)$
- $H(\Lambda)$ periodic: limit cycle

 $\Lambda \to \Lambda \, e^{n\pi/s_0} \approx \Lambda(22.7)^n$

Discrete scale invariance

- Efimov effect (Efimov, 1970)
- Observed in ultracold atoms (Krämer et al., 2006; ...)
- Relevant in halo nuclei?







Efimov Physics in Halo Nuclei



Efimov effect in halo nuclei? (Fedorov, Jensen, Riisager, 1994)

 \implies excited states obeying scaling relations

• Correlation plot: $E_{nn} \leftrightarrow E_{nc}$ (Am

(Amorin, Frederico, Tomio, 1997)



Canham, HWH, Eur. Phys. J. A 37 (2008) 367

Efimov Physics in ^{22}C



- Matter radius from ${}^{22}C + p$ & Glauber: $\langle r_0^2 \rangle^{1/2} = 5.4(9)$ fm (Tanaka et al., Phys. Rev. Lett. **104** (2010) 062701)
- Halo EFT analysis of impact on other observables in ²²C (Acharya, Ji, Phillips, Phys. Lett. B 723 (2013) 196)



Plots for $\langle r_0^2 \rangle^{1/2} = 4.5, 5.4, 6.3$ fm

Excited Efimov states in ²²C appear to be ruled out



(G. Hagen, P. Hagen, HWH, Platter, Phys. Rev. Lett. **111** (2013) 132501)

- The Many and the Few: emergence of effective halo degrees of freedom
- Coupled cluster calculations of ⁶⁰Ca and ⁶¹Ca using chiral N2LO two-body force and schematic three-body force:

⁶¹Ca is a weakly bound S-wave state (or virtual state)

- Quantitative estimate: $S_n = B_{nc} = 5...8 \text{ keV}$
- Scattering Parameters:

 $a_{cn} = 54(1) \text{ fm}, r_{cn} = 9.0(2) \text{ fm} \implies r_{cn}/a_{cn} \approx 1/6$

- Investigate consequences for ⁶²Ca using halo EFT
- Prospects for excited Efimov states in 62 Ca: $S_{\text{deep}} = 1/(\mu_{cn}r_{cn}^2) \approx 500$ keV, scaling factor $\lambda_0 \approx 16$

 \implies possible if $S_{2n} \gtrsim 230 \text{ keV}$

Efimov Physics in ⁶²Ca



Universal correlations between S_{2n}, ⁶¹Ca-n scattering length, ⁶²Ca matter, and charge radii



(G. Hagen, P. Hagen, HWH, Platter, Phys. Rev. Lett. 111 (2013) 132501)

- Excited Efimov state appears around $S_{2n} \approx 230 \text{ keV}$
- Matter radii of order tens of Fermi possible

Production of Excited Efimov States

- How to study excited Efimov states experimentally?
 (A. Macchiavelli, Few-Body Syst. 56, 773 (2015)
- Consider transfer reactions for candidate nucleus $^{A}Z_{N}$
 - (a) One-neutron transfer: ${}^{(A-1)}Z_{(N-1)}(d,p)^A Z_N$
 - (b) Two-neutron transfer: ${}^{(A-2)}Z_{(N-2)}(t,p)^A Z_N$
- Back-of-the-envelope estimate \implies (a) is most promising
- Reaction calculation in Halo EFT would be useful



EFT for Cold Atoms – p. 15

Hadronic Molecules

- New $c\bar{c}$ states at B factories: X, Y, Z
- **Example:** X(3872) (Belle, CDF, BaBar, D0)
- No ordinary $c\bar{c}$ -state
 - Decays violate isospin
 - Measured mass depends on decay channel



 $m_X = (3871.69 \pm 0.17) \text{ MeV}$ $\Gamma < 1.2 \text{ MeV}$ $J^{PC} = 1^{++}$

- Nature of X(3872)? $\overline{D}^0 D^{0*}$ -molecule, tetraquark, charmonium hybrid, ...
- Molecular nature \Rightarrow interaction of X(3872) with D^0 , \overline{D}^0 , D^{0*} , \overline{D}^{0*} determined by large scattering length



Nature of X(3872)



- Nature of X(3872) not finally resolved
- Assumption: X(3872) is weakly-bound $D^0-\overline{D}^{0*}$ -molecule

 $\implies |X\rangle = (|D^0 \bar{D}^{0*}\rangle + |\bar{D}^0 D^{0*}\rangle)/\sqrt{2} , \qquad B_X = (0.11 \pm 0.21) \text{ MeV}$

 \implies universal properties (cf. Braaten et al., 2003-2008, ...)

- Explains isospin violation in decays of $X(3872) \Rightarrow$ superposition of I = 1 and I = 0
- Different masses due to different line shapes in decay channels
- Large scattering length to LO determines interaction of X(3872) with D^0 and D^{0*}
- Higher orders: EFT with perturbative pions \Rightarrow XEFT

(Fleming, Kusunoki, Mehen, van Kolck, 2007; Fleming, Mehen, 2008) (Braaten, HWH, Mehen, 2010; ...) EFT for X(3872)



Effective Lagrangian

$$\mathcal{L} = \sum_{\substack{j=D^0, D^{*0}, \bar{D}^0, \bar{D}^{*0}}} \psi_j^{\dagger} \left(i\partial_t + \frac{\nabla^2}{2m_j} \right) \psi_j + \Delta X^{\dagger} X - \frac{g}{\sqrt{2}} \left(X^{\dagger} (\psi_{D^0} \psi_{\bar{D}^{*0}} + \psi_{D^{*0}} \psi_{\bar{D}^0}) + \mathsf{H.c.} \right) + \dots,$$

• Propagator of the X(3872)



Three-body integral equation









Canham, HWH, Springer, Phys. Rev. D 80, 014009 (2009)

• Three-body scattering lengths: $a_{D^0X} = -9.7a \approx -85$ fm

 $a_{D^0X} = -9.7a \approx -85 \text{ fm}$ $a_{D^{*0}X} = -16.6 a \approx -146 \text{ fm}$

Experimental Observation ?



- Behavior of X(3872) produced in isolation should be distinguishable from its behavior when in the presence of $D^0, D^{*0}, \bar{D}^0, \bar{D}^{*0}$
- Final state interaction of D, D^* mesons in B_c -decays
- Example: quark-level B_c decay yielding three charmed/anticharmed quarks in final state



Process in principle accessible at the LHC

Summary



- Cluster EFT for halo nuclei
 - \implies large scattering length/shallow states
 - Controlled, systematic approach \implies error estimates
 - Straightforward inclusion of external currents
- Universal theory has applications in atomic, nuclear, and particle physics
- Universality predicts correlations between observables
 input from theory or experiment
- Excited Efimov state possible in ⁶²Ca
- Calculations of EM structure and reactions