

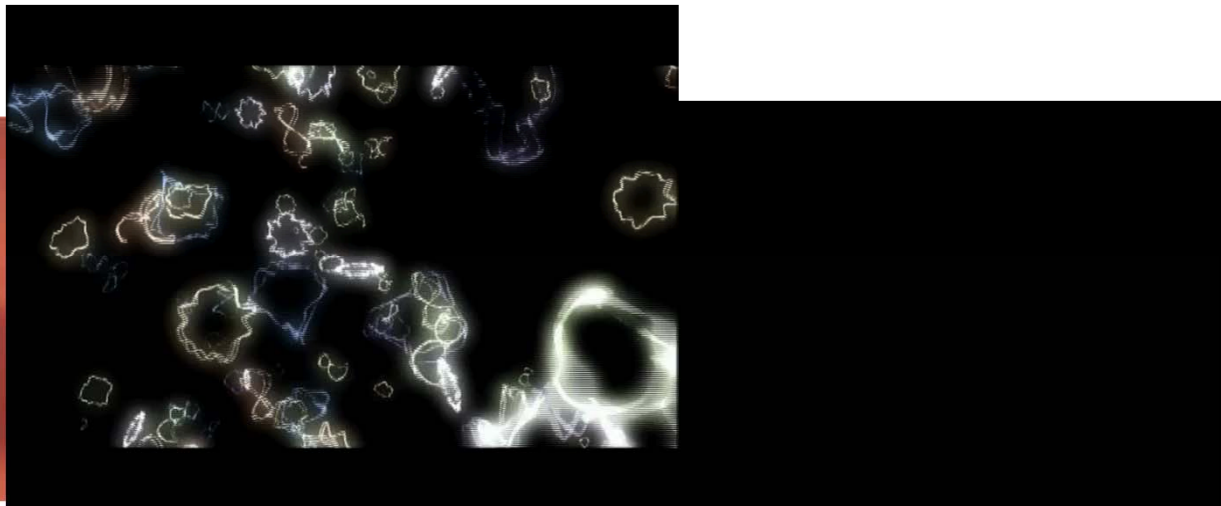
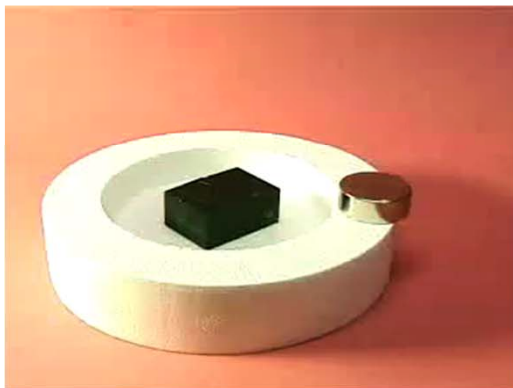
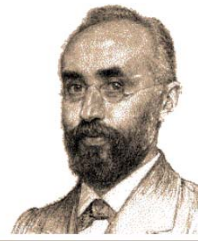
Holographic duality, strange metals and entanglement.

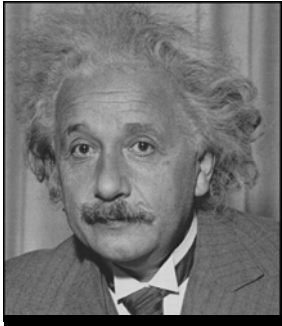
Jan Zaanen



Universiteit
Leiden

Instituut-Lorentz
for theoretical physics

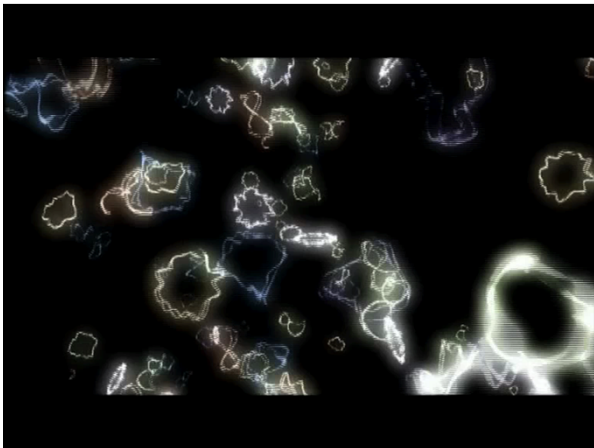




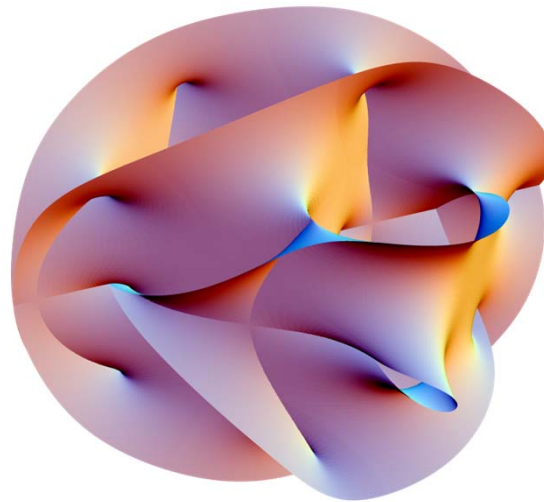
String theory



$$\mathcal{S} = \frac{T}{2} \int d^2\sigma \sqrt{-h} h^{ab} g_{\mu\nu}(X) \partial_a X^\mu(\sigma) \partial_b X^\nu(\sigma)$$



Quantum physics of strings



Beautiful mathematics
(Calabi-Yau, ...)



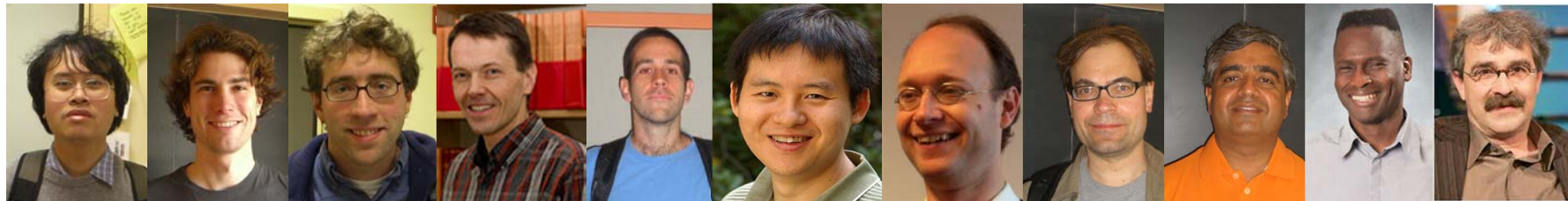
Quantum space-time
(big bang, ...)

String theory: what is it really good for?

- **Quantum matter: heavy fermion systems, high T_c superconductors, CMR manganites...**

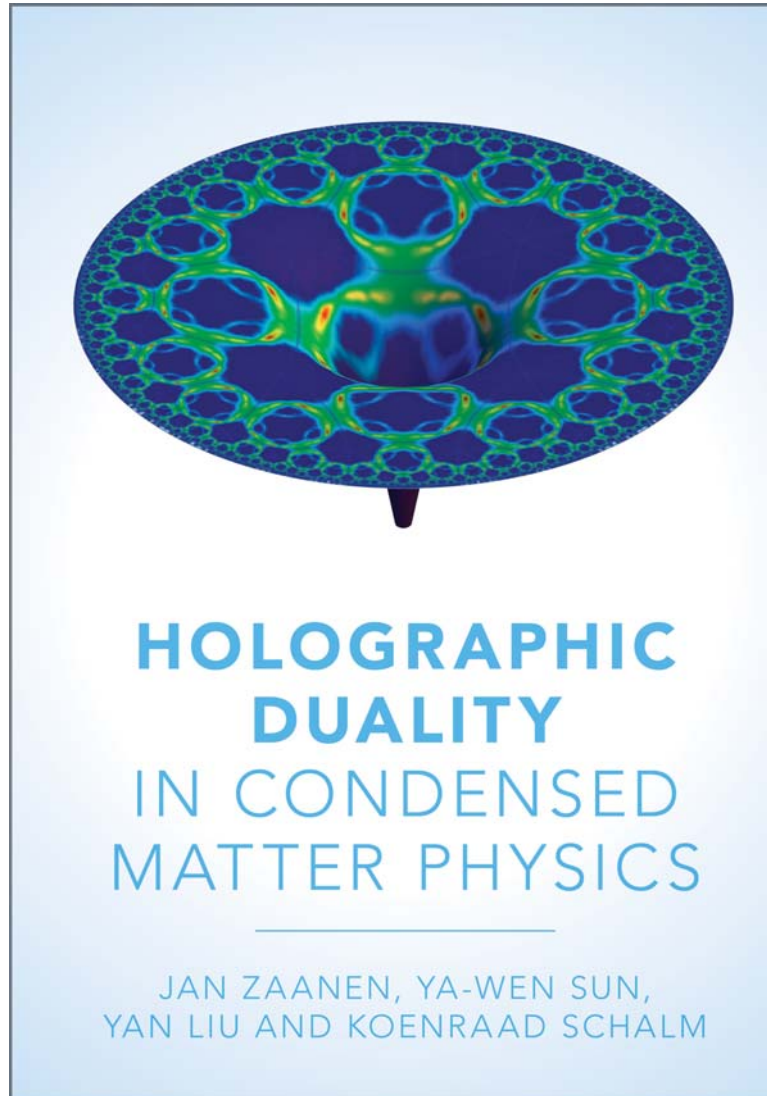


Polchinski Kachru Starinets Erdmenger Gubser Horowitz Kiritsis Gauntlett Policastro Tong



Son Hartnoll Herzog Thoracius McGreevy Liu Schalm Karch Sachdev Phillips Zaanen

Book sales ...



Cambridge University Press

Release: October 28 2015.

It is 600 pages and only € 80!

Quantum field theory = Statistical physics.

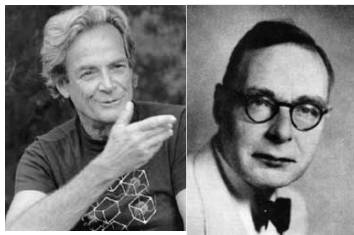


$$Z = \sum_{\text{configs.}} e^{-\frac{E_{\text{config}}}{k_B T}}$$

Path integral mapping

“Thermal QFT”, Wick rotate:

$$t \rightarrow i\tau$$



$$Z_{\hbar} = \sum_{\text{worldhistories}} e^{-\frac{S_{\text{history}}}{\hbar}}$$

But generically: the quantum partition function is not probabilistic: “sign problem”, no mathematical control!

$$Z_{\hbar} = \sum_{\text{worldhistories}} (-1)^{\text{history}} e^{-\frac{S_{\text{history}}}{\hbar}}$$

Fermions at a finite density: the sign problem.

Imaginary time first quantized path-integral formulation

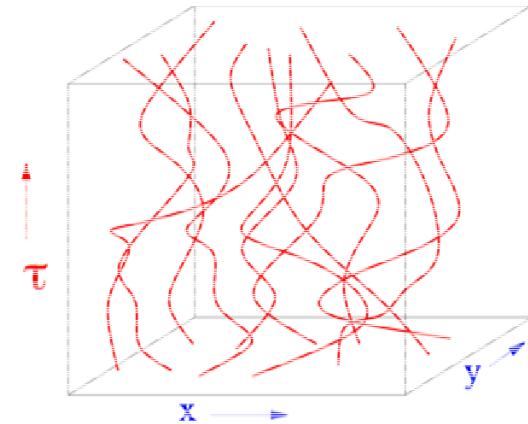


$$\begin{aligned}\mathcal{Z} &= \text{Tr} \exp(-\beta \hat{\mathcal{H}}) \\ &= \int d\mathbf{R} \rho(\mathbf{R}, \mathbf{R}; \beta)\end{aligned}$$

$$\mathbf{R} = (\mathbf{r}_1, \dots, \mathbf{r}_N) \in \mathbb{R}^{Nd}$$

$$\rho_{B/F}(\mathbf{R}, \mathbf{R}; \beta) = \frac{1}{N!} \sum_{\mathcal{P}} (\pm 1)^{\mathcal{P}} \rho_D(\mathbf{R}, \mathcal{P}\mathbf{R}; \beta)$$

$$= \frac{1}{N!} \sum_{\mathcal{P}} (\pm 1)^{\mathcal{P}} \int_{\mathbf{R} \rightarrow \mathcal{P}\mathbf{R}} \mathcal{D}\mathbf{R}(\tau) \exp \left\{ -\frac{1}{\hbar} \int_0^{\hbar/T} d\tau \left(\frac{m}{2} \dot{\mathbf{R}}^2(\tau) + V(\mathbf{R}(\tau)) \right) \right\}$$



Boltzmannons or Bosons:

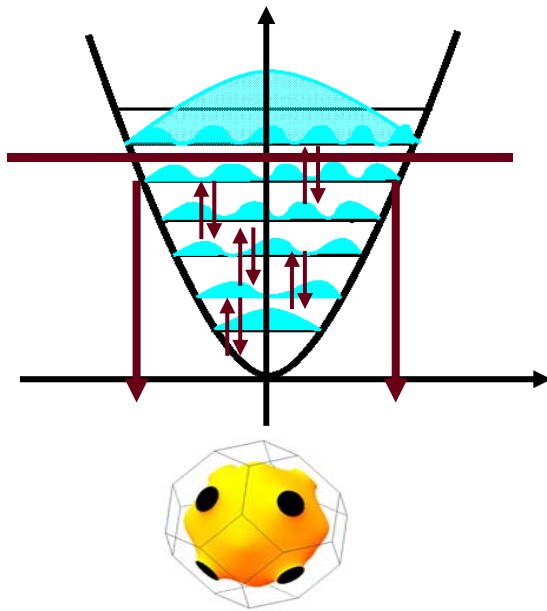
- integrand non-negative
- probability of equivalent classical system: (crosslinked) ringpolymers

Fermions:

- negative Boltzmann weights
- non probabilistic: NP-hard problem (Troyer, Wiese)!!!

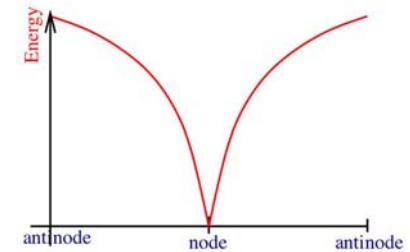
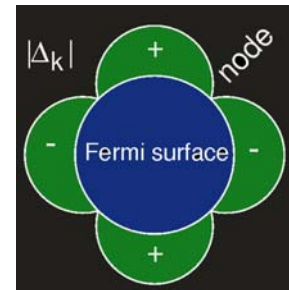
Fermions: the tiny repertoire ...

Fermiology

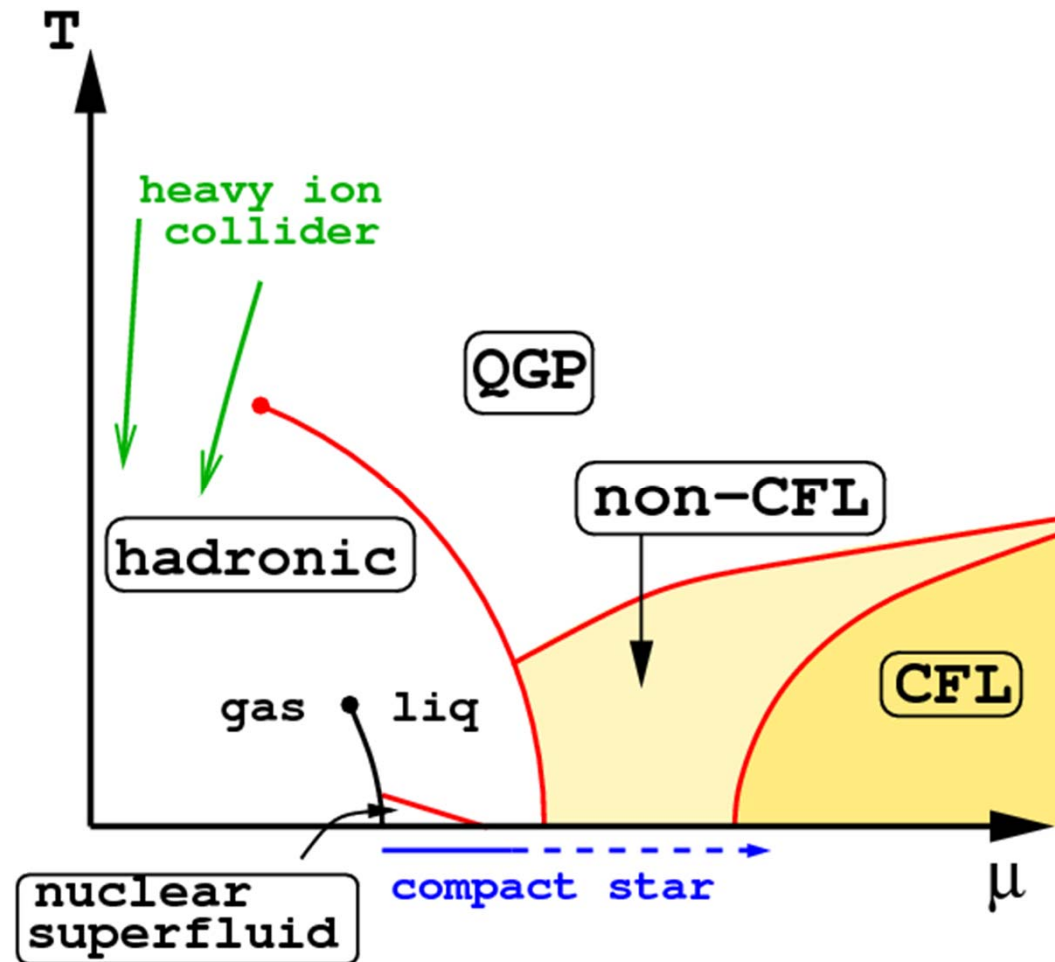


BCS superconductivity

$$\Psi_{BCS} = \prod_k (u_k + v_k c_{k\uparrow}^+ c_{-k\downarrow}^+) |vac.\rangle$$



QCD at intermediate density ...



29 years later: the
“consensus document”.



REVIEW

doi:10.1038/nature14165

From quantum matter to high-temperature superconductivity in copper oxides

B. Keimer¹, S. A. Kivelson², M. R. Norman³, S. Uchida⁴ & J. Zaanen⁵

The discovery of high-temperature superconductivity in the copper oxides in 1986 triggered a huge amount of innovative scientific inquiry. In the almost three decades since, much has been learned about the novel forms of quantum matter that are exhibited in these strongly correlated electron systems. A qualitative understanding of the nature of the superconducting state itself has been achieved. However, unresolved issues include the astonishing complexity of the phase diagram, the unprecedented prominence of various forms of collective fluctuations, and the simplicity and insensitivity to material details of the ‘normal’ state at elevated temperatures.

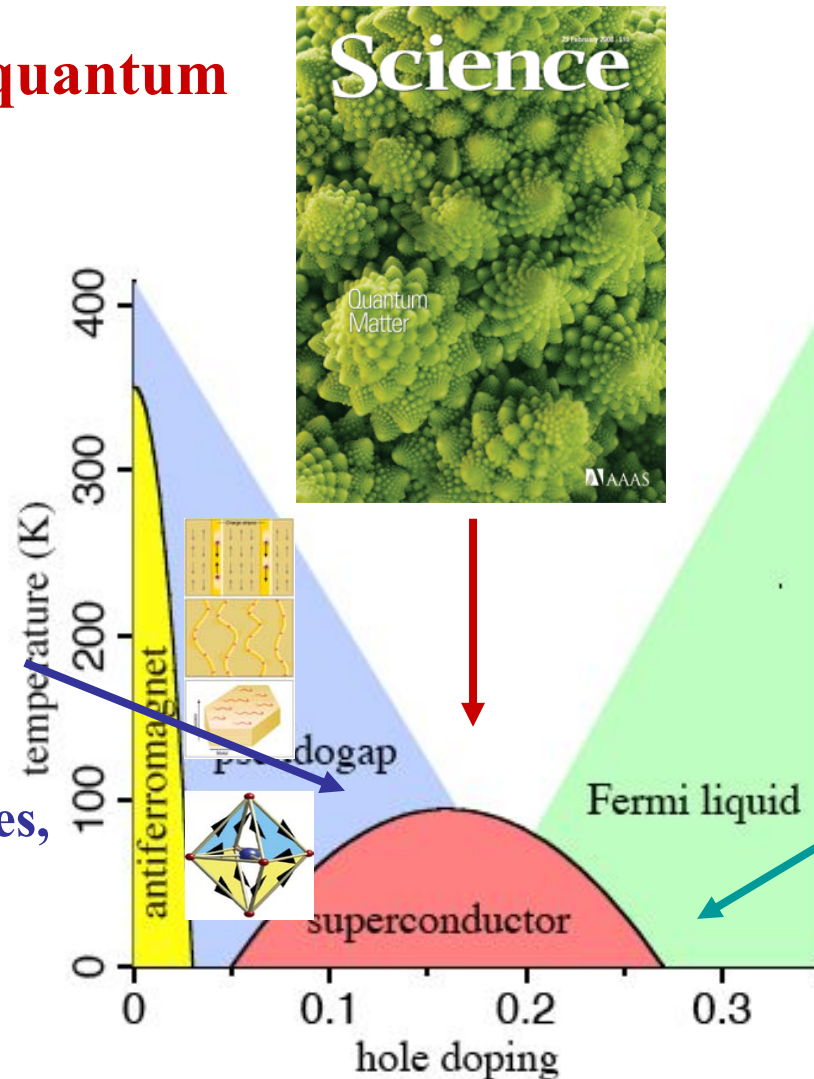
Phase diagram high T_c superconductors (Nature 518, 181, 2015)

The clash: the quantum critical metal

The quantized traffic jam



Exotic orders: stripes, orbital currents, nematics ...

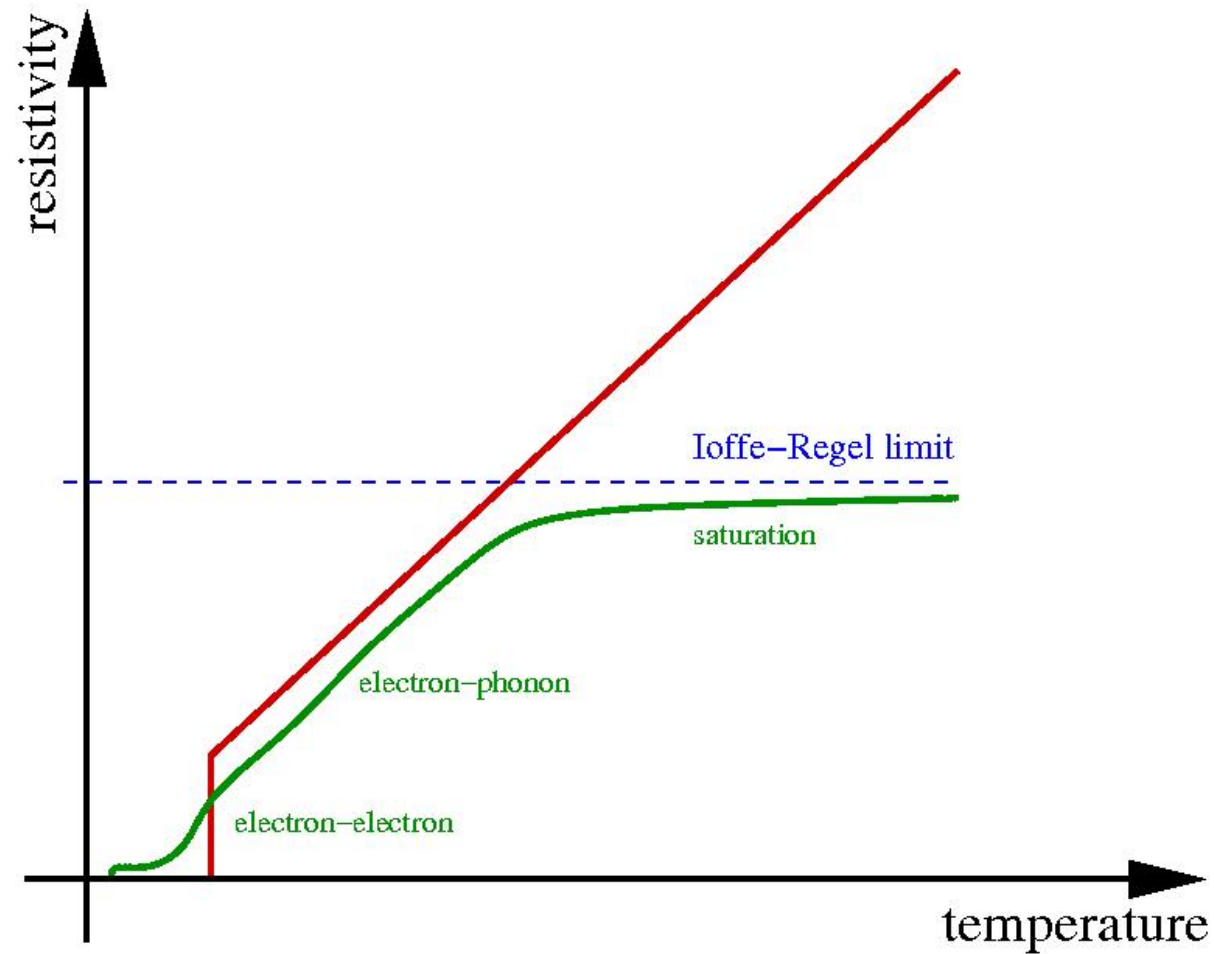


... which is good for superconductivity!

The quantum fog (Fermi gas) returns

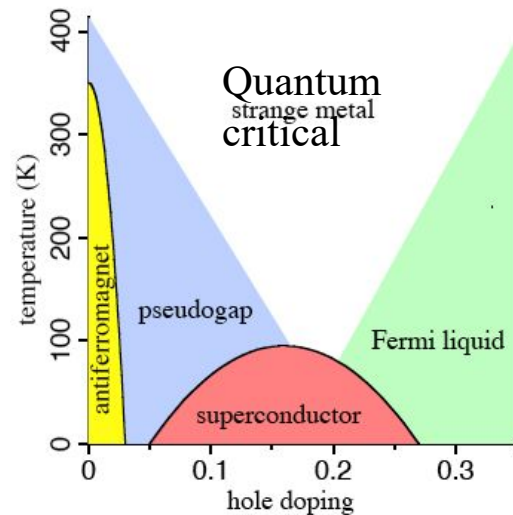


Divine resistivity

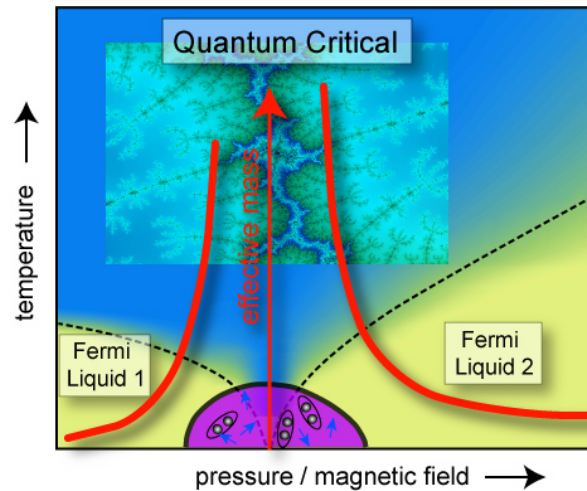


A universal phase diagram

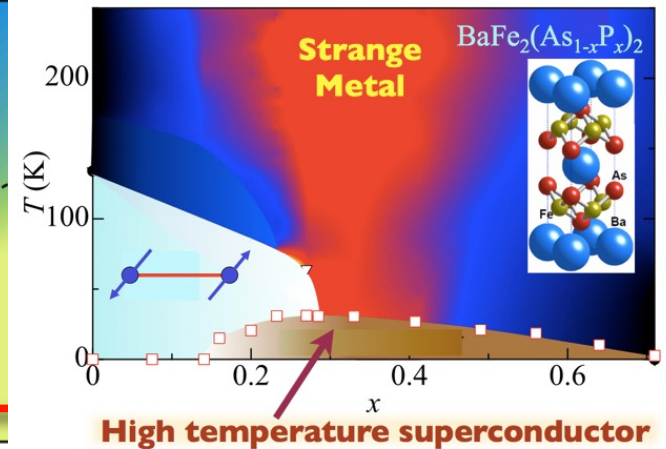
High T_c
superconductors



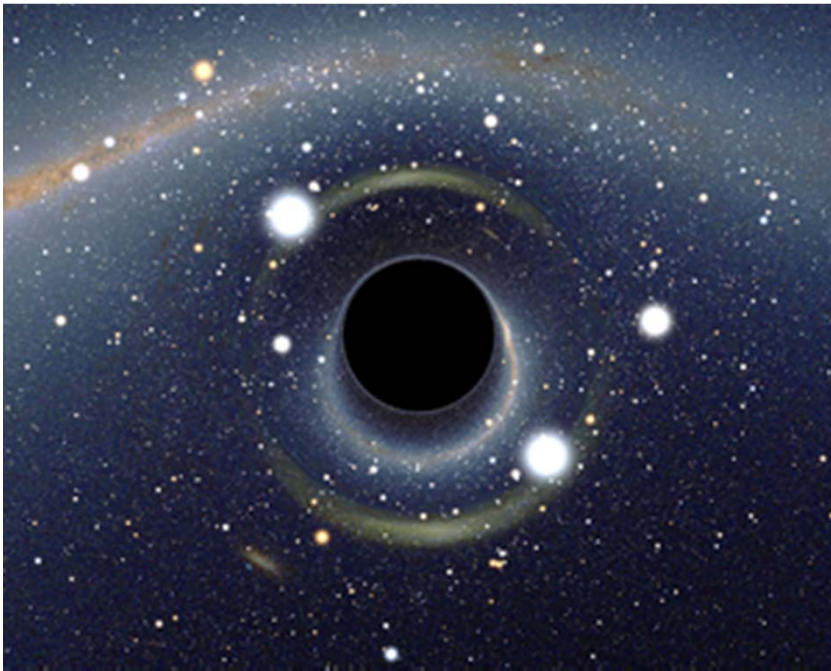
Heavy fermions



Iron
superconductors (?)

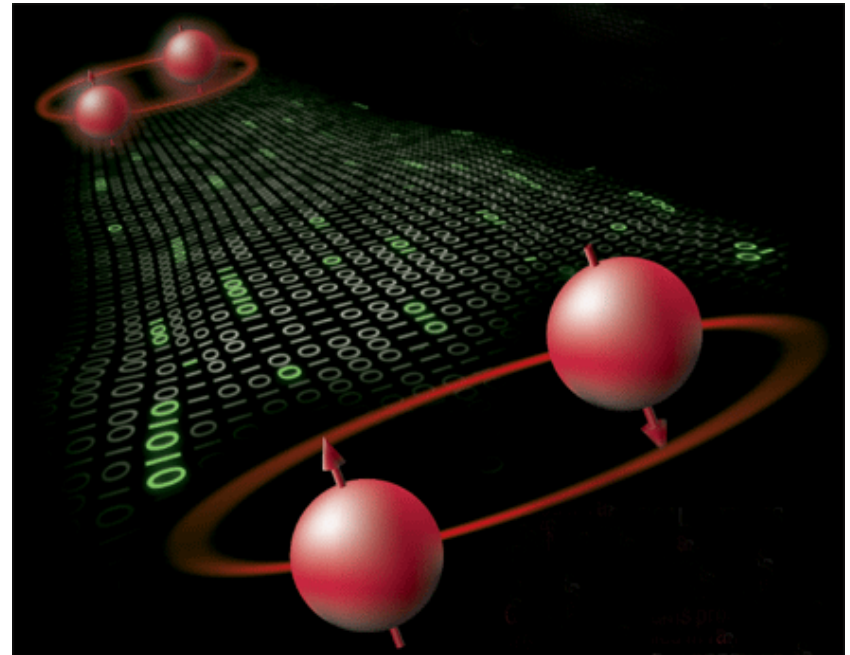


Black holes as “quantum matter computers” !?



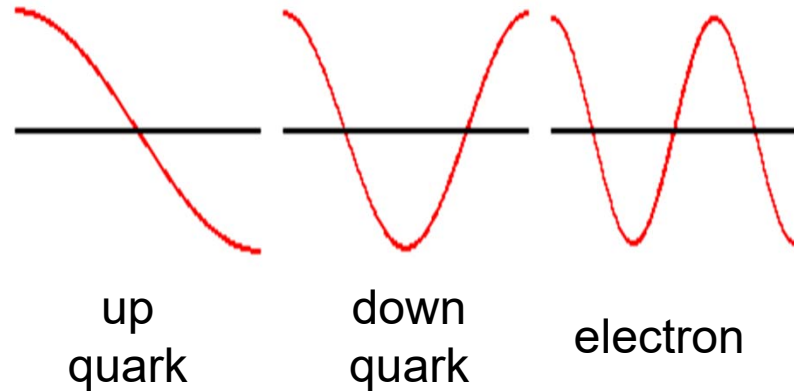
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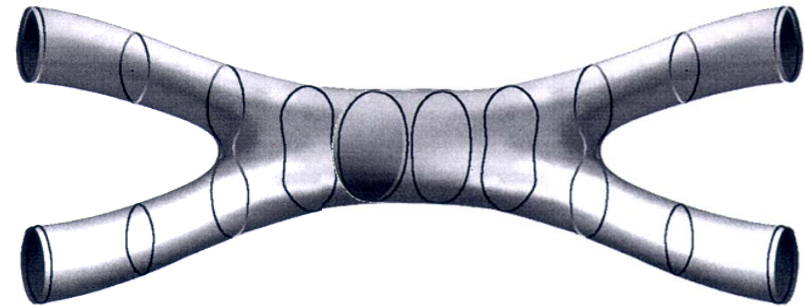


Particles as string vibrations (1980' s)

String vibration modes →
Different particles



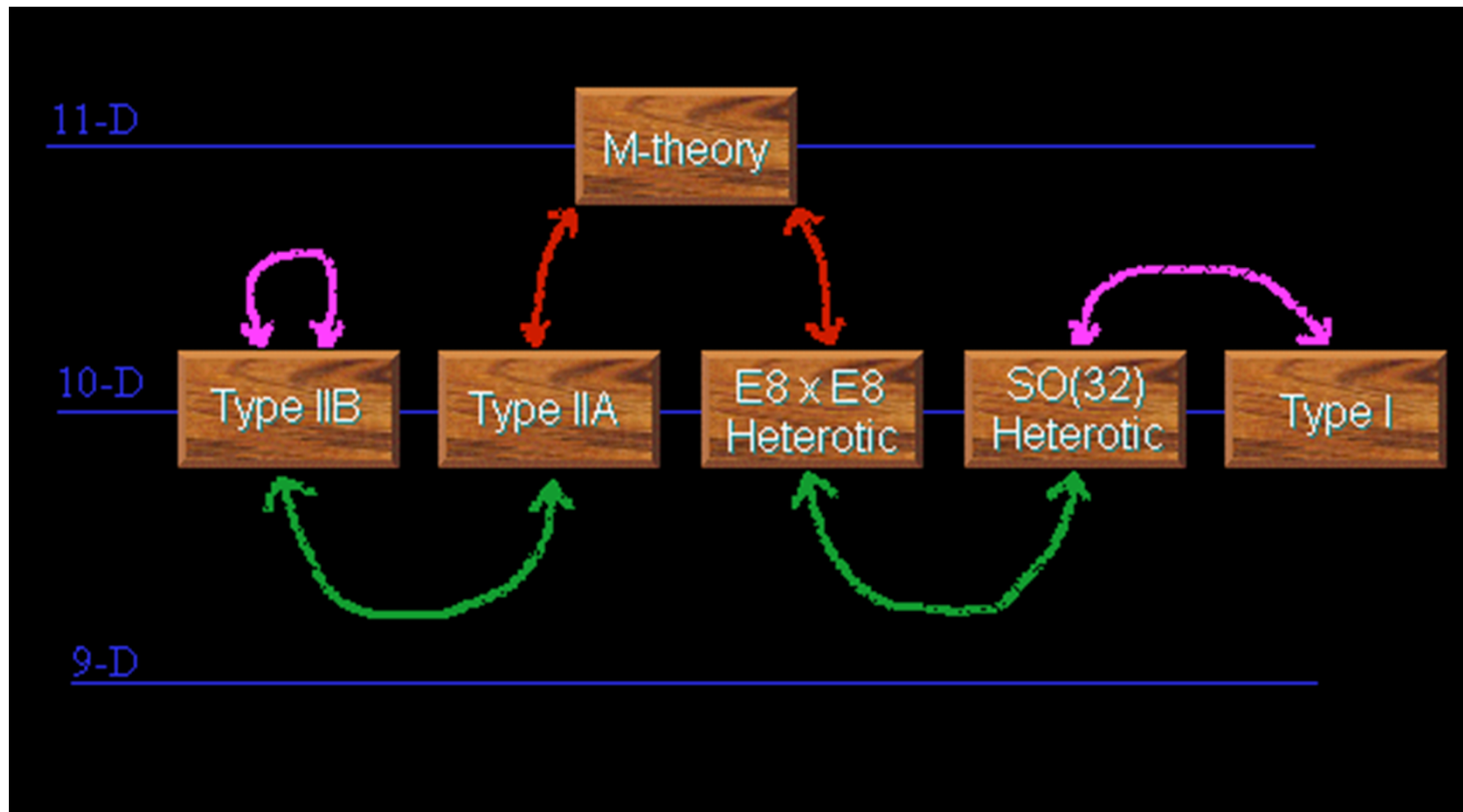
Morphing strings →
Particle interactions



=> Unified theory: one string = all particles

=> Vibrations of “closed strings” describe gravitons
(quantum particles carrying gravitational force).

The “second string revolution” (1995)

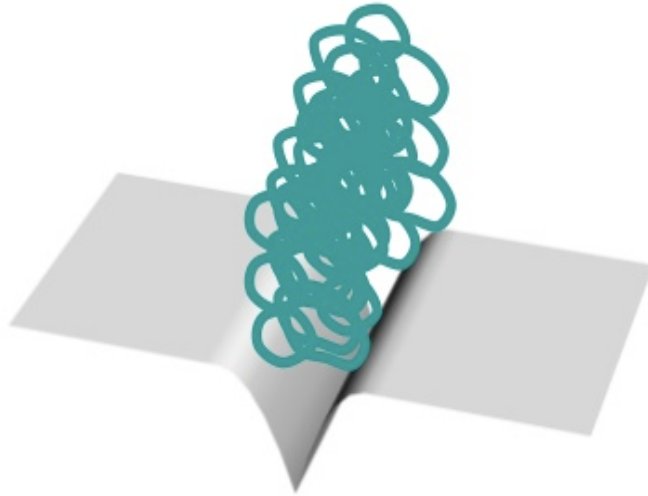


Dualities

AdS/CFT correspondence

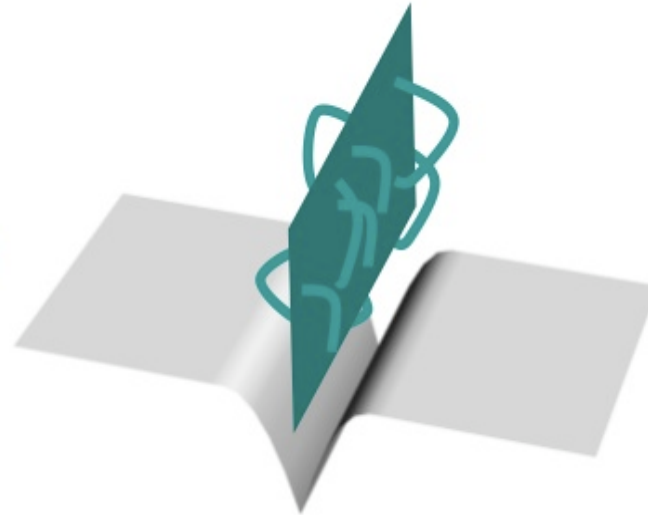
We have two different descriptions for same object!

Closed string
description



=

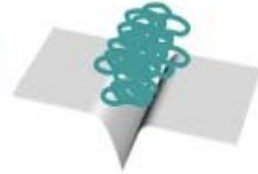
Open string
description



Especially, in the case of **D3-brane**, at low energy these two description will be approximated by

AdS/CFT correspondence

Closed string
description



=

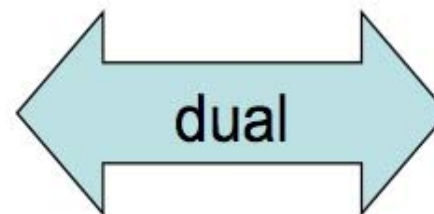
Open string
description



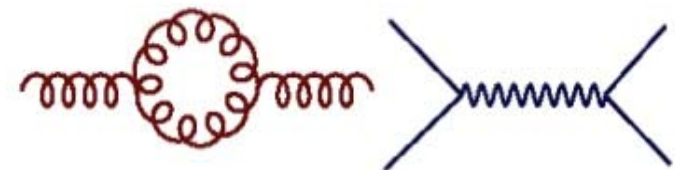
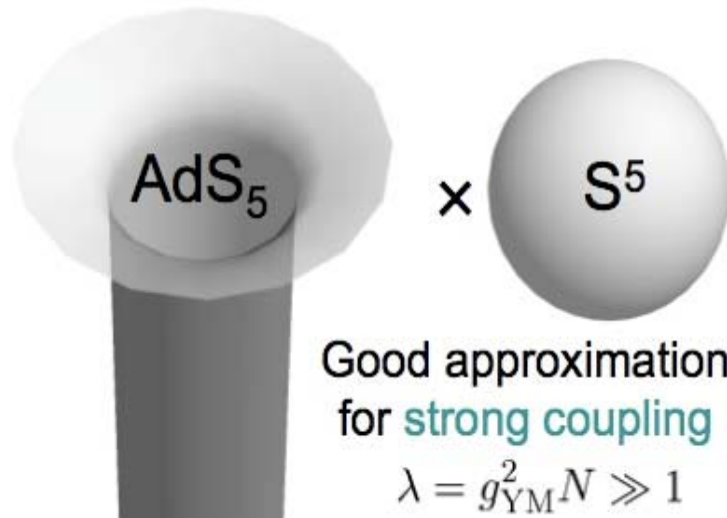
Low energy limit
with large N



Supergravity
on $\text{AdS}_5 \times S^5$



$D=4$ $\mathcal{N}=4$ $U(N)$
Super Yang-Mills

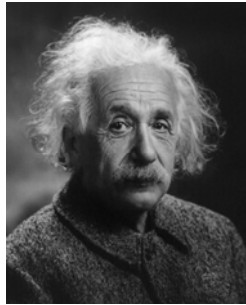


Good approximation
for **weak coupling**

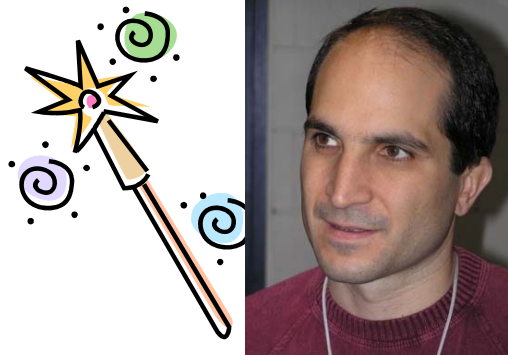
$$\lambda = g_{\text{YM}}^2 N \ll 1$$

General relativity “=” quantum field theory

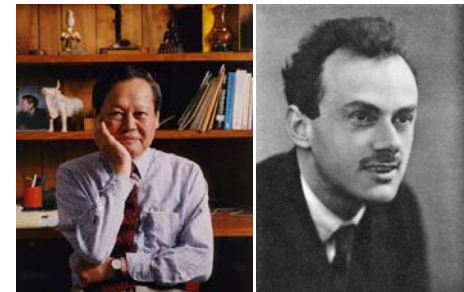
General relativity



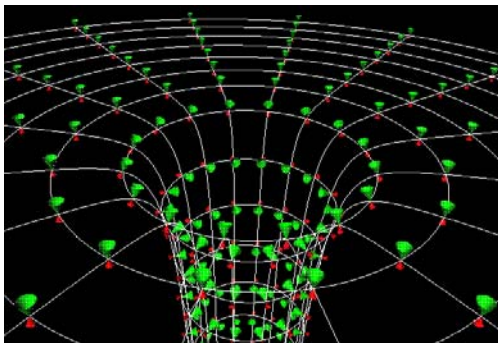
‘AdS/CFT’



Quantum fields



Maldacena 1997



=



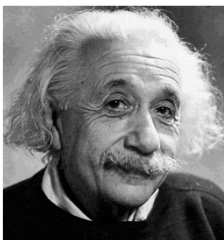
Holographic gauge-gravity duality

Einstein Universe “AdS”

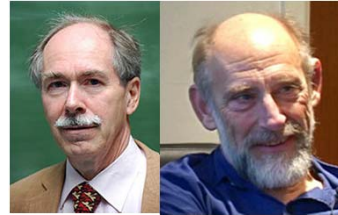


Classical general relativity

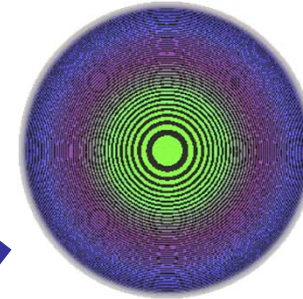
Uniqueness of GR solutions



Quantum field world “CFT”

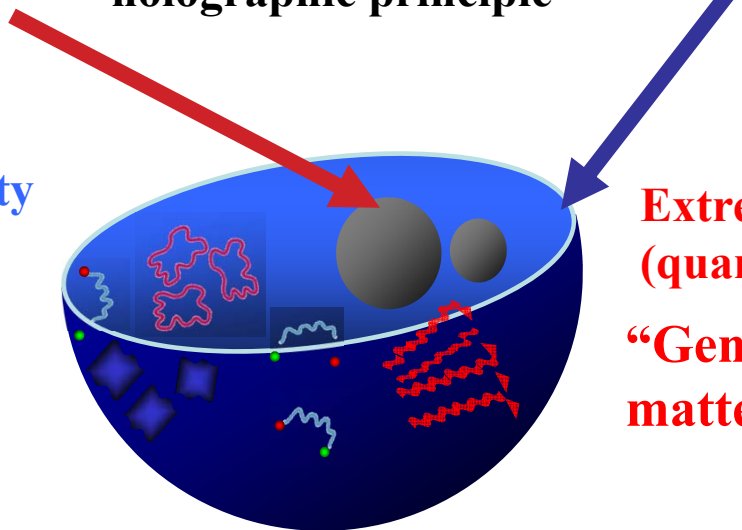


‘t Hooft-Susskind holographic principle



Extremely strongly coupled (quantum) matter

“Generating functional of matter emergence principle”



General Relativity = Renormalization Group

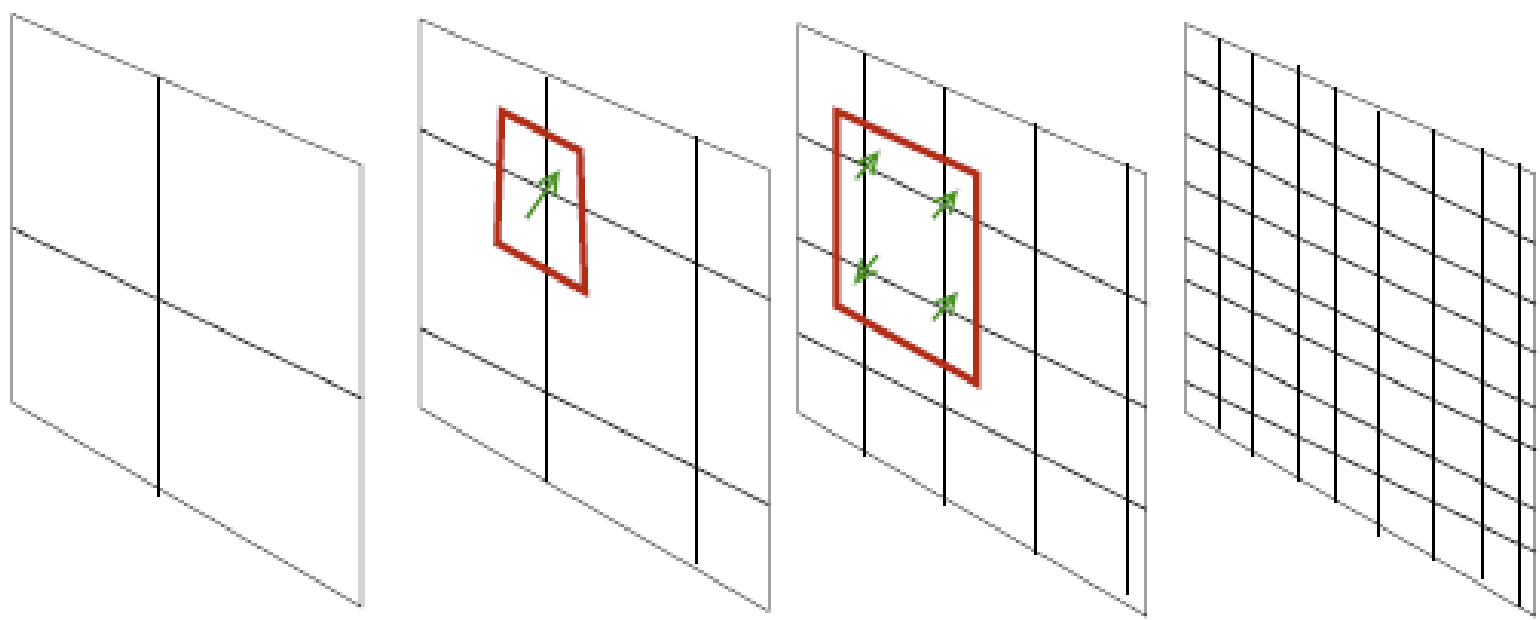


Extra radial dimension
of the bulk \Leftrightarrow scaling
“dimension” in the
field theory

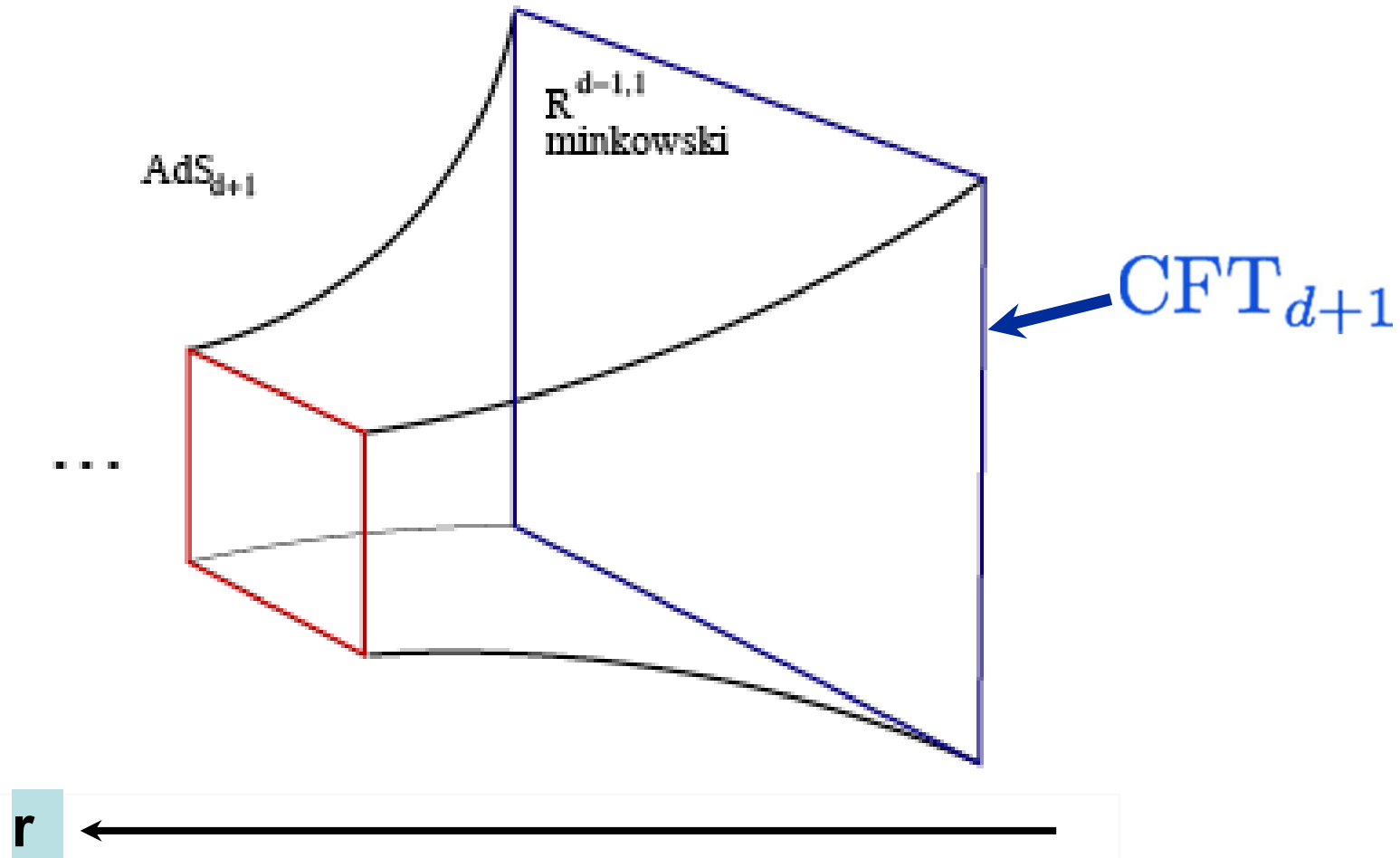
Bulk AdS geometry =
**scale invariance of
the field theory**

$$dr^2 = -F(r)dt^2 + \frac{dr^2}{F(r)} + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

$$F(r) = -\Lambda r^2 + 1, \quad \Lambda < 0$$



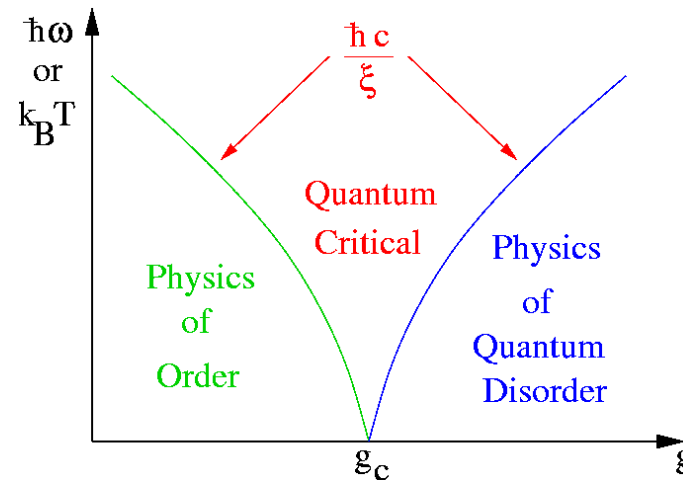
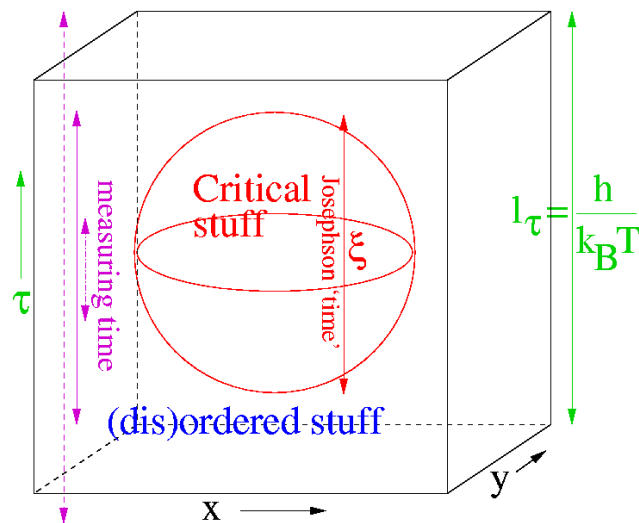
r ←



Quantum criticality.

Sachdev's book "quantum phase transitions"

Scale invariance of the quantum dynamics (in space and time) is dynamically generated, as emergent phenomenon.



In the higher dimensional (bosonic) quantum field theories which are understood this only happens at **isolated points in coupling constant space**.

GKPW rule: propagators in QFT are classical waves in AdS

$$Z_{\text{CFT}}(N) = \int \mathcal{D}\phi e^{iN^2 S_{\text{AdS}}(\phi)}$$

$$\langle e^{\int d^{d+1}x J(x)\mathcal{O}(x)} \rangle_{\text{QFT}} = \int \mathcal{D}\phi e^{iS_{\text{bulk}}(\phi(x,r))|_{\phi(x,r=\infty)=J(x)}}$$

$$g_{YM}^2 N = \frac{R^4}{\alpha} \quad g_{YM}^2 = g_s$$



Only in the *large N* limit the strongly coupled boundary field theory becomes dual to *classical* gravity!

WELCOME TO HELL

MATRIX LARGE N

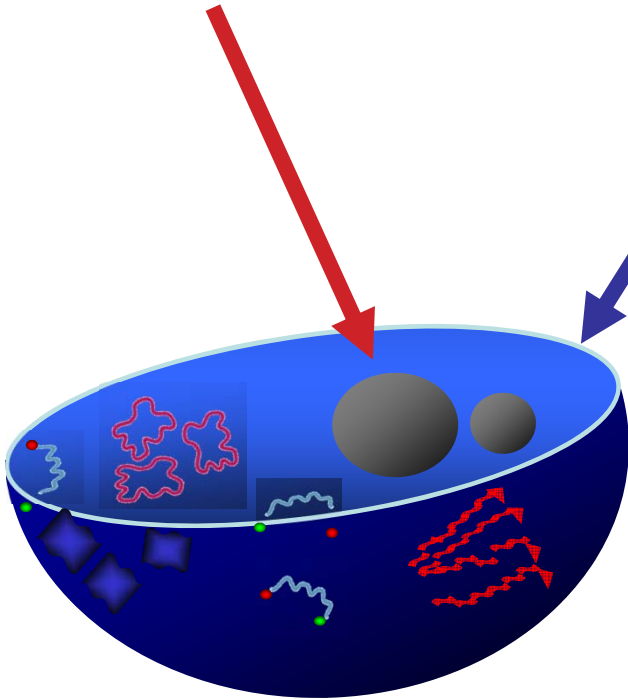
MATRIX LARGE N



UV INDEPENDENCE

The triumph: gravitational encoding of all thermal physics!

**Schwarzschild black hole
in the bulk**



**Boundary: the emergence theories of
finite temperature matter.**

- **All of thermodynamics!** Caveat: phase transitions are mean field (large N limit).

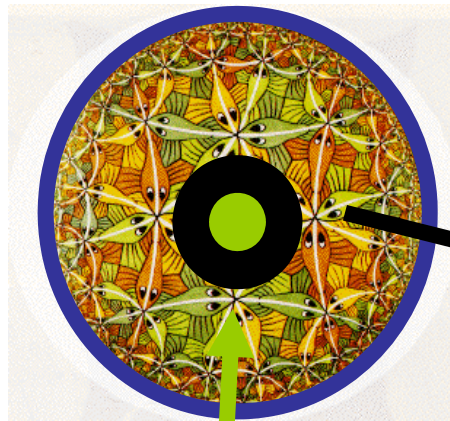
- **Precise encoding of Navier-Stokes hydrodynamics!** Right now used to debug complicated hydrodynamics (e.g. superfluids).

- **For special “Planckian dissipation” values of parameters** (quantum criticality):

$$\tau_{\hbar} = \text{const.} \frac{\hbar}{k_B T}, \quad \text{const.} = O(1)$$

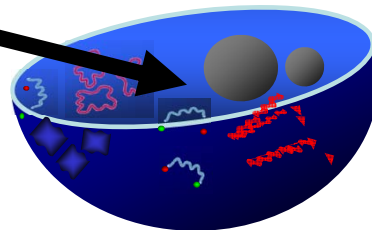
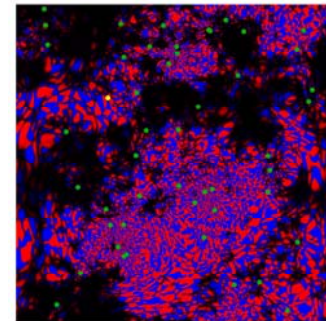
The charged black hole encoding for finite density (2008 - ????)

Anti de Sitter universe.

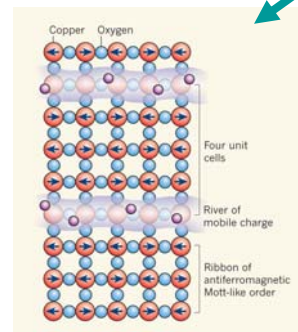


Charged black hole in the middle

Finite density **quantum matter:**



Holographic strange metals



Stripy pseudogap orders



High Tc superconductors



Emergent Fermi liquids

Turbulence and fractal black hole horizons.

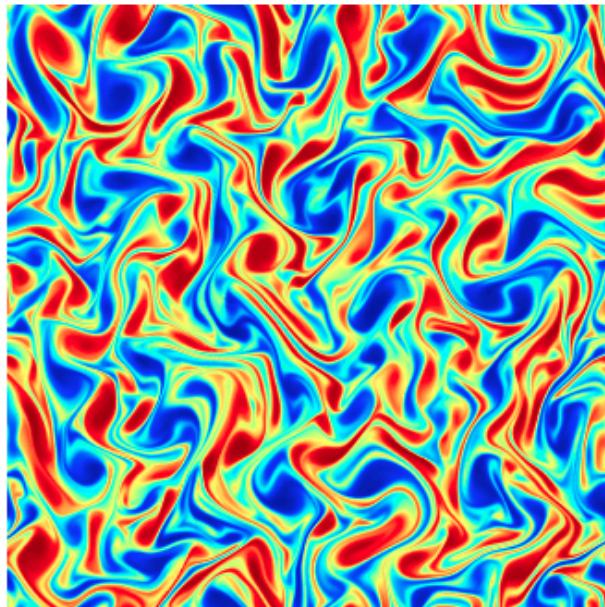


Chesler

Yaffe

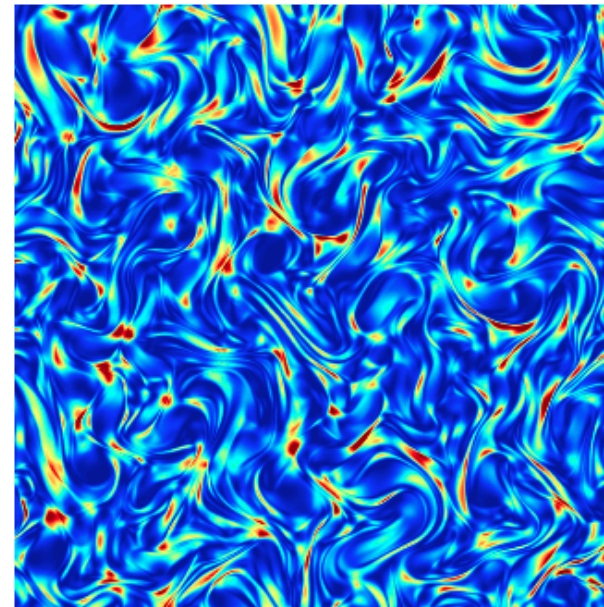
Holography and numerical GR

ω



**Vorticity in the liquid
(Kolmagorov scaling)**

$\partial A/\partial t$



**Near horizon geometry
(fractal)**

Dissipation = absorption of classical waves by Black hole!



Policastro-Son-Starinets (2002):

Viscosity: absorption cross section of gravitons by black hole

$$\eta = \frac{\sigma_{abs}(0)}{16\pi G}$$

= area of horizon (GR theorems)

**Entropy density s: Bekenstein-Hawking
BH entropy = area of horizon**

**Universal viscosity-entropy ratio for CFT' s
with gravitational dual limited in large N by:**

$$\frac{\eta}{s} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

Planckian dissipation ...

Scaling form dynamical susceptibility:

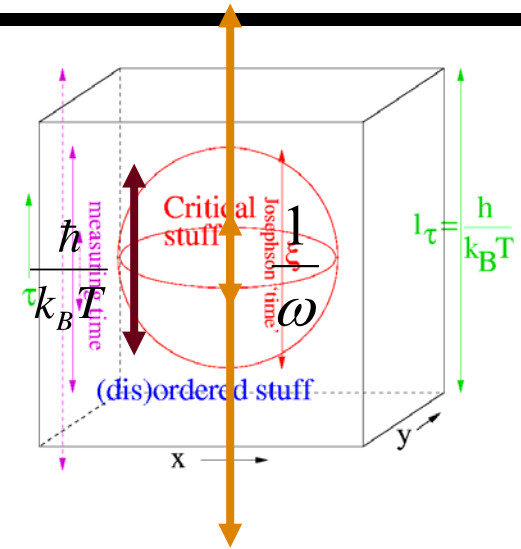
$$\chi(\omega) \propto \frac{1}{T^{2-\eta}} \Psi\left(\frac{\hbar\omega}{k_B T}\right)$$

Quantum critical regime $k_B T \gg \frac{\hbar c}{\xi}$

$$\hbar\omega \gg k_B T: \chi(\omega) \propto \frac{e^{i\frac{\pi}{2}(2-\eta)}}{|\omega|^{2-\eta}}$$

$$\hbar\omega \ll k_B T: \chi(\omega) \propto \frac{1}{T^{2-\eta}} \frac{1}{1 - i\omega\tau_{\hbar}}$$

Planckian dissipation: $\tau_{\hbar} = \text{const.} \frac{\hbar}{k_B T}, \text{const.} = O(1)$



Quantum criticality and the dimension of viscosity ...

Viscosity: $\eta = f\tau_K$

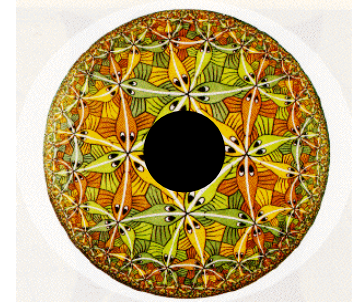
Free energy density QC system: $f = sT$

Planckian dissipation: $\tau_K = A \frac{\hbar}{k_B T}$

$$\frac{\eta}{s} = AT \frac{\hbar}{k_B T} = A \frac{\hbar}{k_B}$$

Large N SUSY Yang Mills: $A = \frac{1}{4\pi}$

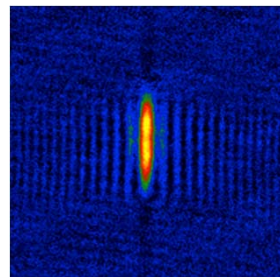
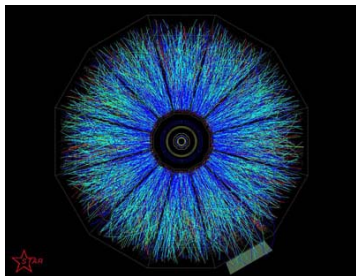
Planckian dissipation



Universal entropy production time in QC system: $\tau = \tau_{\hbar} \approx \frac{\hbar}{k_B T}$

Observed in Quark gluon plasma (heavy ion colliders RIHC, LHC) and cold atom “unitary fermi gas”:

$$\frac{\eta}{s} = T \tau_{\hbar} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$



Science March 4 2016:

PHYSICS

Electrons go with the flow in exotic material systems

Electronic hydrodynamic flow—making electrons flow like a fluid—has been observed

By Jan Zaanen

REPORTS

ELECTRON TRANSPORT

Negative local resistance caused by viscous electron backflow in graphene

D. A. Bandurin,¹ I. Torre,² R. Krishna Kumar,^{1,3} M. Ben Shalom,^{1,4} A. Tomadin,⁵ A. Principi,⁶ G. H. Auton,³ E. Khestanov,^{1,4} K. S. Novoselov,¹ L. V. Grigorieva,⁷ L. A. Ponomarenko,^{1,3} A. K. Geim,^{1,4} M. Polini^{1*}

the space they are moving in is made in-

ELECTRON TRANSPORT

Observation of the Dirac fluid and the breakdown of the Wiedemann-Franz law in graphene

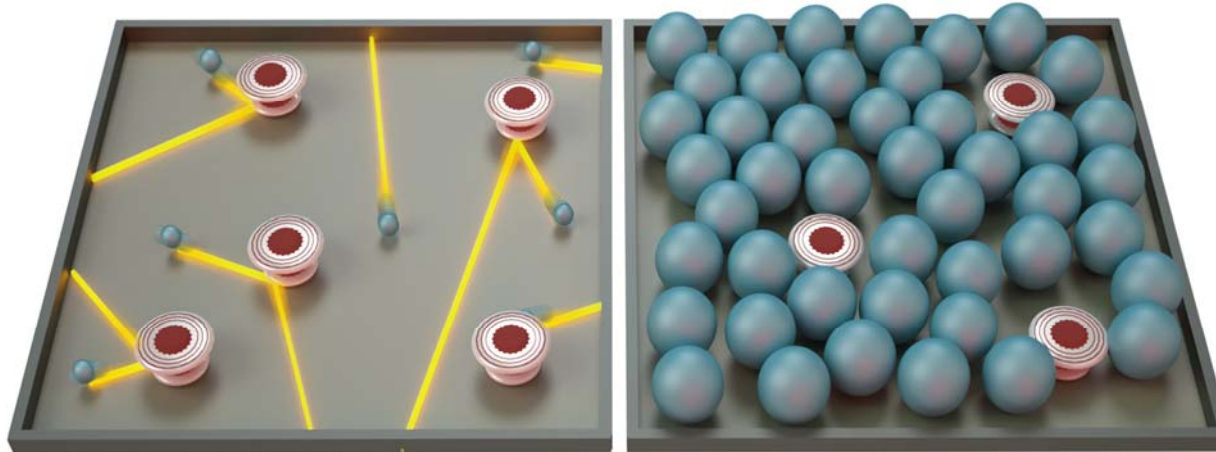
Jesse Crossno,^{1,2} Jing K. Shi,¹ Ke Wang,¹ Xiaomeng Liu,¹ Achim Harzheim,¹ Andrew Lucas,¹ Subir Sachdev,^{1,2} Philip Kim,^{1,2*} Takashi Taniguchi,⁴ Kenji Watanabe,⁴ Thomas A. Ohki,² Kin Chung Fong^{2*}

ELECTRON TRANSPORT

Evidence for hydrodynamic electron flow in PdCoO₂

Philip J. W. Moll,^{1,2,3} Pallavi Kushwaha,³ Nabhanila Nandi,³ Burkhard Schmidt,³ Andrew P. Mackenzie^{3,4*}

Hydro-electricity exhibition in Science ..



Science first march issue:

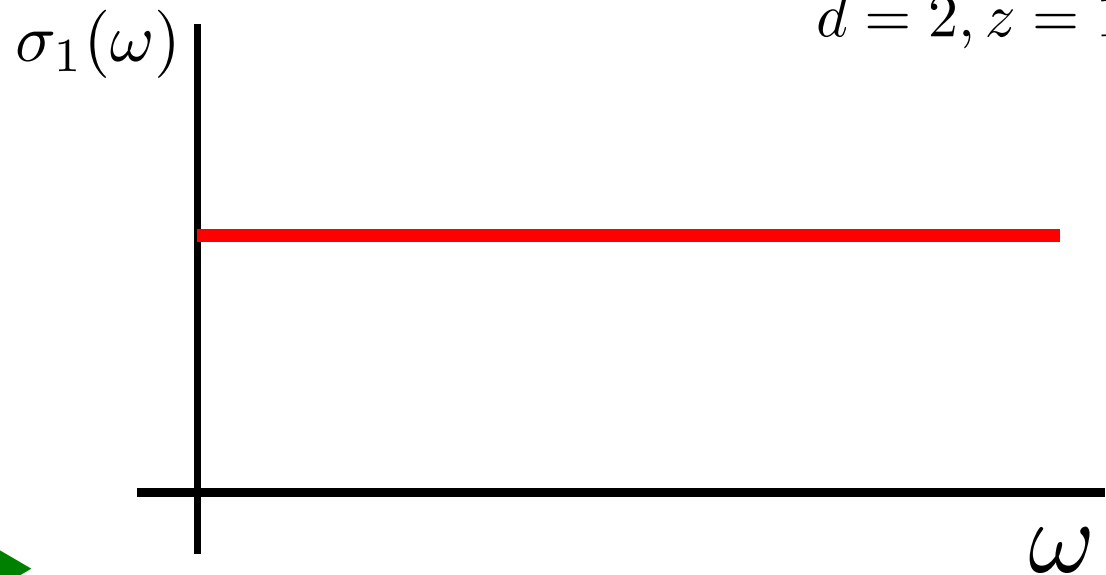
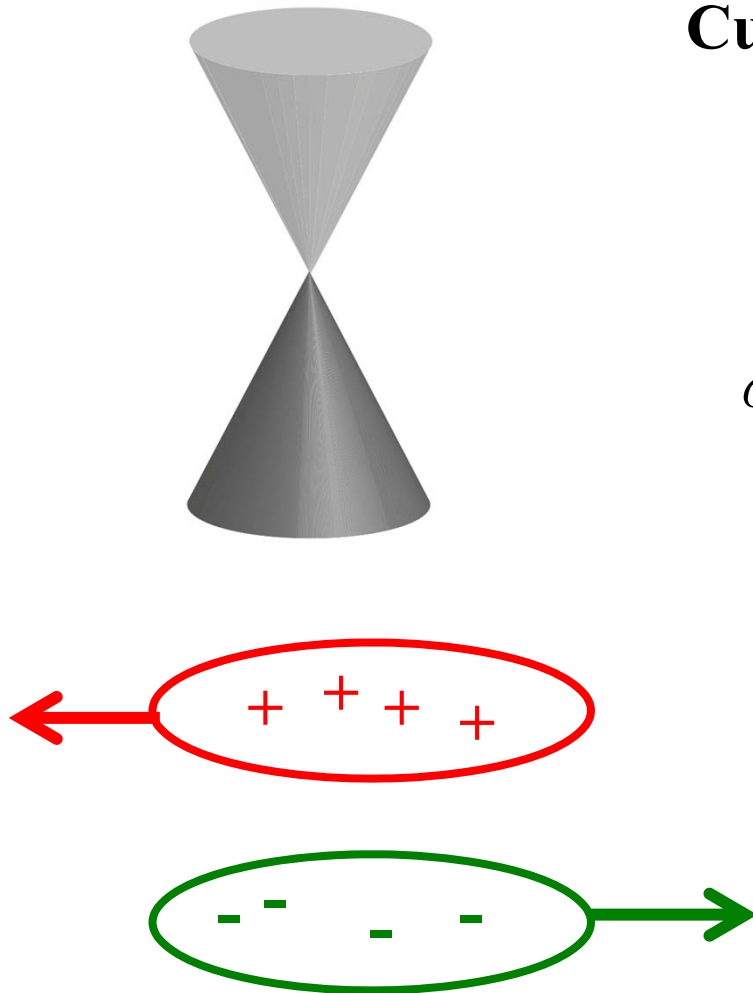
- 1. JZ: Perspective**
- 2. Crossno et al. (Harvard): WF violation in graphene.**
- 3. Bandurin et al. (Manchester): “whirls” in graphene.**
- 4. Moll et al. (Dresden): electron “pipes” in PdCoO₂**

Optical conductivity: zero density CFT.

Current does not carry momentum

$$\sigma_1(\omega) \sim \omega^{(d-2)/z}$$

$$d = 2, z = 1$$



Graphene as an interacting Dirac fluid.

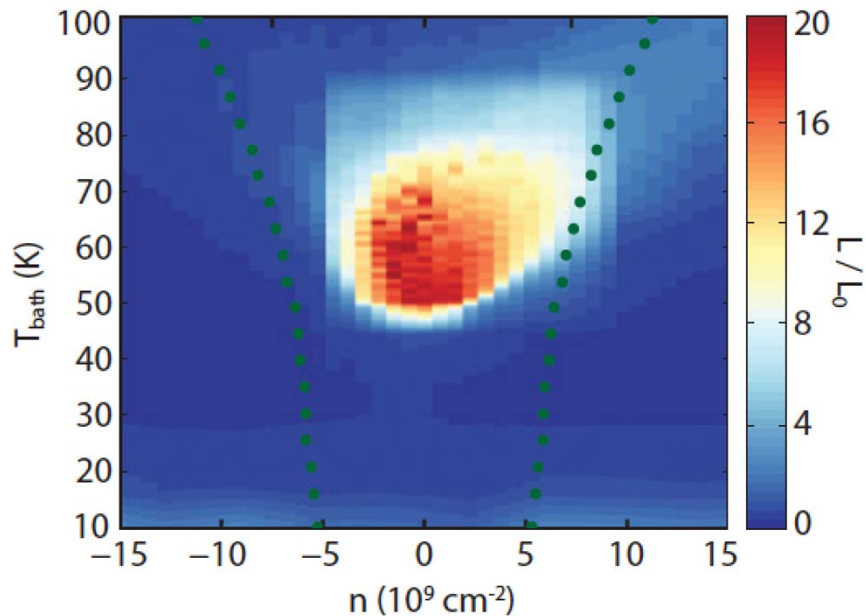
Crossno et al., arXiv:1509.04713

Wiedemann-Franz ratio:

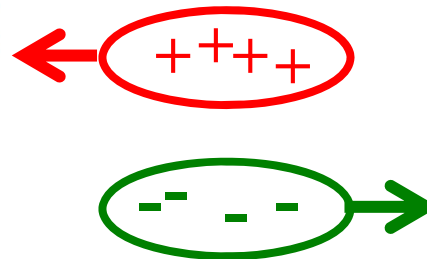
$$\mathcal{L} = \frac{\kappa_e}{\sigma T^2}$$

Lorentz-ratio (quasiparticles):

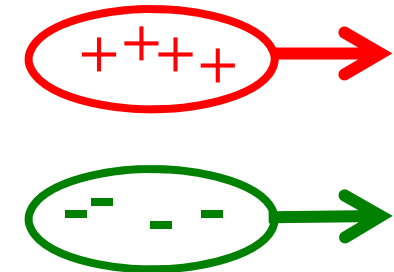
$$\mathcal{L}_0 = \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2$$



Electric field



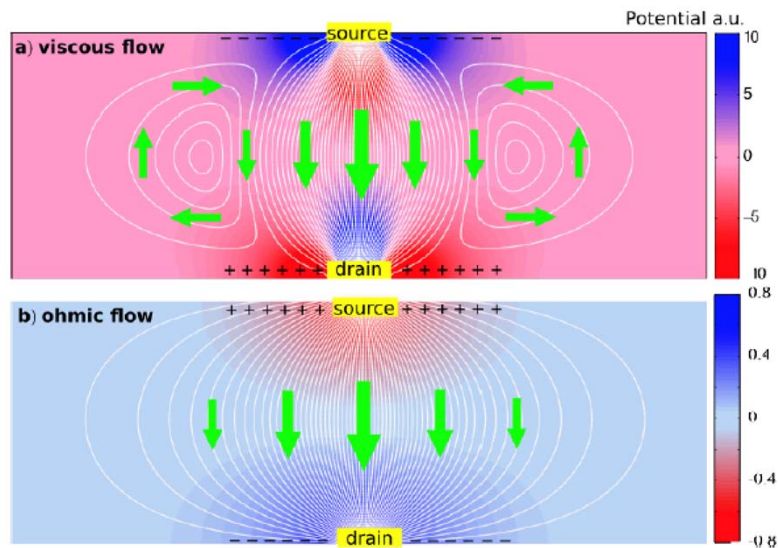
Temp. gradient



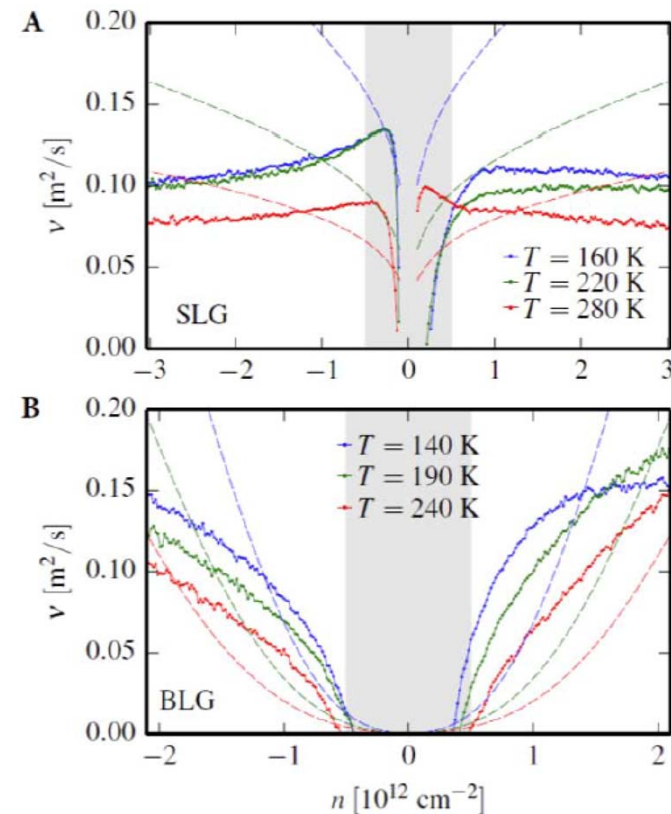
Lucas et al. (arXiv:1510.01738): “minimalish” viscosity

Graphene's Hydrodynamical whirls and the negative local resistance .

Prediction: Levitov, Falkovich, arXiv:1508.00836

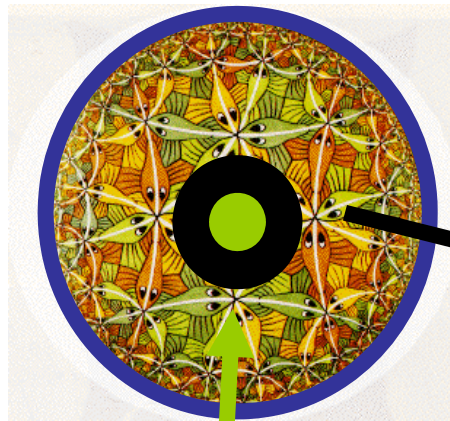


Confirmation: Geim group, arXiv:1509.04165



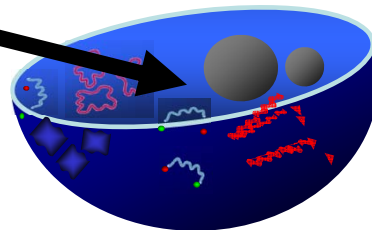
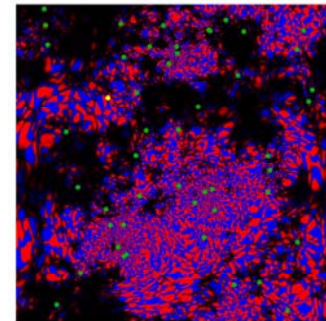
The charged black hole encoding for finite density (2008 - ????)

Anti de Sitter universe.

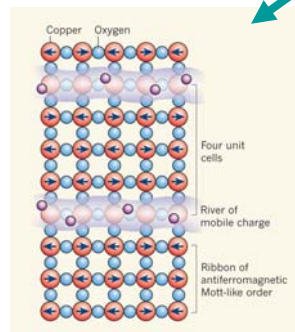


Charged black hole in the middle

Finite density **quantum matter:**



Holographic strange metals



Stripy pseudogap orders



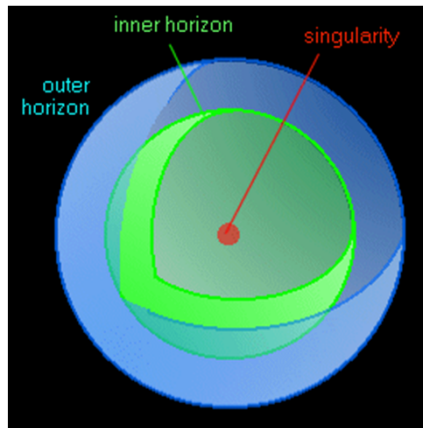
High Tc superconductors



Emergent Fermi liquids

The holographic stable states: uncollapsing in an AdS “star”.

(Reissner-Nordstrom)
“Black hole like object”



“fractionalized”, “unstable”:
strange metal

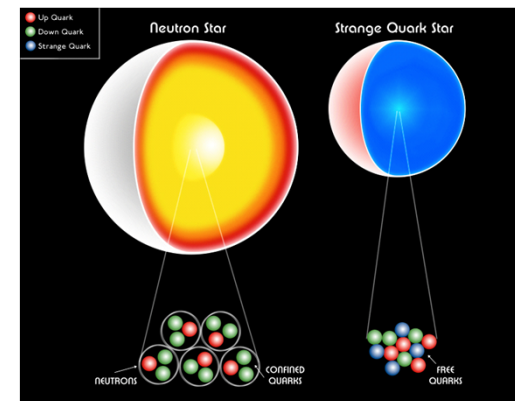
“uncollapse”



Phase transition



“star like object”



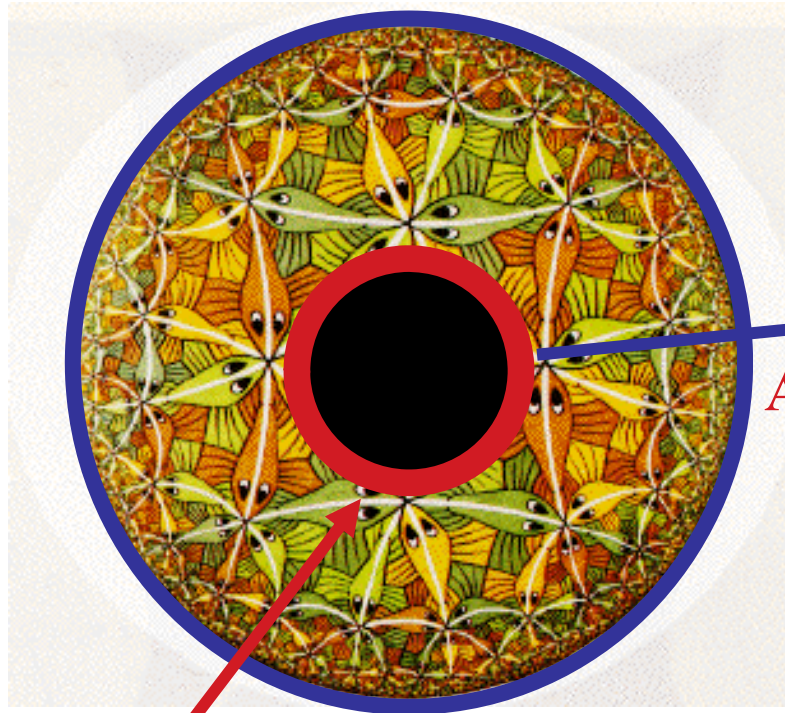
“Cohesive state”:

Symmetry breaking:
superconductor, crystal
 (“scalar hair”)

Fermi-liquid (“electron star”)

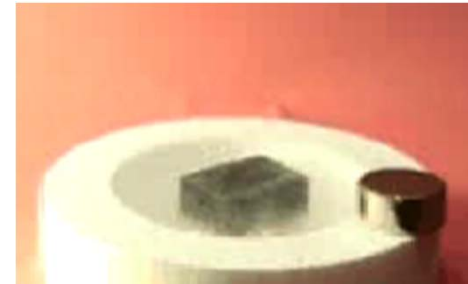
The holographic superconductor

Hartnoll, Herzog, Horowitz, arXiv:0803.3295



(Scalar) matter
'atmosphere'

Condensate (superconductor,
...) on the boundary!



AdS-CFT

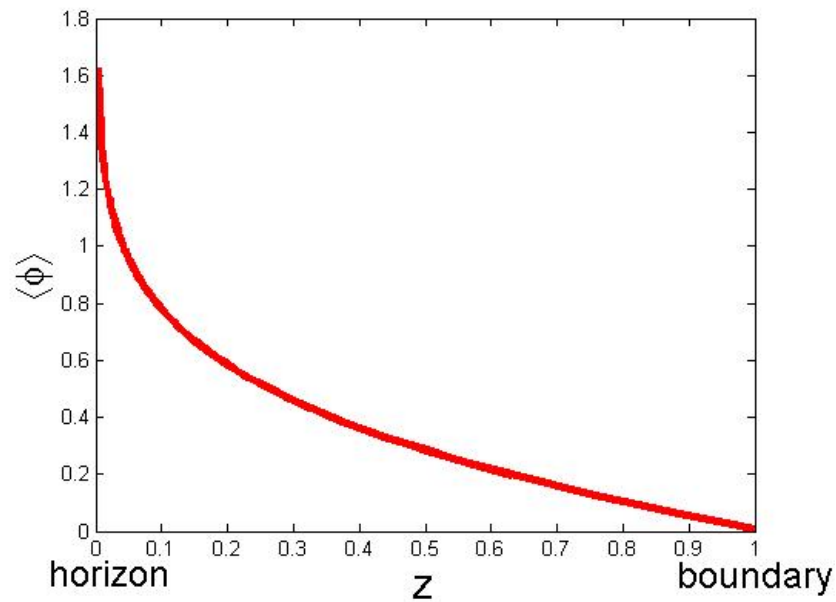
'Super radiance' : in the
presence of matter the
extremal BH is unstable =>
zero T entropy always
avoided by low T order!!!

The Bose-Einstein Black hole hair

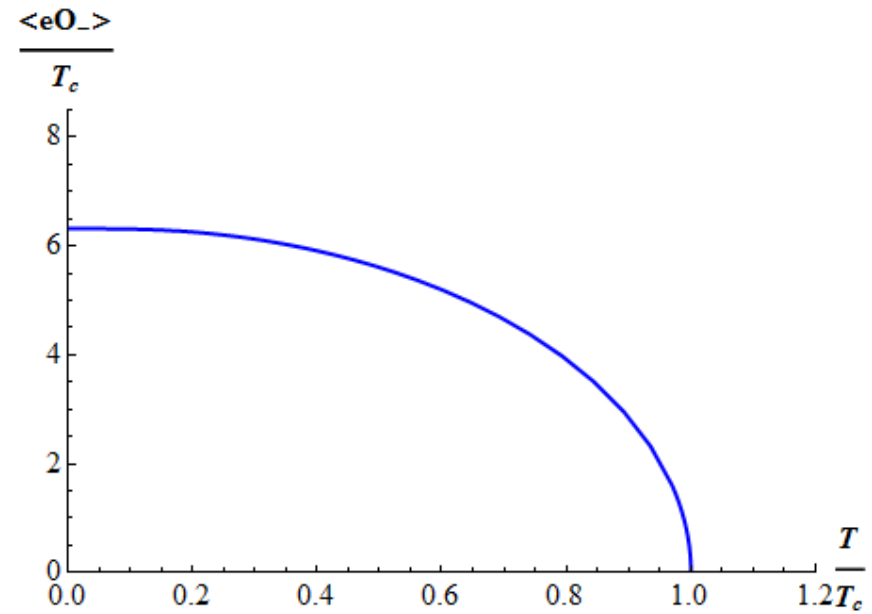


Hartnoll Herzog Horowitz

Scalar hair accumulates at the horizon

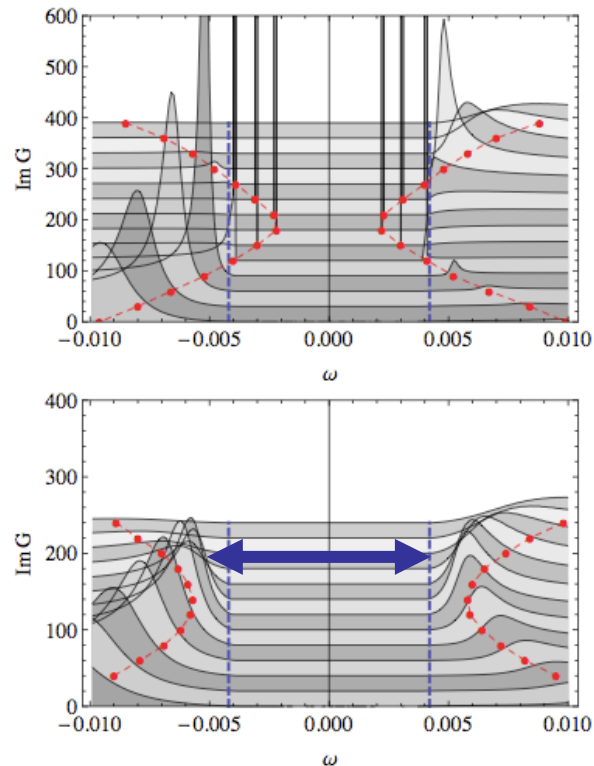


Mean field thermal transition.

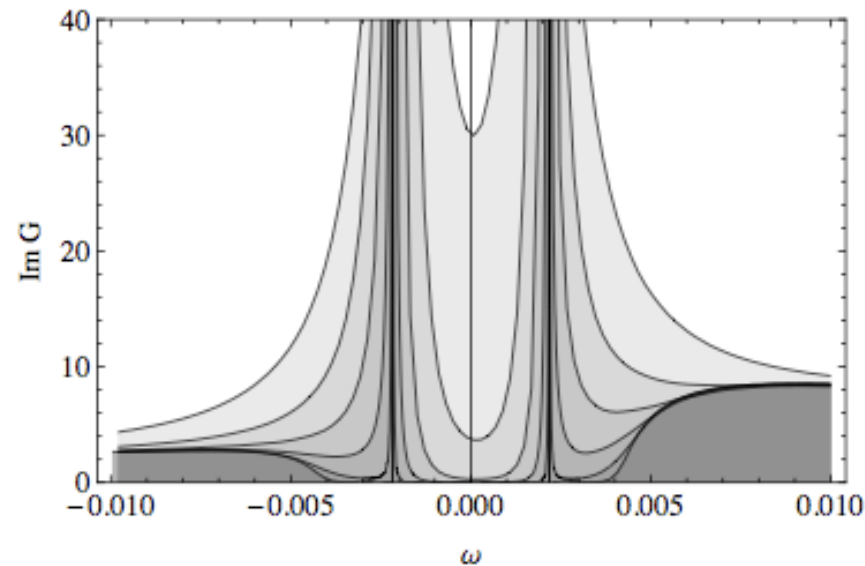


Holographic superconductivity: stabilizing quasiparticles.

Fermion spectrum for scalar-hair black hole (Faulkner et al., 911.340):

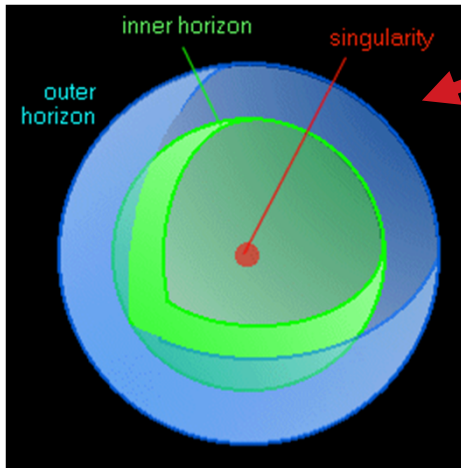


‘BCS’ Gap in fermion spectrum !!



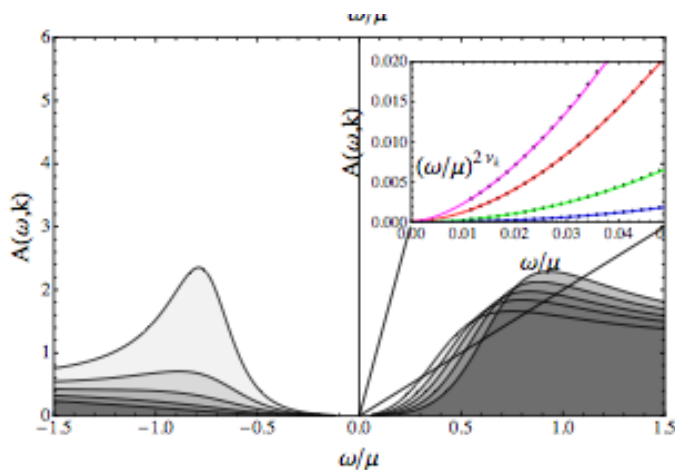
“Pseudogap” Temperature dependence

Finite density: the Reissner-Nordstrom strange metals (Liu et al.).



Near-horizon geometry of the extremal RN black hole:

- **Space directions: flat**, codes for **simple Galilean invariance** in the boundary.
- **Time-radial(=scaling) direction: emergent AdS₂**, codes for **emergent temporal scale invariance!**



Fermion spectral functions:

$$A(k, \omega) \propto G'_{AdS_2}(k, \omega) \propto \omega^{2\nu_k}$$

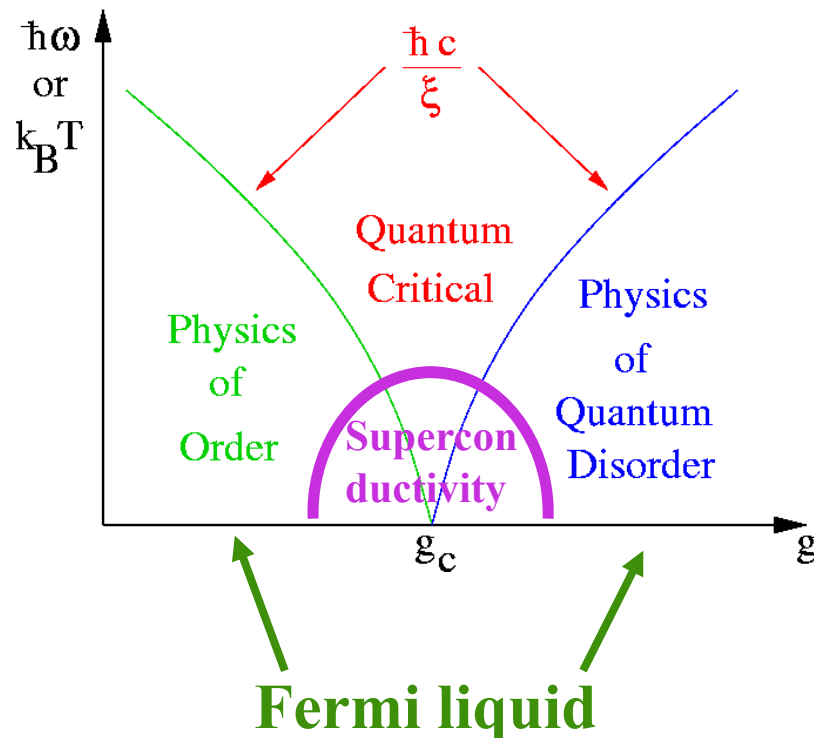
$$\nu_k = \frac{1}{\sqrt{6}} \sqrt{k^2 + \frac{1}{\xi^2}}$$

“Un-particle physics!”

Hertz-Millis metallic “Quantum Critical Point”

Assertion: in the “UV” a *Fermi-liquid* is formed co-existing with an electronic *order parameter* (e.g. magnet) interacting via a Yukawa coupling.

The *order parameter* is subjected to a *bosonic quantum phase transition*: always isolated unstable fixed point (stat phys rule book).



Electrons: *Fermi-gas* = heat bath *damping bosonic critical fluctuations*.

Lingering singularities of QP *on the Fermi surface* due to critical bosons.

“Critical Fermi-surface”: Fermi surface survives but branch cut fermionic propagators.

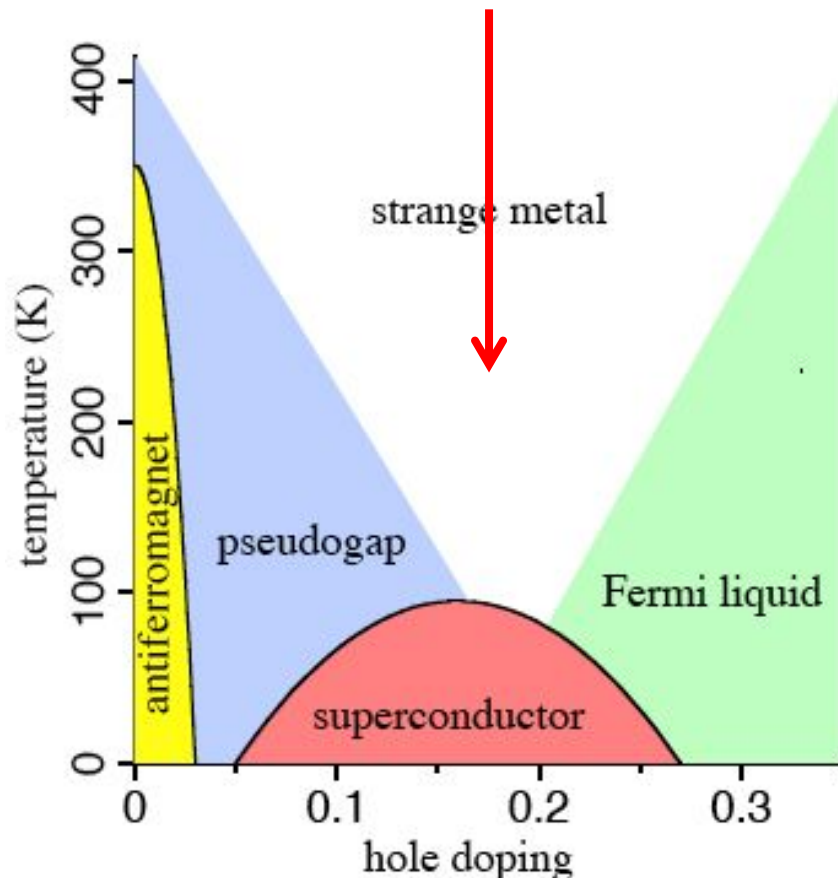
Critical fluctuations acting as pairing glue: the “superconducting domes”.

The unstable conformal metal of finite density holography.



Hong Liu

Conformal metal: *quantum critical fermionic phase not requiring fine tuning!*



Characterised by ***non-bosonic scaling properties***: large (infinite) z , hyperscaling violation, ...

Intrinsically ***very unstable***: “**mother**” of Fermi-liquids, superconductors, stripes, CDW’s, loop currents, electronic nematics, ...

The finite temperature state: governed by ***Planckian dissipation***.

“Scaling atlas” of holographic quantum critical phases.



Kiritsis

Deep interior geometry sets the scaling behavior in the emergent deep infrared of the field theory. Uniqueness of GR solutions:

1. “Cap-off geometry” = confinement: conventional superconductors, Fermi liquids

2. Geometry survives: “hyperscaling violating geometries” (Einstein – Maxwell – Dilaton – Scalar fields – Fermions).

$$ds^2 = \frac{1}{r^2} \left(-\frac{dt^2}{r^{2d(z-1)/(d-\theta)}} + r^{2\theta/(d-\theta)} dr^2 + dx_i^2 \right)$$

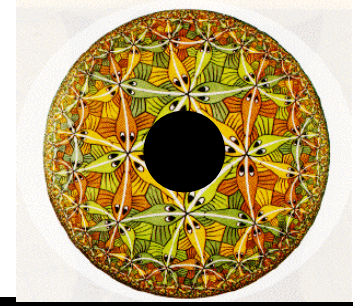
$$x_i \rightarrow \zeta x_i, \quad t \rightarrow \zeta^z t, \quad ds \rightarrow \zeta^{\theta/d} ds$$

$$S \propto T^{(d-\theta)/z}$$

Quantum critical phases with unusual values of:

- $z =$ Dynamical critical exponent
- $\theta =$ Hyperscaling violation exponent
- $\zeta =$ Charge exponent

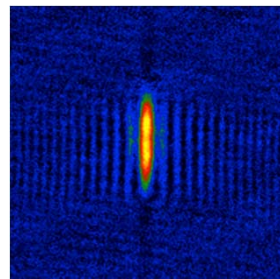
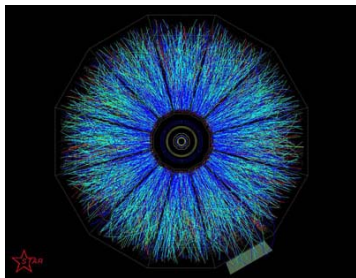
Planckian dissipation



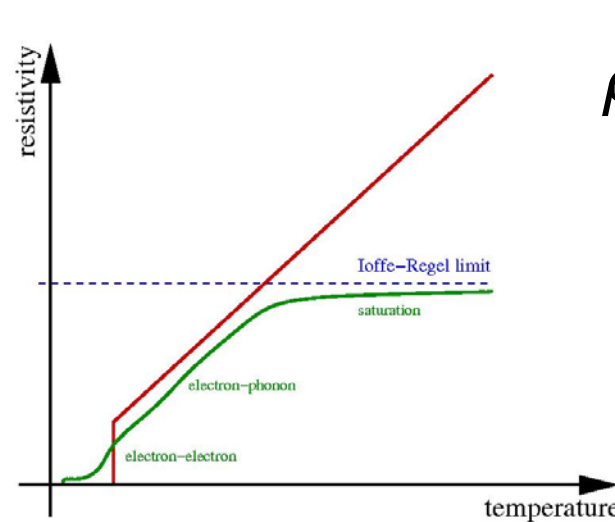
Universal entropy production time in QC system: $\tau = \tau_{\hbar} \approx \frac{\hbar}{k_B T}$

Observed in Quark gluon plasma (heavy ion colliders RIHC, LHC) and cold atom “unitary fermi gas”:

$$\frac{\eta}{s} = T \tau_{\hbar} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

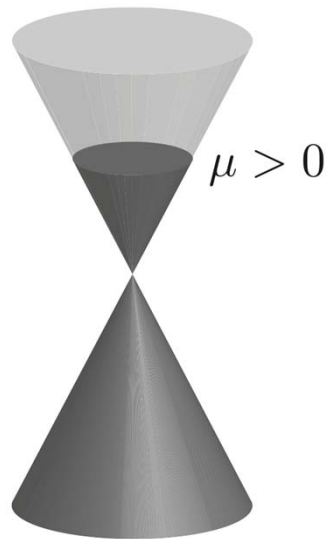


Since early 1990’ s recognized as responsible for strange metal properties, e.g. linear resistivity high Tc metals:

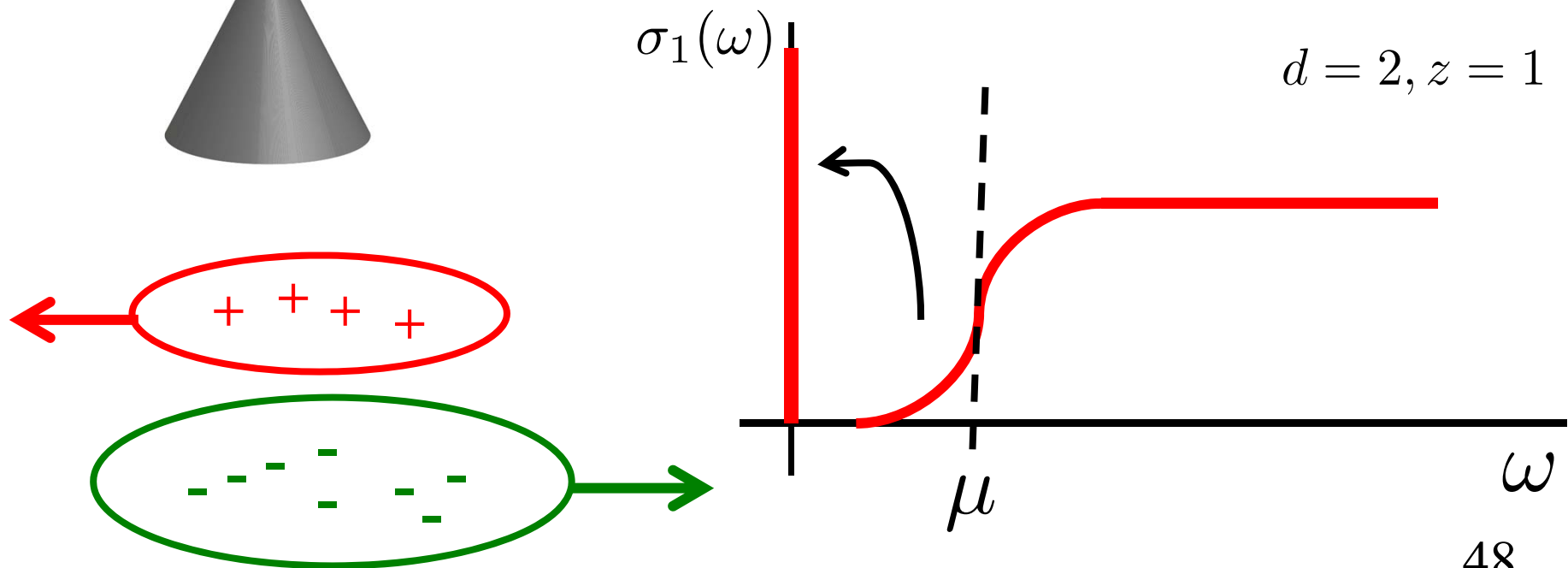


$$\rho \propto \frac{1}{\tau_{\hbar}} \propto k_B T$$

Optical conductivity: finite density *particle* system.



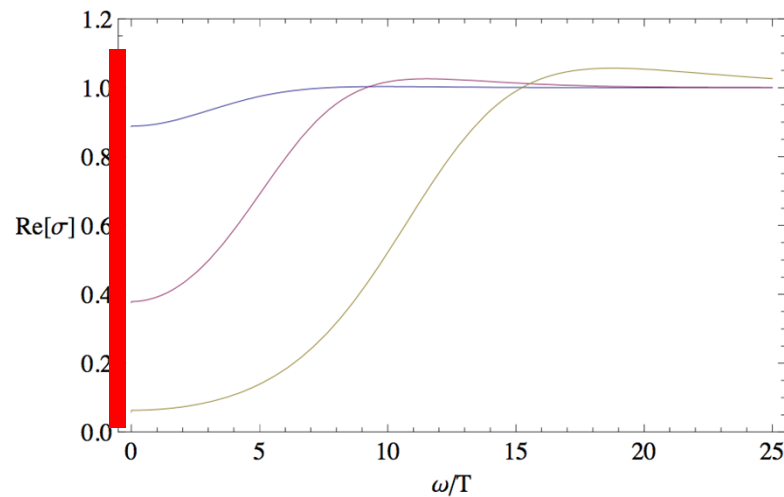
Current carries momentum: in Galilean continuum the weight below the chemical potential accumulates in the “perfect metal”
Drude peak.



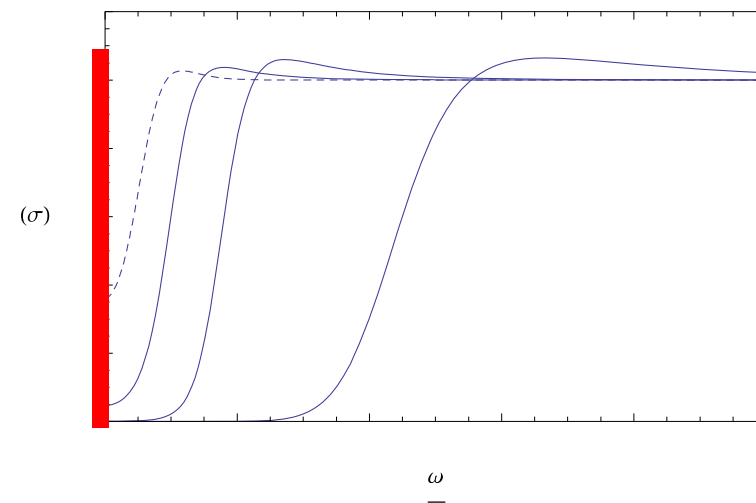
Holographic optical conductivity (2+1D).

Optical conductivity in finite density systems:

Strange metal.

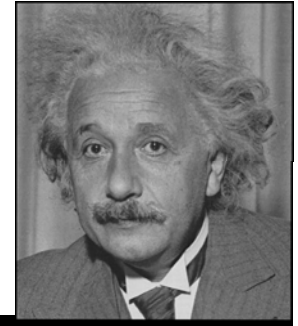


Holographic superconductor.



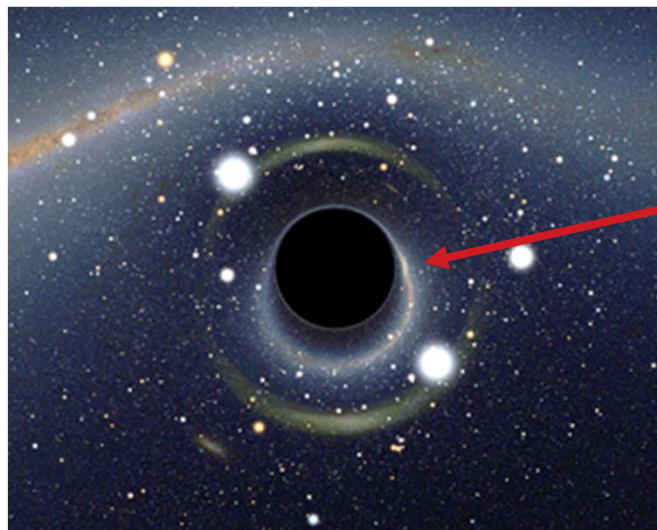
“Gapping” as for a BCS superconductor!

Black holes with a corrugated horizon



Charged Black Hole: describes finite density strange metal .

Breaking translational symmetry in the boundary:



Corrugate the black hole horizon

Not a favorite thing of general relativity -- hard work, still in progress!

Holographic quenched disorder.



David Vegh

Dictionary entry “**number one**”:

Global **translational invariance** in the boundary (energy-momentum conservation)



General covariance in the bulk (Einstein theory)

Breaking of Galilean invariance in the boundary = elastic scattering (?)



Fix the (spatial) frame in the bulk = “**Massive gravity**”

$$S = \frac{1}{2\kappa_4^2} \int d^4x \sqrt{-g} \left(\mathcal{R} + \frac{6}{L^2} - \frac{L^2}{4} F_{\mu\nu} F^{\mu\nu} + m^2 \left(\alpha \text{Tr}(\mathcal{K}) + \beta \left(\text{Tr}(\mathcal{K})^2 - \text{Tr}(\mathcal{K}^2) \right) \right) \right)$$

$$\mathcal{K}_{\alpha}^{\mu} \mathcal{K}_{\nu}^{\alpha} = g^{\mu\alpha} f_{\alpha\nu}$$

Couple the metric g_{ab} to a fixed metric $f_{xx}=f_{yy}=1$

Holographic linear resistivity

(PRB 89, 2451161, 2014).



Richard Davison



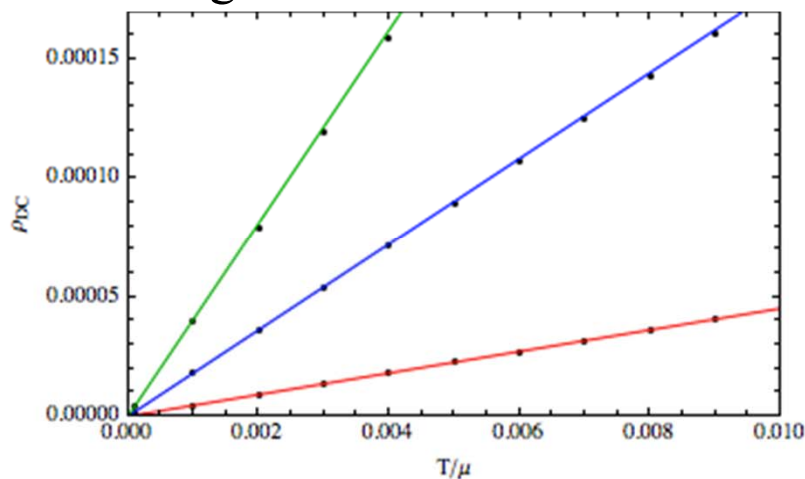
Steve Gubser

“Champion” strange metal: Einstein-Maxwell-dilaton (consistent truncation), **local quantum critical, marginal Fermi-liquid (3+1D)**, susceptible to **holo. superconductivity**, healthy thermodynamics: unique ground state, **Sommerfeld thermal entropy**.



David Vegh

Breaking of Galilean invariance (finite conductivities) due to **quenched disorder**: “massive gravity” = “**Higgsing**” **space-like diffeomorphisms in the bulk !?**



$$S = \frac{1}{2\kappa_4^2} \int d^4x \sqrt{-g} \left[\mathcal{R} - \frac{1}{4} e^\phi F_{\mu\nu} F^{\mu\nu} - \frac{3}{2} \partial_\mu \phi \partial^\mu \phi + \frac{6}{L^2} \cosh \phi - \frac{1}{2} m^2 \left(\text{Tr}(\mathcal{K})^2 - \text{Tr}(\mathcal{K}^2) \right) \right]$$

Explicit holographic construction explaining linear resistivity!

The secret of the linear resistivity (PRB 89, 2451161, 2014).



Davison



Planckian dissipation = very rapid local equilibration: a hydrodynamical fluid is established before it realizes that momentum is non conserved due to the lattice potential (not true in Fermi-gas: Umklapp time is of order collision time).

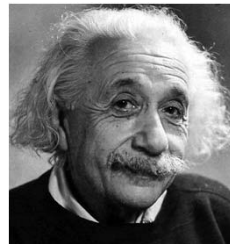
Hartnoll



Stokes

Resistivity in hydrodynamics

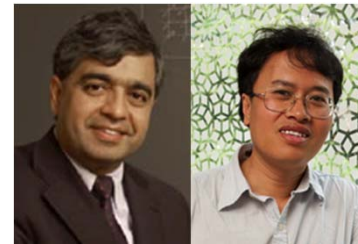
$$\rho(T) \propto \frac{1}{\tau_{rel}} = \frac{D}{l^2}$$



Einstein

Einstein relation:

$$D = \frac{\eta}{m_e n_e}$$



Sachdev Son

Planckian viscosity

$$\eta = A \frac{\hbar}{k_B} S$$

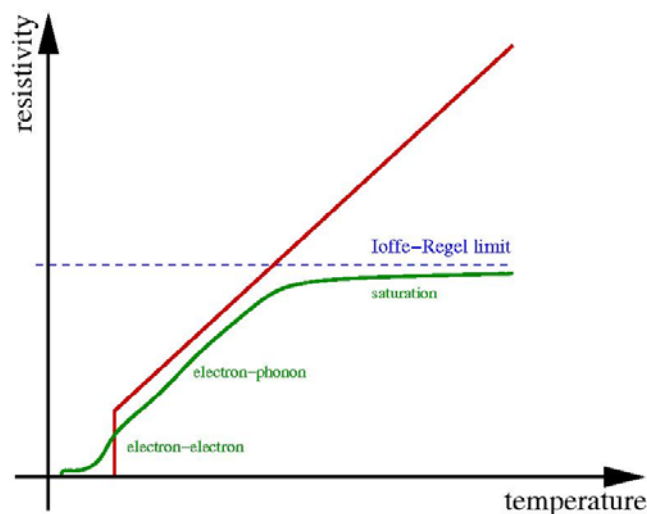
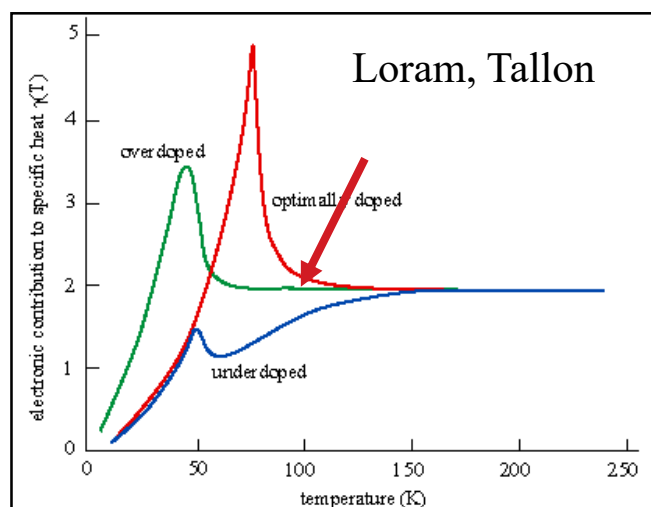
$$\rho(T) = \frac{1}{\omega_p^2 \tau_{rel}} = A \frac{\hbar}{\omega_p^2 l^2 m_e} \frac{S}{k_B}$$

Entropy versus transport: optimal doping

Optimally doped

$$C = \gamma T \Rightarrow S = T/\mu$$

$$\rho \propto \frac{1}{\tau_{rel}} \propto S \propto T$$



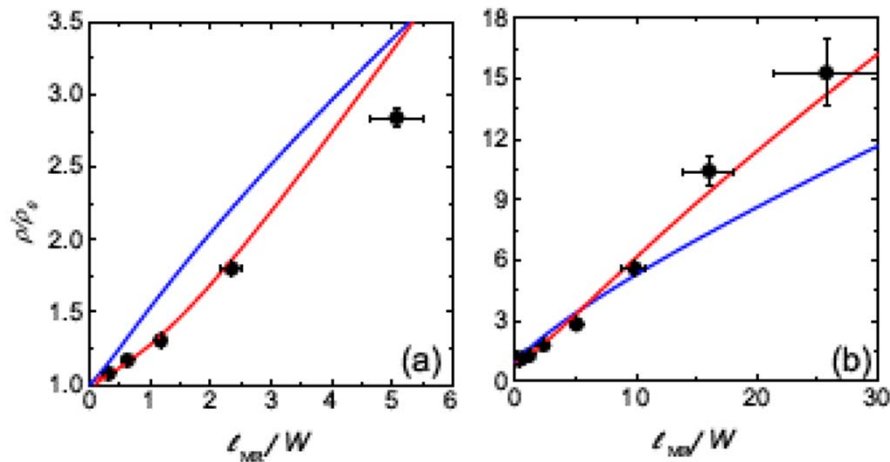
Plugging in numbers: “mean-free path” $l \approx 10^{-9} m$

Quite dirty but no residual resistivity since the fluid becomes perfect at $T = 0$!

Evidence for hydrodynamic electron flow in PdCoO₂

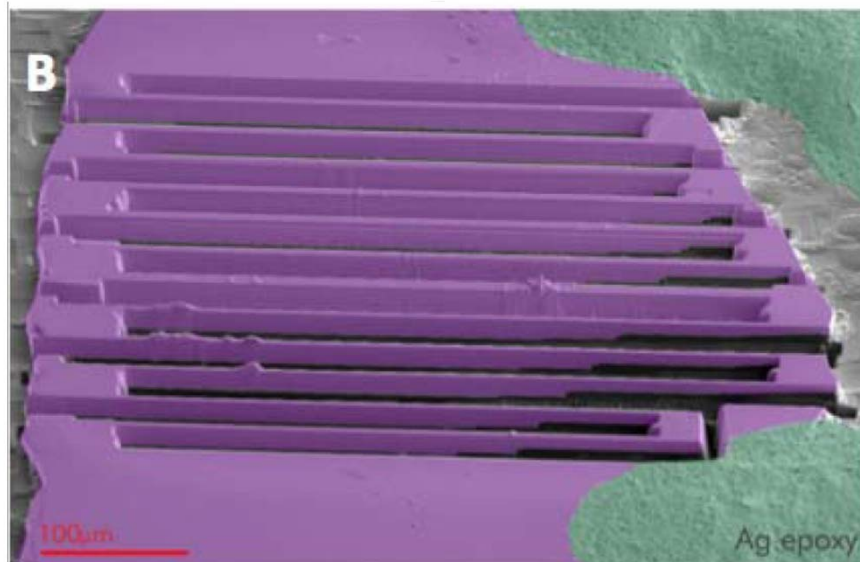
P. J. W. Moll, P. Kushwaha, N. Nandi, B. Schmidt, and A. P. Mackenzie

arXiv:1509.05691



Hydrodynamic effect on transport.

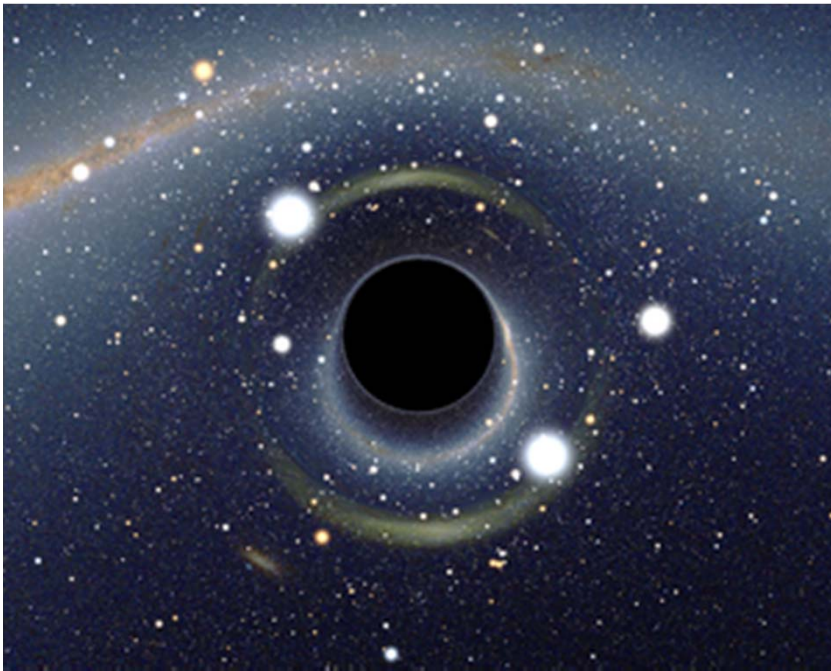
The measured resistivity of PdCoO₂ channels normalised to that of the widest channel (ρ_0), plotted against the inverse channel width $1/W$ multiplied by the bulk momentum-relaxing mean free path.



The tip of the empirical iceberg ...

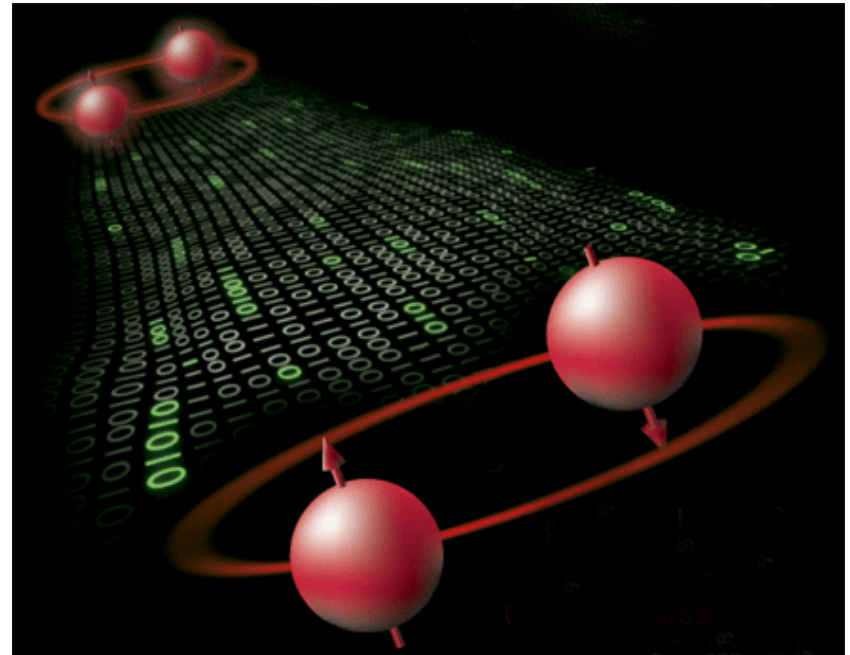
- **“Magneto-hydrodynamical” effects in microwave and optical response (vd Marel).**
- **Direct evidence for a quantum critical phase in cuprates: doping dependent transport, the Hall angle, transversal WF (Hussey, ...)**
- **Strange metal behaviour showing up in Photoemission and scanning tunneling spectroscopy (Dessau, ...)**
- **(Dis)proving holographic superconductivity: measuring the conformal metal in the pair susceptibility**
- **“Charge renormalization”: “fractional” Ahronov-Bohm effect (Phillips et al).**

Black holes as “quantum matter computers” !?



?

=



Bell pairs and the “spooky action at a distance”.

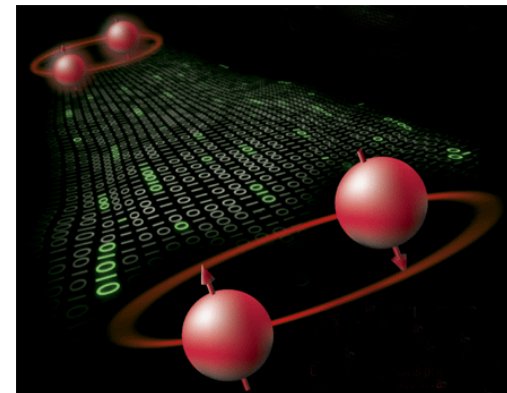
Classical computers live in tensor product space:

$$| \textit{product} \rangle = |0\rangle_A \otimes |1\rangle_B \textit{ or } |1\rangle_A \otimes |0\rangle_B$$



Quantum computers exploit entangled states capable of “spooky action at a distance” (EPR paradox).

$$| \textit{Bell} \rangle = \frac{1}{\sqrt{2}} \left(|0\rangle_A \otimes |1\rangle_B - |1\rangle_A \otimes |0\rangle_B \right)$$



“The classical condensates: from crystals to Fermi-liquids.”

States of matter that we understand are “short range entangled product”!

$$|\Psi_0\{\Omega_i\}\rangle = \prod_i \hat{X}_i^+(\Omega_i) |vac\rangle$$

- Crystals: put atoms in real space wave packets $X_i^+(R_i^0) \propto e^{(R_i^0 - r)^2 / \sigma^2} \psi^+(r)$

- Magnets: put spins in generalized coherent state

$$X_i^+(\vec{n}_i) \propto e^{i\varphi_i/2} \cos(\theta_i/2) c_{i\uparrow}^+ + e^{-i\varphi_i/2} \sin(\theta_i/2) c_{i\downarrow}^+$$

- Superconductors/superfluids: put bosons/Cooper pairs in coherent superposition

$$X_{k/i}^+ \propto u_k + v_k c_{k\uparrow}^+ c_{-k\downarrow}^+, \quad u_i + v_i e^{i\varphi_i} b_i^+$$

- Fermi gas/liquid special, but only “Fermi-Dirac entanglement”

$$|\Psi_{FL}\rangle = \prod_k^{k_F} c_k^+ |vac\rangle$$

Quantum matter.

“Macroscopic stuff that can quantum compute all by itself”

$$|\Psi\rangle = \sum_{\text{configs}} A_{\text{configs}} |\text{configs}\rangle$$

- **Topological incompressible systems, no low energy excitations but the whole carries quantum information: fractional quantum Hall, top.**

Superconductors/insulators (Majorana's, theta vacuum, ..)

- **Compressible systems: strongly interacting bosonic quantum critical states have dense long range entanglements (Planckian dissipation)**

- **Compressible systems: are the strange metals of this kind??**

Strongly interacting fermions at finite density: the fermion signs as entanglement resource!

Turning on the backflow

Kruger, JZ, PRB 78,
035104 (2008)

Nodal surface has to
become fractal !!!



Try backflow wave functions

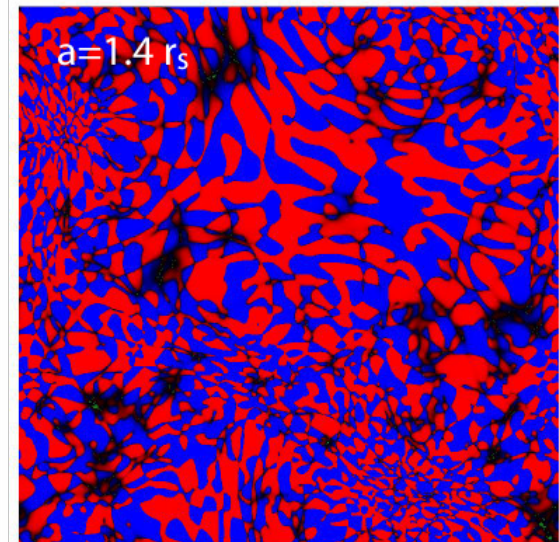
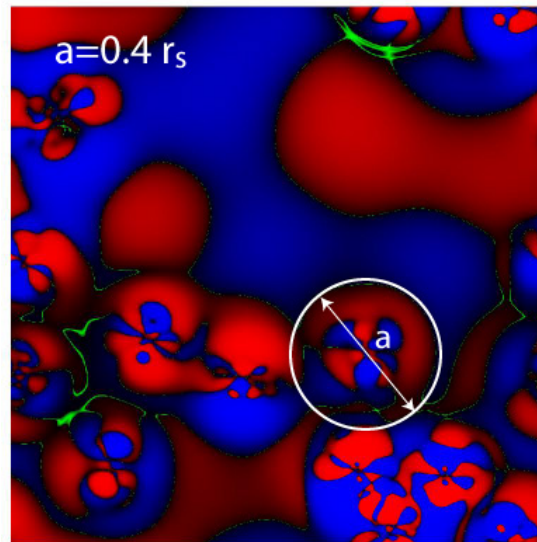
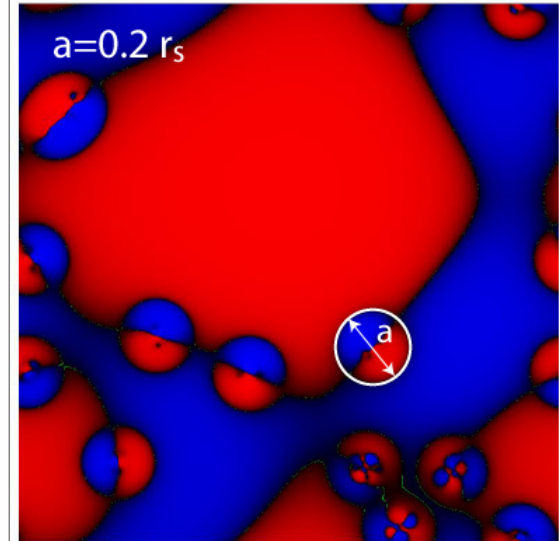
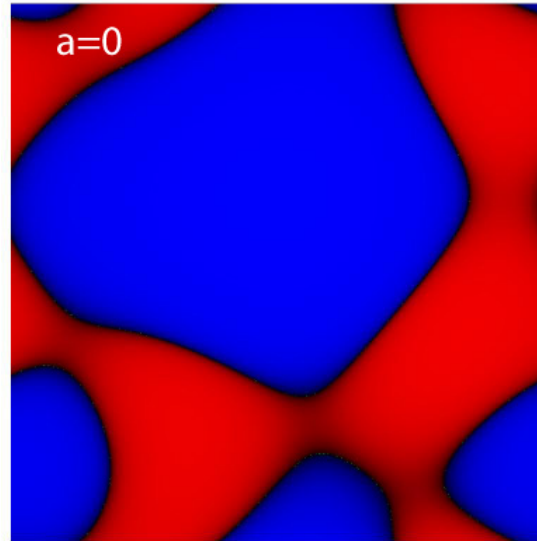
$$\psi_{bf}(\mathbf{R}) \sim \text{Det} (e^{i\mathbf{k}_i \tilde{\mathbf{r}}_j})_{ij}$$

$$\tilde{\mathbf{r}}_j = \mathbf{r}_j + \sum_{l(\neq j)} \eta(r_{jl})(\mathbf{r}_j - \mathbf{r}_l)$$

$$\eta(r) = \frac{a^3}{r^3 + r_0^3}$$

Collective (hydrodynamic)
regime:

$$a \gg r_s$$

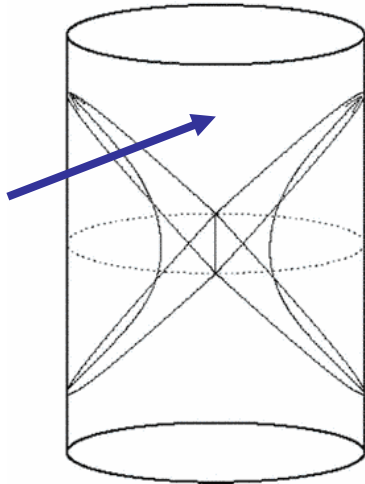


Its from quantum bits ?



Van Raamsdonk

Classical
space time
in bulk ...



Entangled state of
pair of H^d CFTs

$$|\Psi\rangle = \sum_i e^{-\pi R_H E_i} |E_i\rangle_{H^d} \otimes |E_i\rangle_{H^d}$$

Encoded by
quantum info
(entanglement
spectrum) in
boundary

“Rindler” description of pure global AdS

also

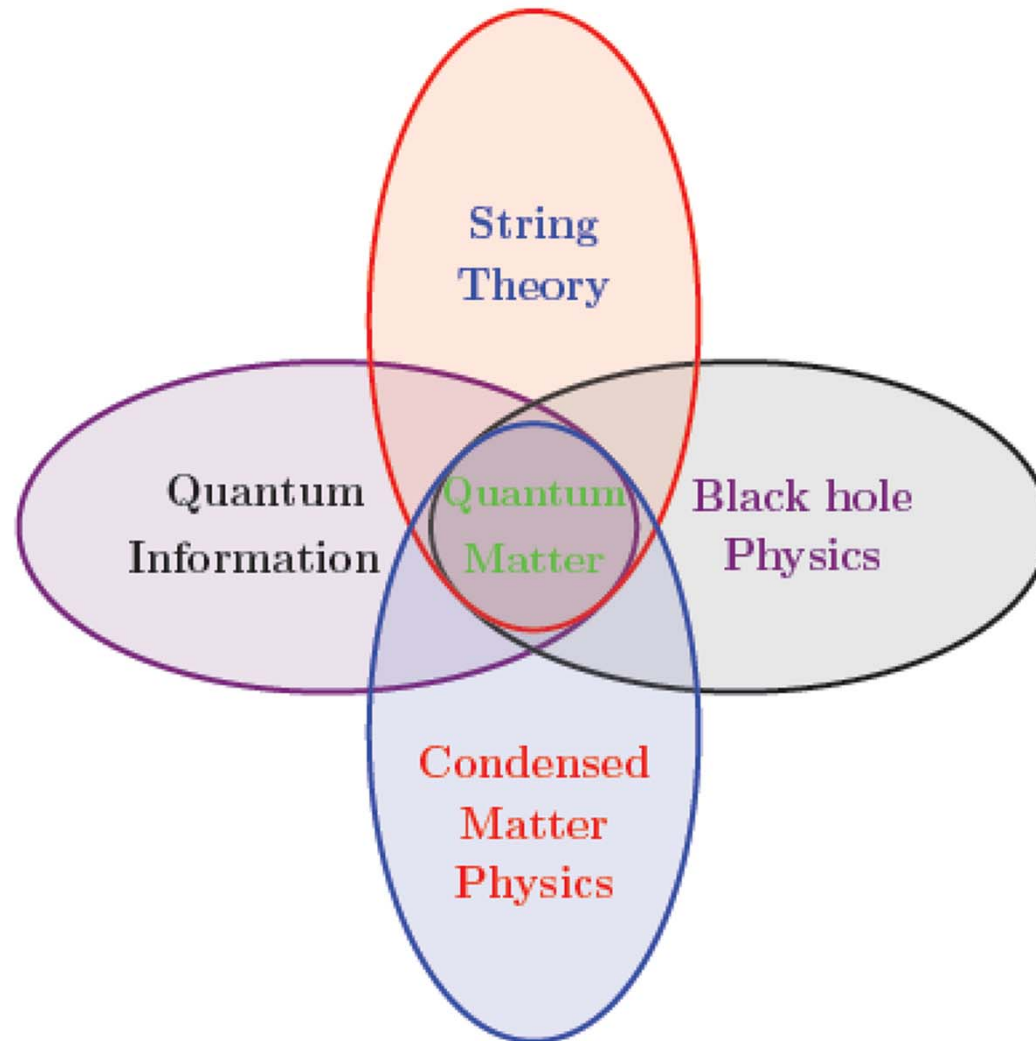


Hubeny Myers Swingle

The tip of the iceberg.

- **Electron systems in solids: the glass ceiling of quantum field theory = *sign problem*.**
- **Holographic duality = the generating functional of the mathematical theories describing **strong emergence in quantum systems**.**
- The extended **conventional textbook** of condensed matter physics is rewritten in terms of the **equations of general relativity**.
- Predicts “**beauty behind the fermion brick wall**”: strange metals as self organized **quantum critical phases** with novel scaling properties, presently chased in the cond. mat. laboratories
- These are extremely **densely entangled states of quantum matter**.
- **What does this all mean for the greater *quantum gravity* agenda?**

Koenraad's cloverleaf



Entanglement entropy



Bipartite von Neumann entropy: measures entanglement = quantum information of Bell pairs.

Trace the full density matrix over B:

$$\rho_A = \text{Tr}_B \rho$$

Compute the entropy associated with the reduced density matrix:

$$S_{VN,A} = \text{Tr}[\rho_A \ln \rho_A]$$

Universal measure of two bit entanglement:

$$\begin{aligned} |Bell\rangle &= \frac{1}{\sqrt{2}} (|0\rangle_A \otimes |1\rangle_B - |1\rangle_A \otimes |0\rangle_B) & S_{VN,A} &= \sqrt{2} \\ |Prod\rangle &= \frac{1}{2} (|0\rangle + |1\rangle)_A \otimes (|0\rangle + |1\rangle)_B & S_{VN,A} &= 0 \end{aligned}$$

“Space bipartite” Von Neumann entropy: entanglements and fields

Divide **space in two**:

$$\rho_A = \text{Tr}_B \rho, \quad S_A = -\text{Tr}[\rho_A \ln \rho_A]$$

1+1 D CFT's (Wilczek et al., Calabrese-Cardy):

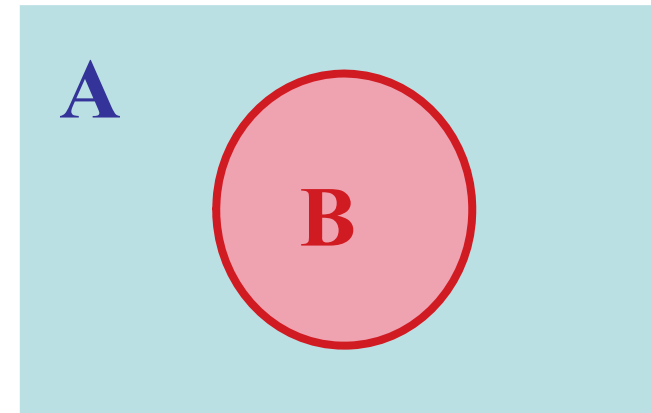
$$S_A = \frac{c}{3} \log L \quad \mathbf{c} = \text{central charge}$$

Topological insulators (Kitaev-Preskil, Levin-Wen):

$$S_A = \alpha \Sigma - \gamma + \dots, \quad \Sigma = L^{d-1}, \quad \gamma = \log D \quad \mathbf{D} = \text{total quantum dimension.}$$

Conformal fields in higher (e.g. odd) d (Myers, Klebanov, ...):

$$S_A = \alpha \Sigma + \dots + (-1)^{(d-1)/2} S_d \quad \mathbf{S}_d \text{ is universal}$$



Bipartite entanglement entropy and quantum field theory.

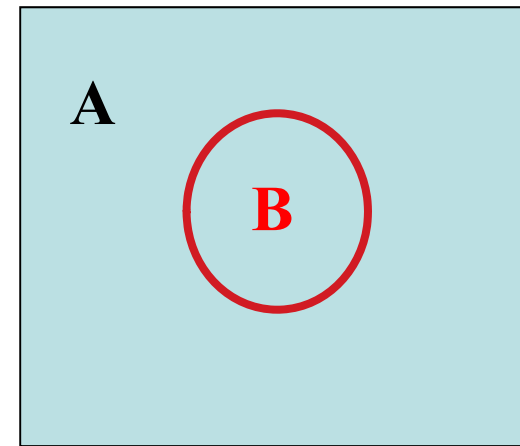


Wilczek

$$\rho_A = \text{Tr}_B[\rho]$$

$$S_{vN} = -\text{Tr}[\rho_A \ln \rho_A]$$

Measure of entanglement of degrees of freedom in spatial volume B with those in A.



Generic energy eigenstates: S_{vN} scales with volume L^d of B.

Ground states of bosonic systems: S_{vN} scales with the area L^{d-1} of B.

Fermi gas: longer ranged “signful” entangled $S_{vN} \sim L^{d-1} \ln(L^{d-1})$

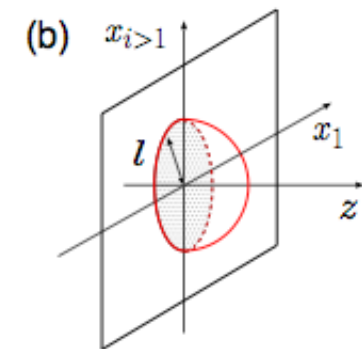
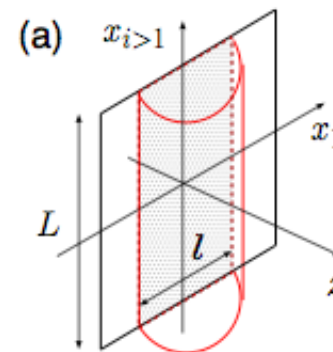
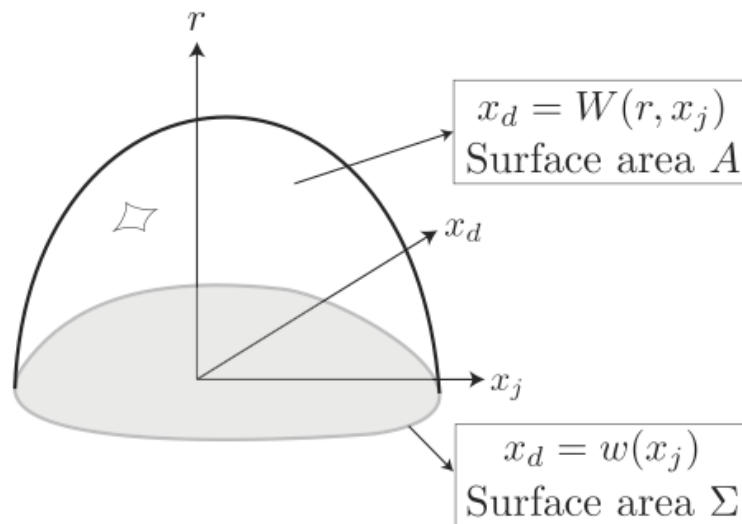
Entanglement entropy versus AdS/CFT.



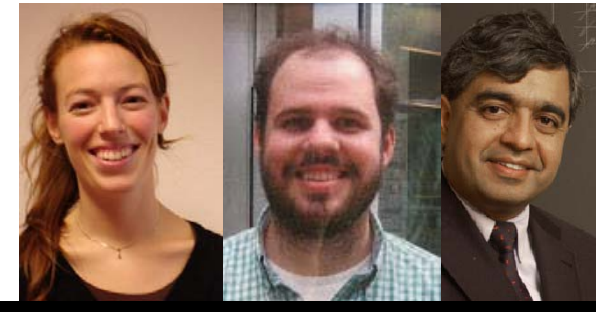
Takayanagi Ryu

$$\rho_A = \text{Tr}_B[\rho] \quad S_{VN} = -\text{Tr}[\rho_A \ln \rho_A]$$

The spatial bipartite entanglement entropy in the boundary is dual to the area of the minimal surface in the bulk, bounded by the cut in the space of the boundary



Holographic strange metal entanglement entropy.



Huijse Swingle Sachdev

Einstein-Maxwell-Dilaton bulk => “hyperscaling violating geometry” (Kiritsis et al.): **PRD** 85, 035121 (2012)

$$ds^2 = \frac{1}{r^2} \left(-\frac{dt^2}{r^{2d(z-1)/(d-\theta)}} + r^{2\theta/(d-\theta)} dr^2 + dx_i^2 \right)$$

Boundary: interpolating between “normal” and RN strange metals.

$$x_i \rightarrow \zeta x_i, \quad t \rightarrow \zeta^z t, \quad ds \rightarrow \zeta^{\theta/d} ds \quad S \propto T^{(d-\theta)/z}$$

Entanglement entropy:

$$S_{VN} \propto L^{d-1}, \quad \theta < d-1$$

$$S_{VN} \propto L^{d-1} \ln L^{d-1}, \quad \theta = d-1$$

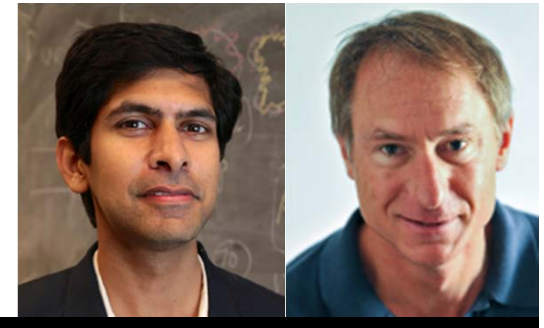
$$S_{VN} \propto L^\theta, \quad d-1 < \theta < d$$

Bosonic fields

Fermi liquid-like

But this is longer ranged !

Fermion signs and dense entanglement ...



Grover Fisher
arXiv:1412.3534

S_{vN} area law: ground states of “sign-free” systems (bosons, tensor product states ..)



Energy eigenstates:

$$|\Psi_i\rangle = \sum_{conf} (-1)^{i,conf} |A_{conf}^i| |conf\rangle$$

$(-1)^{i,conf} =$ - Antisymmetrization \Rightarrow area log area S_{vN} (Fermi gas)
- Random \Rightarrow Volume S_{vN} (typical excited states)

The *quantum critical metallic phases* of holography are characterized by dense “sign driven” entanglement as characterized by the hyperscaling violation exponent!

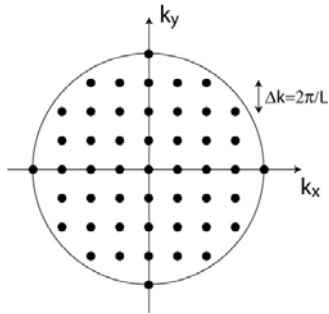
The nodal hypersurface

Antisymmetry of the wave function

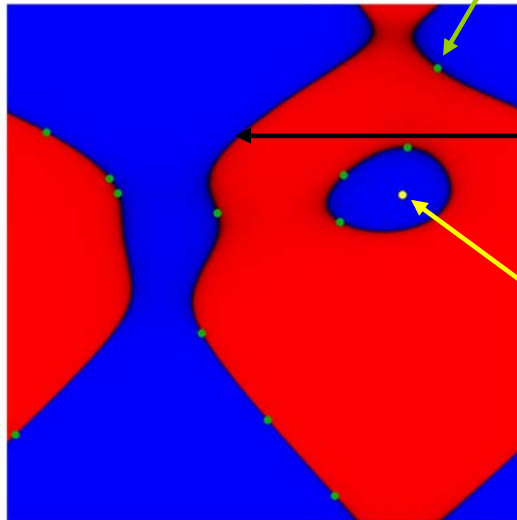
$$\Psi(\mathbf{r}_1, \dots, \mathbf{r}_i, \dots, \mathbf{r}_j, \dots, \mathbf{r}_N) = -\Psi(\mathbf{r}_1, \dots, \mathbf{r}_j, \dots, \mathbf{r}_i, \dots, \mathbf{r}_N)$$

Free Fermions

$$\Psi_0(\mathbf{R}) \sim \text{Det} (e^{i\mathbf{k}_i \mathbf{r}_j})_{ij}$$



d=2



Pauli hypersurface

$$P = \bigcup_{i \neq j} P_{ij}$$

$$P_{ij} = \{\mathbf{R} \in \mathbb{R}^{Nd} | \mathbf{r}_i = \mathbf{r}_j\}$$

$$\dim P = Nd - d$$

Nodal hypersurface

$$\Omega = \{\mathbf{R} \in \mathbb{R}^{Nd} | \Psi(\mathbf{R}) = 0\}$$

$$\dim \Omega = Nd - 1$$

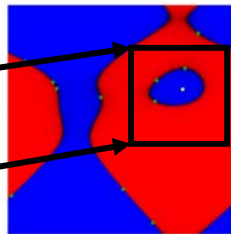
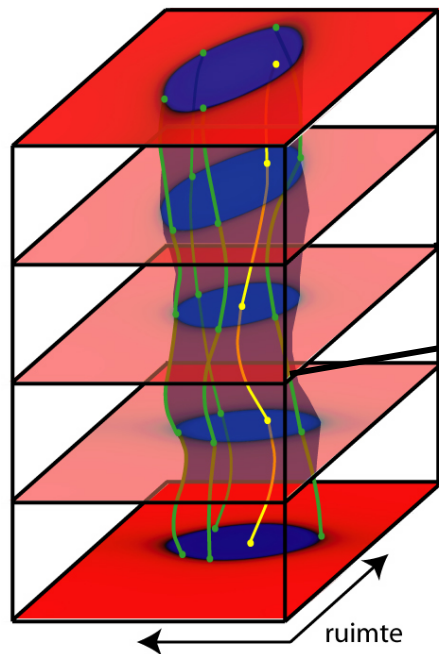
Test particle

Constrained path integrals

Formally we can solve the sign problem!!

$$\rho_F(\mathbf{R}, \mathbf{R}; \beta) = \frac{1}{N!} \sum_{\mathcal{P}, \text{even}} \int_{\gamma: \mathbf{R} \rightarrow \mathcal{P}\mathbf{R}}^{\gamma \in \Gamma(\mathbf{R}, \mathcal{P}\mathbf{R})} \mathcal{D}\mathbf{R}(\tau) \exp \left\{ -\frac{1}{\hbar} \int_0^{\hbar/T} d\tau \left(\frac{m}{2} \dot{\mathbf{R}}^2(\tau) + V(\mathbf{R}(\tau)) \right) \right\}$$

$$\Gamma(\mathbf{R}, \mathbf{R}') = \{ \gamma : \mathbf{R} \rightarrow \mathbf{R}' \mid \rho_F(\mathbf{R}, \mathbf{R}(\tau); \tau) \neq 0 \}$$



Ceperley, *J. Stat. Phys.* (1991)

Self-consistency problem:
Path restrictions depend on ρ_F !

Reading the worldline picture

Fermi-energy: confinement energy imposed by local geometry

$$l^2(\tau) = \langle (\mathbf{r}_i(\tau) - \mathbf{r}_i(0))^2 \rangle = 2dD\tau = 2d\frac{\hbar}{2m}\tau$$

$$l^2(\tau_c) \simeq r_s^2 \rightarrow \tau_c \simeq \frac{1}{2d} \frac{2m}{\hbar} n^{-2/d}$$

$$\hbar\omega_c = \frac{\hbar}{\tau_c} \simeq d \frac{\hbar^2}{2m} n^{2/d} \simeq E_F$$

Fermi surface encoded globally: $\rho_F = \text{Det}(e^{ik_i r_j}) = 0$

Change in coordinate of one particle changes the nodes everywhere

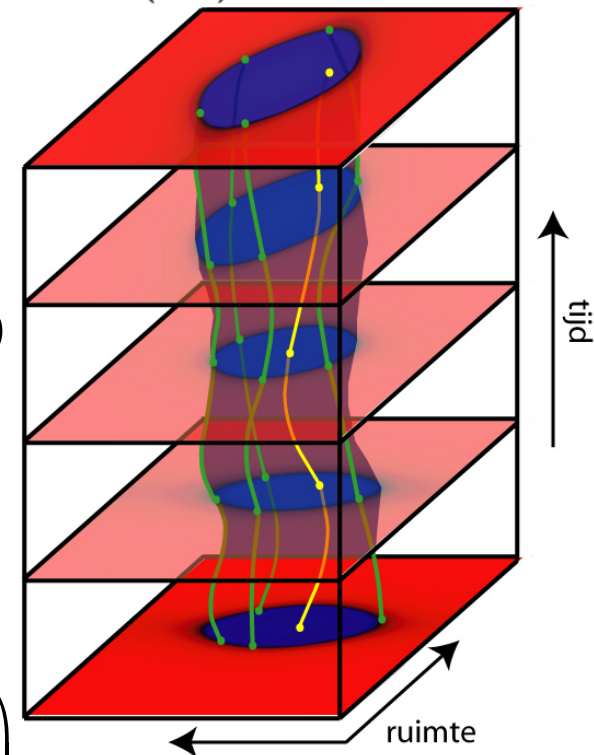
Finite T: $\rho_F = (4\pi\lambda\beta)^{-dN/2} \text{Det} \left[\exp\left(-\frac{(r_i - r_{j0})^2}{4\lambda\tau}\right) \right]$
 $\lambda = \hbar^2 / (2M)$

Non-locality length:

$$\lambda_{nl} = v_F \tau_{inel} = v_F \left(\frac{E_F}{k_B T} \right) \left(\frac{\hbar}{k_B T} \right)$$

Average node to node spacing

$$\sim r_s = \left(\frac{V}{N} \right)^{1/d} = n^{-1/d}$$



Key to fermionic quantum criticality



Kruger

JZ

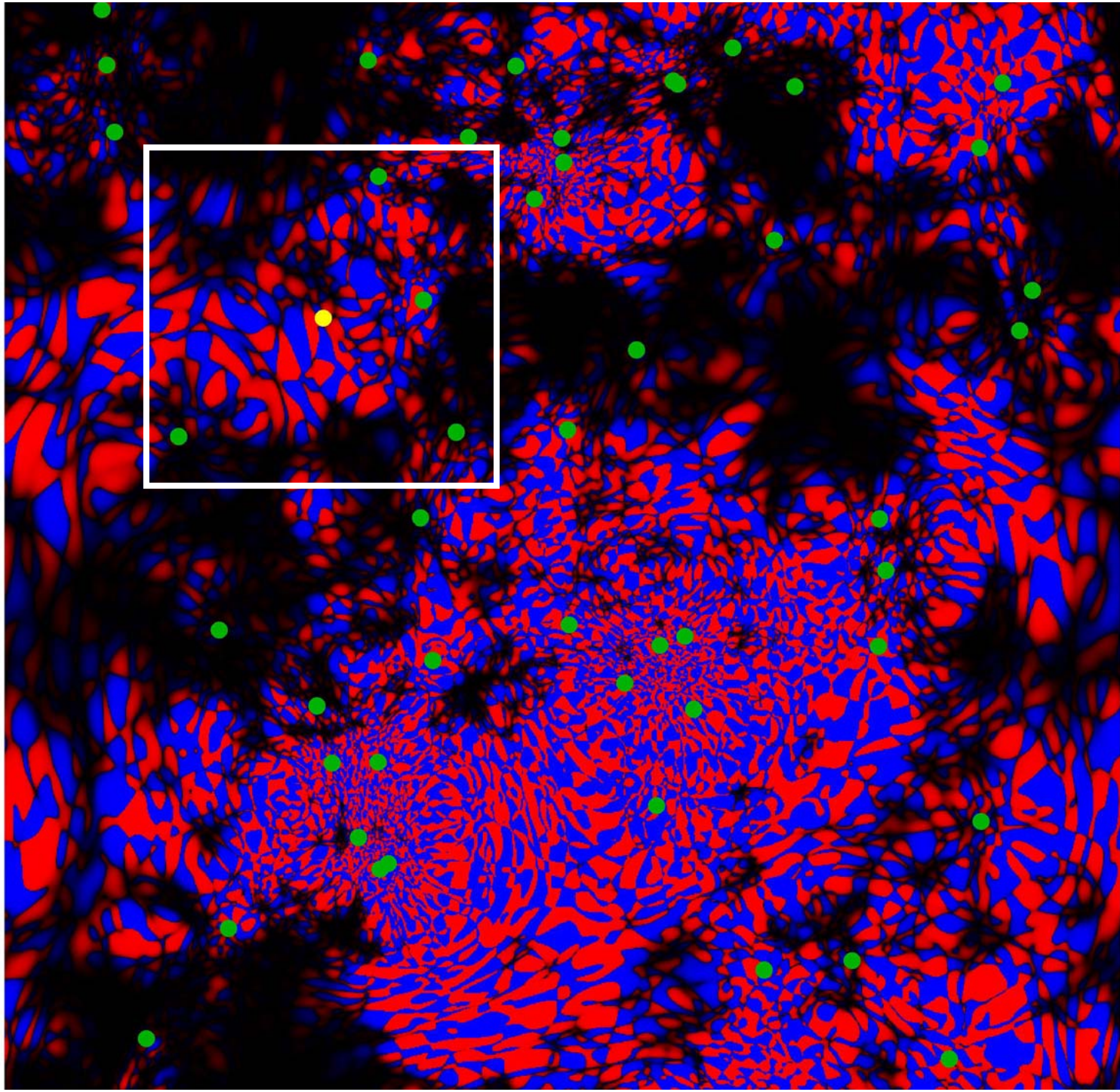
Phys. Rev. B **78**, 035104 (2008)

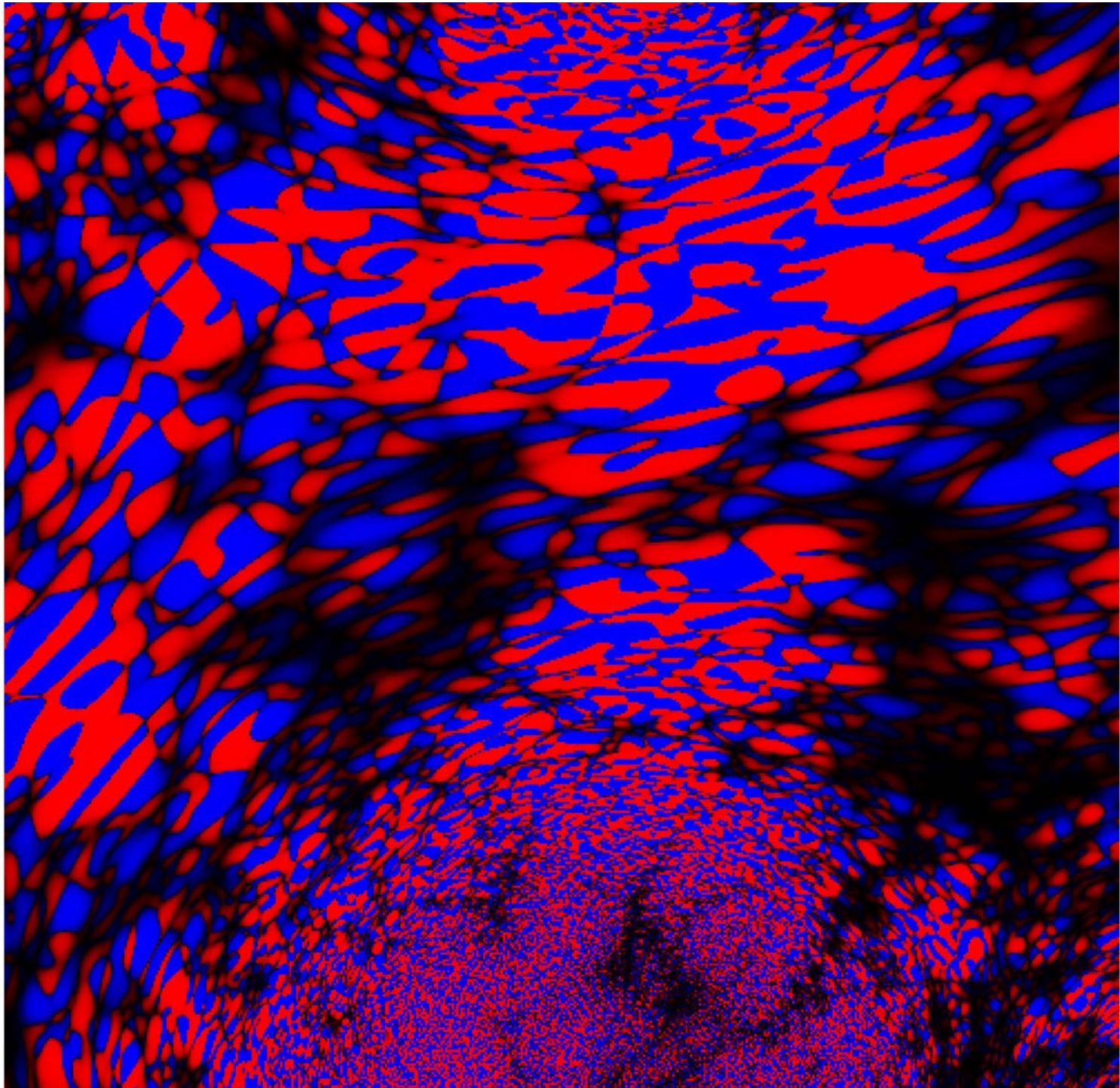
At the QCP scale invariance, no E_F



Nodal surface has to become fractal !!!







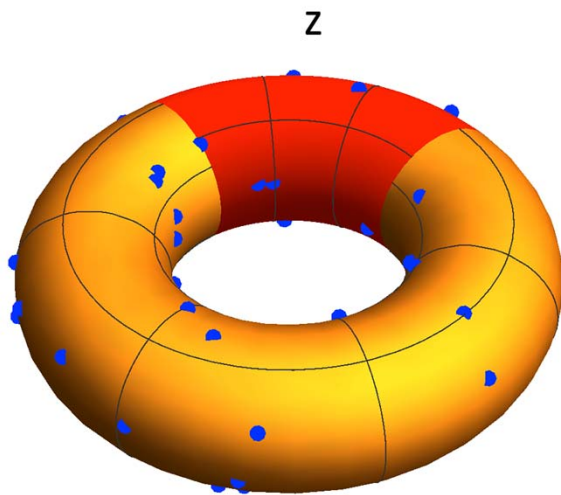
Fractal nodes and entanglement entropy.



Kaplis Kruger

Second Renyi entropy: leading contribution scales like vN entropy.

$$S^q(z) = \frac{\ln(\text{Tr} \rho_A^q)}{1 - q}, \quad q = 2$$



$$\rho(\mathbf{R}, \mathbf{R}') = \psi^*(\mathbf{R})\psi(\mathbf{R}')$$

$$\psi_{bf}(\mathbf{R}) \sim \text{Det} (e^{i\mathbf{k}_i \tilde{\mathbf{r}}_j})_{ij}$$

$$\tilde{\mathbf{r}}_j = \mathbf{r}_j + \sum_{l(\neq j)} \eta(r_{jl})(\mathbf{r}_j - \mathbf{r}_l)$$

Backflow range exponent “eta” (=3 for hydro backflow): $\eta(r) = \frac{a^\eta}{r^\eta + r_0^\eta}$

Fractal nodes and entanglement entropy.

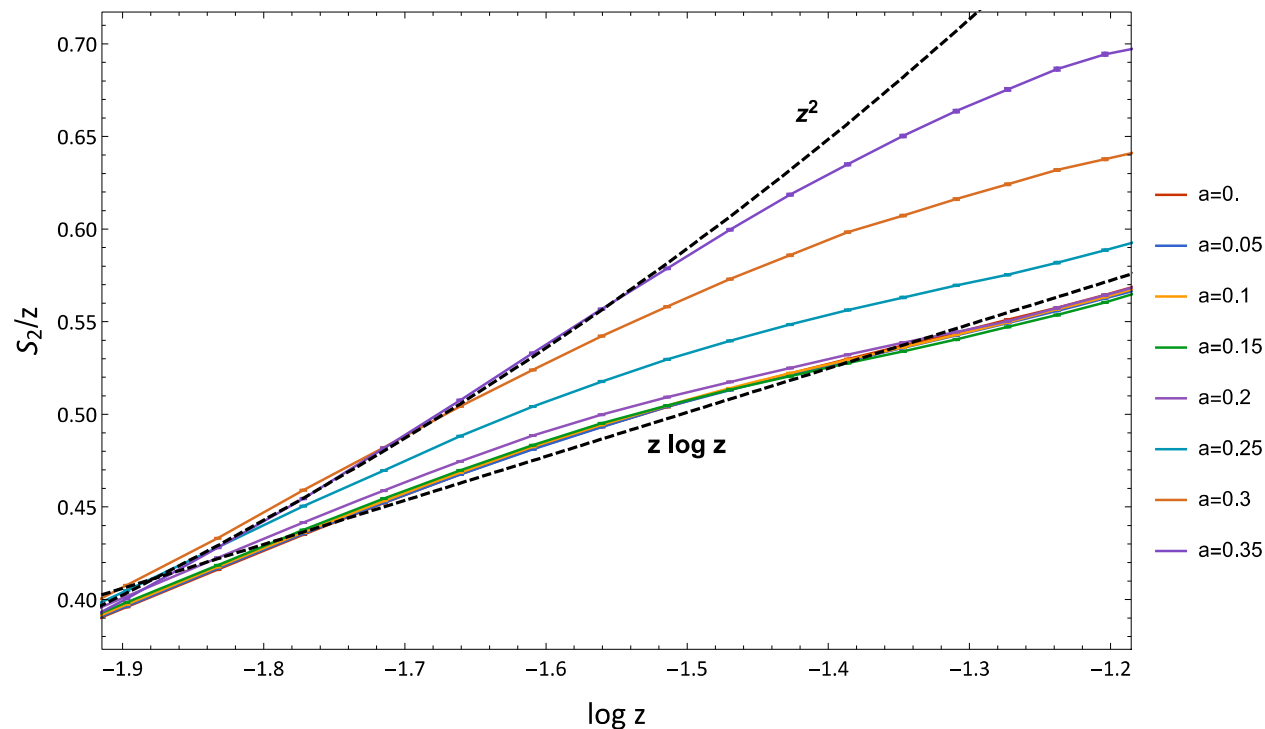


Kaplis

Kruger

Second Renyi entropy: $S^q(z) = \frac{\ln(\text{Tr} \rho_A^q)}{1 - q}$, $q = 2$

Hydrodynamical backflow, for increasing backflow length a ($a_c = 0.5$):



Empty.

Empty.

Empty.

Empty.
