

# Discovering electroweakinos in (mini)split supersymmetry.

(BASED ON BADZIAK, DELGADO, OLECHOWSKI, SP, SAKURAI, JHEP, 2015)

BEFORE „FORGETTING” SUPERSYMMETRY, LET’S  
EXPLORE MORE  
IT’S POSSIBLE PARAMETER RANGE:

LIGHT GAUGINOS, HEAVY SFERMIONS - (MINI) SPLIT?

$$\Omega h^2 \leq 0.12 \quad (\text{MULTI - COMPONENT DARK MATTER?})$$

## (MINI)SPLIT SUPERSYMMETRY

SCALAR MASSES AND  $B_\mu$  (CORRESPONDING TO DIM 2 OPERATORS)

ARE CHARACTERIZED BY  $\tilde{m}$  AND GAUGINO/HIGGSINO MASSES

AND A-TERMS (CORRESPONDING TO DIM 3 OPERATORS)

ARE ASSUMMED TO BE  $\mathcal{O}(1TeV)$

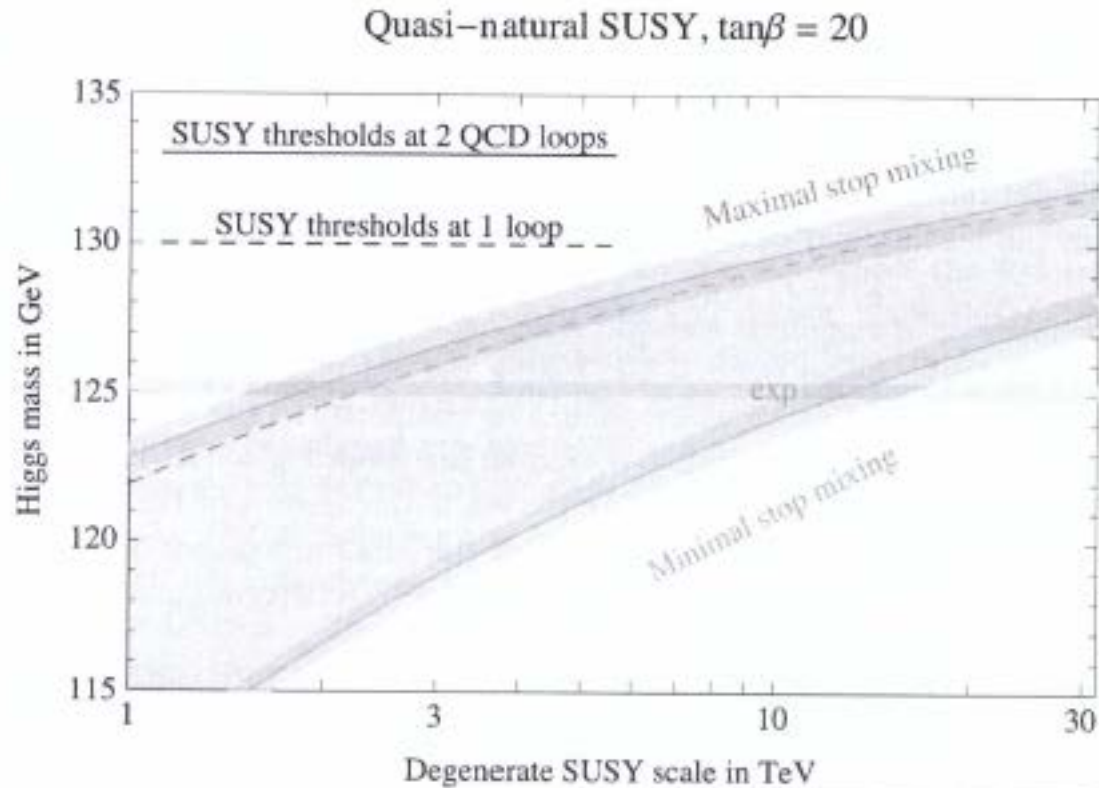
- A VARIETY OF GAUGE MEDIATION MECHANISMS, WITH DIFFERENT PATTERNS OF ELECTROWEAKINO MASSES, ANOMALY MEDIATION, MIRAGE MEDIATION

THE SPLITTING IN SPLIT SUPERSYMMETRY IS CONSTRAINED  
BY THE HIGGS MASS

(  $\tan \beta$  DEPENDENT )

# Summing $\ln(M_{\text{susy}}/M_Z) = \text{MATCHING SM TO MSSM AT } M_{\text{susy}}$

Bagnaschi, Giudice, Slavich, Strumia (2014)



Superpartner mass mass dependence can be described (in LL approx) by a single effective parameter  $T_{\text{susy}}$

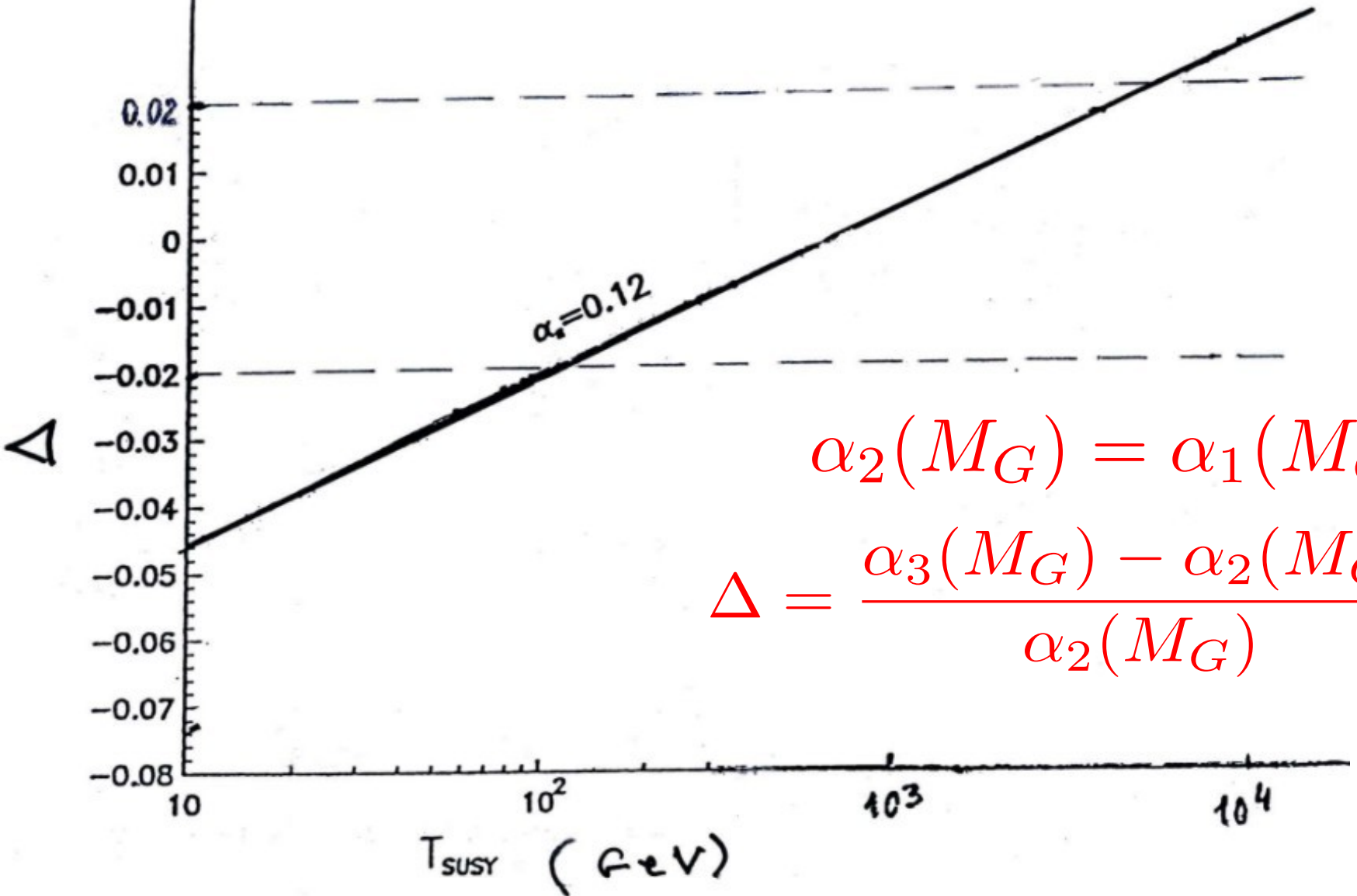
$$T_{\text{susy}} = |\mu| \left( \frac{m_{\tilde{w}}}{m_{\tilde{g}}} \right)^{3/2} \left( \frac{M_{\tilde{t}}}{M_{\tilde{q}}} \right)^{3/19} \left( \frac{M_{A^0}}{|\mu|} \right)^{3/19} \left( \frac{m_{\tilde{w}}}{|\mu|} \right)^{3/19}$$

$$\approx |\mu| \left[ \frac{m_{\tilde{w}}}{m_{\tilde{g}}} \right]^{3/2}$$

CARENA, SP, WAGNER '93

A HINT FOR EITHER HEAVY OR COMPRESSED SPECTRUM BECAUSE...

$$m_t = 180 \text{ GeV}, \quad \frac{1}{g\beta} = 10$$





For instance, for universal gaugino masses at the GUT scale

$$T_{\text{susy}} \sim |\mu| \left[ \frac{\alpha_2(M_2)}{\alpha_3(M_2)} \right]^{3/2} \sim \frac{1}{7} |\mu|$$

**A HINT FOR EITHER HEAVY O(1) TeV OR COMPRESSED SPECTRUM**

- ELECTROWEAK SECTOR MAY PLAY LEADING ROLE IN DISCOVERING SUPERSYMMETRY
- ACCEPTING THERMAL HISTORY OF UNIVERSE, IN MODELS WITH STABLE NEUTRALINO (R-PARITY), ITS MASS IS BOUNDED FROM ABOVE BY

$$\Omega h^2 \leq 0.12$$

- **DECOUPLED SFERMIONS** (actually, it's enough to have them 20-30% heavier than gauginos/higgsinos)

PARAMETERS

$$M_1, M_2, \mu, \tan \beta$$

APPLIES TO MSSM AND ALL MODELS WHERE THE ADMIXTURE OF ADDITIONAL STATES TO LSP AND NLSP IS SMALL

DIRECT DETECTION + LHC -> RELEVANT ELECTROWEAKINO PARAMETER SPACE BEGINS TO BE TESTED

- WHAT IS THE RANGE OF VALUES OF INTEREST FOR  $\Omega h^2$
- DIRECT DETECTION
- COLLIDER
- THEIR INTERPLAY

ANALYTICAL QUALITATIVE GUIDE AND  
NUMERICAL SCANS

It is convenient to organize the discussion according to the LSP composition

BINO-HIGGSINO MIXING  $(M_2 = m_{sf} = m_A = 7 \text{ TeV})$

HIGGSINO-WINO MIXING  $(M_1 = m_{sf} = m_A = 7 \text{ TeV})$

BINO-WINO MIXING WITH  $\mu = 1.1 \min(M_1, M_2)$

AND MORE GENERAL BINO-HIGGSINO-WINO MIXING

## PURE HIGGSINO

$$\Omega_h h^2 = 0.10 \left( \frac{\mu}{1 \text{ TeV}} \right)^2$$

$$\Omega h^2 < 0.12 \quad \text{for}$$

$$\mu < 1 \text{ TeV}$$

## PURE WINO

$$\Omega_w h^2 = 0.13 \left( \frac{M_2}{2.5 \text{ TeV}} \right)^2$$

$$\Omega h^2 < 0.12 \quad \text{for}$$

$$M_2 < 2.2(2.8) \text{ TeV}$$

LOWER BOUND ON CHARGINO MASS GIVES LOWER BOUND FOR  $\Omega$ ; ADDING BINO COMPONENT GIVES LARGER  $\Omega$  (WHEN SFERMIONS ARE HEAVY)

# SPIN INDEPENDENT SCATTERING CROSS SECTION

## EXCLUSION LIMITS (LUX) AND FUTURE PROSPECTS

REMEMBER: FOR NEUTRALINOS WITH  $\Omega_\chi$

AND CROSS SECTION  $\sigma_{SI}$

THE EXCLUSION LIMIT IS  $\frac{\sigma_{SI}^\chi}{\sigma_{SI}^{LUX}} \frac{\Omega_\chi}{\Omega_{DM}} > 1$

$$\sigma_{SI} = 8 \times 10^{-45} \text{cm}^2 \left( \frac{C_{h\chi\chi}}{0.1} \right)^2$$

where

$$L = \frac{C_{h\chi\chi}}{2} h(\chi\chi + \chi^\dagger\chi^\dagger)$$

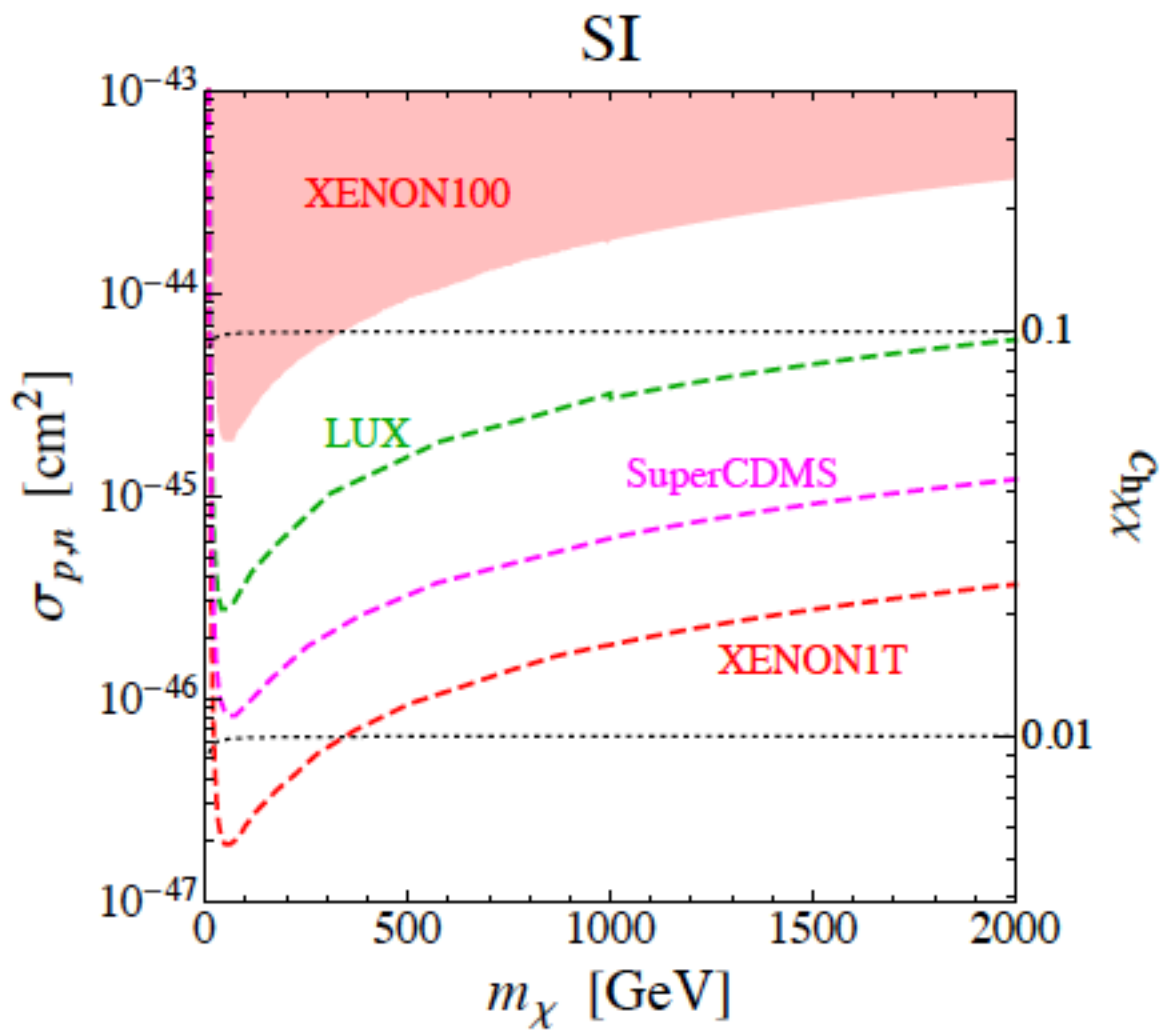
and (approximately)

$$C_{h\chi\chi} \sim \theta \quad (\text{MIXING})$$

Gaugino/higgsino  $\theta = \frac{(\sin \beta \pm \cos \beta)}{\sqrt{2}} \left( \frac{M_Z}{(\mu \mp M_i)} \right)$

Bino/wino  $\theta = \frac{(\sin 2\beta \sin 2\theta_W)}{2} \left( \frac{M_Z^2}{(M_2 - M_1)\mu} \right)$

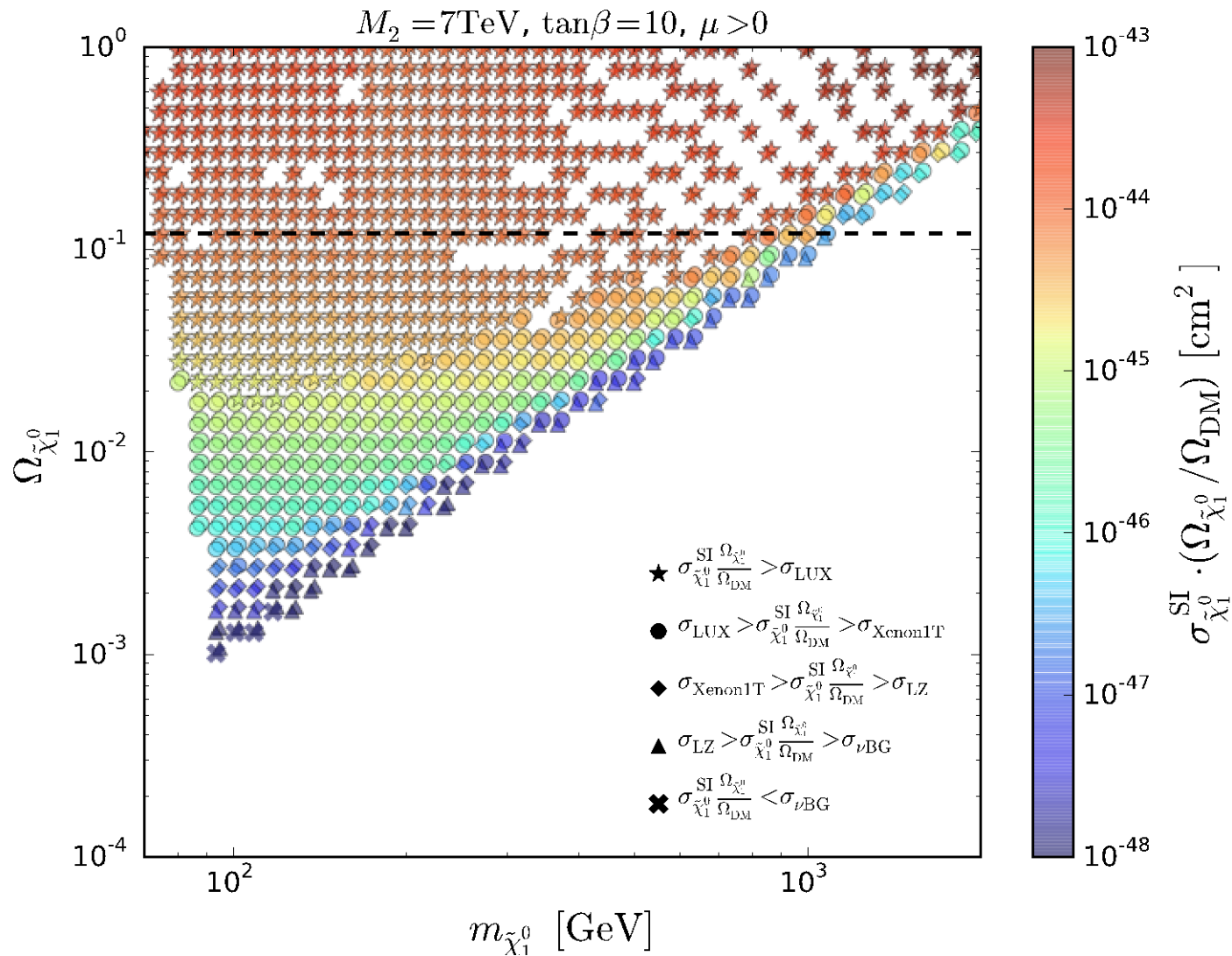


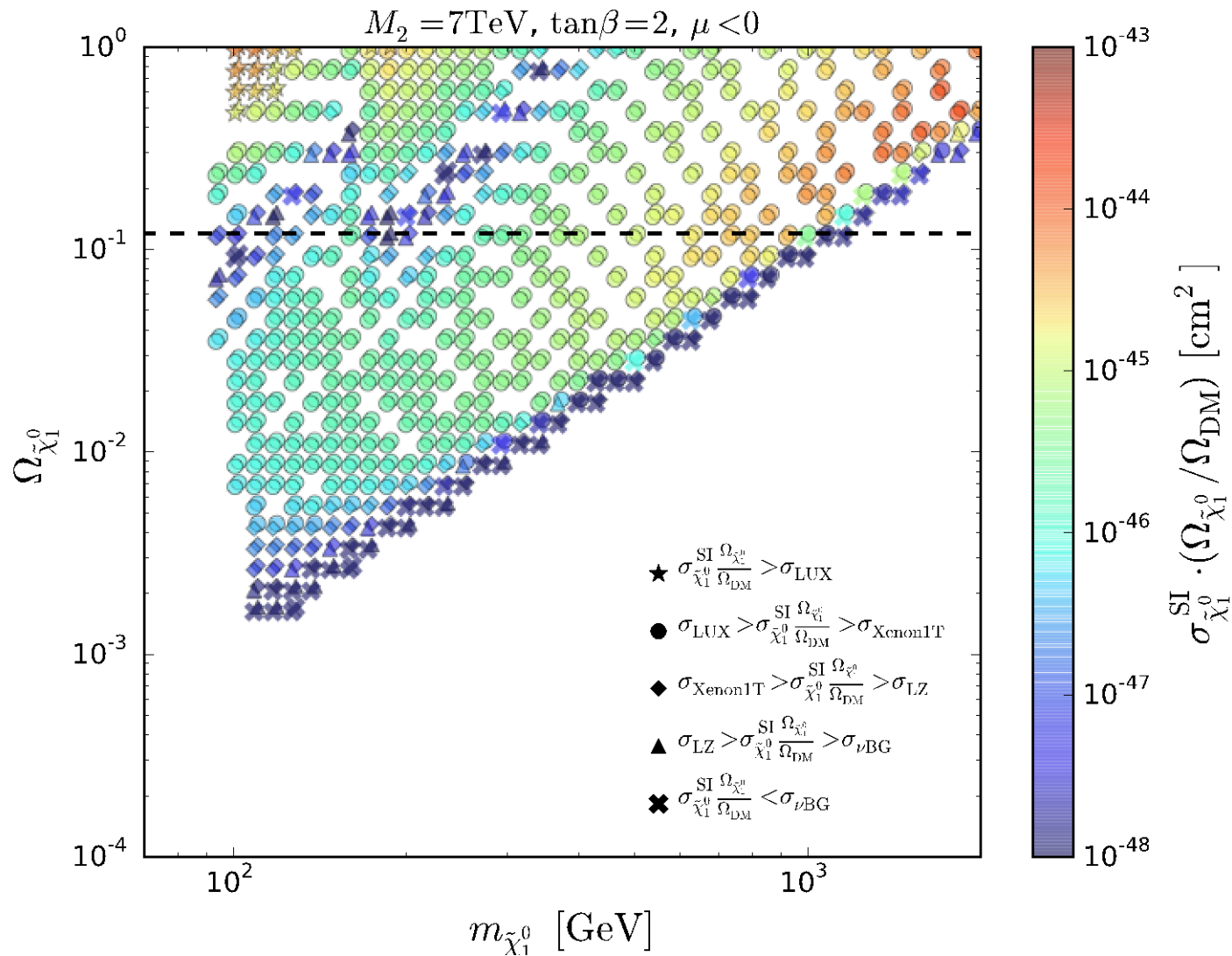


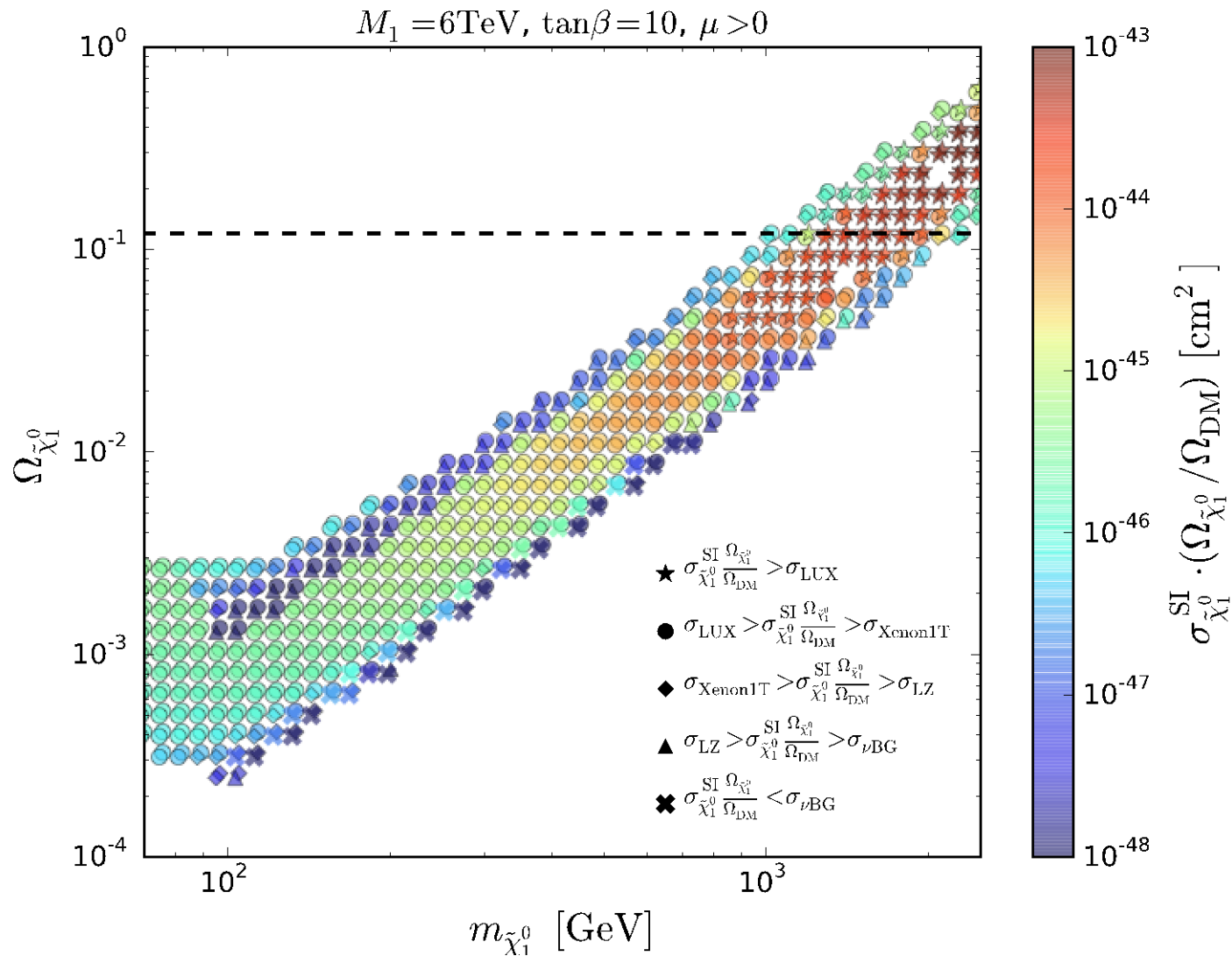
## FIRST IMPORTANT CONCLUSION:

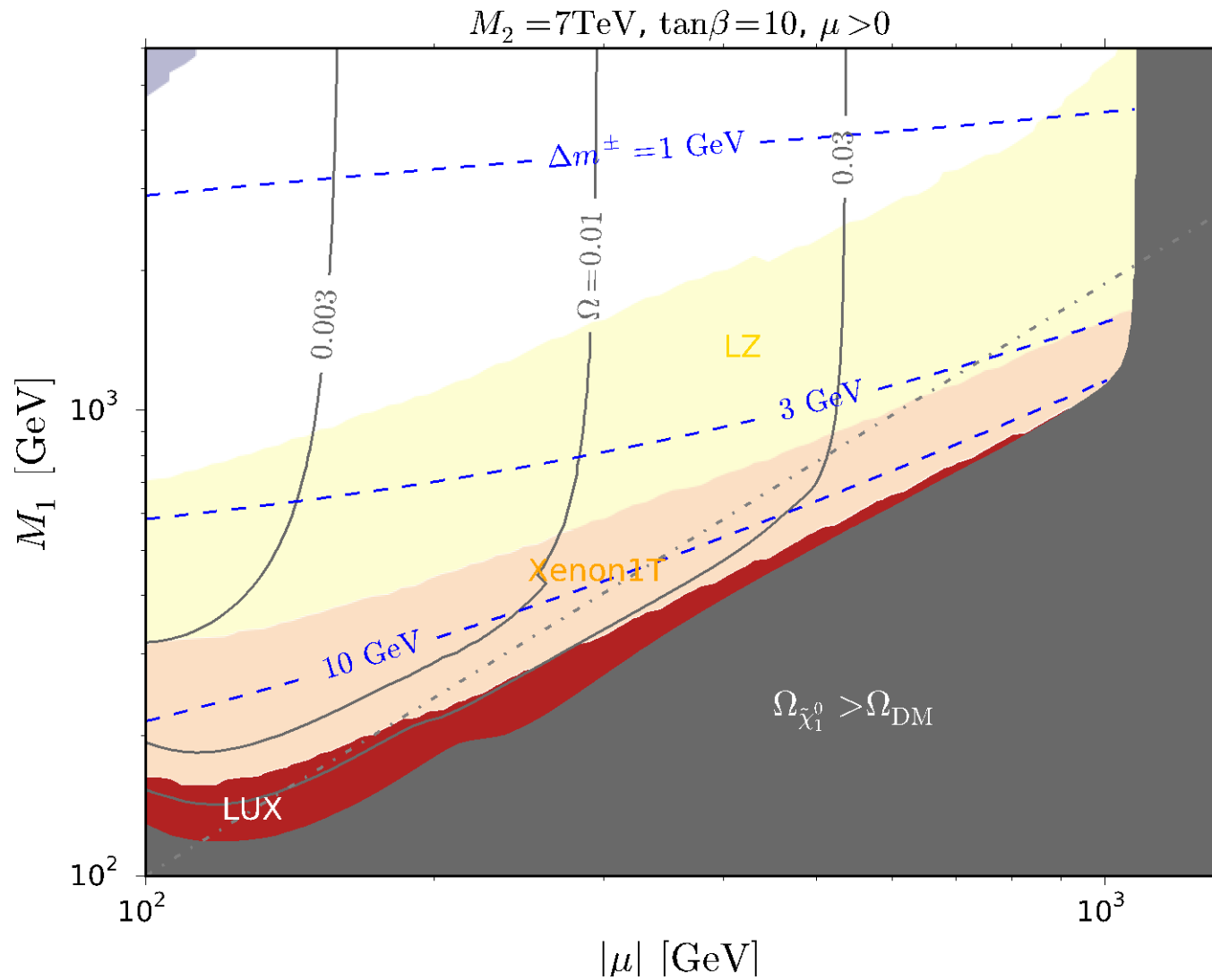
IN SPITE OF THE (SMALL) FLUX FACTOR,  
LUX AND FUTURE DD EXPERIMENTS ARE/WILL  
BE SENSITIVE TO NEUTRALINOS EVEN WITH  
VERY SMALL  $\Omega_\chi$

(A SIGNAL WOULD NOT NECESSARILY MEAN THE  
DISCOVERY OF THE MAIN DM COMPONENT)







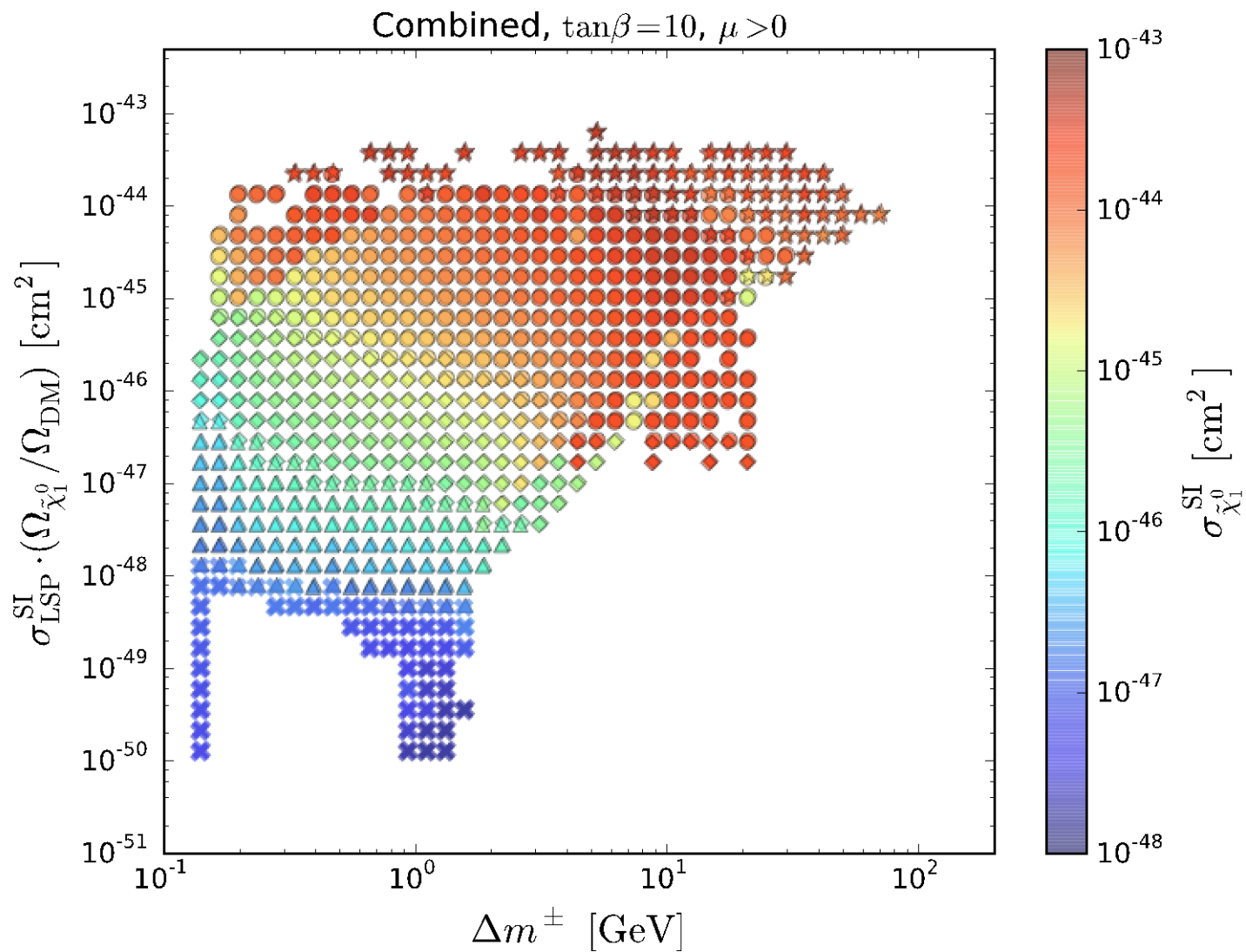


ANOTHER IMPORTANT GENERAL CONCLUSION,  
WITH STRONG IMPACT ON THE COLLIDER SEARCHES:

**DD LIMITS IMPLY SMALL NLSP-LSP MASS SPLITTINGS**

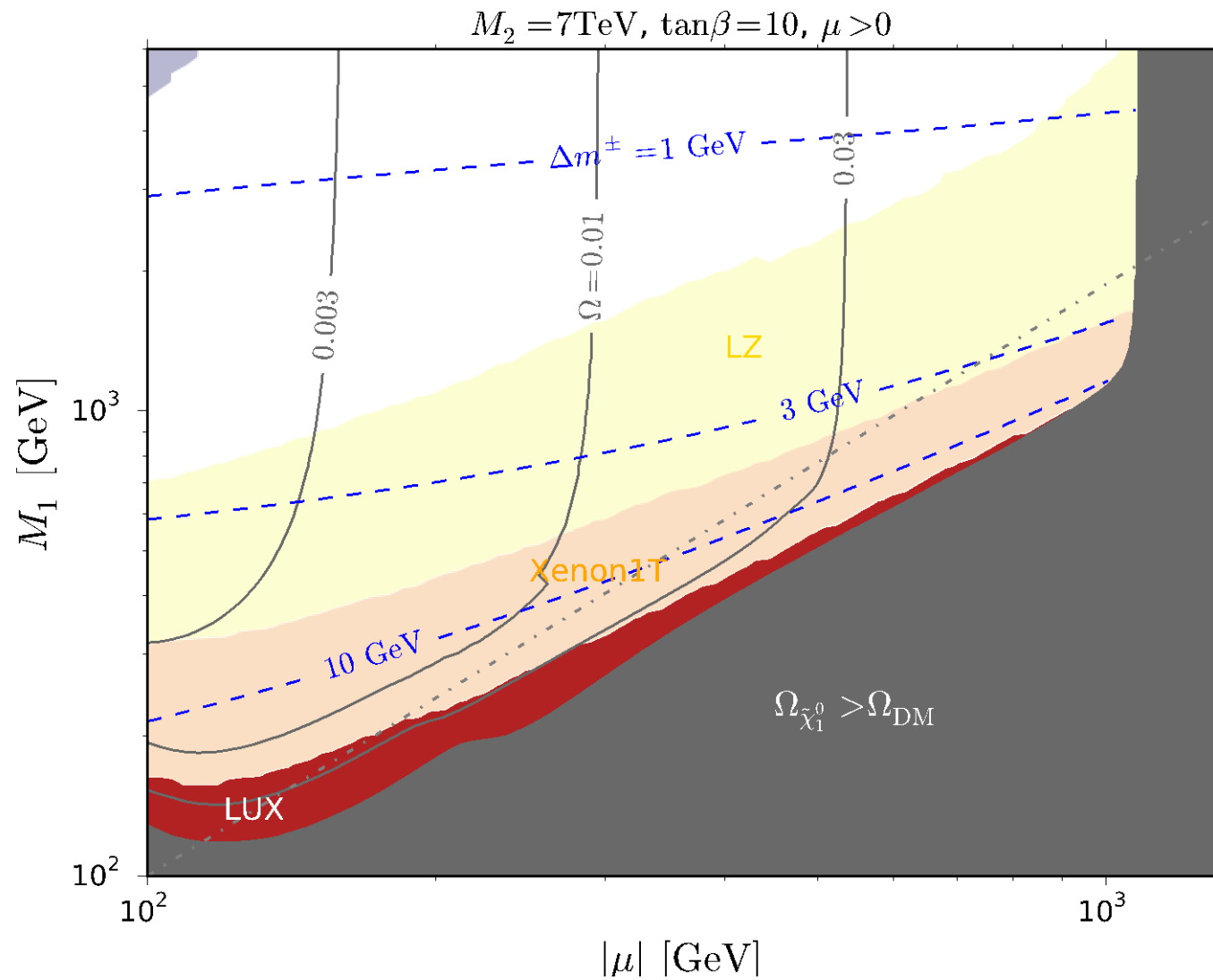
LUX-→ <50 GeV, XENON1T-→ <10GeV

FOLLOWS FROM THE STRONGLY CONSTRAINED FROM  
ABOVE BINO COMPONENT IN THE LSP BY THE DD LIMITS





General correlation: the smaller the spin independent cross section the smaller the mass differences and longer life times of the NLPS



# Collider

Drell-Yan production of  
a pair of gauginos + a hard jet



$$\chi_1^+ \chi_1^-, \chi_2^0 \chi_1^0, \\ \chi_1^+ \chi_1^0, \chi_2^0 \chi_1^+ \text{ etc}$$

$$pp \rightarrow \text{jet} \cancel{E_T} + X$$

depends on mass differences

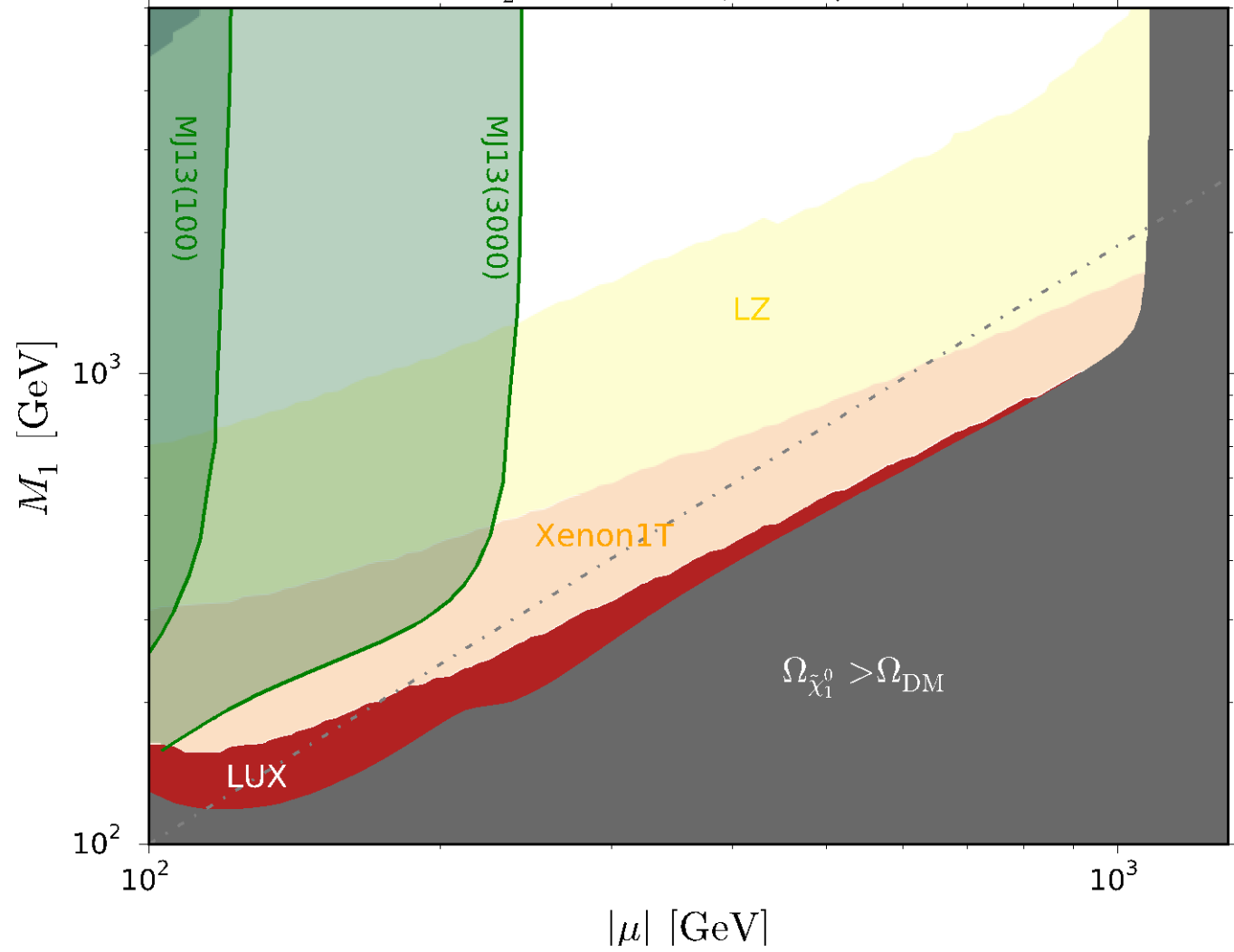
Mass differences for  $\Omega h^2 \leq 0.12$

$\Delta m = 20 - 40 \text{ GeV}$  (rare) soft leptons

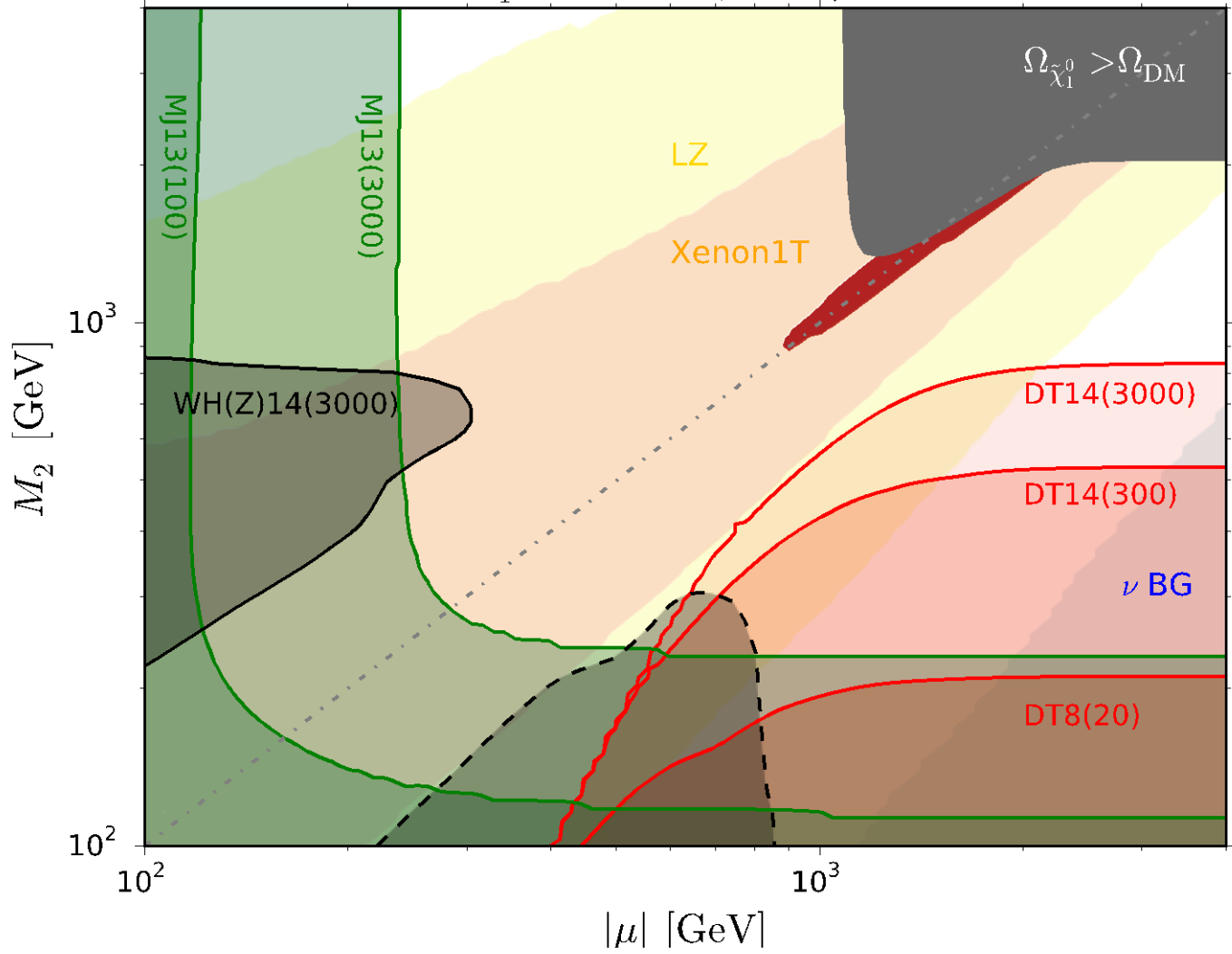
$\Delta m \sim O(1 \text{ GeV})$  (most frequent, monojets,  
mono-Z, mono-gamma)

$\Delta m \sim O(200 - 300 \text{ MeV})$  (disappearing  
tracks)

$M_2 = 7 \text{ TeV}, \tan\beta = 10, \mu > 0$



$M_1 = 7 \text{ TeV}, \tan\beta = 10, \mu > 0$



# CONCLUSIONS (WITH HEAVY SFERMIONS)

DD EXPERIMENTS ARE SENSITIVE TO NEUTRALINOS WITH

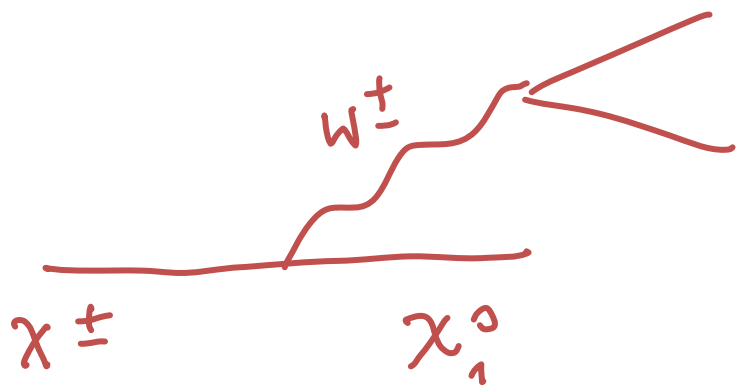
$$\Omega h^2 \approx (10^{-4} - 0.12)$$

THE BOUNDS FROM DD EXPERIMENTS PUSH TOWARDS  
SQUEEZED ELECTROWEAKINO SPECTRUM, WITH STRONG  
IMPACT ON THE COLLIDER SIGNATURES

DD EXPERIMENTS + LHC WILL SIGNIFICANTLY EXPLORE THE  
ELECTROWEAKINO PARAMETER SPACE  
(SOME DEGREE OF COMPLEMENTARITY AND SOME  
OVERLAP IN THE PARAMETER SPACE)

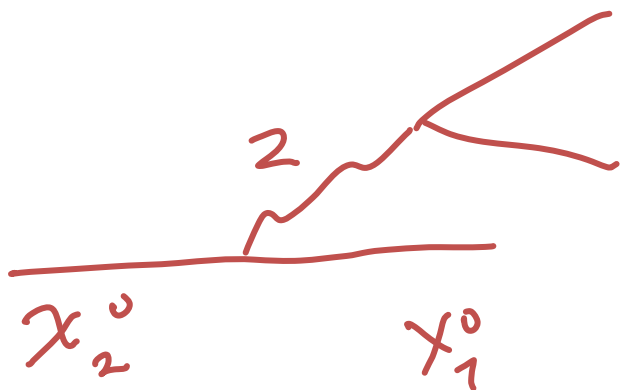
BACKUP





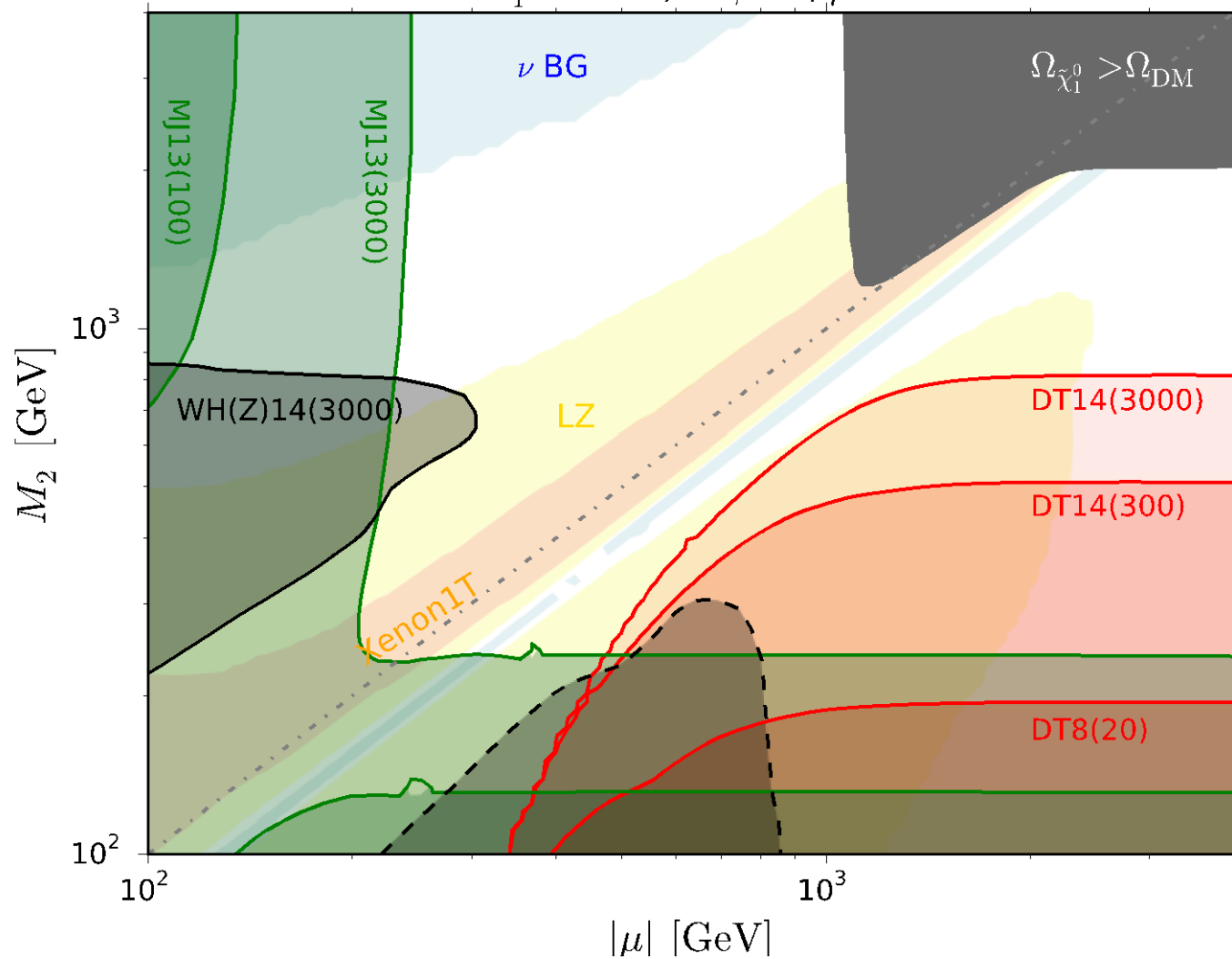
$$\Delta m = \chi_1^+ - \chi_1^0$$

(\* boost)



$$\Delta m = \chi_2^0 - \chi_1^0$$

$M_1 = 7 \text{ TeV}, \tan\beta = 2, \mu < 0$



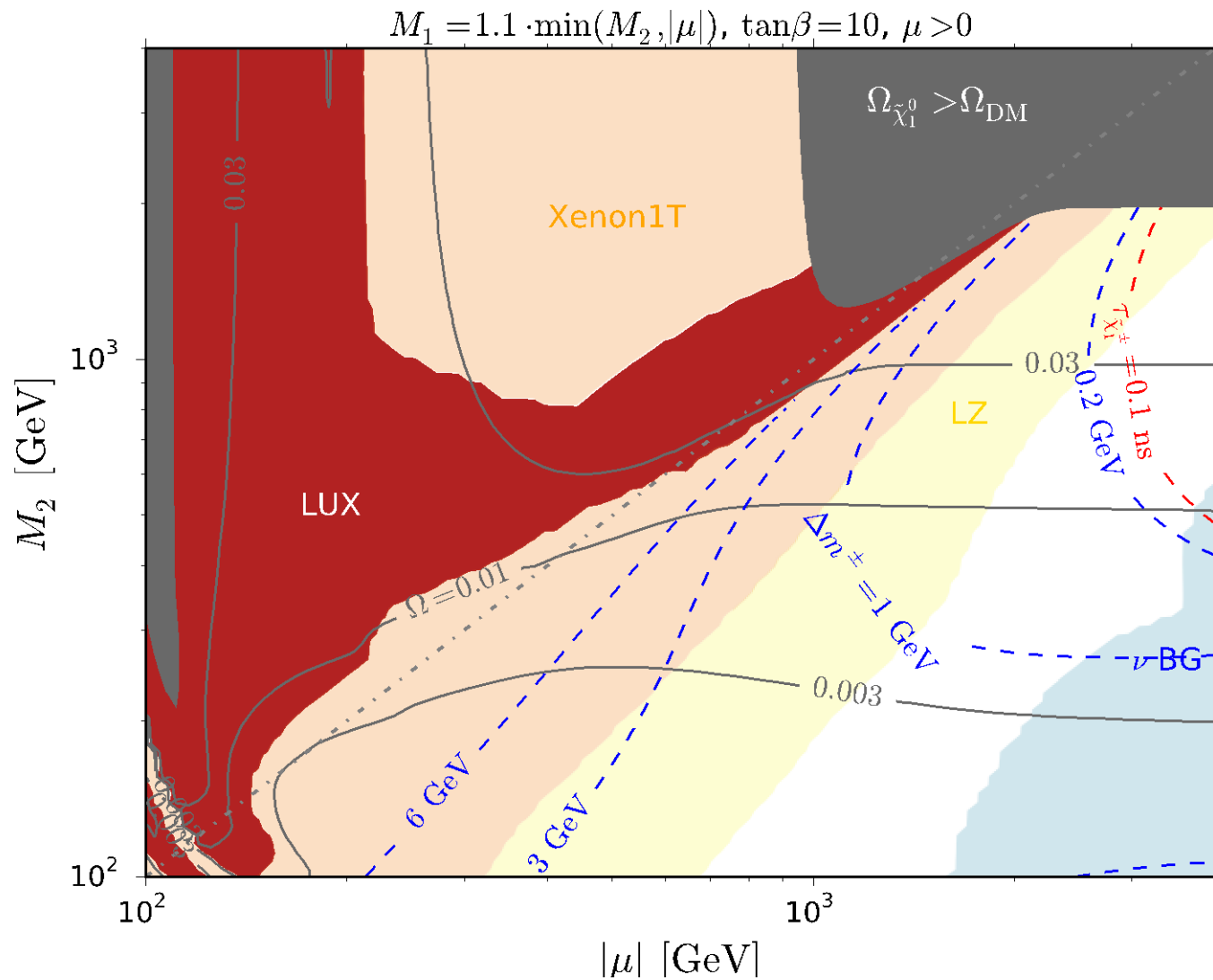
## MASS DIFFERENCES (E.G.)

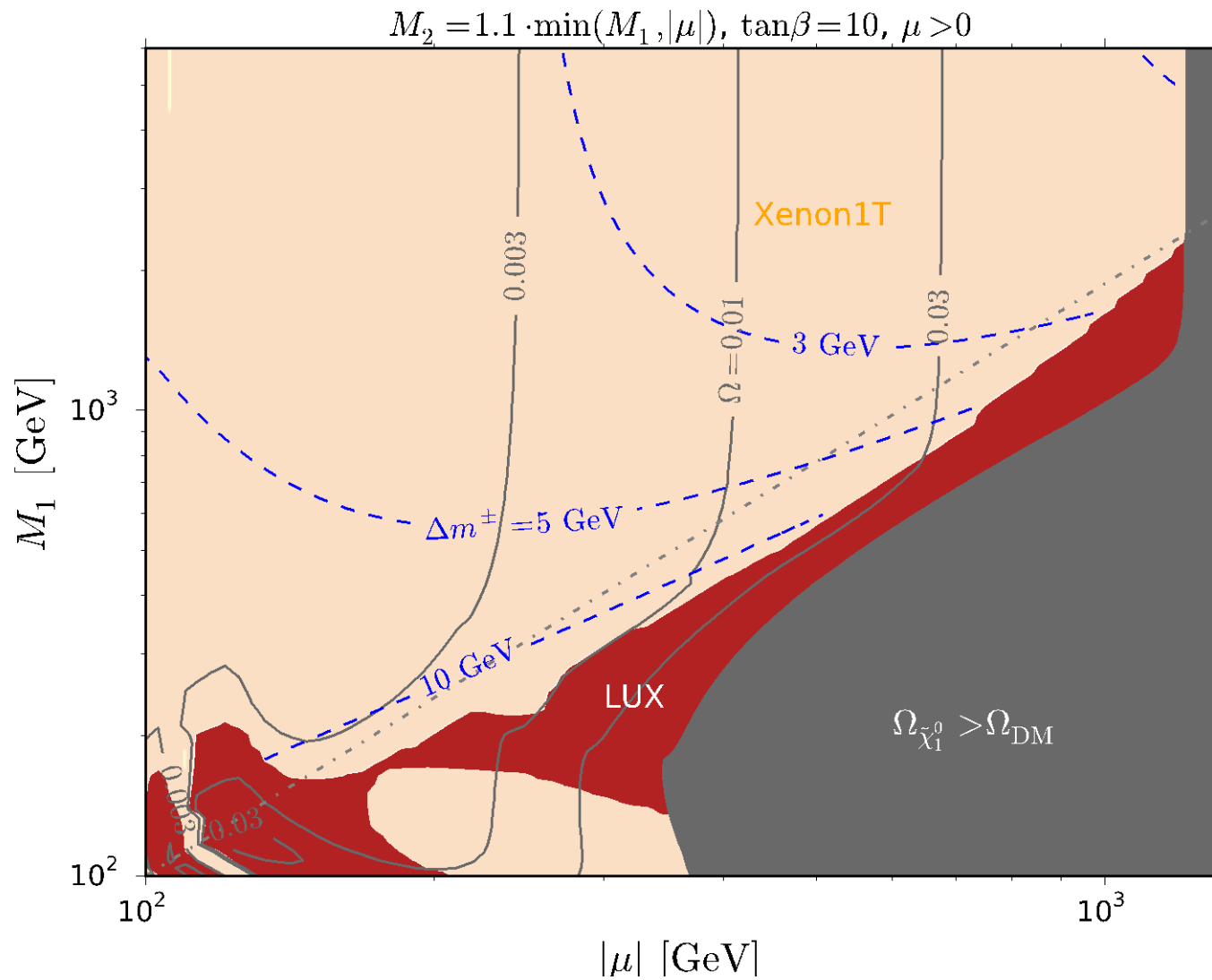
$$|\mu| < |M_1| \ll |M_2| \quad \text{HIGGSINO-BINO}$$

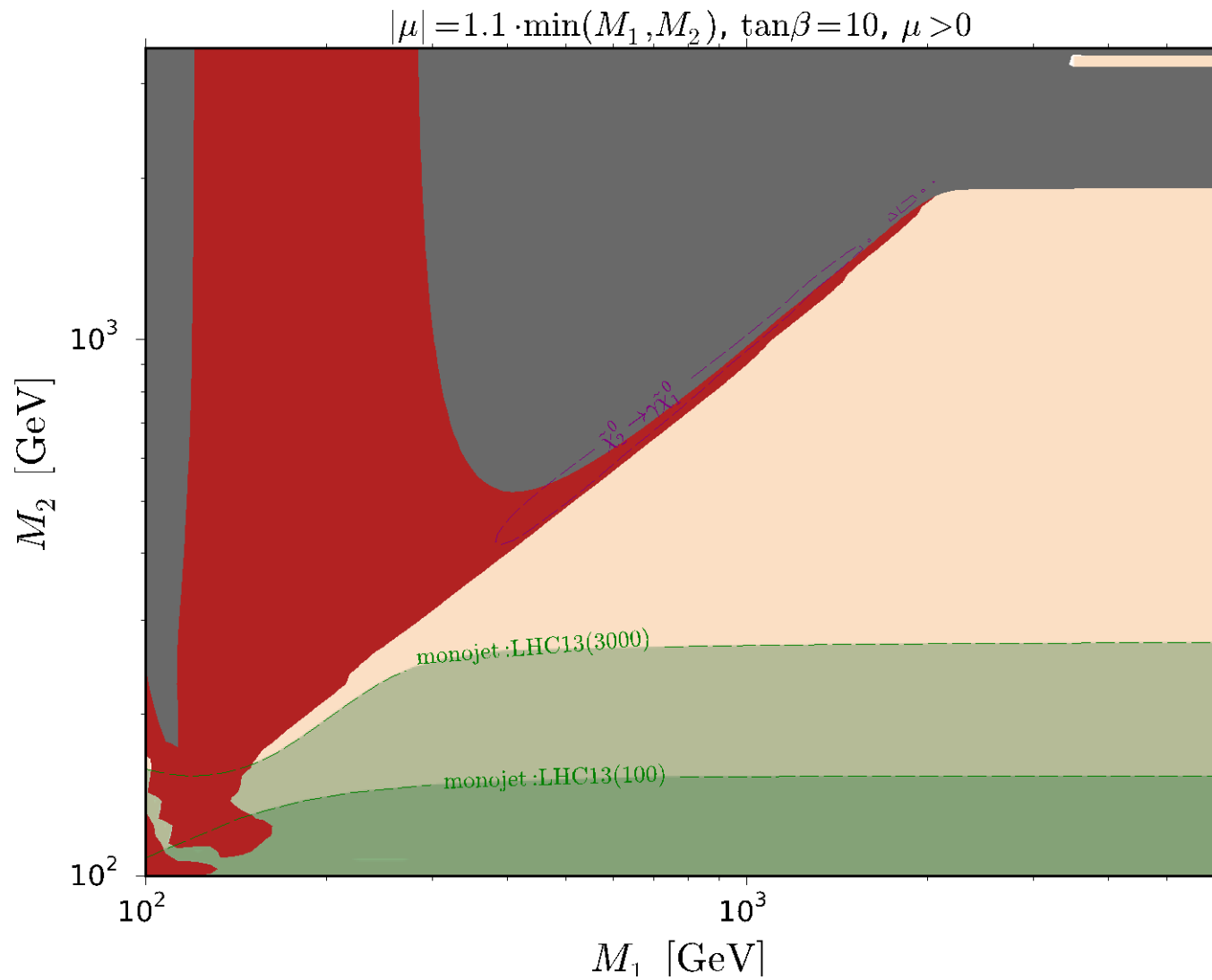
$$m_{\chi_2^0} - m_{\chi_1^0} \approx 2(m_{\chi_1^+} - m_{\chi_1^0}) \approx \frac{1}{2} M_Z^2 \left( \frac{\cos^2 \theta_W}{M_2} + \frac{\sin^2 \theta_W}{M_1} \right)$$

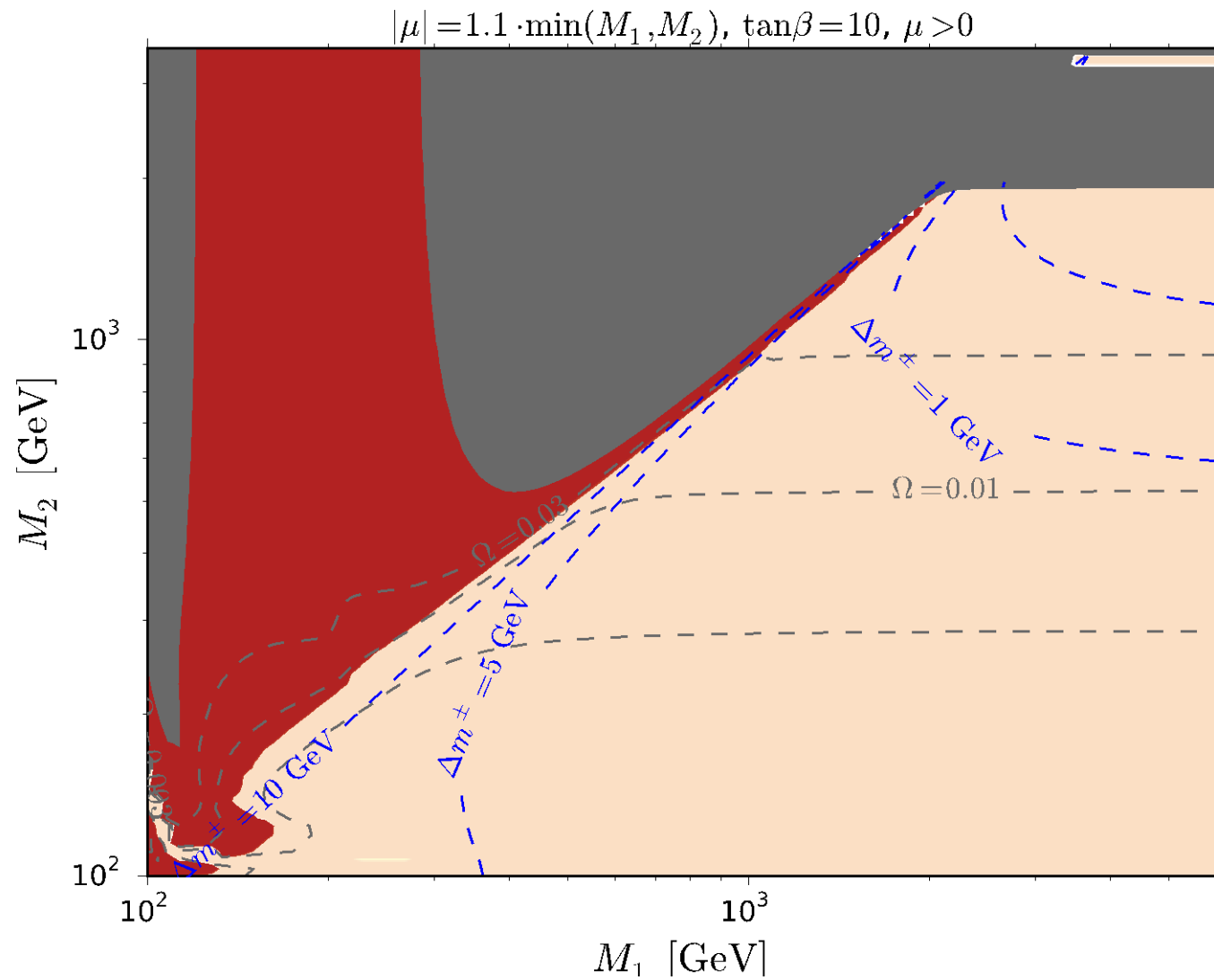
$$|M_2| < |\mu| \ll M_1 \quad \text{HIGGSINO-WINO}$$

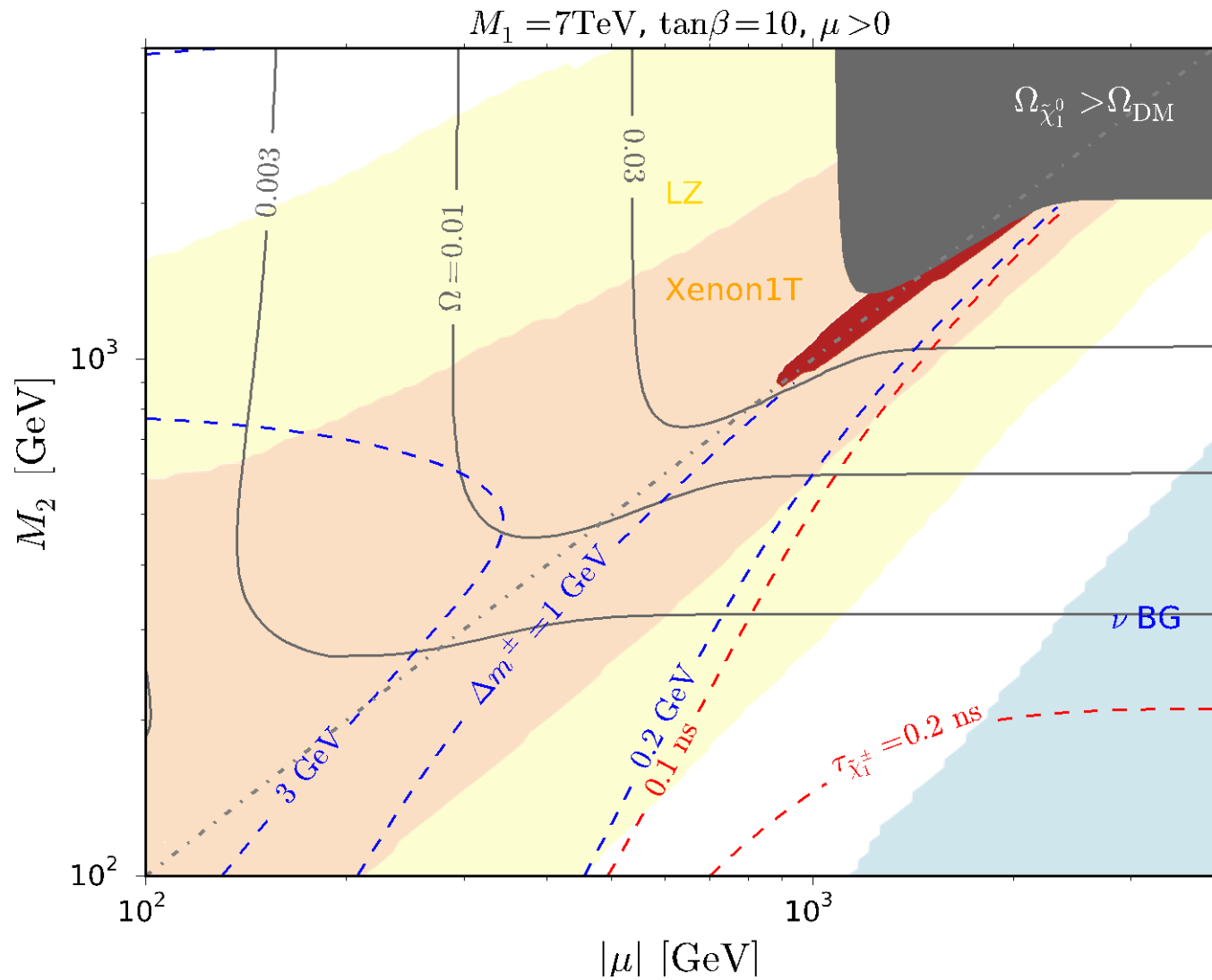
$$m_{\chi_1^+} - m_{\chi_1^0} = \frac{1}{2} \frac{|M_2| M_Z^2}{\mu^4} \cos^4 \theta_W \cos^2(2\beta)$$



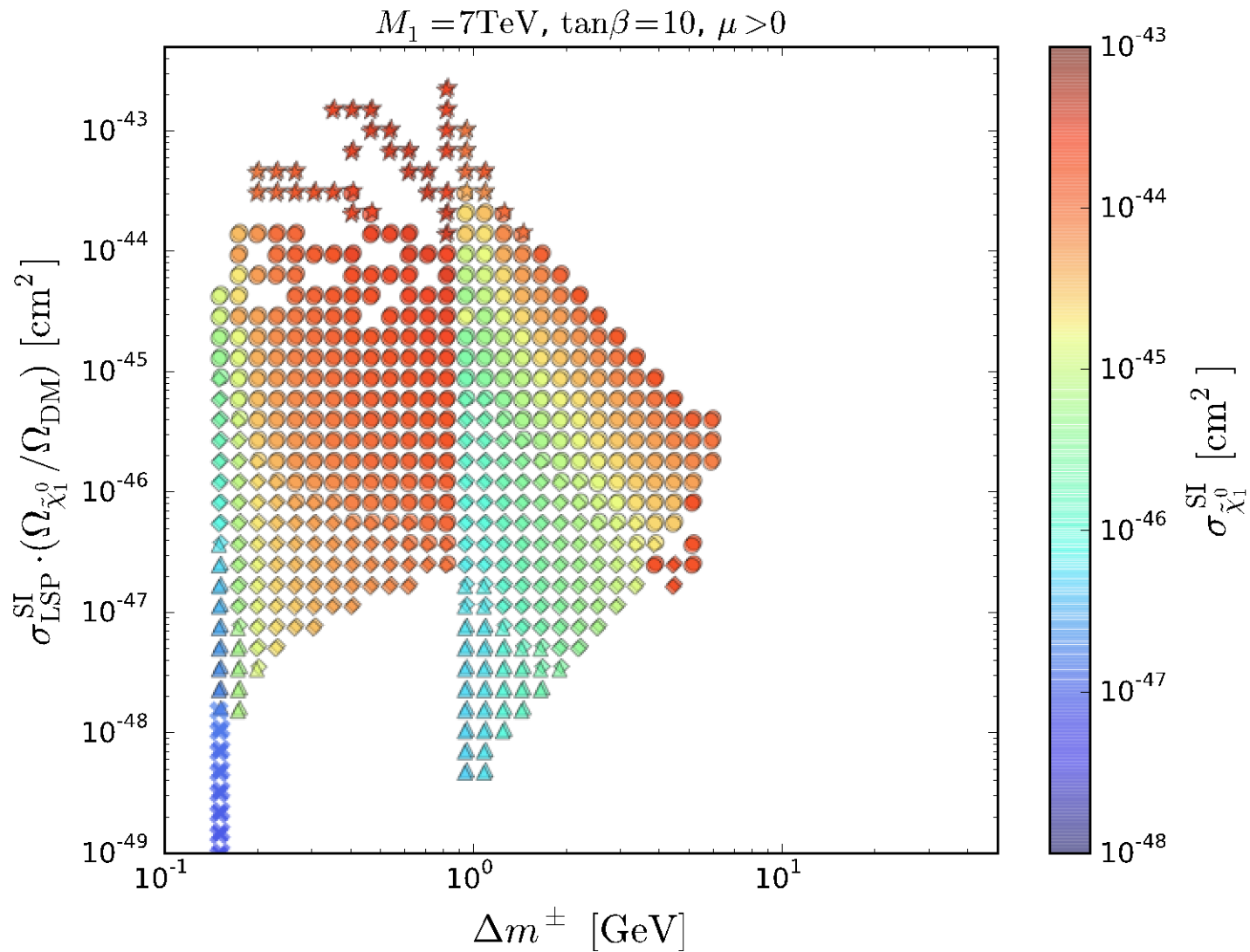




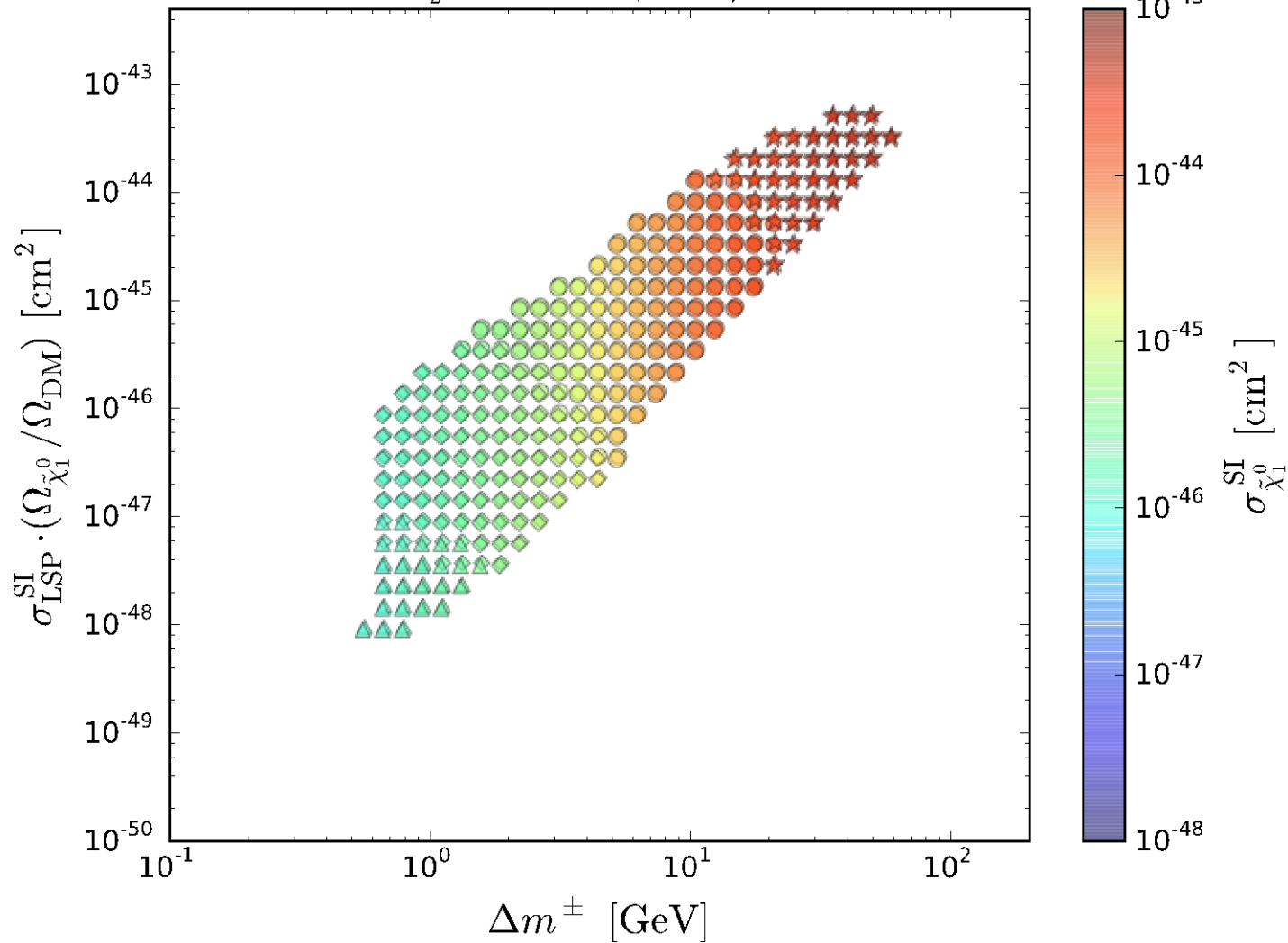


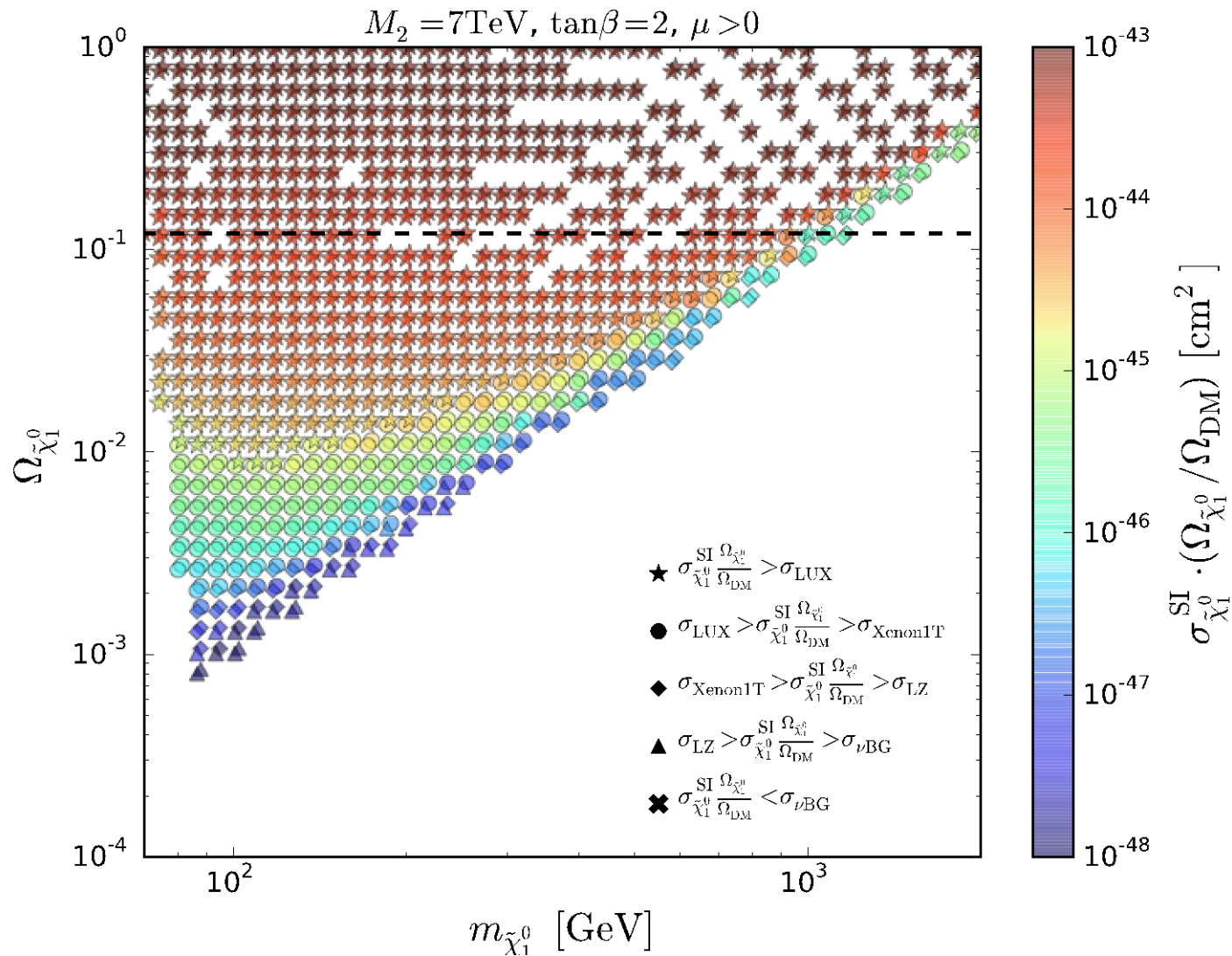




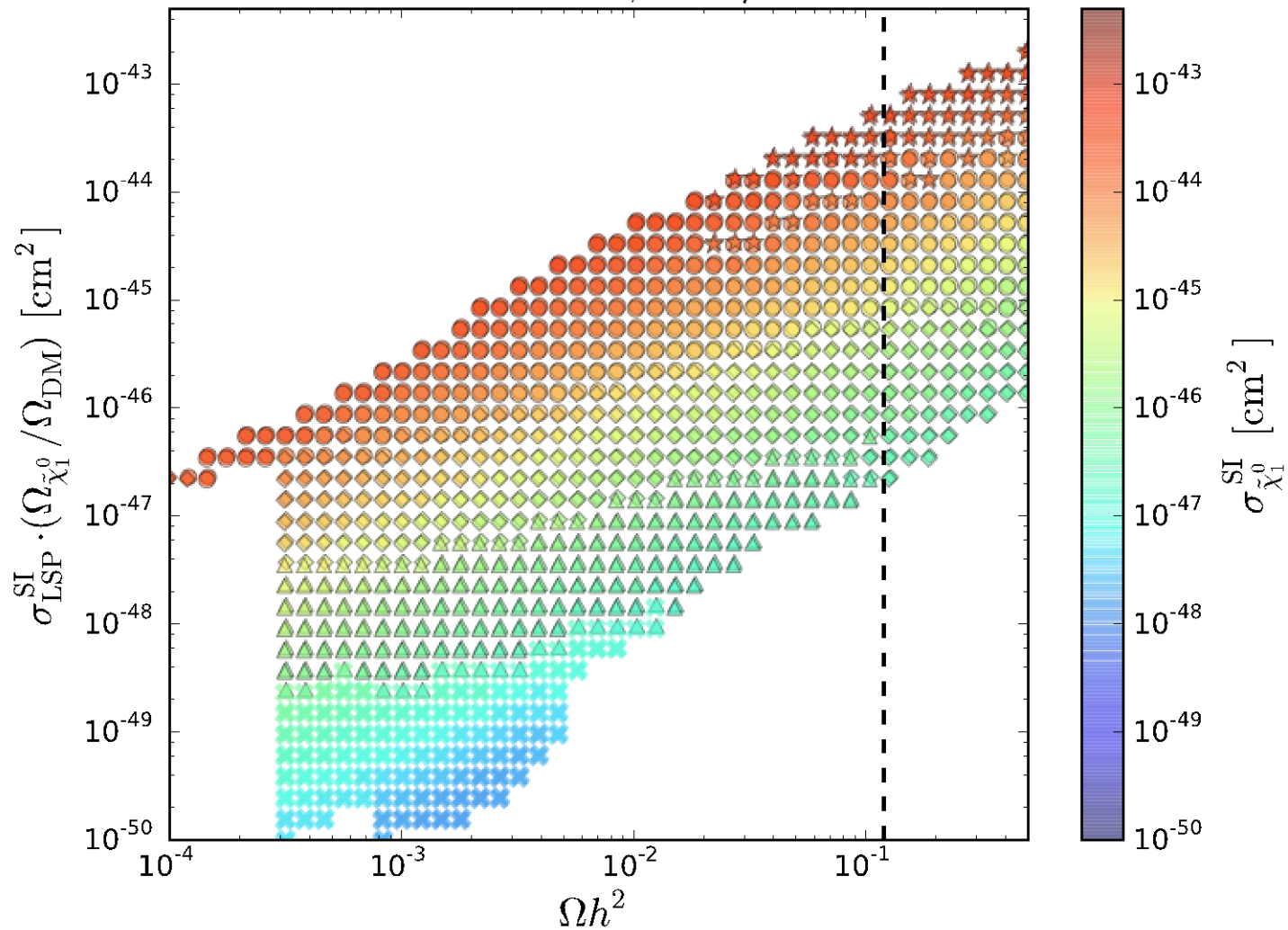


$M_2 = 7\text{TeV}, \tan\beta = 2, \mu > 0$

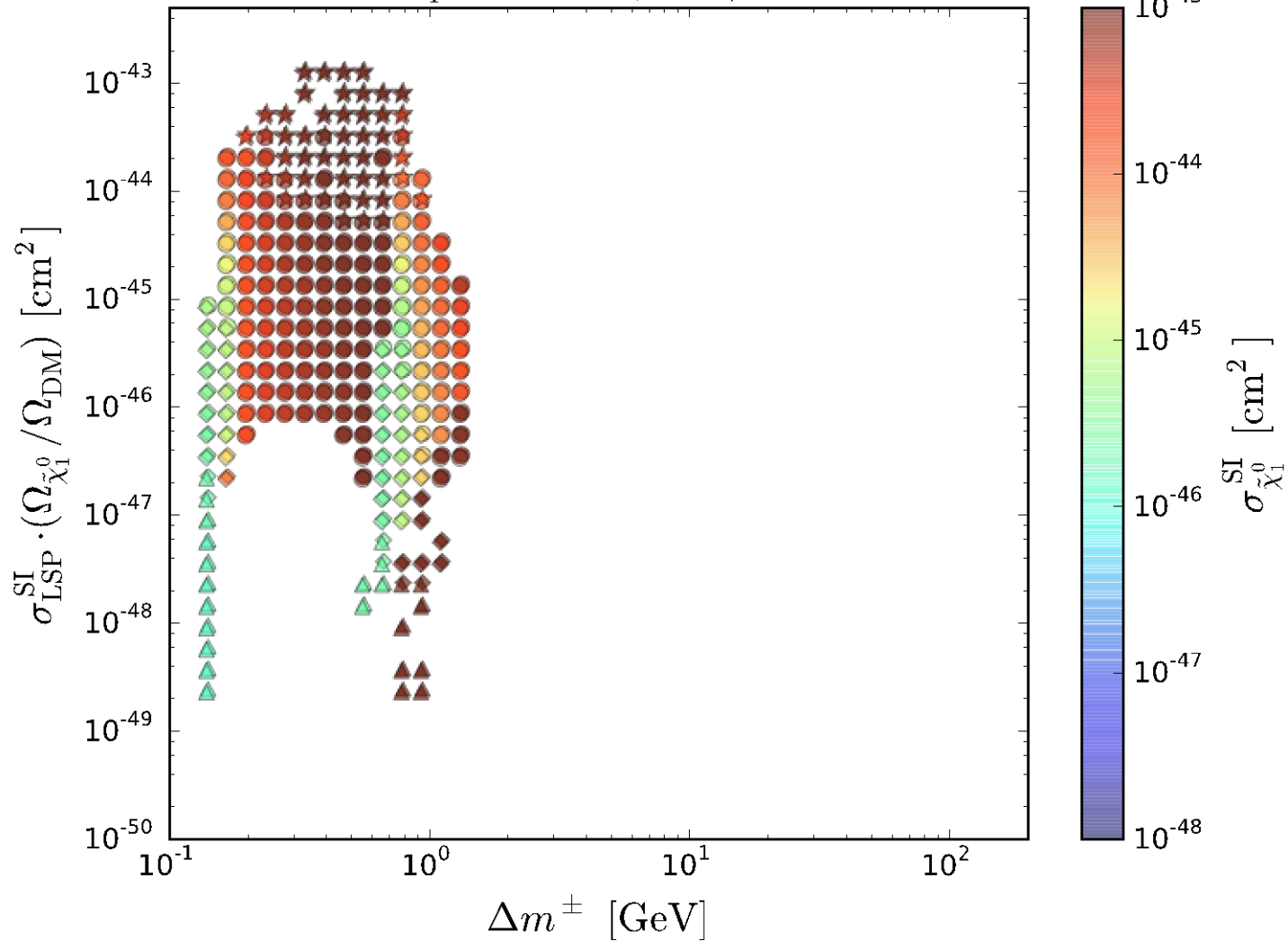


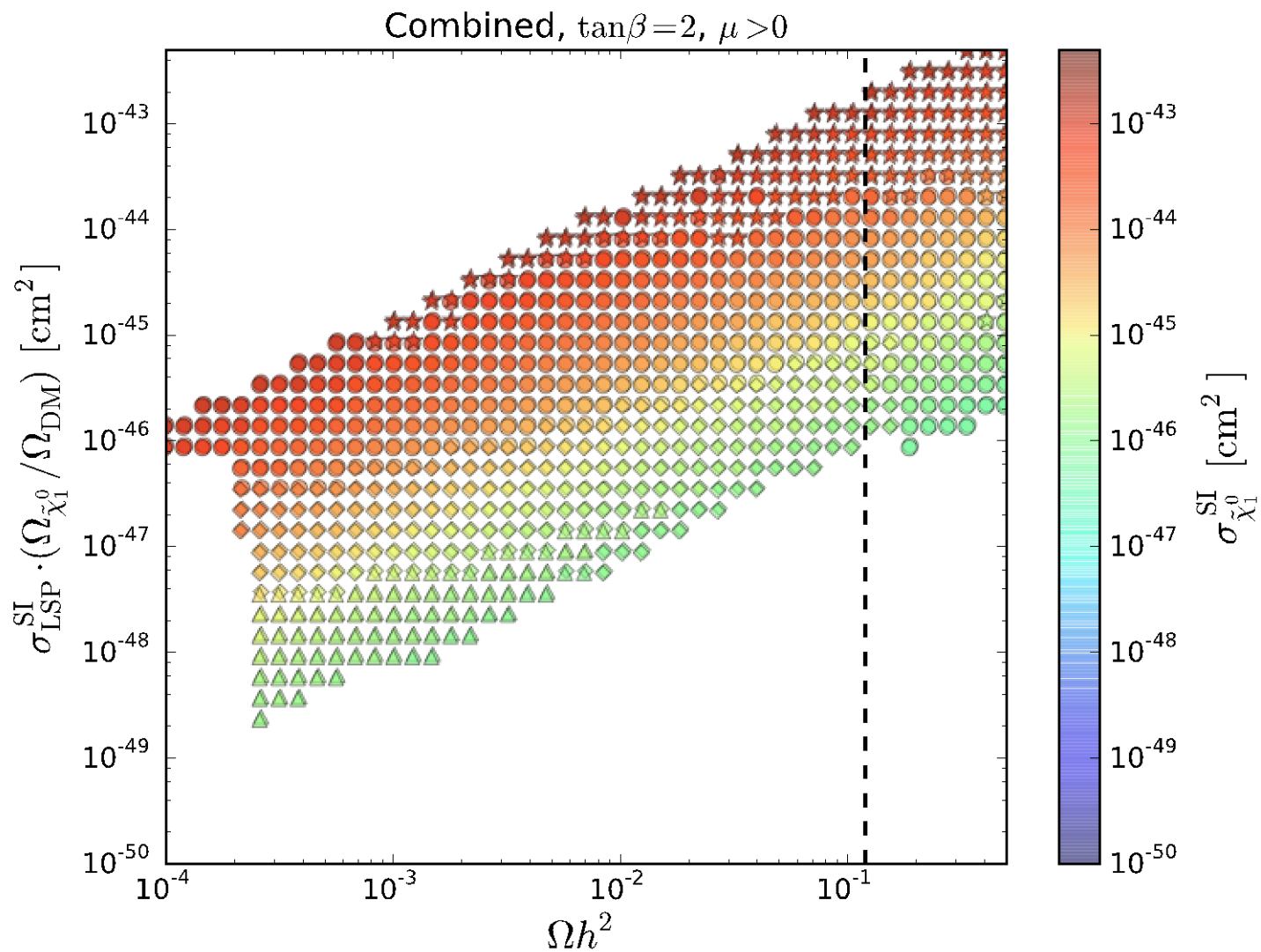


Combined,  $\tan\beta=10, \mu > 0$

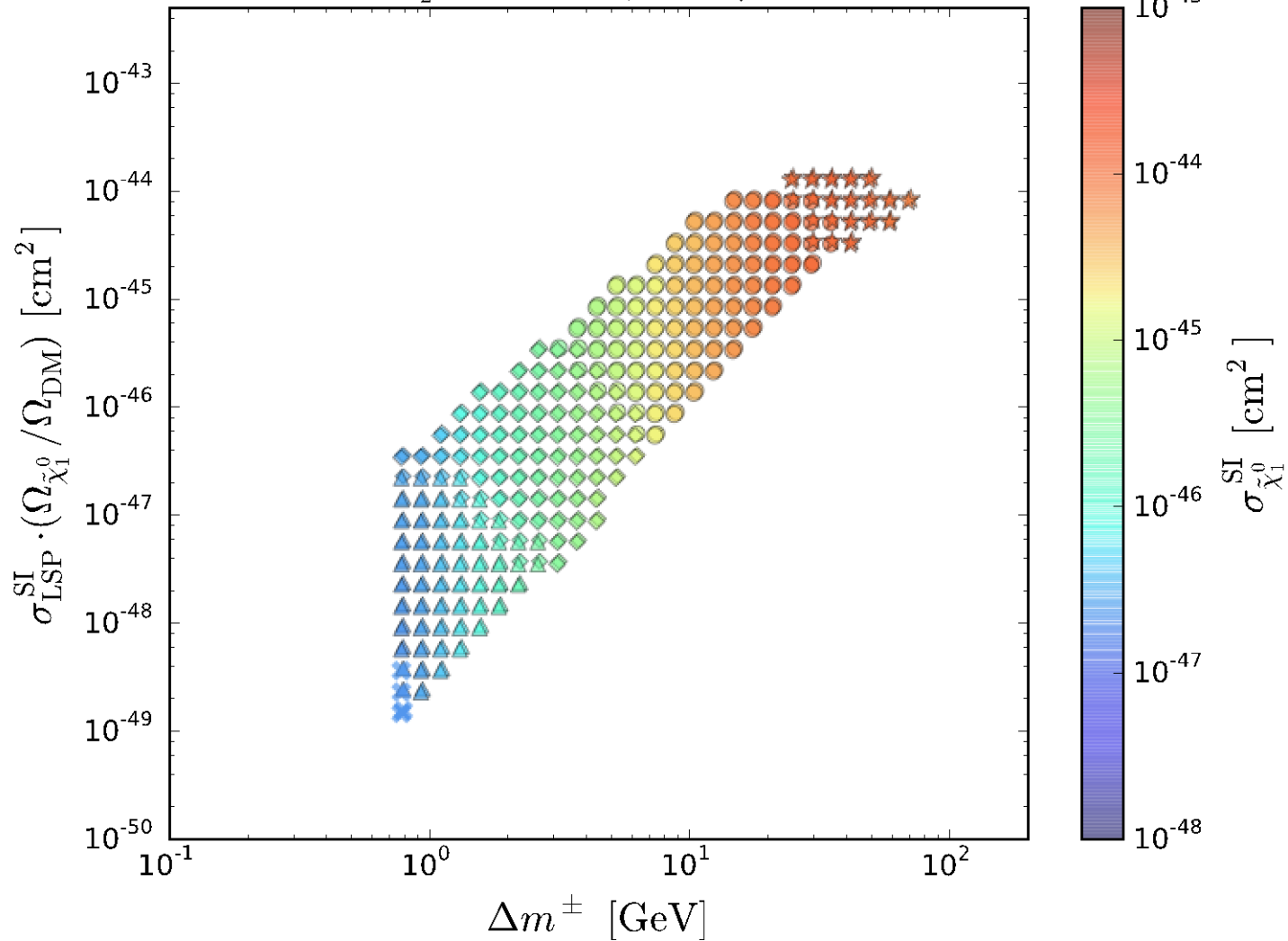


$M_1 = 7\text{TeV}, \tan\beta = 2, \mu > 0$





$M_2 = 7\text{TeV}, \tan\beta = 10, \mu > 0$



$M_2 = 7 \text{ TeV}, \tan\beta = 10, \mu > 0$

