Sharpening cosmological constraints with EFT methods



Kudos to Rafael Porto

Nick Kokron (KIPAC/Stanford) w/ Jérôme Gleyzes, Pierre Zhang, Leonardo Senatore, Dida Markovic, Florian Beutler, Héctor Gil-Marín, others

EFT?

- *Effective Field Theory* (EFT for short) is a general paradigm
 Valid when hierarchy between length scales exists
- Wild success in particle theory, condensed matter/stat mech, biophysics, gravitational waves, physics of inflation, hydrodynamics
- Baumann et al. (2010) proposed using EFT to study dark matter as a cosmological fluid • E.g., Large-scale structure formation should not depend on the microphysics of stellar structure
- Allows us to import particle theory technology into cosmology
- Provides an "analytic", physically well motivated, approach to computing quasilinear corrections to cosmological observables
 - \circ This is a good thing.

Doesn't SPT do this already?

- Other approaches to cosmological perturbation theory have the same objective
- Several have pitfalls
 - For example, SPT gives formally infinite answers already at one loop
- EFT is as an extension of SPT
 - Uses the tools of particle theory to control for these infinities
- Also consistently includes:
 - Tracer biasing from gravitational physics
 - Stochastic biases
 - Higher-order RSD
 - "IR-resummation", for BAO purposes
 - Baryons



• Normally the first and second moments of the collisionless Boltzmann eqn. are solved sequentially:



- Not zero in EFT, expanded in combinations of large-scale quantities allowed by Galilean + rotational symmetry
- Higher order contributions are always smaller, due to the scale hierarchy mentioned above
 Like a matryoshka doll! (I told you it was relevant)
- Explicitly:

$$\tau^{ij} = p_b \delta^{ij} + \rho \left[c_s^2 \delta^{ij} \delta - \frac{c_{bv}^2}{Ha} \delta^{ij} \vec{\nabla} \cdot \vec{v} - \frac{3}{4} \frac{c_{sv}^2}{Ha} \left(\partial_j v^i + \partial^i v^j - \frac{2}{3} \delta^{ij} \partial_k v^k \right) \right] \\ + \Delta \tau^{ij} + \cdots$$

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• Combined into one free parameter, correction will look like:

$$P_{\rm ct}(k) = -\tilde{c}_s^2 \frac{k^2}{k_{\rm NL}^2} P_{\rm lin}(k)$$

• Quality of fit improves dramatically by fitting for this sound speed simultaneously



The Universe isn't "Fair and Balanced"TM: Some bias required

- Formalism so far only treats dark matter density
- We observe galaxies, which trace DM
 - DM halos themselves are biased tracers of the dark matter density
- EFT principles guide us in how to describe biasing procedure
 - Desjacques, Jeong & Schmidt 2018 is the most comprehensive review on this
- Biasing can be split into two contributions: $\delta_{h} = b_{1}\delta + b_{\nabla^{2}\delta}\nabla^{2}\delta + \frac{1}{2}b_{2}\delta^{2} + b_{K^{2}}(K_{ij})^{2} + \frac{1}{6}b_{3}\delta^{3} + b_{\delta K^{2}}\delta(K_{ij})^{2} + b_{K^{3}}(K_{ij})^{3} + b_{td}O_{td}^{(3)}$ Deterministic $+ \varepsilon + \varepsilon_{\delta}\delta + \varepsilon_{\delta^{2}}\delta^{2} + \varepsilon_{K^{2}}(K_{ij})^{2} \longrightarrow \text{Stochastic}$



Wechsler & Tinker 2018

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$$\delta_r(\vec{k}) = \delta(\vec{k}) + \int d^3x e^{-i\vec{k}\cdot\vec{x}} \left(\exp\left[-i\frac{k_z}{aH}v_z(\vec{x})\right] - 1 \right) (1 + \delta(\vec{x}))$$

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• To cubic order, will give $\delta_r(\vec{k}) = \delta(\vec{k}) + \mu^2 \frac{\dot{\delta}(\vec{k})}{H} - \epsilon^{zij} \frac{\mu k_i}{k} \frac{\pi_{V,j}}{H} - \frac{1}{2} \left(\frac{k \mu}{aH}\right)^2 \left([v_z^2]_{\vec{k}} + [v_z^2 \delta]_{\vec{k}}\right) + \frac{i}{6} \left(\frac{k \mu}{aH}\right)^3 [v_z^3]_{\vec{k}}$

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Jump forward in time ~4 years (and around 30 papers)

"(...) long-wavelength statistics that are routinely observed in LSS surveys can be finally computed in the EFTofLSS. This formalism thus is ready to start to be compared directly to observational data." 1610.09321 (Perko et al)

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Jump forward in time another 2 years (10 more papers)

A model fit for a survey I: A cornucopia of free parameters

- With the tools outlined above we may write a model for the anisotropic power spectrum at one loop (third order in δ , v)
- It looks like this

$$\begin{split} P_g(k,\mu) &= Z_1(\mu)^2 P_{11}(k) \\ &+ 2 \int d^3 q \, Z_2(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q})^2 P_{11}(|\boldsymbol{k}-\boldsymbol{q}|) P_{11}(q) + 6 Z_1(\mu) P_{11}(k) \int d^3 q Z_3(\boldsymbol{q},-\boldsymbol{q},\boldsymbol{k}) P_{11}(q) \\ &+ 2 Z_1(\mu) P_{11}(k) \left(c_{\rm ct} \frac{k^2}{k_{\rm NL}^2} + c_{r,1} \mu^2 \frac{k^2}{k_{\rm M}^2} + c_{r,2} \mu^4 \frac{k^2}{k_{\rm M}^2} \right) + \frac{1}{\bar{n}_g} \left(c_{\epsilon,1} + c_{\epsilon,2} \frac{k^2}{k_{\rm M}^2} + c_{\epsilon,3} f \mu^2 \frac{k^2}{k_{\rm M}^2} \right) \end{split}$$

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There are actually 12 bias parameters a priori. They can be re-cast into a basis of four less-degenerate parameters.

A model fit for a survey II: window functions and likelihoods

- Window functions in 3D surveys are applied in real space
 - EFT predictions diverge at large k. This is a bad thing.

$$P_{\ell}(k) \stackrel{\text{\tiny FFTlog}}{\to} \xi_{\ell}(r) \to \xi_{\ell}^{(W)}(r) = Q_{\ell,\ell'}(r) \cdot \xi_{\ell}(r) \stackrel{\text{\tiny FFTlog}}{\to} P_{\ell}^{(W)}(k)$$

- Lessons from CMB: mode-coupling window functions à la pseudo-Cl $P_{\ell}^{W}(k) = \sum_{k',\ell'} W_{\ell,\ell'}(k,k') P_{\ell'}^{W}(k')$
- 6 of the 10 bias parameters only enter linearly in the prediction
 - They factor quadratically in the likelihood
- Analytically marginalized, leading to only 3+3 parameters being varied
 - Strong, physical priors can be used on marginalized params

Validating on simulations

- Extensive validation across various numerical suites of simulations
 - You can ask me for a list
- Of specific interest are 'Lettered Challenge Boxes' from Hand et al. (2017)
 - (Their model has 13 free parameters)
- High resolution, varying HODs for each catalog
 - Covariance matrices scaled to BOSS DR12 volume



Validating on simulations

- Errors decrease well as we increase kmax
- Posteriors contain the truth within 1σ, normally
 - Looking at product of posteriors reveals slight biases in A_s and h, ~0.5 σ
- Looking at bias contours reveals we can tell the HOD apart
 - Without significantly affecting the cosmology



Validating on simulations

- Other suites show similar results
 Modulo their quality at small scales
- Our model can also self-consistently includes bispectrum predictions
 - Apparently not much cosmology improvement
- Are the simulations to blame?



Validating on simulations: a heavenly anchor

- A "light" way to combine with CMB information: sound horizon at decoupling
- r_d is set by early universe physics
 - Imposes strong constraints on $h \Omega_m$ plane.
- Constraints aren't biased by this prior



Onto data!

- Confident in our validation we take another look at BOSS DR12
- Simulated constraints are already exciting • ~Plank 13 for data released in 2015!
- Error bars in data comparable to
- simulated chains
- Stay tuned. arXiv: 1812.XXXX
 - hopefully

Data

An aside

- Barriers exist between community working on this and survey cosmologists
- Sociology aside, I think these techniques can be valuable for DES/DESI/LSST
 Please don't shoot the messenger
- We now have (moderately) fast codes and analysis pipelines set up to do this for spectroscopic surveys
- Extending to photometric observables is feasible
 - We can push the modelling harder for lensing with no additional parameters

Wrapping up

- Extensive tests of EFT modelling for the anisotropic power spectrum have been performed
- We can wring out even more constraining power from existing data
- Combining these techniques with survey data is tantalizing
- Or perhaps we should just go to the field level as Marko explained earlier? :)



Thank you for your time!





An extra slide: planck shift

