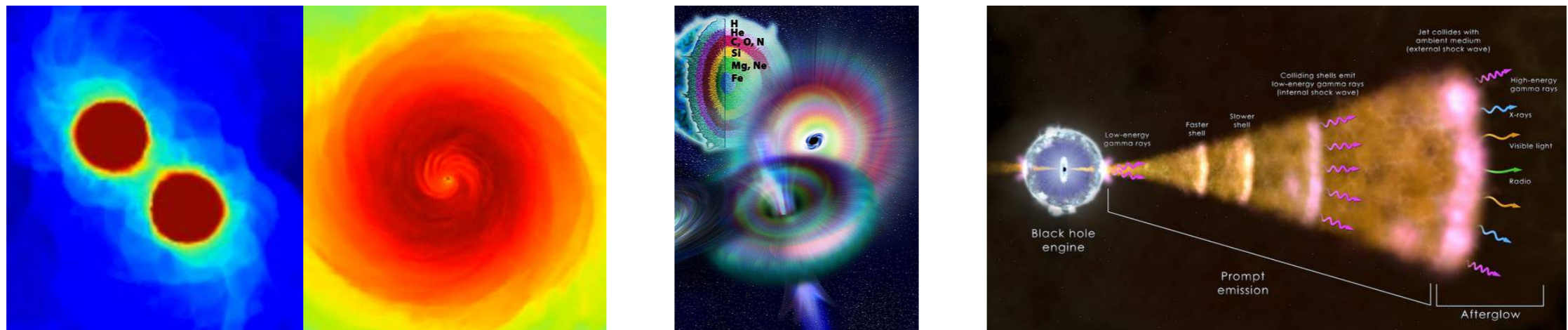


# ICTP school: GRB central engine slides



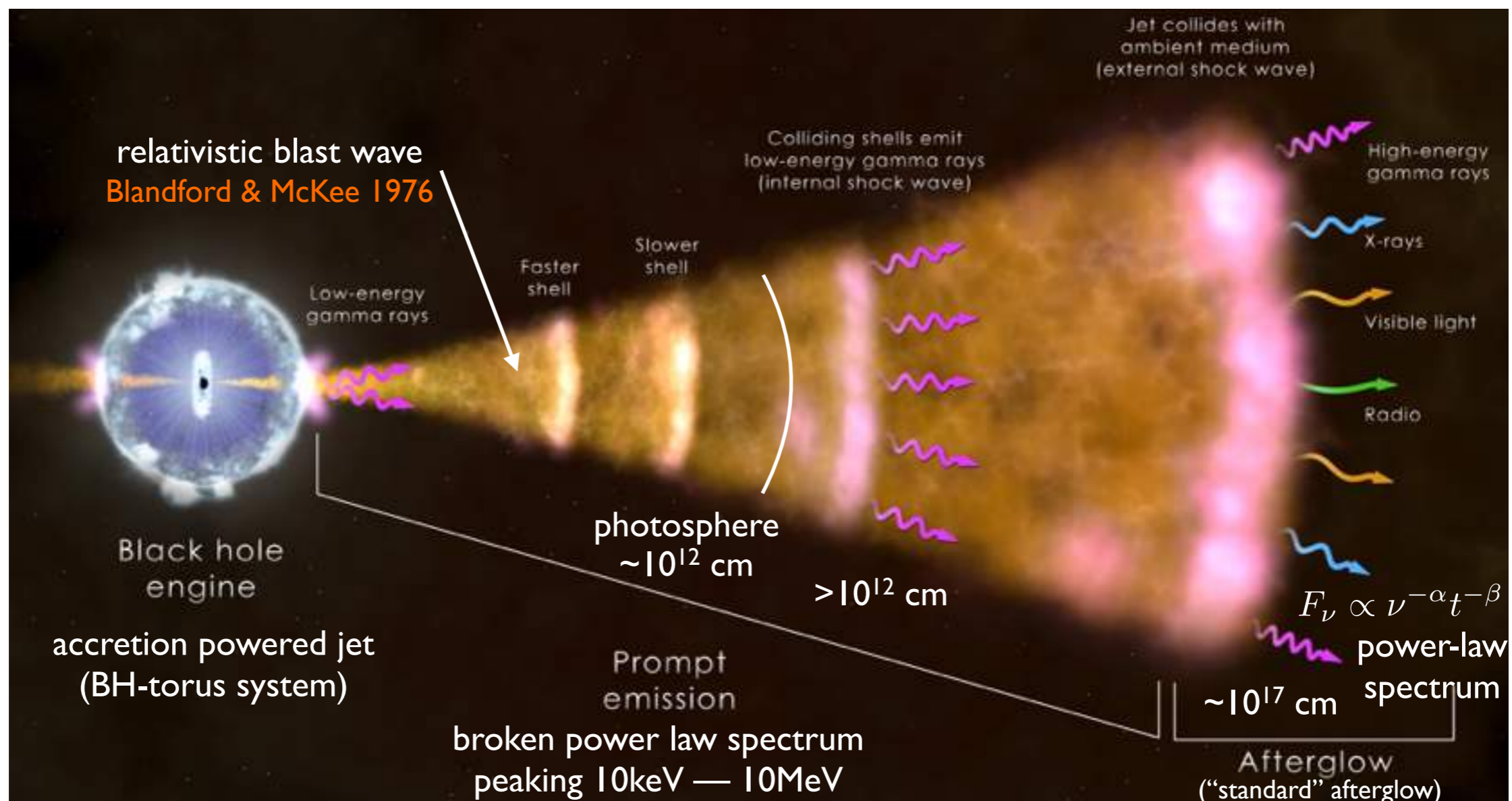
Daniel M. Siegel

*Center for Theoretical Physics & Columbia Astrophysics Laboratory*

*Columbia University*

ICTP school *The Sound of Space-time: The Dawn of Gravitational Wave Science*,  
Sao Paulo, Dec 10-14, 2018

# The canonical GRB picture



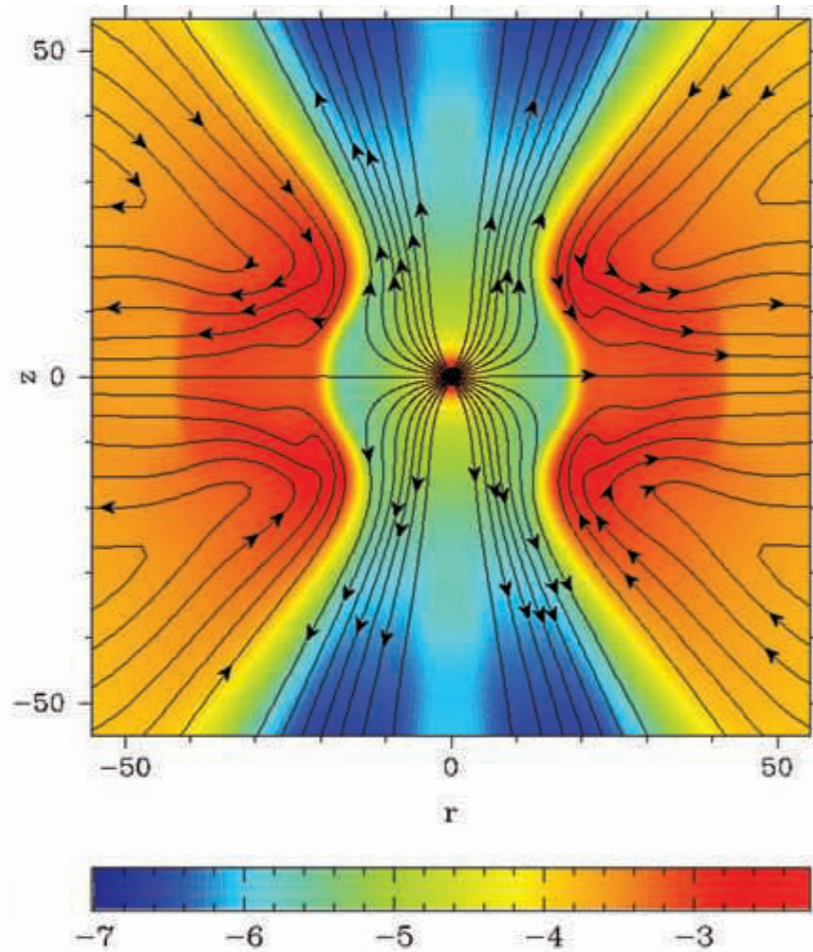
What actually powers these explosions?  
Where does the energy come from?

# Preliminaries

# What is the central engine of GRBs?

Rapidly rotating black hole?

Blandford-Znajek mechanism

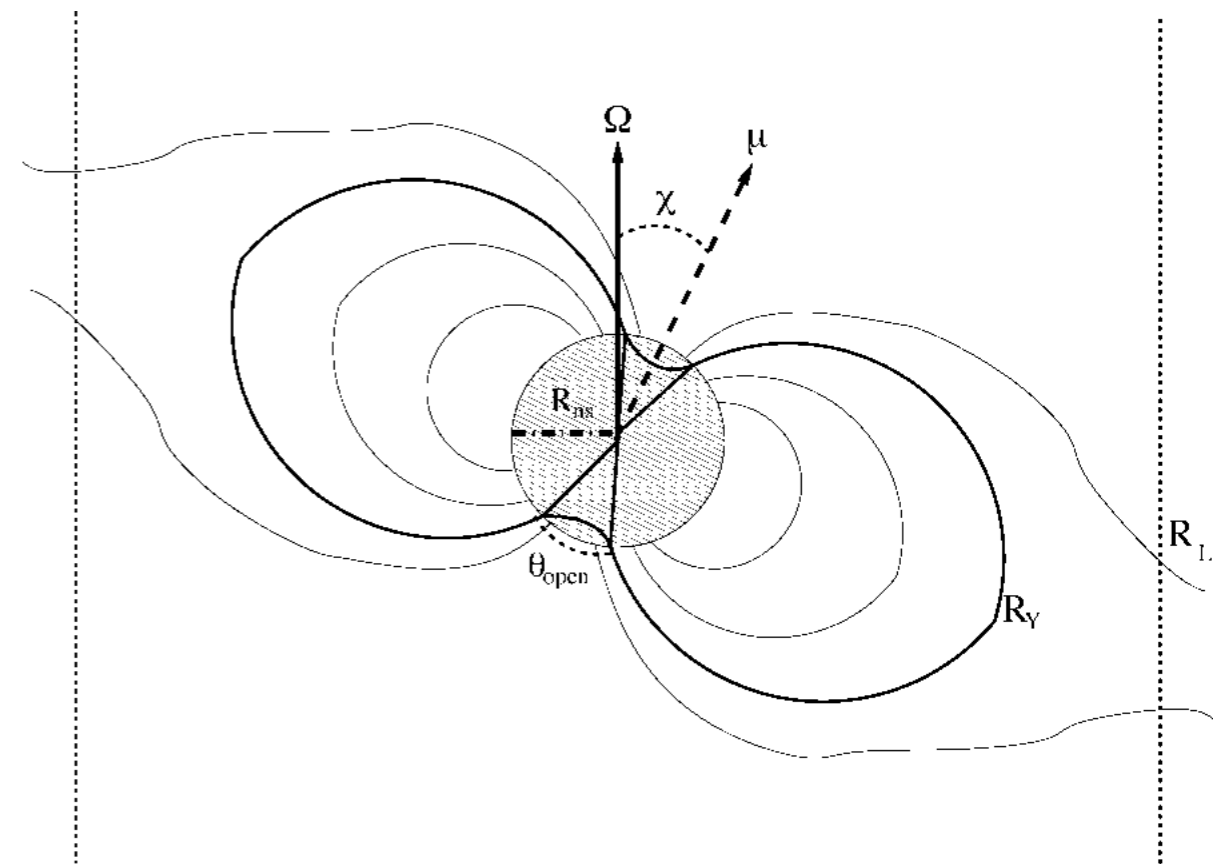


Komissarov & Barkov 2009

$$E_{\text{rot}} = 10^{54} f \left( \frac{a}{0.9} \right) \left( \frac{M_{\text{BH}}}{4M_{\odot}} \right) \text{ erg}$$

Rapidly rotating, highly magnetized neutron stars (ms-magnetars?)

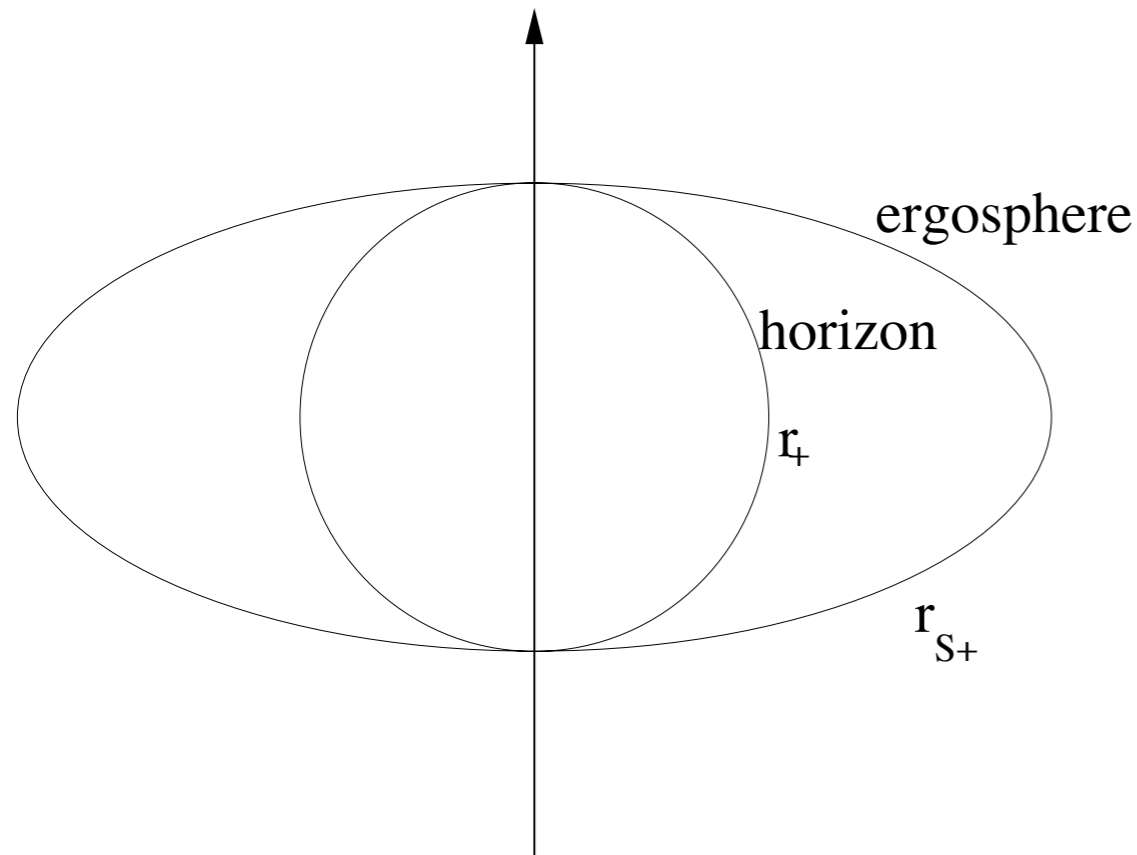
Magnetar model for GRBs



Metzger+ 2011

$$E_{\text{rot}} \approx 3 \times 10^{52} \text{ erg} \left( \frac{M_{\text{ns}}}{1.4 M_{\odot}} \right) \left( \frac{R_{\text{ns}}}{12 \text{ km}} \right)^2 \left( \frac{P}{\text{ms}} \right)^{-2}$$

# Basic Kerr-BH anatomy



Boyer-Lindquist coordinates:

$$ds^2 = g_{tt}dt^2 + 2g_{t\phi}dtd\phi + \gamma_{rr}dr^2 + \gamma_{\phi\phi}d\phi^2 + \gamma_{\theta\theta}d\theta^2$$

$$\begin{aligned} g_{tt} &= (2Mr/\Sigma) - 1, & \Sigma &= r^2 + a^2 \cos^2 \theta, \\ g_{t\phi} &= -2aMr \sin^2 \theta / \Sigma, & \Delta &= r^2 - 2Mr + a^2, \\ \gamma_{rr} &= \Sigma / \Delta, & A &= (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta, \\ \gamma_{\phi\phi} &= A \sin^2 \theta / \Sigma, & a &= J/M, \quad -1 < a/M < +1. \\ \gamma_{\theta\theta} &= \Sigma, \end{aligned}$$

Writing as a **3+1 split** (remember from NR lecture):

$$ds^2 = -(\alpha^2 - \beta^2)dt^2 + 2\beta_i dx^i dt + \gamma_{ij} dx^i dx^j$$

$$\begin{aligned} \alpha^2 &= -1/g^{tt} = \Delta\Sigma/A, \\ \beta^2 &= \alpha^2 + g_{tt} = 4a^2 r^2 \sin^2 \theta / A, \\ \beta^\phi &= \alpha^2 g^{t\phi} = -2aMr/A, \quad \beta^r = \beta^\theta = 0. \end{aligned}$$

**Horizon location:**  $\Delta = 0$

$$r_{\pm} = M \pm \sqrt{M^2 - a^2}.$$

**Ergosphere:**  $g_{tt} = \beta^2 - \alpha^2 = 0$

$$r_{S\pm} = M \pm \sqrt{M^2 - a^2 \cos^2 \theta}.$$

Inside, the killing vector  $t^\alpha = \partial_t$  not timelike anymore:

$$g_{\mu\nu} t^\mu t^\nu = g(t^\mu, t^\nu) = g_{tt} > 0$$

Identify using general form of metric components:

$$g_{\mu\nu} = \begin{pmatrix} -\alpha^2 + \beta_k \beta^k & \beta_i \\ \beta_j & \gamma_{ij} \end{pmatrix}, \quad g^{\mu\nu} = \begin{pmatrix} -\alpha^{-2} & \alpha^{-2} \beta^i \\ \alpha^{-2} \beta^i & \gamma^{ij} - \alpha^{-2} \beta^i \beta^j \end{pmatrix}$$

# Black hole electrodynamics

Covariant Maxwell equations:

$$\begin{aligned} \nabla_{\nu} F^{*\mu\nu} &= 0, \\ \nabla_{\nu} F^{\mu\nu} &= I^{\mu} \end{aligned}$$

Faraday tensor
Maxwell tensor

Define **electric** and **magnetic field** as seen by normal (Eulerian) observer in 3+1 split:

$$\begin{aligned} D^{\mu} &= -F^{\mu\nu} n_{\nu}, \\ B^{\mu} &= -F^{*\mu\nu} n_{\nu} \end{aligned}$$

Define **auxiliary fields** (only physical meaning at infinity):

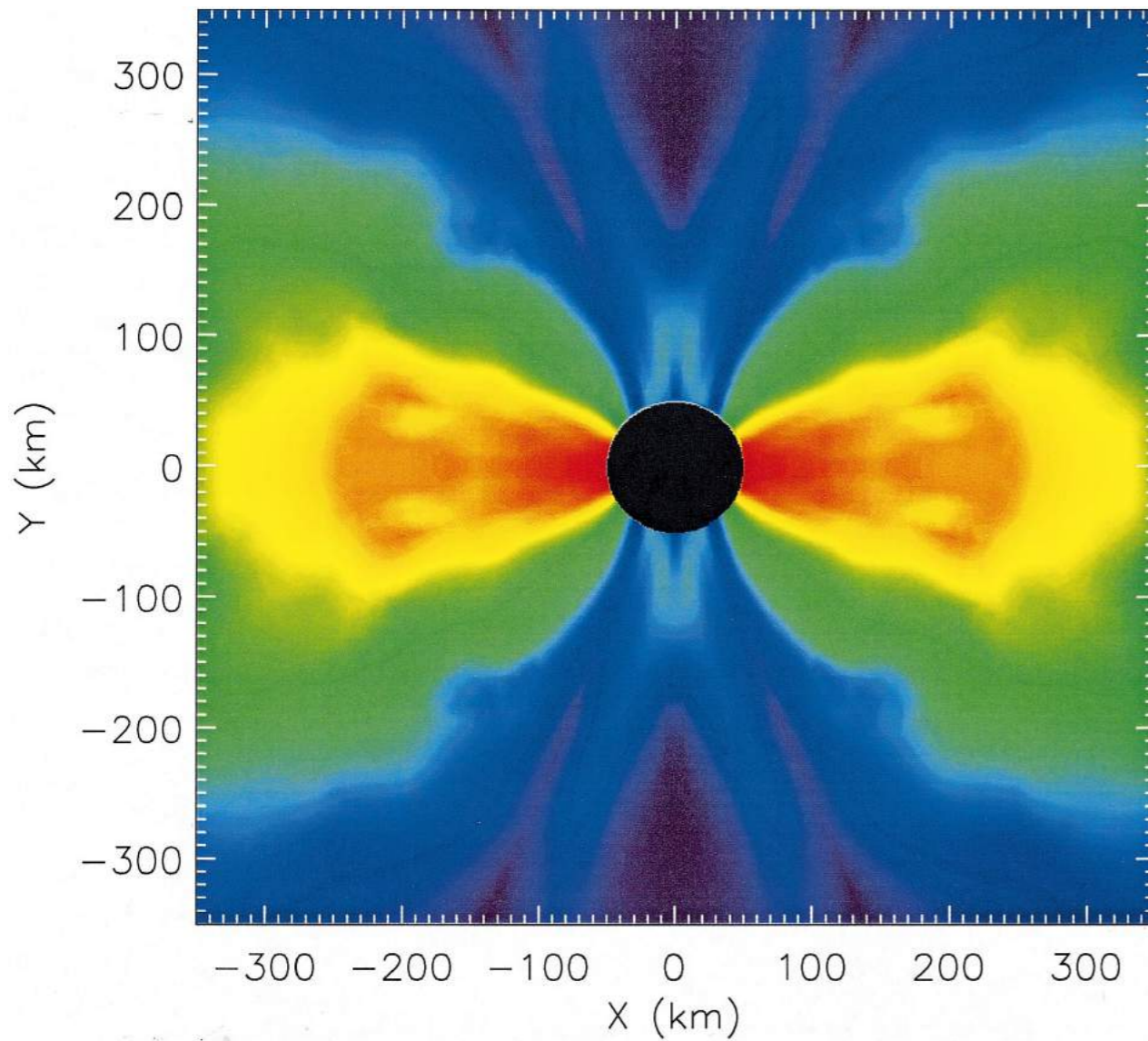
$$\begin{aligned} E^{\mu} &= -\frac{1}{2} \gamma^{\mu\nu} \eta_{\nu\alpha\beta\gamma} k^{\alpha} F^{*\beta\gamma}, \\ H^{\mu} &= -\frac{1}{2} \gamma^{\mu\nu} \eta_{\nu\alpha\beta\gamma} k^{\alpha} F^{\beta\gamma}, \\ \text{electric current} \longrightarrow J^{\mu} &= 2I^{[\nu} k^{\mu]} n_{\nu}, \\ \text{electric charge} \longrightarrow \rho_q &= -I^{\nu} n_{\nu}. \end{aligned}$$

Exercise: all these fields are **purely spatial**:  $X^{\mu} n_{\mu} = 0$

(i.e., they live on the spatial hyper-surfaces of the 3+1 split)

# Activation of the BZ mechanism

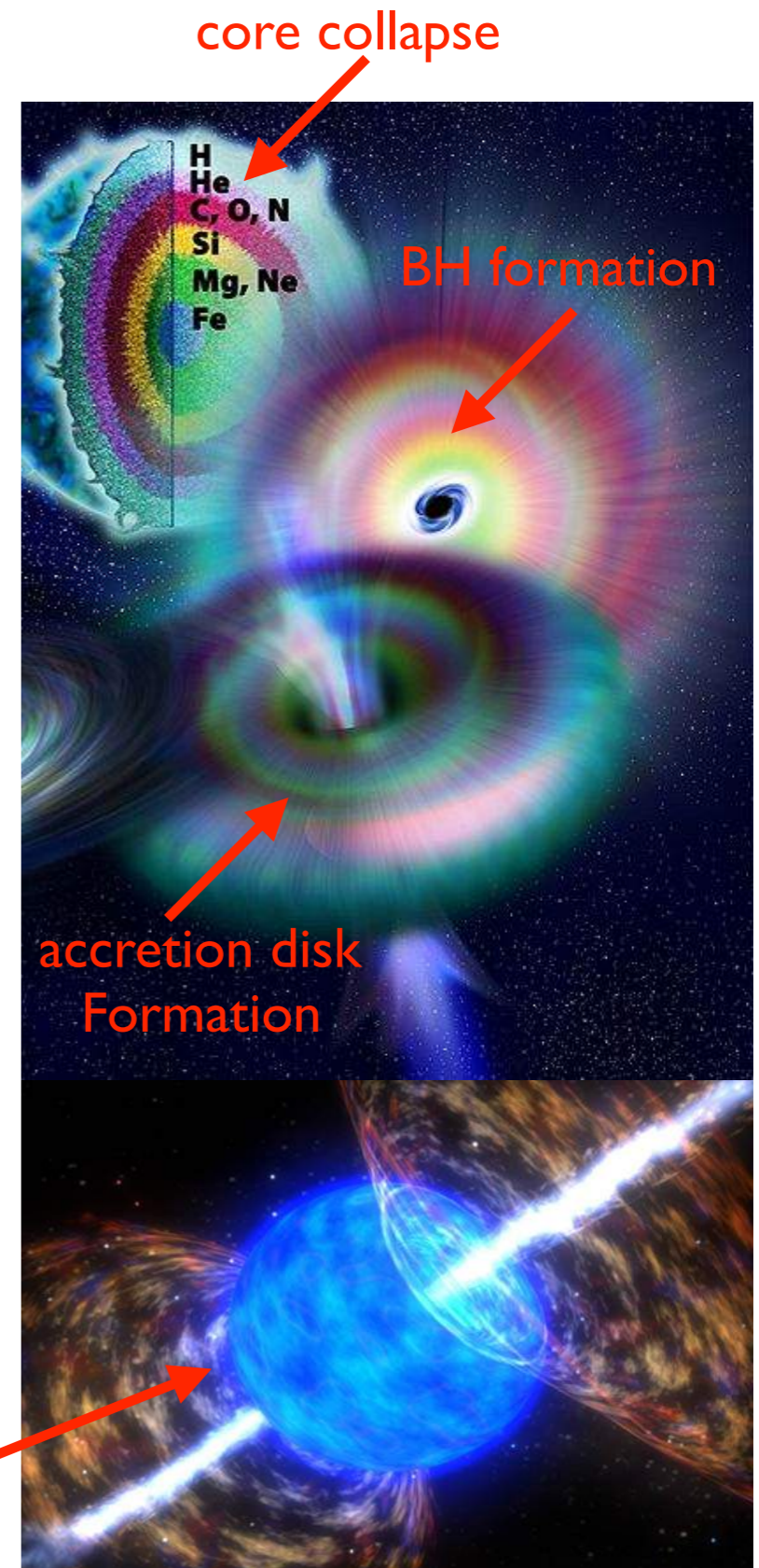
# Long GRBs: collapsars



MacFadyen & Woosley 1999

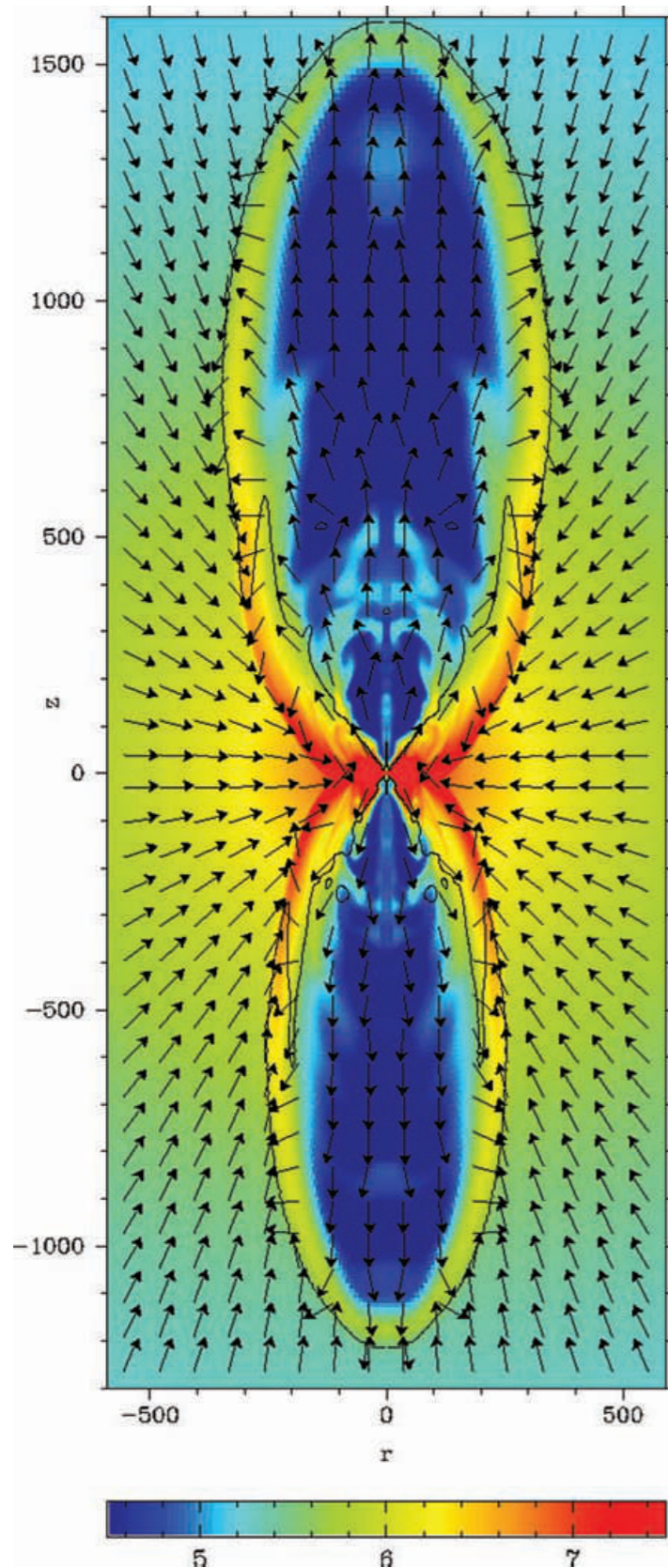
BH-accretion disk from collapse of *rapidly* rotating massive stars ( $M > 20 M_{\text{sun}}$ )

jet punches through infalling material, generates GRB





# Long GRBs: collapsars



- Initial conditions: free-fall material with angular momentum, roughly approximating collapse of massive, rapidly rotating star
- Initial dipole B-field superimposed on infalling material
- Find that accretion onto BH predominantly occurs through accretion disk
- Find activation of BZ mechanism if

$$\beta_\rho = \frac{4\pi\rho c^2}{B^2} < 1$$

develops in polar region near BH horizon

Komissarov & Barkov 2009

# Short GRBs: binary neutron star mergers

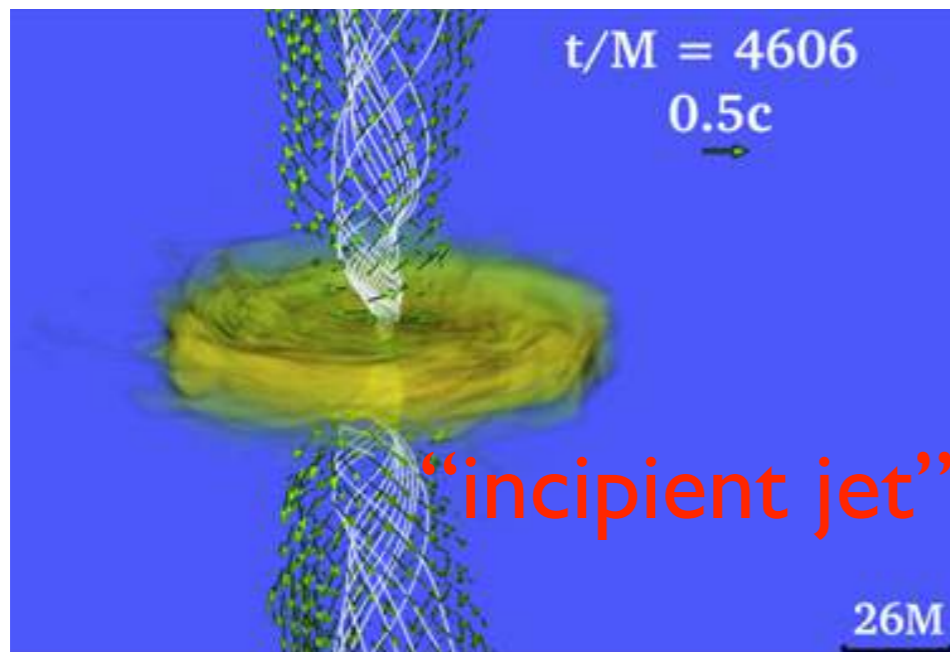


Fig.: Magnetic funnel (“incipient jet”) emerging from a BH-torus system (BNS merger)

Ruiz+ 2016  
Paschalidis et al. 2015

- jet formation in NS mergers not understood yet
- successful case only for prompt collapse

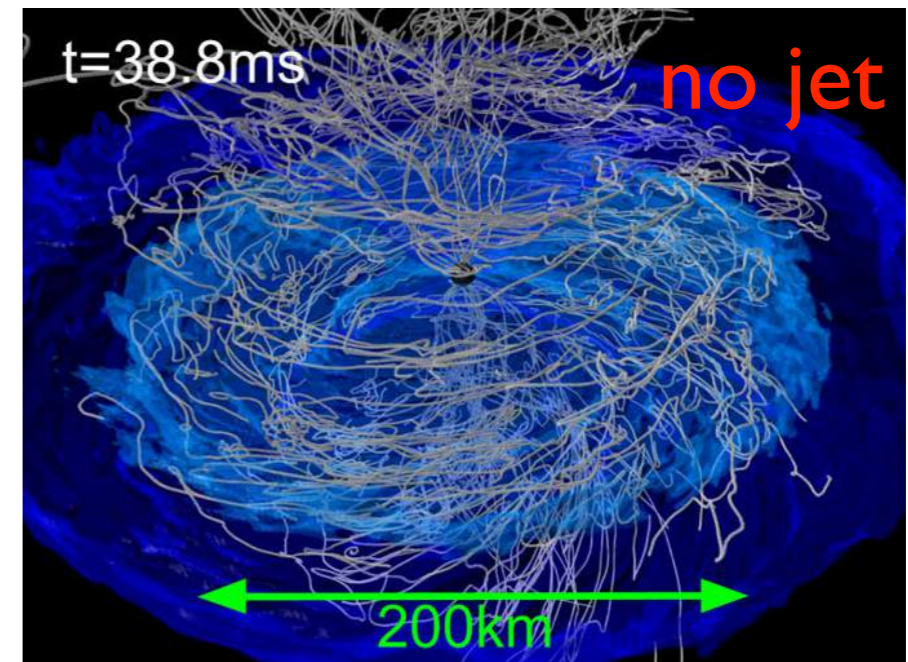


Fig.: Magnetic configuration from the highest resolution BNS simulations: no jet

Kiuchi+ 2017

no jet

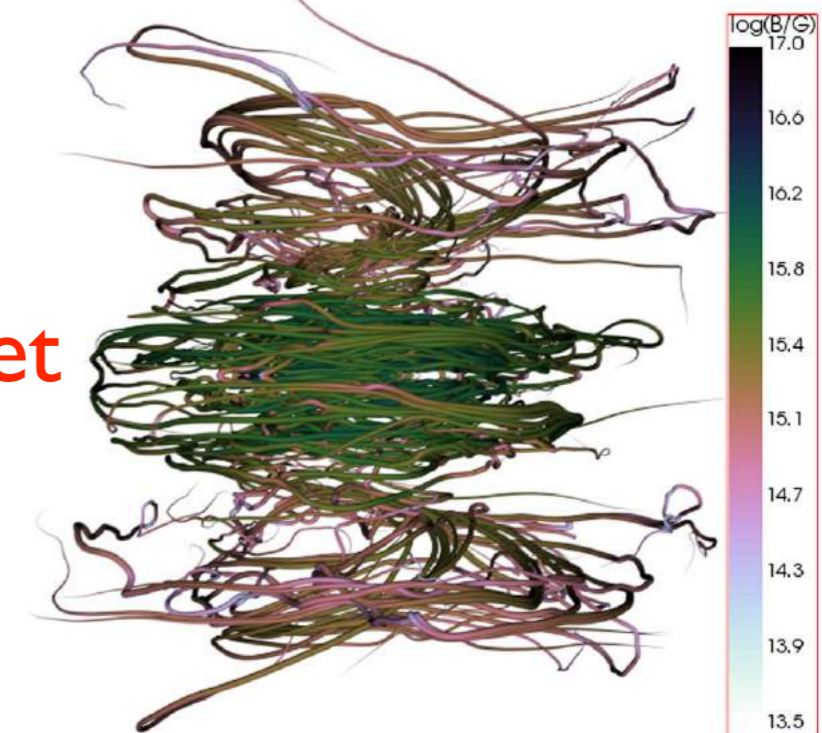
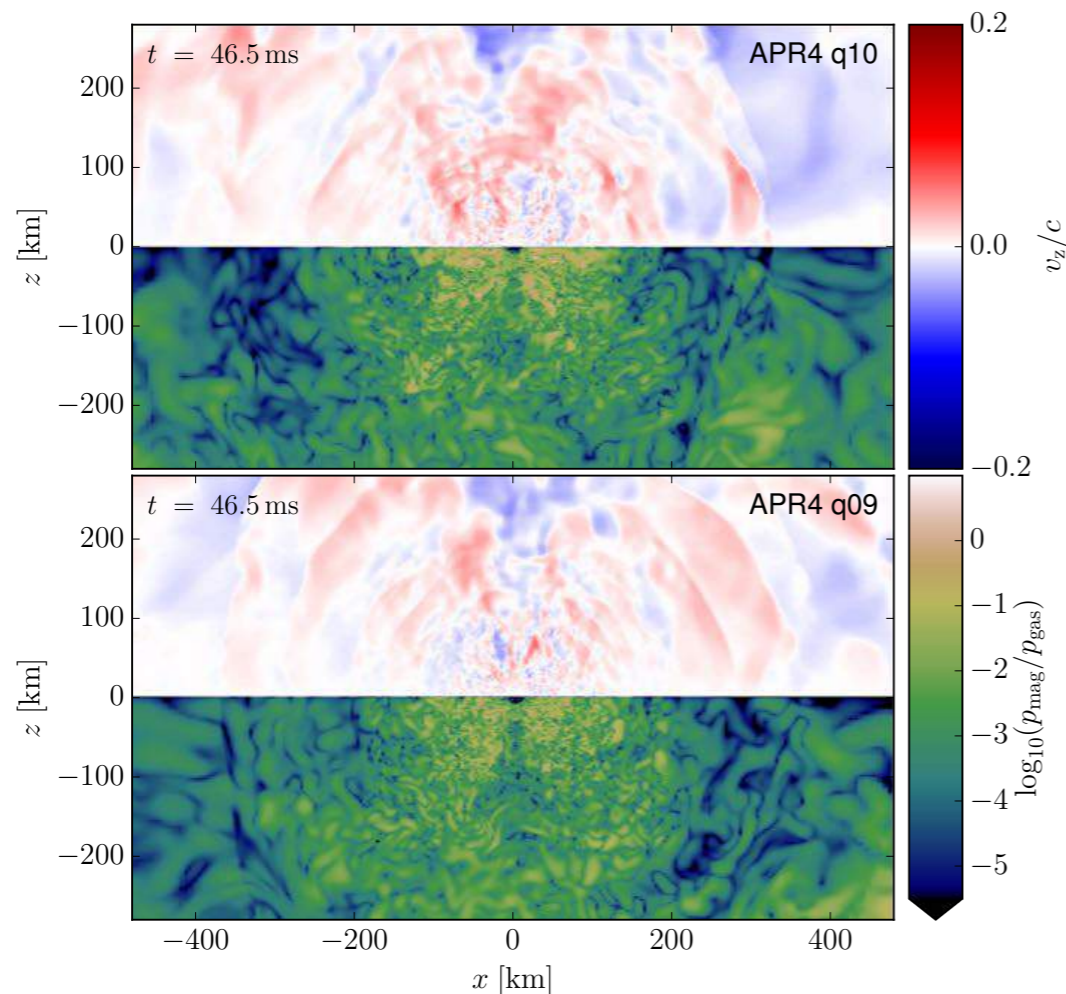


Fig.: Magnetic configuration from the latest BNS simulations: no jet

Cioffi, Siegel+ 2017

# Baryon pollution in BNS mergers

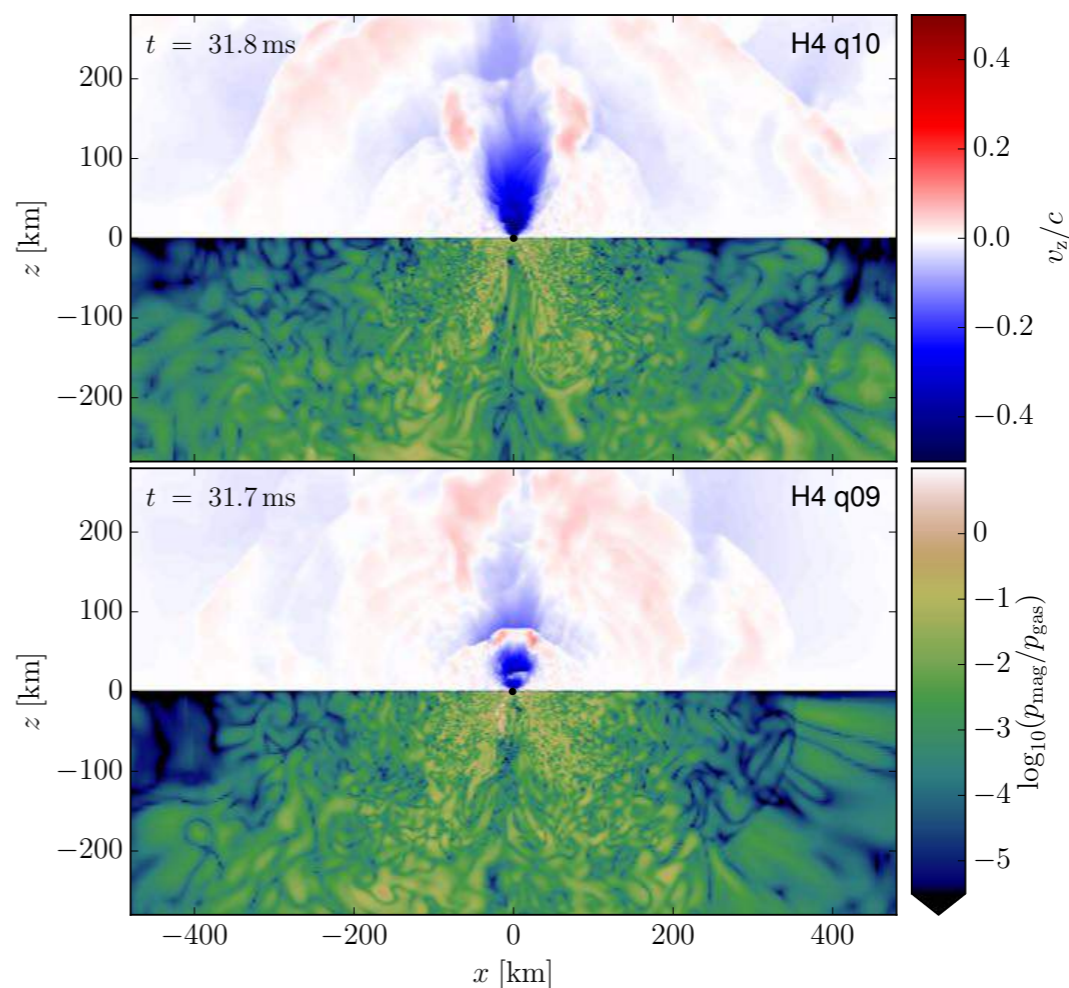


BNS → NS  
isotropic baryon pollution



relativistic outflows likely choked  
or prevented to form

Murguia-Berthier+ 2015,2016  
Nagakura+ 2014  
Just+ 2016



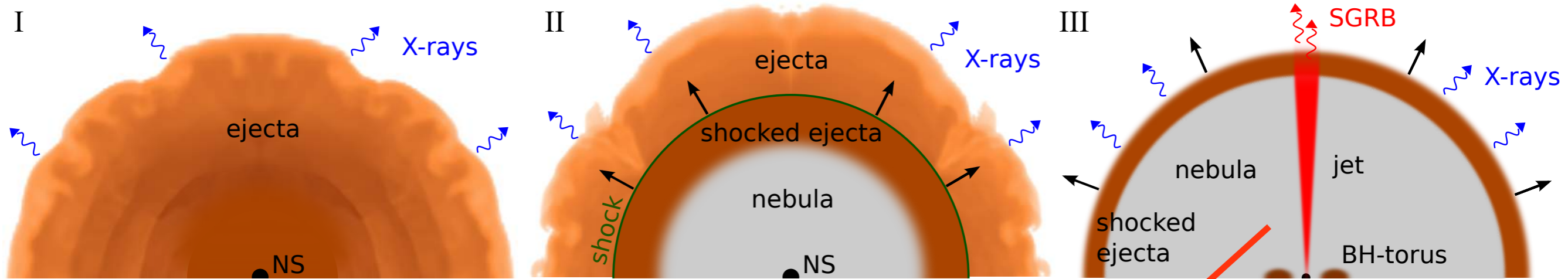
BNS → BH  
low-density polar funnel

Ciolfi+ 2017

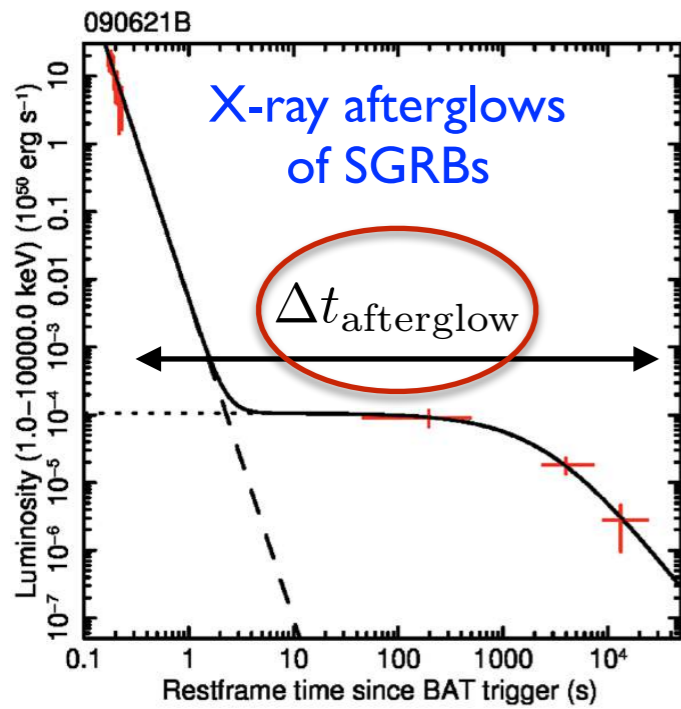
→ how can GRB jets be produced at all in BNS → NS events?

**How to avoid baryon pollution?**

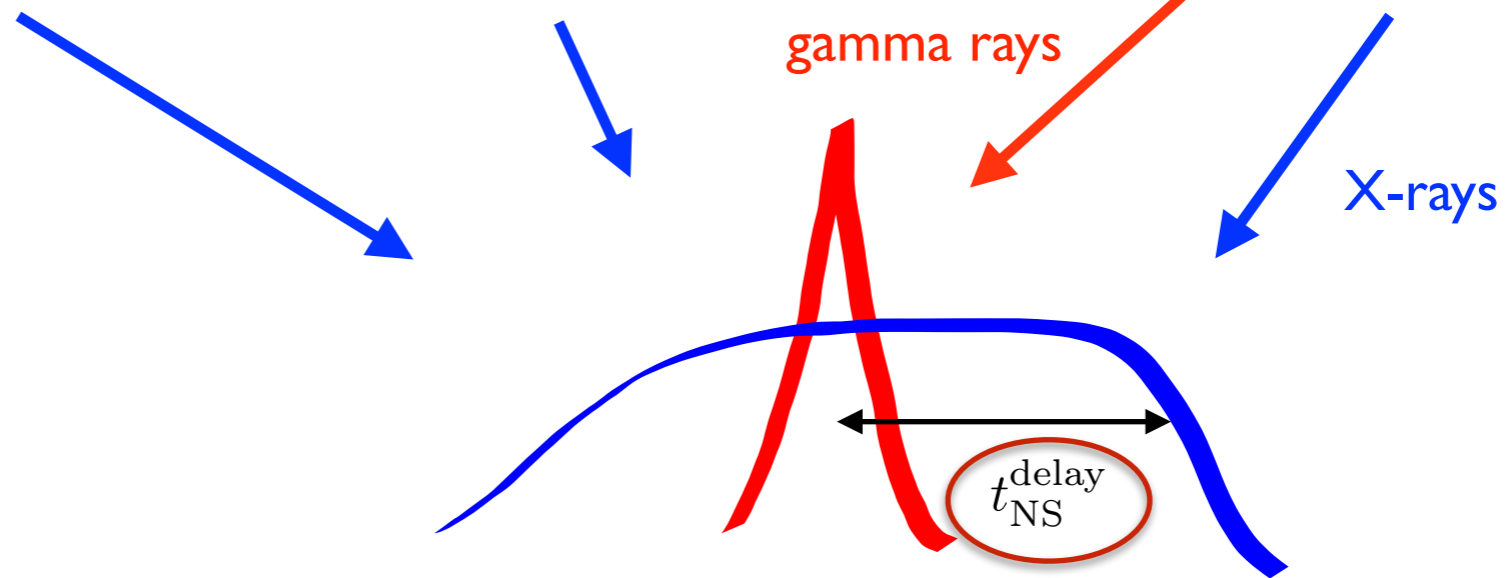
# Time-reversal scenario Cioffi & Siegel 2015a,b



Cioffi & Siegel 2015a,b

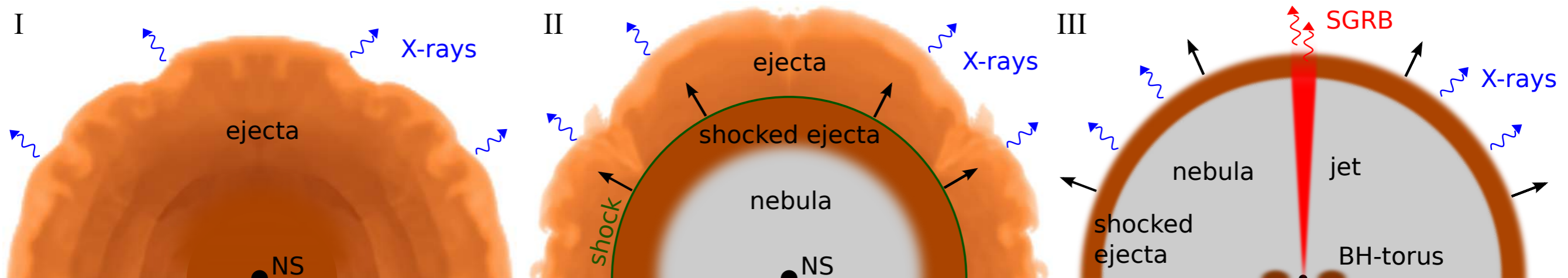


Rowlinson+2013

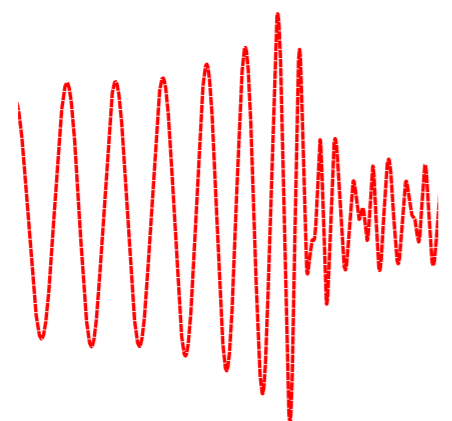


→ SGRB at time of NS collapse  
X-ray precursor and afterglow

# Time-reversal scenario Cioffi & Siegel 2015a,b



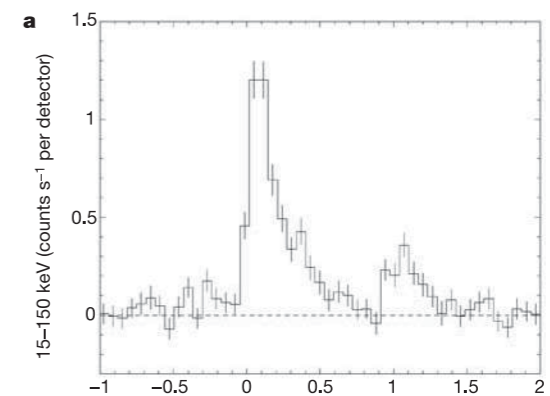
Cioffi & DMS 2015a,b



peak amplitude GW signal

$$t_{\text{coll}} \sim t_{\text{sd}} \sim 10^3 \text{ s}$$

lifetime of the NS



EM SGRB signal

→ GW-EM multimessenger observations ideal to reveal the central engine