Implications of the Higgs discovery

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• Before the 4th of July
• The 4th of July and after
• Implications for the Standard Model
• Implications for the MSSM
• Perspectives
1. Before the 4th of July

A longstanding and most crucial problem in particle physics: how to generate particle masses in an SU(2) × U(1) gauge invariant way? in the Standard Model ⇒ the Higgs–Englert–Brout mechanism

Introduce a doublet of scalar fields \( \Phi = (\Phi^0, \Phi^0) \) with \( \langle 0 | \Phi^0 | 0 \rangle \neq 0 \): fields/interactions symmetric under SU(2) × U(1) but vaccum not.

\[
\mathcal{L}_S = D_\mu \Phi^\dagger D^\mu \Phi - \mu^2 \Phi^\dagger \Phi - \lambda (\Phi^\dagger \Phi)^2
\]

\[
v = \left(-\frac{\mu^2}{\lambda}\right)^{1/2} = 246 \text{ GeV}
\]

⇒ three d.o.f. for \( M_{W^\pm} \) and \( M_Z \).

For fermion masses, use same \( \Phi \): \[
\mathcal{L}_{\text{Yuk}} = -f_e (\bar{e}, \bar{\nu}) L \Phi e_R + ...\]

Residual d.o.f corresponds to spin–0 H particle.

- The scalar Higgs boson: \( J^{PC} = 0^{++} \) quantum numbers (CP–even).
- Masses and self–couplings from \( V \) : \( M_H^2 = 2\lambda v^2, g_{H^3} = 3M_H^2/v, ... \)
- Higgs couplings \( \propto \) particle masses: \( g_{Hff} = m_f/v, g_{HVV} = 2M_V^2/v \)

Since \( v \) is known, the only free parameter in the SM is \( M_H \) (or \( \lambda \)).
1. Before the 4th of July

Once $M_H$ known, all properties of the Higgs are fixed (modulo QCD).

Example: Higgs decays in the SM

- As $g_{HPP} \propto m_P$, H will decay into heaviest particle phase-space allowed:
  - $M_H \lesssim 130 \text{ GeV}$:
    - $H \rightarrow b\bar{b}$: dominant decay
    - $H \rightarrow cc, \tau^+\tau^-, gg = \mathcal{O}(\text{few } \%)$
    - $H \rightarrow \gamma\gamma, Z\gamma = \mathcal{O}(0.1\%)$
  - $M_H \gtrsim 130 \text{ GeV}$:
    - $H \rightarrow WW, ZZ$ dominant
    - decays into $t\bar{t}$ for heavy Higgs

- Total Higgs decay width:
  - very small for a light Higgs
  - comparable to mass if heavy

HDECAY ⇒
1. Before the 4th of July

Higgs production rates also fixed (modulo QCD):

Large production cross sections
with $gg \rightarrow H$ by far dominant process
$1 \text{ fb}^{-1} \Rightarrow O(10^4)$ events@LHC
$\Rightarrow O(10^3)$ events@Tevatron
but eg $\text{BR}(H \rightarrow \gamma\gamma, ZZ \rightarrow 4\ell) \approx 10^{-3}$
... a small # of events at the end...
with a huge QCD-jet background.
see G. Salam talk
$\Rightarrow$ an extremely challenging task!

Main sensitive channels:
$gg \rightarrow H \rightarrow \gamma\gamma$
$gg \rightarrow H \rightarrow ZZ \rightarrow 4\ell, 2\ell2\nu, 2\ell2b$
$gg \rightarrow H \rightarrow WW \rightarrow \ell\nu\ell\nu + 0, 1j$
also help from other channels:
– VBF+$gg \rightarrow H \rightarrow \tau\tau$
– $q\bar{q} \rightarrow HV \rightarrow b\bar{b}\ell X$
1. Before the 4th of July

But a major problem in the SM: the hierarchy/naturalness problem

Radiative corrections to $M_H^2$ in SM with a cut–off $\Lambda = M_{NP} \sim M_{Pl}$

$$\Delta M_H^2 \equiv \frac{H}{f} \frac{H}{f} \propto \Lambda^2 \approx (10^{18} \, \text{GeV})^2$$

$M_H$ prefers to be close to the high scale than to the EWSB scale...

Three main avenues for solving the problem: see Diaz-Cruz talk

**Supersymmetry:** a set of new/light SUSY particles cancel the divergence.
- MSSM $\equiv$ two Higgs doublet model $\Rightarrow$ 5 physical states $h, H, A, H^\pm$
- very predictive: only two free parameters at tree–level ($\tan \beta, M_A$)
- upper bound on light Higgs $M_h \lesssim 130 \, \text{GeV}$ and $M_{H,H^\pm} \approx M_A \lesssim \text{TeV}$

**Extra dimensions:** there is a cut–off at TeV scale where gravity sets in.
- in most cases: SM–like Higgs sector but properties possibly affected
- but also: scenarios with Higgs–gauge unification and Higgsless models....

**Strong interactions/compositness:** the Higgs is not an elementary scalar.
- $H$ is a bound state of fermions like for the pions in QCD...
- $H$ emerges as a Nambu–Goldstone of a strongly interacting sector.
1. Before the 4th of July

and along the avenues, many possible streets, paths, corners...

Which scenario chosen by Nature? The LHC was supposed to tell!
2. The 4th of July and after

After 48 years of postulat, 30 years of search (and a few heart attacks),
the Higgs is discovered at LHC on the 4th of July: Hi(gg)storical day!

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**ATLAS 2011 - 2012**

- Obs.
- Exp.

![Graph](Image)

**CMS Preliminary**

- Observed
- Expected (68%)
- Expected (95%)

![Graph](Image)

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The particle decays into $\gamma\gamma$ states
- not spin–1: Landau–Yang
- could be spin–2 like graviton?
  – miracle that rates/distributions fit that of a scalar Higgs boson,
  $\Rightarrow$ “prima facie” evidence against it.
  Many theoretical analyses...

Is it a CP–even state or CP–odd?

$HV^\mu V^\mu$ versus $H\epsilon^{\mu\nu\rho\sigma}Z_{\mu\nu}Z_{\rho\sigma}$

$\Rightarrow \frac{d\Gamma(H\rightarrow ZZ^*)}{dM_*}$ and $\frac{d\Gamma(H\rightarrow ZZ)}{d\phi}$

CMS: 2.5$\sigma$ for CP–even...

Problem: if $H$ is CP mixture, only $0^+$ component is projected out!
(or very large $0^-$VV loop coupling).
$\Rightarrow$ better probe: $\hat{\mu}_{ZZ} = 0.95 \pm 0.3$?
2. The 4th of July and after

Rates compatible with those expected in the SM

See ATLAS/CMS talks this morning...

Higgs couplings to gauge bosons and fermions as dictated by unitarity:
- fermiophobic, gauge-phobic scenarios ruled out
- still two solutions for fermion cplgs: non–SM–like is non unitary...

SM particle spectrun now complete: no 4th generation fermions
- Rates in $ZZ, WW, \gamma\gamma, b\bar{b}$ incompatible with SM4
- direct searches and precision data against it

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3. Implications for the SM

So it's looks like expected in SM ⇒ a triumph for high-energy physics!

Indirect constraints from EW data

H contributes to RC to W/Z masses:

\[ H \propto \frac{\alpha}{\pi} \log \frac{M_H}{M_W} + \cdots \]

Fit the EW precision measurements, one obtains

\[ M_H = 92^{+34}_{-26} \text{ GeV}, \]

or

\[ M_H \lesssim 160 \text{ GeV at 95\% CL} \]

compared with the measured mass

\[ M_H \approx 126 \text{ GeV}. \]

A very non–trivial consistency check!

(remember the stop of the top quark!).

The SM is a very successfull theory!

\[ ^a \text{Still some problems with } A^0_{FB} \text{ (LEP), } A^t_{FB} \text{ (TeV) and } g - 2 \text{ but not severe...} \]

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3. Implications in the SM

- The theory preserves unitarity:
  without $H$: $|A_0(\text{VV} \to \text{VV})| \propto E^2$
  including $H$: $|A_0| \propto M_H^2 / v^2$

- Extrapolable up to highest scales.

Stability of the EW vacuum?

- $\lambda = M_H^2 / 2v^2$ evolves with $Q$: $\frac{\lambda(Q^2)}{\lambda(v^2)} \approx 1 + 3 \cdot \frac{2M_W^4 + M_Z^4 - 4m_t^4}{16\pi^2 v^4} \log \frac{Q^2}{v^2}$
  tops make $\lambda(0) < \lambda(v)$: unstable vacuum

- SM valid only if $v \equiv \text{EW-min}$, ie $\lambda(Q^2) > 0$
  $\Lambda_C \sim M_{\text{Planck}} \Rightarrow M_H \gtrsim 129$ GeV!
  for $m_t = 173$ GeV; but what is $m_t^{\text{TEV}}$ ??

- Unambiguous $m_t$ only from $\sigma(\bar{t}t)$:
  but value at TEV/LHC not precise...

- Standardissimo=TOE? Maybe not (?):
  $m_\nu$, DM, GUT, hierarchy pb...

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**4. Implications for BSM**

Rates compatible with those expected in the SM

Some beyond the SM scenarios are ruled out:
- Higgsless models, extreme Technicolor and composite scenarios, ..
- fermiophobic Higgs, gauge-phobic Higgs, 4th generation, ...

Some beyond the SM scenarios are in “hospital” (no names..)
Other BSM scenarios are strongly constrained...

Here, I discuss the example of Supersymmetry and the MSSM:

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4. Implications for the MSSM

In MSSM with two Higgs doublets: \( H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix} \) and \( H_2 = \begin{pmatrix} H_2^0 \\ H_2^+ \end{pmatrix} \),

- to cancel the chiral anomalies introduced by the new \( \tilde{h} \) field,
- give separately masses to d and u fermions in SUSY invariant way.

After EWSB (which can be made radiative: more elegant than in SM):
three dof to make \( W_L^\pm, Z_L \Rightarrow 5 \) physical states left out: \( h, H, A, H^\pm \)

Only two free parameters at the tree level: \( \tan \beta, M_A \); others are:

\[
M_{h,H}^2 = \frac{1}{2} \left[ M_A^2 + M_Z^2 \mp \sqrt{(M_A^2 + M_Z^2)^2 - 4M_A^2M_Z^2\cos^22\beta} \right]
\]

\[
M_{H^\pm}^2 = M_A^2 + M_W^2
\]

\[
\tan2\alpha = \tan2\beta (M_A^2 + M_Z^2)/(M_A^2 - M_Z^2)
\]

We have important constraint on the MSSM Higgs boson masses:

\[
M_h \leq \min(M_A, M_Z) \cdot |\cos2\beta| \leq M_Z, \ M_{H^\pm} > M_W, M_H > M_A \ldots
\]

\[M_A \gg M_Z: \text{decoupling regime, all Higgses heavy except for } h:\]

\[
M_h \sim M_Z|\cos2\beta| \leq M_Z!, \ M_H \sim M_{H^\pm} \sim M_A , \ \alpha \sim \frac{\pi}{2} - \beta
\]

\( \Rightarrow \) Inclusion of radiative corrections to \( M_h \) important and necessary.
4. Implications for MSSM: mass

The mass value 126 GeV is rather large for the MSSM h boson, ⇒ one needs from the very beginning to almost maximize it...

Maximizing $M_h$ is maximizing the radiative corrections; at 1-loop:

$$M_h \sim M_A \gg M_Z \rightarrow M_Z |\cos 2\beta| + \frac{3\tilde{m}_t^4}{2\pi^2 v^2 \sin^2 \beta} \left[ \log \frac{M_S^2}{m_t^2} + \frac{X_t^2}{M_S^2} \left( 1 - \frac{X_t^2}{12M_S^2} \right) \right]$$

- decoupling regime with $M_A \sim O(\text{TeV})$;
- large values of $\tan \beta \gtrsim 10$ to maximize tree-level value;
- maximal mixing scenario: $X_t = \sqrt{6}M_S$;
- heavy stops, i.e. large $M_S = \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}$;

we choose at maximum $M_S \lesssim 3$ TeV, not to have too much fine-tuning....

- Do the complete job: two-loop corrections and full SUSY spectrum
- Use RGE codes (Suspect) with RC in $\overline{\text{DR}}$/compare with FeynHiggs (OS).

Perform a full scan of the phenomenological MSSM with 22 free parameters
- determine the regions of parameter space where $123 \leq M_h \leq 129$ GeV (3 GeV uncertainty includes both “experimental” and “theoretical” error)
- require h to be SM–like: $\sigma(h) \times \text{BR}(h) \approx H_{SM}$ ($H = H_{SM}$ later)

Many analyses! Here, the one from Arbey et al. 1112.3028+1207.1348
4. Implications for pMSSM: mass

Main results:
- **Large** $M_S$ values needed:
  - $M_S \approx 1\,\text{TeV}$: only maximal mixing
  - $M_S \approx 3\,\text{TeV}$: only typical mixing.
- **Large** $\tan \beta$ values favored
  but $\tan \beta \approx 3$ possible if $M_S \approx 3\,\text{TeV}$

**How light sparticles can be with the constraint** $M_h = 126\,\text{GeV}$?
- $1s/2s$ gen. $\tilde q$ should be heavy...
  But not main player here: the stops:
  \[ m_{\tilde t_1} \lesssim 500\,\text{GeV} \] still possible!

- $M_1$, $M_2$ and $\mu$ unconstrained,
- non-univ. $m_{\tilde f}$: decouple $\tilde \ell$ from $\tilde q$

EW sparticles can be still very light but watch out the new limits.
4. Implications for the cMSSM

Constrained MSSMs are interesting from model building point of view:

- concrete schemes: SSB occurs in hidden sector gravity → MSSM fields
- provide solutions to some MSSM problems: CP, flavor, etc..
- parameters obey boundary conditions ⇒ small number of inputs...

• mSUGRA: $\tan \beta$, $m_{1/2}$, $m_0$, $A_0$, $\text{sign}(\mu)$
• GMSB: $\tan \beta$, $\text{sign}(\mu)$, $M_{\text{mes}}$, $\Lambda_{SSB}$, $N_{\text{mess}}$ fields
• AMSB: $m_0$, $m_3/2$, $\tan \beta$, $\text{sign}(\mu)$

Full scans of the model parameters with $123 \text{ GeV} \leq M_h \leq 129 \text{ GeV}$

Very strong constraints and some (minimal) models ruled out...

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4. Implications for MSSM: other searches

There are other (stringent) constraints on pMSSM to be included:

- production/decay rates of the observed Higgs particle;
- the observation of heavier Higgses in the ZZ,WW signal channels;
- CMS and ATLAS $pp \rightarrow A/H/(h) \rightarrow \tau\tau$ and $t \rightarrow bH^+$ searches;
- constraints from sparticle searches and eventually Dark Matter,
- constraints from flavor: at least (direct!) limits from $B_s \rightarrow \mu\mu$...

### CMS Preliminary

- $\sqrt{s} = 7$ TeV, $L = 5.1 \text{ fb}^{-1}$ $\sqrt{s} = 8$ TeV, $L = 12.2 \text{ fb}^{-1}$
- $\sqrt{s} = 8$ TeV, $L = 5.1 \text{ fb}^{-1}$
- $\sqrt{s} = 7$ TeV, $L = 4.8 \text{ fb}^{-1}$
- $\sqrt{s} = 8$ TeV, $L = 5.9 \text{ fb}^{-1}$
- Combination

### ATLAS Preliminary

- $\sqrt{s} = 7$ TeV, $L = 4.8 \text{ fb}^{-1}$ $\sqrt{s} = 8$ TeV, $L = 13 \text{ fb}^{-1}$
- $\sqrt{s} = 7$ TeV, $L = 4.6 \text{ fb}^{-1}$
- $\sqrt{s} = 8$ TeV, $L = 13 \text{ fb}^{-1}$
- Combination

$\tan\beta = 1.3 \pm 0.3$ for $\sqrt{s} = 8$ TeV, $L = 13 \text{ fb}^{-1}$

### Observed

- 95% CL limit on $\sigma/\sigma_{SM}$
- Combined $\sigma/\sigma_{SM}$
- $\sigma/\sigma_{SM}$ from $WZ$ and $Z^0$ decay

### Expected

- 95% CL limit on $\sigma/\sigma_{SM}$
- Expected $\sigma/\sigma_{SM}$
- Expected $\sigma/\sigma_{SM}$ (68%)
- Expected $\sigma/\sigma_{SM}$ (95%)

### 95% CL Excluded Regions

- Observed
- Expected
- $z\rightarrow \mu\mu$ expected
- $z\rightarrow \tau\tau$ expected
- LEP

### MSSM $m_h^{\max}$ scenario

- $M_{\text{SUSY}} = 1 \text{ TeV}$

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4. Implications for MSSM: other searches

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- CMS and ATLAS $pp \rightarrow A/H/(h) \rightarrow \tau\tau$ and $t \rightarrow bH^+$ searches;
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- constraints from flavor: at least (direct!) limits from $B_s \rightarrow \mu\mu$. 

![Graph showing constraints on M(A) vs. tanβ with various regions and lines indicating different limits and searches.](image)
4. Implications for MSSM: rates

Sets stringent constraints on pMSSM regimes/benchmark scenarios?

- Heavier \( \tilde{H} \) being the observed Higgs is now excluded.
- Close \( h, H, A, H^\pm \) (intense coupling regime) excluded.
- Small \( \alpha_{\text{eff}} \) scenario with \( g_{hbb} \approx 0 \) and thus small \( \Gamma_h \):
  - ruled out by LHC/Tevatron data: ex: loose \( Wh \rightarrow \ell \nu b \bar{b} \) signal.
- gluophobic \( h \) with \( g_{hgg} \ll g_{H_{SM}gg} \) due to squark loops?
  - ruled out by \( ZZ, WW, \gamma\gamma \) signals at LHC (and also the \( h \) mass)

But some difference with the SM!
- \( a > 2\sigma \) excess in \( H \rightarrow \gamma\gamma \).
- Statistical fluctuation?
- Systematics problem?
- Maybe QCD uncertainties?

Baglio+Godbole+AD \( \Rightarrow \)

Hope it is due to SUSY!
- total Higgs width suppressed?
- SUSY effects in \( h\gamma\gamma \) loop?
4. Implications for pMSSM: rates

Pretty hard to change tree-level Higgs couplings and loop hgg vertex

Can SUSY contributions significantly enhance the $h \rightarrow \gamma\gamma$ rate?

- light stau’s and large $\mu \tan \beta$
  Carena+Gori+Shah+Wagner
- light $\tilde{\chi}_1^\pm$ in non-univ MSSM
  Driesen+Illana+Hollik+AD
- possibility of light $\tilde{t}$:
  $\Rightarrow$ max-mixing: $\sigma(gg \rightarrow h)$ suppressed.
  $\Rightarrow$ no mixing: yes, but stops too heavy.
  Arvanitaki+Villadoro, AD

- BMSSM? One example is the NMSSM:
  Ellwanger etal, King etal., Gunion etal,..
  - stops lighter as $M_{h_{\text{max}}}$ larger,
  - additional singlet for couplings,
  - less severe non-H constraints.

  Common features: some light sparticles are around the corner!
  
  Data also OK with non SUSY BSM; ex: 2HDM, triplets, new fermions,..
5. Conclusions: MSSM

A 126 GeV Higgs provides information on BSM and SUSY in particular:

- $M_H = 119$ GeV would have been a boring value: everybody OK..
- $M_H = 145$ GeV would be a devastating value: mass extinction..
- $M_H \approx 126$ GeV is Darwinian: (natural) selection among models..

SUSY spectrum heavy; except maybe for weakly interacting sparticles and also stops $\Rightarrow$ more focus on them in SUSY searches!

One has to include other Higgs/SUSY searches in particular:

- $H/A/H^\pm$ searches at the LHC are becoming very constraining..
- SUSY searches and flavor constraints are to be taken into account.
- Little room for other Higgs searches at the LHC.
- Need to start thinking bout changing the benchmark scenarios....

But let’s not get desperate yet! (maybe just slightly nervous..)
- light stops and charginos/neutralinos/sleptons still possible...
- watch out the full 2012 data for some possible (weak) signal....
- still have to look to Higgs coupling mesurements such as $H\gamma\gamma$...

7–8 TeV LHC for the Higgs and 13–14 TeV LHC for SUSY?
5. Conclusions: SM

Now that Higgs is found (and nothing else yet): is Particle Physics “closed”? No! Need to check that H is indeed responsible of sEWSB (and SM-like?)

Measure its fundamental properties in the most precise way:

- its mass and total decay width (invisible width due to dark matter?),
- its spin–parity quantum numbers and check SM prediction for them,
- its couplings to fermions and gauge bosons and check that they are indeed proportional to the particle masses (fundamental prediction!),
- its self–couplings to reconstruct the potential $V_H$ that makes EWSB.

Possible for $M_H \approx 126$ GeV as all production/decay channels useful!
5. Conclusion

Now, this is not the end.
It is not even the beginning to the end.
But it is, perhaps, the end of the beginning.
Sir Winston Churchill, November 1942

We hope that at the end we finally understand the EWSB mechanism, but there is a long way until then.... and there might be many surprises!
Backup: implications for other scenarios

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Backup: implications for MSSM: is H observed?

It looks like in decoupling regime. True 😃

- are small values of $M_A$ allowed?
- can H be the SM-like Higgs boson? YES!, if no other constraints than:
  - $M_H \approx 126 \pm 3$ GeV
  - $g_{HVV} \approx g_{H_{SM}VV}$

Heinemeyer+Stal+Weiglein

$M_A \approx 100$ GeV, $\tan \beta \approx 6 - 10$, $M_S \approx \mu \approx 1$ TeV, $X_t \approx \sqrt{6}M_S$,

$\Rightarrow M_H \approx 126$ GeV ; $M_h \approx 98$ GeV!

[ABDM scan: only few points, $10^{-6}$ OK but they are all ruled out by flavor data

$\Rightarrow$ only h SM–like is likely...

With new CMS update, $\tan \beta \lesssim 5$:

$\Rightarrow H \equiv$ observed is now excluded...