

# In search of general relativity

or: How I learned to stop worrying and understood that not everything is what it seems to be

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# Context and ADM decomposition

## Context: quantum gravity (QG)

- A standalone theory of gravity in  $3 + 1$  dimensions, valid at all scales.
- Particular candidate: **Hořava-Lifshitz gravity**<sup>a</sup>:
  - Built from anisotropic UV fixed point, defined s.t.  $[g_N] = 1$ ,
  - EFT point of view is then applicable,
  - “IR limit” is different from Einstein-Hilbert action.  
 $\implies$  searching for GR in this limit.

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## Ingredients: ADM decomposition

- 3+1 decomposition of the metric,
 
$$g_{ij} \equiv {}^{(4)}g_{ij}, \quad N \equiv (-{}^{(4)}g^{00})^{-1/2}, \quad N_i \equiv {}^{(4)}g_{0i}.$$
- All time derivatives encoded in the extrinsic curvature,
 
$$K_{ij} = \frac{1}{2N} (\dot{g}_{ij} - \nabla_i N_j - \nabla_j N_i).$$

# GR revisited

## Action + Wheeler-DeWitt metric

- The 3 + 1 version of the Einstein-Hilbert action,

$$\begin{aligned} S &= \int dt \int d^3x \sqrt{g} N (K_{ij} K^{ij} - K^2 + R - 2\Lambda) \\ &= \int dt \int d^3x \sqrt{g} N (K_{ij} G^{ijkl} K_{kl} + R - 2\Lambda), \end{aligned}$$

- $G^{ijkl}$  is the Wheeler-DeWitt metric,

$$G^{ijkl} = \frac{1}{2} (g^{ik} g^{jl} + g^{il} g^{jk}) - g^{ij} g^{jk},$$

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## Total Hamiltonian and constraints

- Performing the Legendre transformation yields the Hamiltonian,

$$H = \int d^3x (N\mathcal{H} + N^i \mathcal{H}_i + \alpha\phi + \alpha^i \phi_i)$$

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- $\mathcal{H}$  and  $\mathcal{H}_i$  denote the Hamiltonian and momentum constraints,

$$\mathcal{H} = \frac{\pi^{ij} G_{ijkl} \pi^{kl}}{\sqrt{g}} - \sqrt{g} (R - 2\Lambda).$$

- Imposing  $(\dot{\phi}, \dot{\phi}_i) = (0, 0)$  yields  $(\mathcal{H}, \mathcal{H}_i) \approx (0, 0)$  and no new constraints arise.

# The $\lambda$ -R model

## Action + generalized Wheeler-DeWitt metric

- Breaking  $\text{Diff}(\mathcal{M})$ , a new dimensionless coupling appears,

$$\begin{aligned} S &= \int dt \int d^3x \sqrt{g} N (K_{ij} K^{ij} - \lambda K^2 + R - 2\Lambda) \\ &= \int dt \int d^3x \sqrt{g} N \left( K_{ij} G_{\lambda}^{ijkl} K_{kl} + R - 2\Lambda \right), \end{aligned}$$

- $G_{\lambda}^{ijkl}$  is the generalized Wheeler-DeWitt metric,

$$G_{\lambda}^{ijkl} = \frac{1}{2} (g^{ik} g^{jl} + g^{il} g^{jk}) - \lambda g^{ij} g^{jk},$$

## Total Hamiltonian and constraints

- Performing the Legendre transformation yields the Hamiltonian,

$$H = \int d^3x (N\mathcal{H}_{\lambda} + N^i \mathcal{H}_i + \alpha\phi + \alpha^i \phi_i)$$

- $\mathcal{H}_i$  and its algebra remain unchanged, with  $\lambda$  present in  $\mathcal{H}_{\lambda}$ ,

$$\mathcal{H}_{\lambda} = \frac{\pi^{ij} G_{ijkl}^{\lambda} \pi^{kl}}{\sqrt{g}} - \sqrt{g} (R - 2\Lambda).$$

- While  $(\mathcal{H}_{\lambda}, \mathcal{H}_i) \approx (0, 0)$ ,  $\dot{\mathcal{H}}_{\lambda} \approx 0$  is not trivial.

# The tertiary constraint

Tertiary constraint  $\sim$  gauge fixing

- General solution:  $\nabla_j \pi \approx 0$ .
- Asymptotically flat spaces:  $\pi = 0^a$ 
  - GR in the maximal slicing gauge ( $\pi = 0$ ) is recovered.
- For compact spaces:  $\frac{\pi}{\sqrt{g}} = a(t)$  (CMC gauge condition),
  - at first glance, the  $\lambda$ -dependence remains.

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## More on the compact case

- Time preservation of tertiary constraint yield two more equations:
  - $\lambda$ -dependent  $N$  (denoted by  $\mathcal{A} \approx 0$ ) and  $\alpha$  fixing equations.
- 1st and 2nd class classification of constraints yields 2 d.o.f.,
  - 2nd class constraints:  $(\mathcal{A}, \phi, \mathcal{H}_\lambda, \pi - a\sqrt{g})$
  - 1st class constraints:  $(\phi_i, \mathcal{H}_i)$ .

# Outlook and references

- Originally, we worked with  $\dot{a} = 0$ , new results soon with fully general  $a(t) = \frac{1}{V} \int d^3x \pi$ .
  - Using York's conformal methods, it seems to be possible to prove equivalence with CMC general relativity for  $\lambda > 1/3$ ,
  - Proof of Solution of Lichnerowicz-York does not seem to generalize for  $\lambda < 1/3$ .
- Would also be interesting to:
  - Check spaces with different boundary conditions,
  - Include (possibly as a perturbation) the term  $\nabla_i \log N \nabla^i \log N$ .

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## Bibliography:

- P. Hořava: arXiv:0901.3775v2 [hep-th];
- D. Giulini, and C. Kiefer: <http://arxiv.org/abs/gr-qc/9405040>;
- J. Bellorin, and A. Restuccia: <http://arxiv.org/abs/1004.0055>;
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Thank you!

Don't make a sound: they're not dead, just sleeping.

