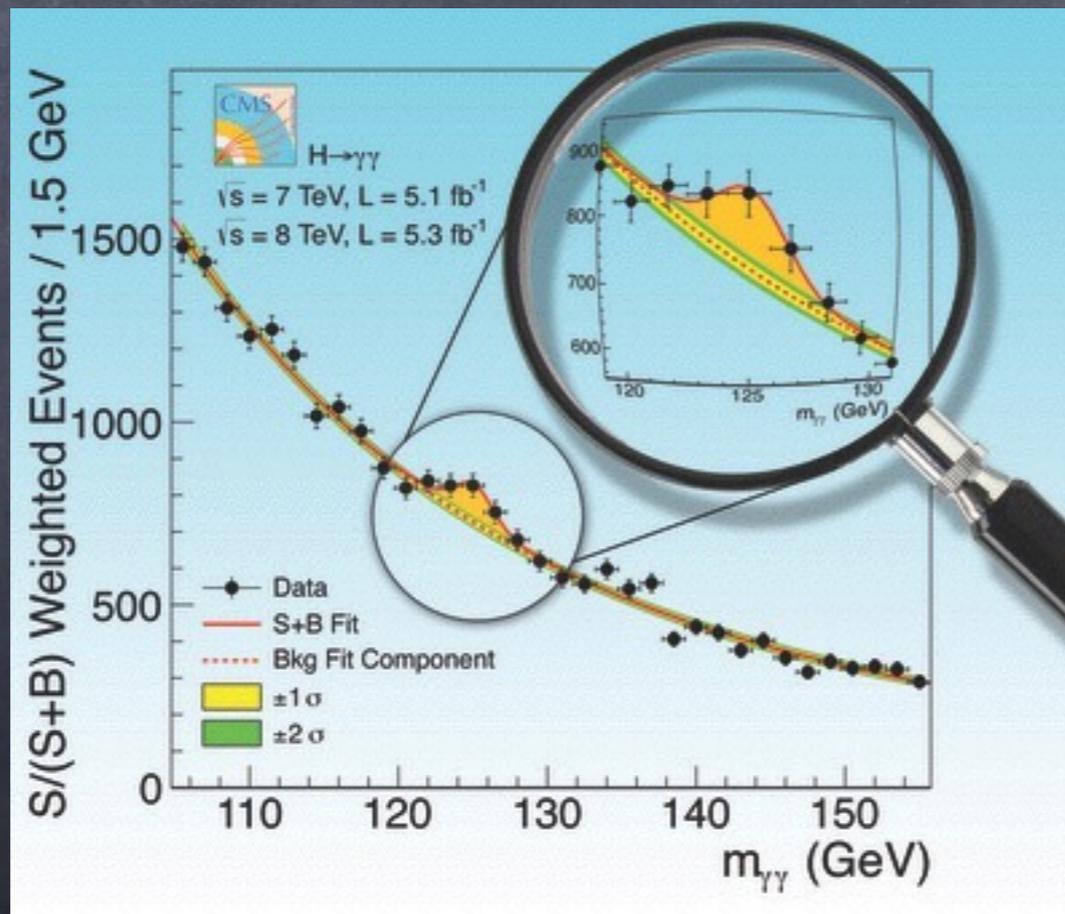


Composite Higgses:
non-minimal, Dark Matter
and light T

G. Cacciapaglia (IPNL)
@ ICTP, São Paulo

Monsieur Le Higgs: enfin!

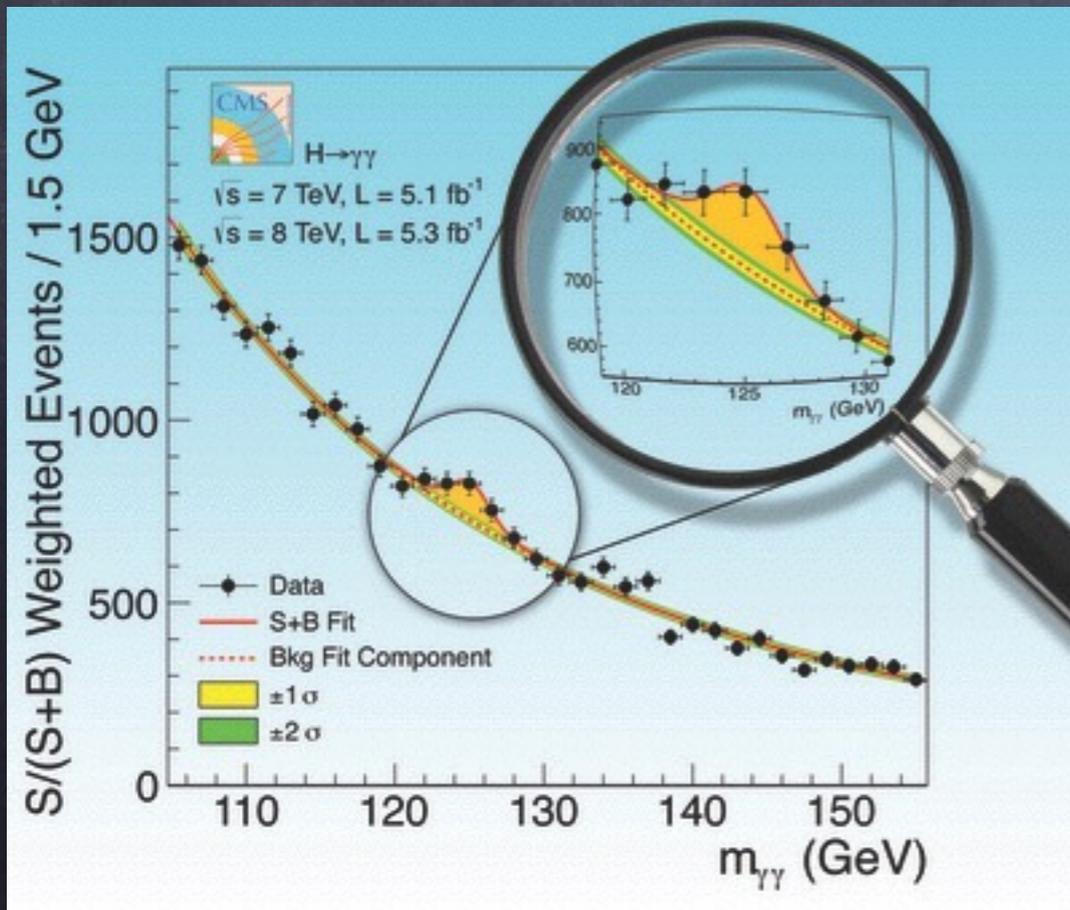


The Standard Model
is finally complete!

(is it?)

Monsieur Le Higgs: enfin!

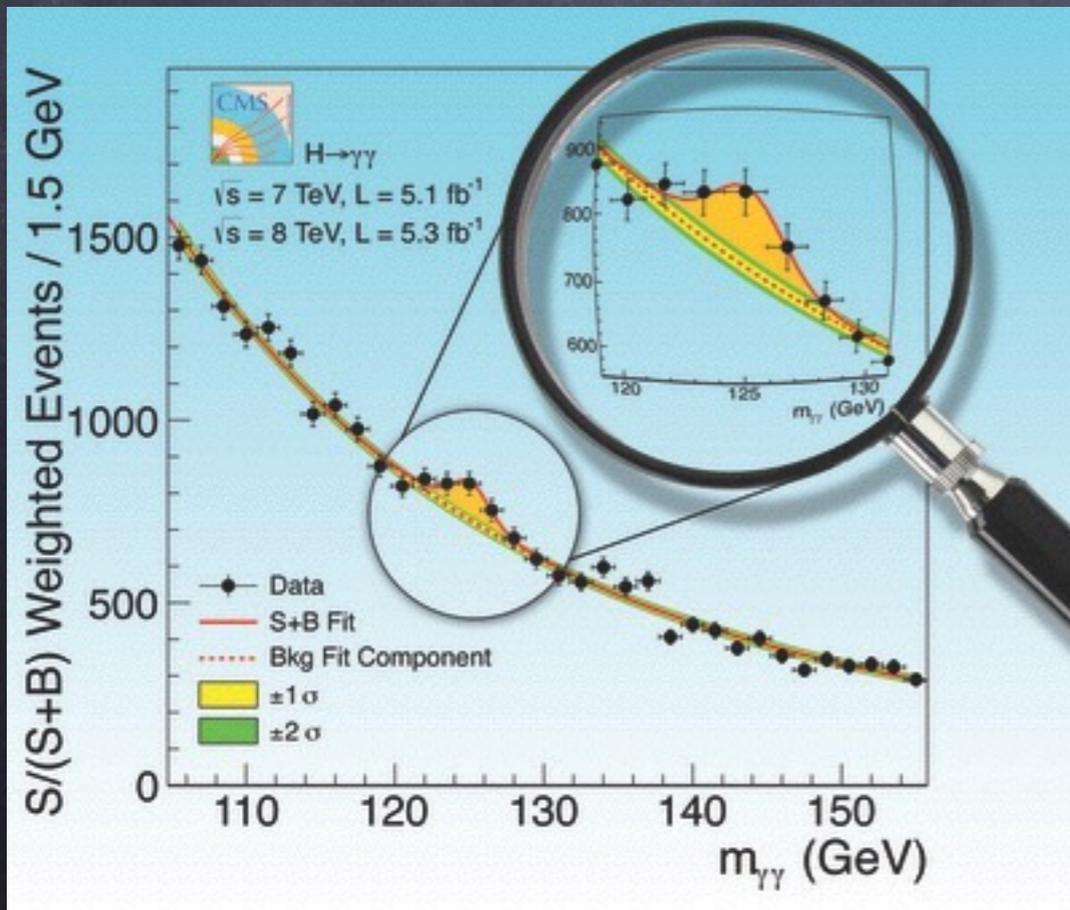
...and the LHC crew
is happy!



Monsieur Le Higgs: enfin!

For a theorist, this is the beginning of a new era!

We have a new toy,



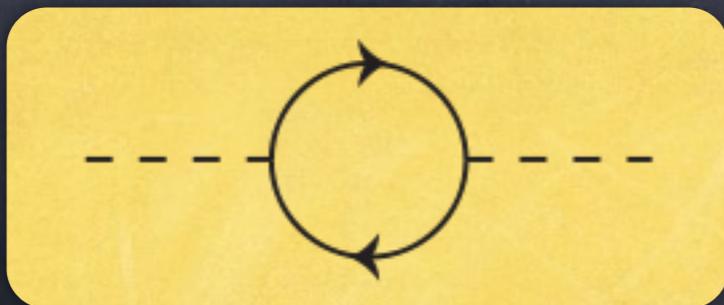
a new probe in the EW sector!

Do we still need BSM?



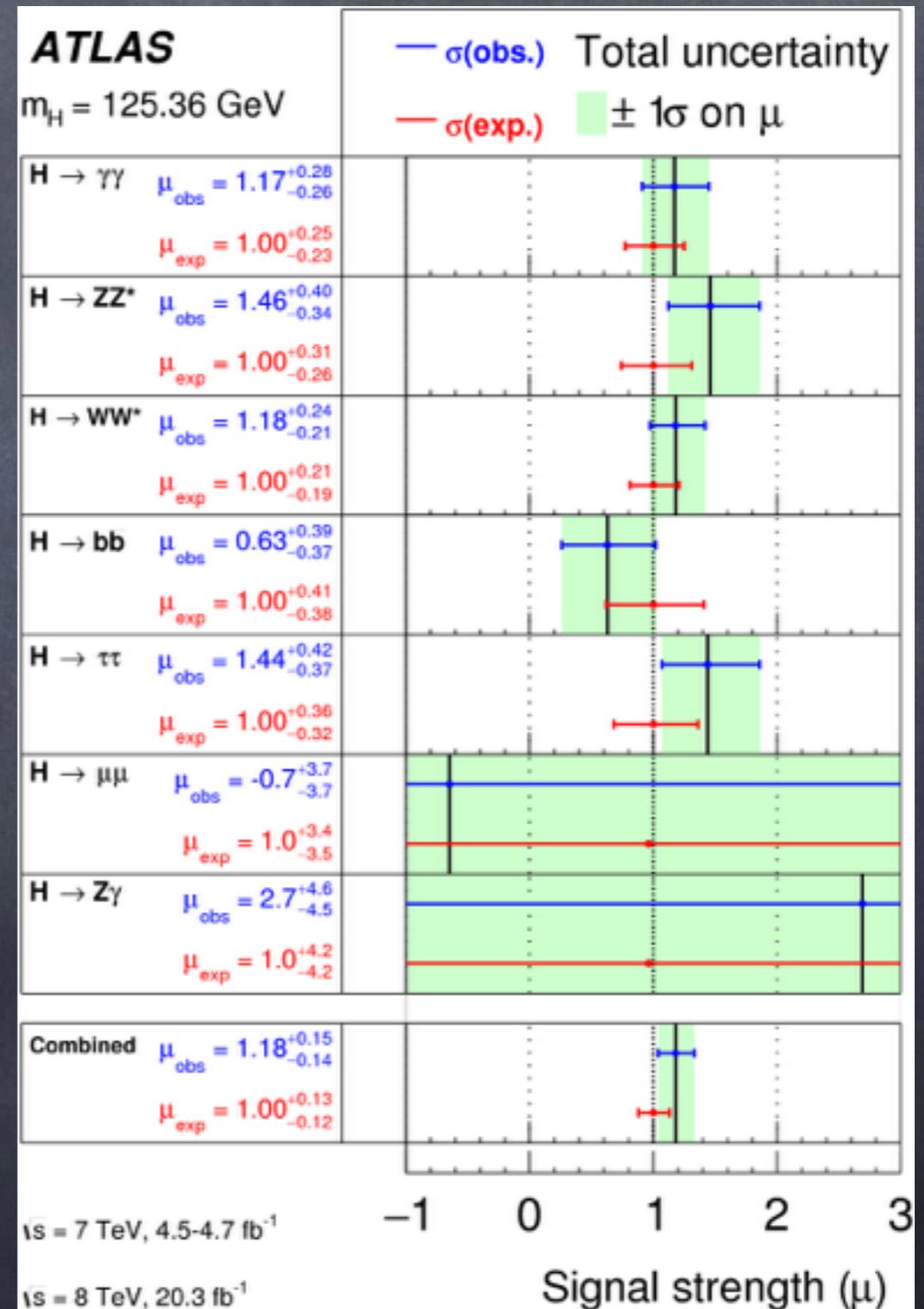
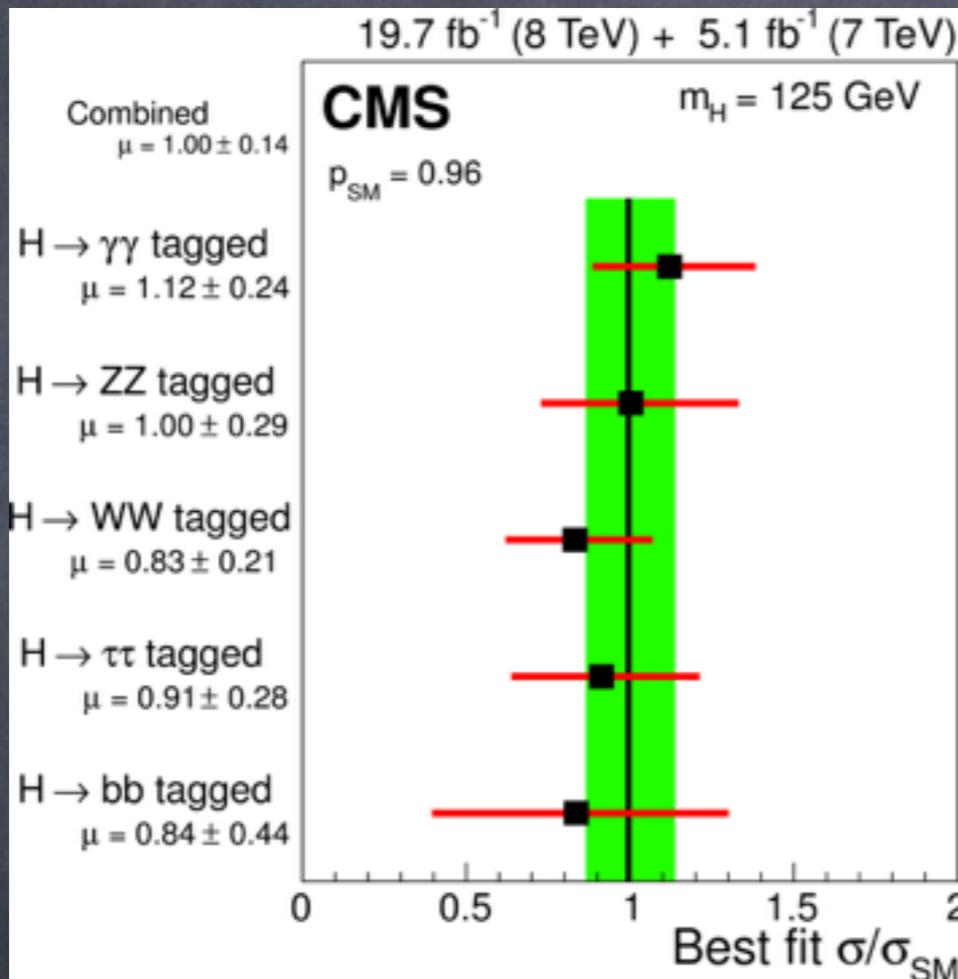
We have a pretty good idea of the mechanism

But, we don't know how to protect it:



$$\delta m_h^2 \sim \frac{g^2}{16\pi^2} M_{\text{NPh}}^2$$

Do we still need BSM?

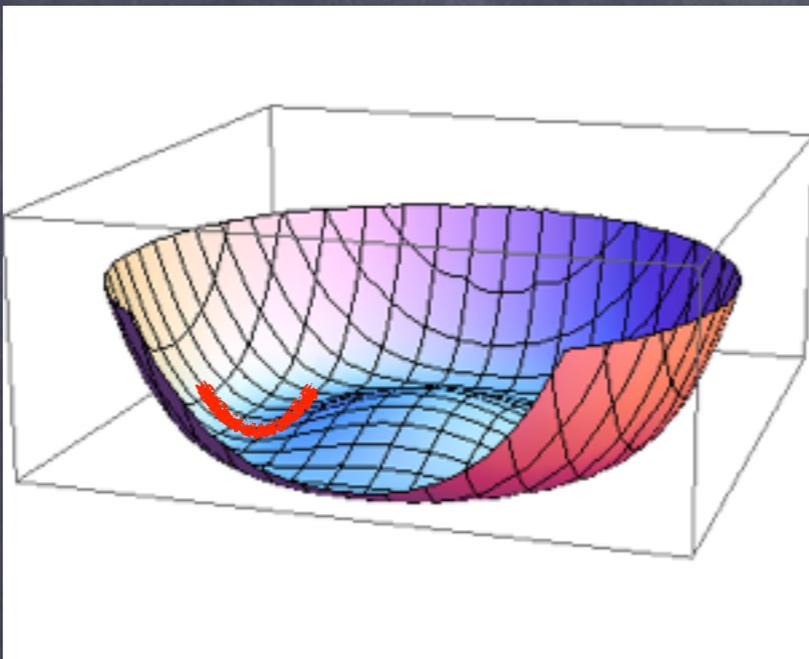


- Still large uncertainties on the Higgs coupling measurements.
- Ample space for BSM!!!!

Outline:

- An old idea: compositeness
- Rediscovery of composite Higgses
- Back to the origins: fundamental composite dynamics (FCD) approach
- Non-minimal models, and Dark Matter

Compositeness, and the Higgs boson



Global symmetry of the system:

$$G \rightarrow H$$

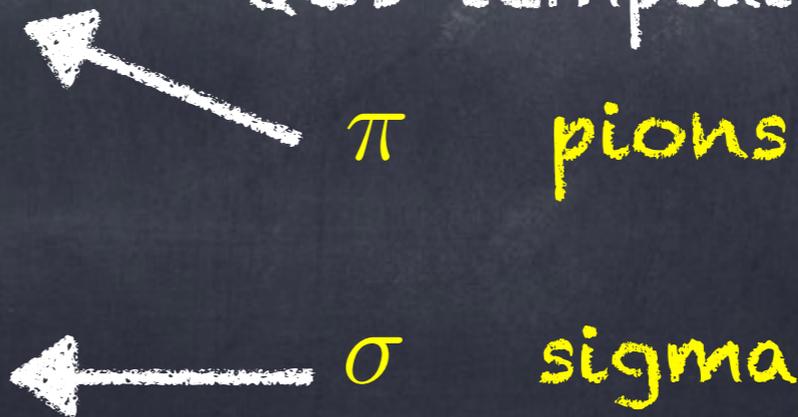


$$SU(2) \times U(1) \rightarrow U(1)_{em}$$

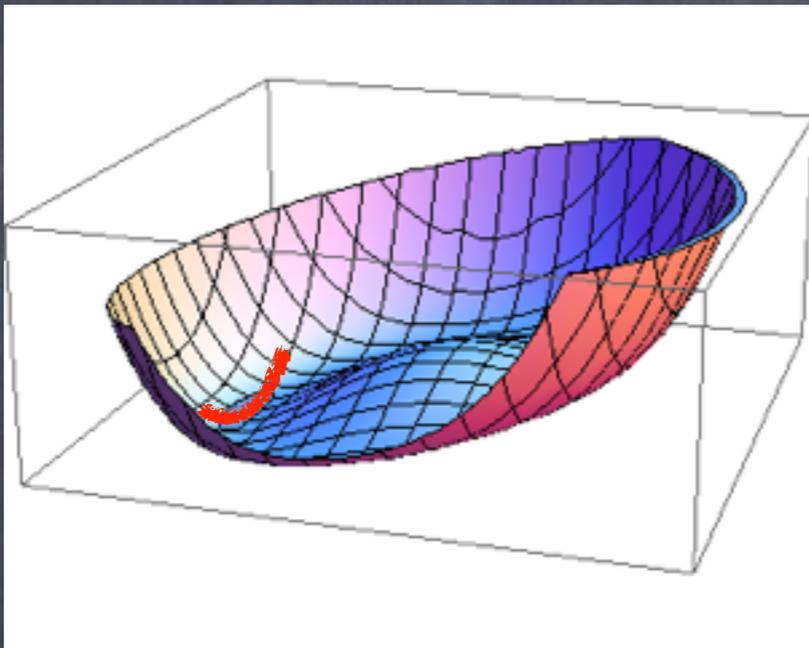
SM gauge symmetry!

- Goldstones include the longitudinal d.o.f. of W and Z
- the Higgs is a heavy bound state (singlet under H)

QCD template:



Compositeness, and the Higgs boson



Global symmetry of the system:

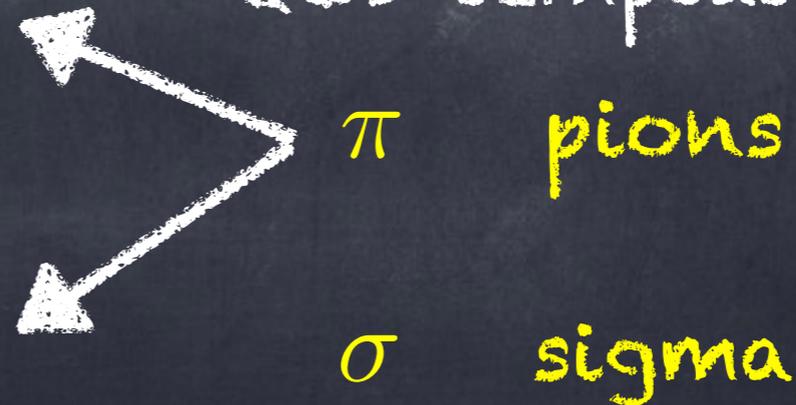
$$G \rightarrow \mathcal{H}$$

$$SU(2) \times U(1) \rightarrow U(1)_{\text{em}}$$

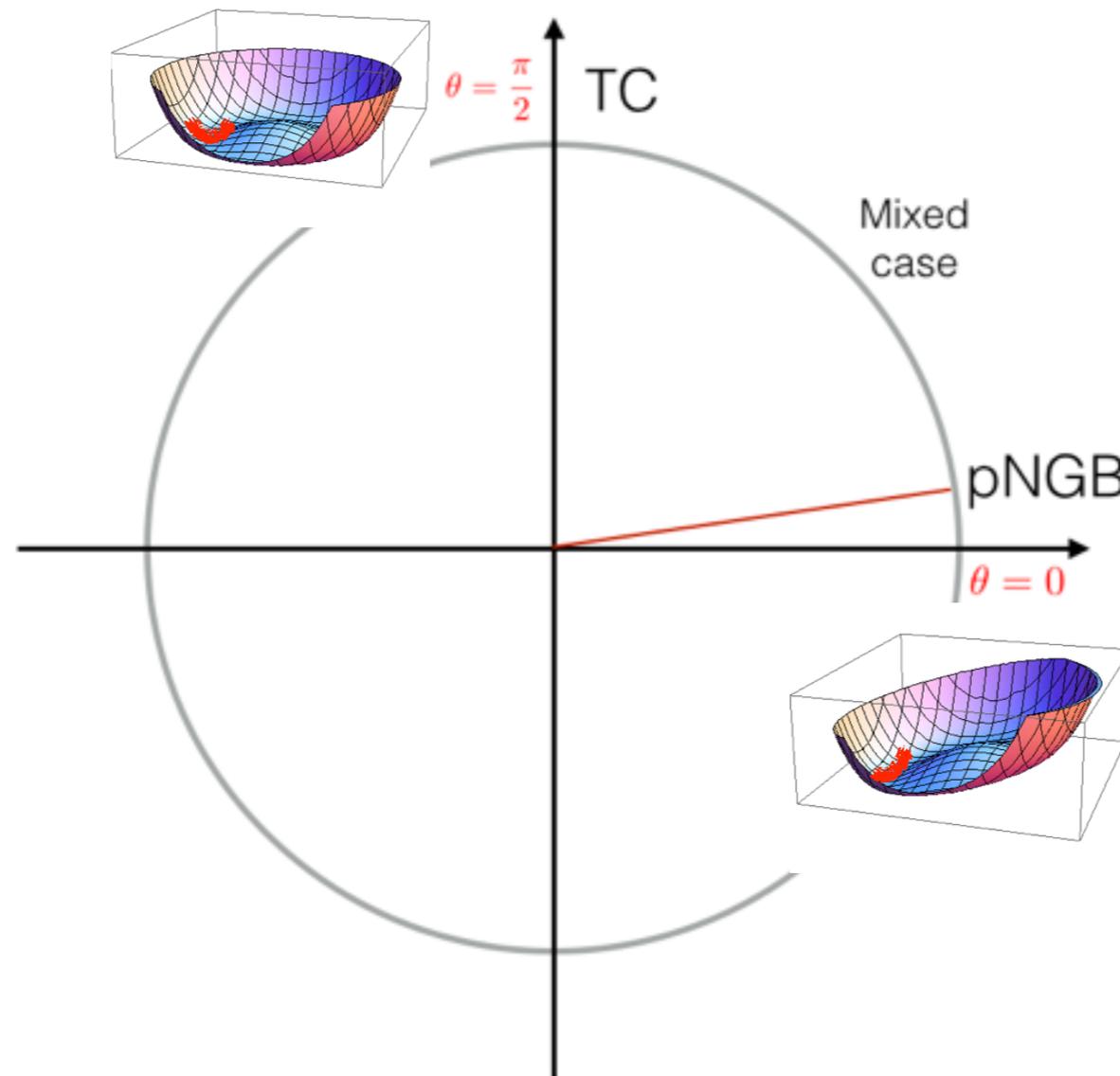
SM gauge symmetry!

- Goldstones include the longitudinal d.o.f. of W and Z
- the Higgs is a pseudo-Goldstone

QCD template:



Compositeness, and the Higgs boson



Higgs: p-Goldstone vs. sigma

👍 Mass is param. lighter than the compositeness scale

👍 The couplings to SM states are naturally close to SM

👎 Tuning the tilt in the potential!

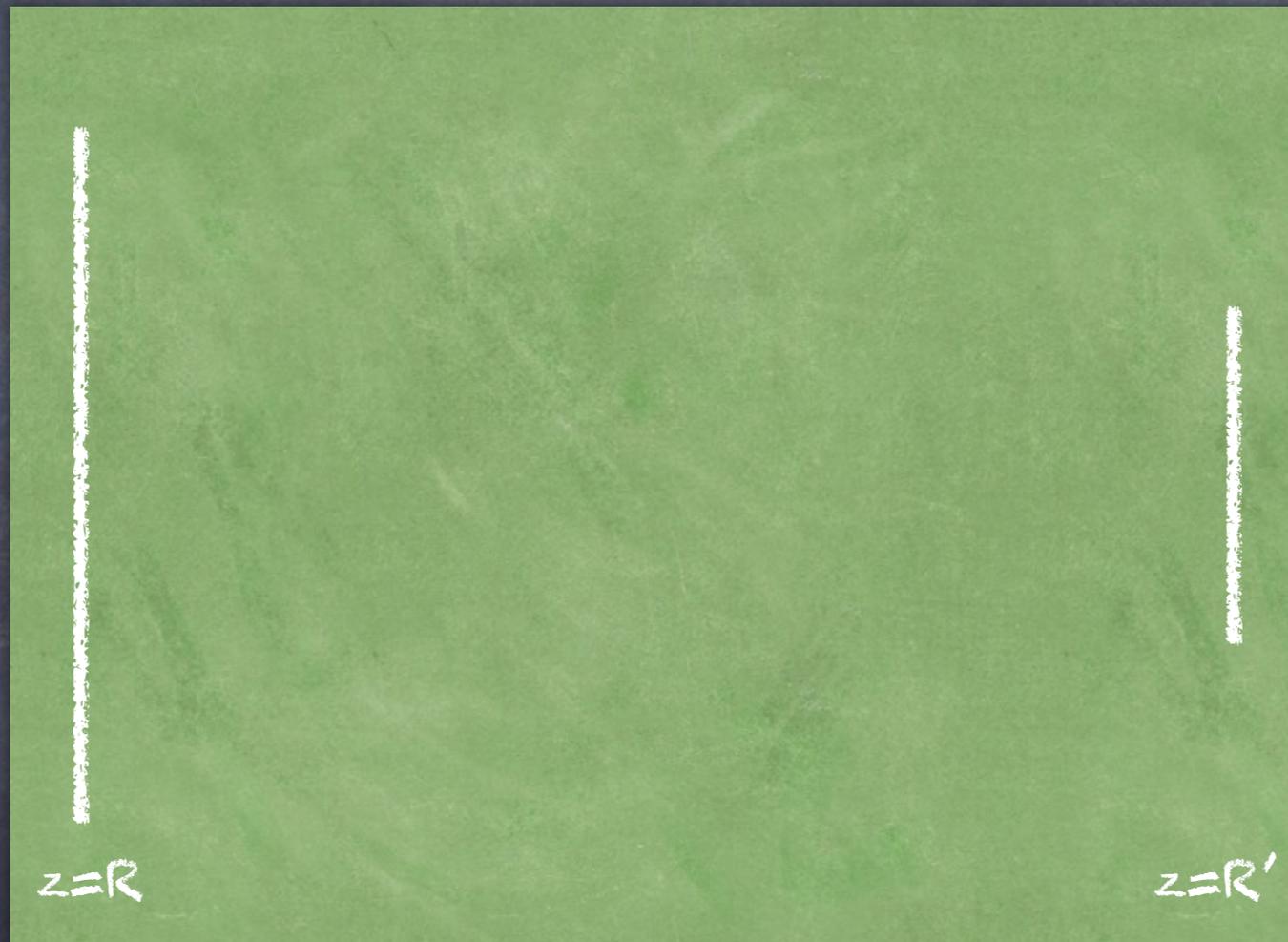
👎 Mass is expected to be heavy (close to the rho)

👐 The couplings to SM states are Unknown!

👍 No tuning is necessary!

Rediscovery of composite Higgses

Randall Sundrum

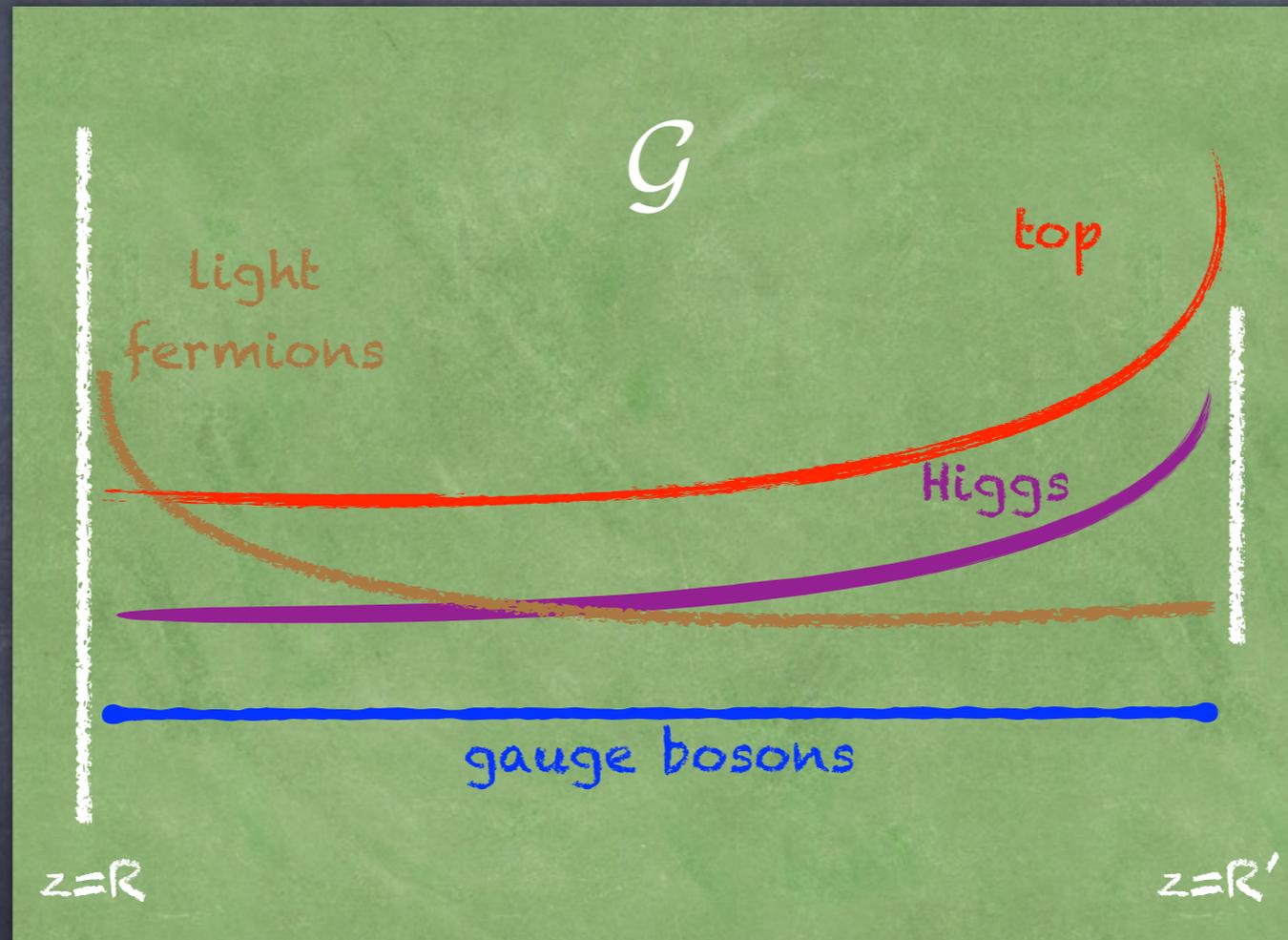


$$ds^2 = \left(\frac{R}{z}\right)^2 (dx^2 - dz^2)$$

Rediscovery of composite Higgses

Randall Sundrum

Elementary
sector



Breaking to H

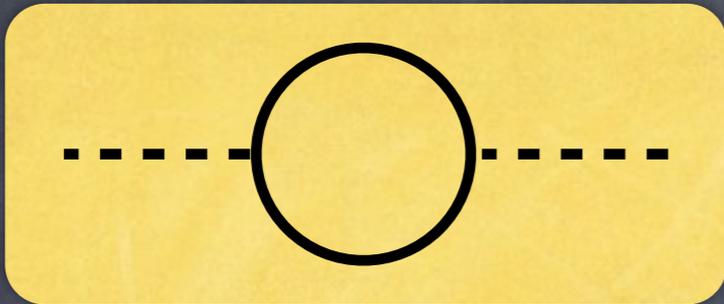
- Conformal dynamics
- UV-insensitive potential
- Fermion partial compositeness
- Flavour under control

Composite Higgses: which model?

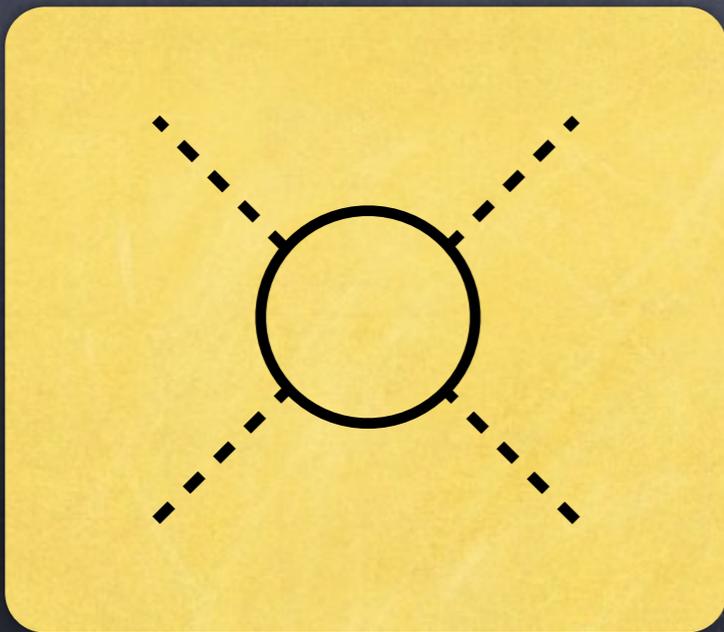
| \mathcal{G} | \mathcal{H} | C | N_G | $\mathbf{r}_\mathcal{H} = \mathbf{r}_{\text{SU}(2) \times \text{SU}(2)} (\mathbf{r}_{\text{SU}(2) \times \text{U}(1)})$ | Ref. |
|----------------------|-----------------------------|-----|-------|---|--------------|
| SO(5) | SO(4) | ✓ | 4 | $\mathbf{4} = (\mathbf{2}, \mathbf{2})$ | [11] |
| SU(3) × U(1) | SU(2) × U(1) | | 5 | $\mathbf{2}_{\pm 1/2} + \mathbf{1}_0$ | [10, 35] |
| SU(4) | Sp(4) | ✓ | 5 | $\mathbf{5} = (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$ | [29, 47, 64] |
| SU(4) | [SU(2)] ² × U(1) | ✓* | 8 | $(\mathbf{2}, \mathbf{2})_{\pm 2} = 2 \cdot (\mathbf{2}, \mathbf{2})$ | [65] |
| SO(7) | SO(6) | ✓ | 6 | $\mathbf{6} = 2 \cdot (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$ | – |
| SO(7) | G ₂ | ✓* | 7 | $\mathbf{7} = (\mathbf{1}, \mathbf{3}) + (\mathbf{2}, \mathbf{2})$ | [66] |
| SO(7) | SO(5) × U(1) | ✓* | 10 | $\mathbf{10}_0 = (\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{3}) + (\mathbf{2}, \mathbf{2})$ | – |
| SO(7) | [SU(2)] ³ | ✓* | 12 | $(\mathbf{2}, \mathbf{2}, \mathbf{3}) = 3 \cdot (\mathbf{2}, \mathbf{2})$ | – |
| Sp(6) | Sp(4) × SU(2) | ✓ | 8 | $(\mathbf{4}, \mathbf{2}) = 2 \cdot (\mathbf{2}, \mathbf{2})$ | [65] |
| SU(5) | SU(4) × U(1) | ✓* | 8 | $\mathbf{4}_{-5} + \bar{\mathbf{4}}_{+5} = 2 \cdot (\mathbf{2}, \mathbf{2})$ | [67] |
| SU(5) | SO(5) | ✓* | 14 | $\mathbf{14} = (\mathbf{3}, \mathbf{3}) + (\mathbf{2}, \mathbf{2}) + (\mathbf{1}, \mathbf{1})$ | [9, 47, 49] |
| SO(8) | SO(7) | ✓ | 7 | $\mathbf{7} = 3 \cdot (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$ | – |
| SO(9) | SO(8) | ✓ | 8 | $\mathbf{8} = 2 \cdot (\mathbf{2}, \mathbf{2})$ | [67] |
| SO(9) | SO(5) × SO(4) | ✓* | 20 | $(\mathbf{5}, \mathbf{4}) = (\mathbf{2}, \mathbf{2}) + (\mathbf{1} + \mathbf{3}, \mathbf{1} + \mathbf{3})$ | [34] |
| [SU(3)] ² | SU(3) | | 8 | $\mathbf{8} = \mathbf{1}_0 + \mathbf{2}_{\pm 1/2} + \mathbf{3}_0$ | [8] |
| [SO(5)] ² | SO(5) | ✓* | 10 | $\mathbf{10} = (\mathbf{1}, \mathbf{3}) + (\mathbf{3}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$ | [32] |
| SU(4) × U(1) | SU(3) × U(1) | | 7 | $\mathbf{3}_{-1/3} + \bar{\mathbf{3}}_{+1/3} + \mathbf{1}_0 = 3 \cdot \mathbf{1}_0 + \mathbf{2}_{\pm 1/2}$ | [35, 41] |
| SU(6) | Sp(6) | ✓* | 14 | $\mathbf{14} = 2 \cdot (\mathbf{2}, \mathbf{2}) + (\mathbf{1}, \mathbf{3}) + 3 \cdot (\mathbf{1}, \mathbf{1})$ | [30, 47] |
| [SO(6)] ² | SO(6) | ✓* | 15 | $\mathbf{15} = (\mathbf{1}, \mathbf{1}) + 2 \cdot (\mathbf{2}, \mathbf{2}) + (\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{3})$ | [36] |

Table 1: Symmetry breaking patterns $\mathcal{G} \rightarrow \mathcal{H}$ for Lie groups. The third column denotes whether the breaking pattern incorporates custodial symmetry. The fourth column gives the dimension N_G of the coset, while the fifth contains the representations of the GB's under \mathcal{H} and $\text{SO}(4) \cong \text{SU}(2)_L \times \text{SU}(2)_R$ (or simply $\text{SU}(2)_L \times \text{U}(1)_Y$ if there is no custodial symmetry). In case of more than two SU(2)'s in \mathcal{H} and several different possible decompositions we quote the one with largest number of bi-doublets.

Anatomy of the potential

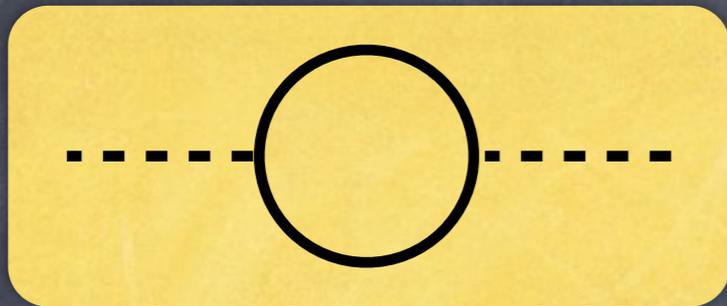


$$\delta m_h^2 \sim \frac{y^2}{16\pi^2} \Lambda^2$$



$$\delta \lambda \sim \frac{y^4}{16\pi^2} \log \Lambda$$

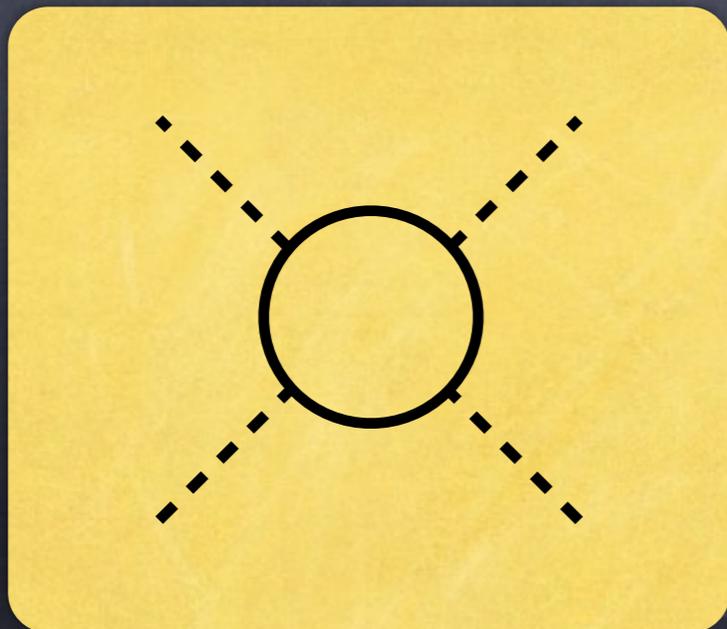
Anatomy of the potential



$$\delta m_h^2 \sim \frac{y^2}{16\pi^2} (4\pi f)^2 \sim y^2 f^2$$

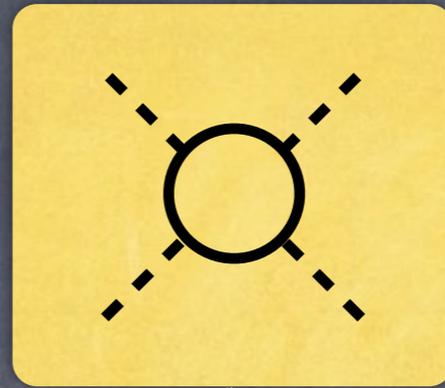
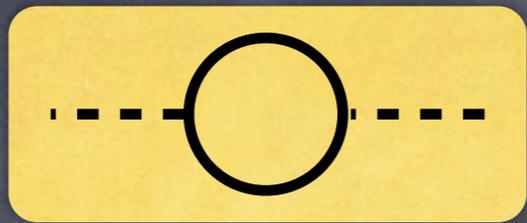
$$\Lambda \sim 4\pi f$$

Strong dynamics
estimate



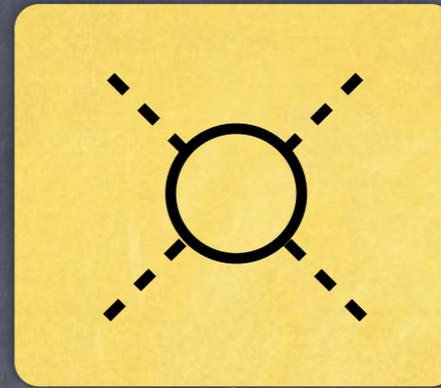
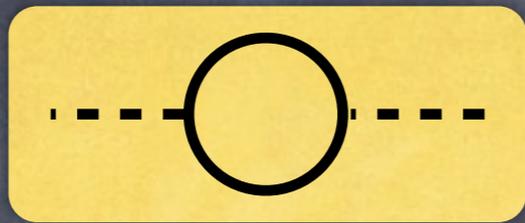
$$\delta\lambda \sim \frac{y^4}{16\pi^2} \log \Lambda$$

Anatomy of the potential



$$V \sim \alpha \sin^2 \theta + \beta \sin^4 \theta$$

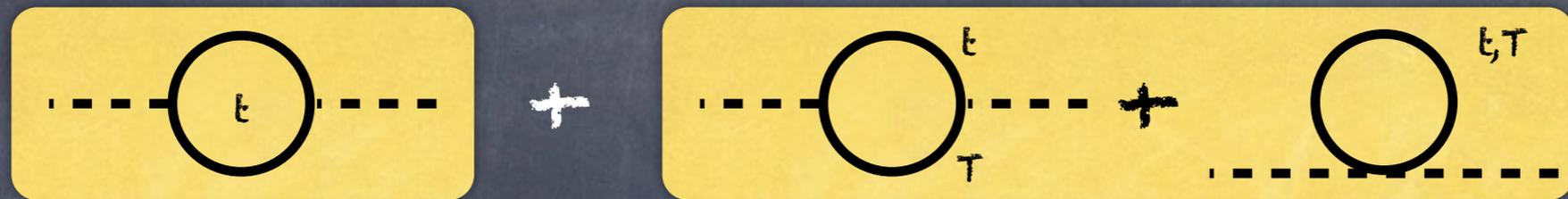
Anatomy of the potential



$V \sim \alpha \sin^2 \theta + \beta \sin^4 \theta$

Minimum: $\theta \sim \frac{\pi}{2}$

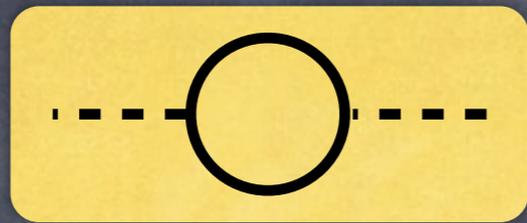
Anatomy of the potential



$$V \sim \alpha \sin^2 \theta + \beta \sin^4 \theta$$

Minimum: $\theta \sim \epsilon$

Anatomy of the potential



Explicit
symmetry
breaking?

$$V \sim \alpha \sin^2 \theta + \gamma \sin \theta$$

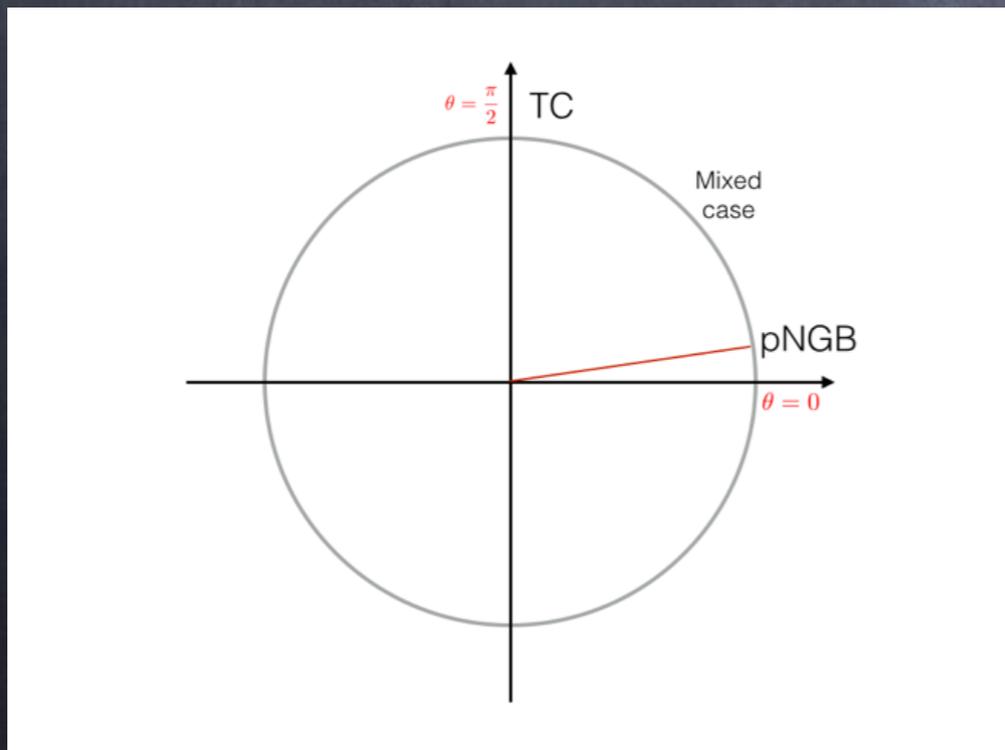
Minimum: $\theta \sim \epsilon$

Possible!

Anatomy of the potential

Higgs mass in the small theta limit:

$$m_h \sim y f \sin \theta \sim y v_{SM}$$



Naturally in the
right ballpark,
without further
fine tuning!

The FCD approach

G.C., F.Sannino

1402.0233

- Define a confining gauge group (GTC)
- Add in N fermions charged under the confining group GTC
- Assign SM quantum numbers to the fermions
- Couple them to SM fermions

The FCD approach

| coset | GTC | TF | doublets | pNGBs |
|------------------------------|----------|------|----------|-------|
| $SU(4)/Sp(4)$ | $Sp(2N)$ | fund | 1 | 5 |
| $SU(5)/SO(5)$ | $SU(4)$ | 6 | 1 | 14 |
| $SU(4) \times SU(4) / SU(4)$ | $SU(N)$ | fund | 2 | 15 |
| $SU(6)/Sp(6)$ | $Sp(2N)$ | fund | 2 | 14 |

← Minimal!

G.C., T.Ma
1508.07014

G.C., M.Lespinasse
in prep.

The minimal case

T.Ryttov, F.Sannino 0809.0713
Galloway, Evans, Luty, Tacchi 1001.1361

Anti-symmetric

$$\langle \psi^i \psi^j \rangle = 6_{\text{SU}(4)} \rightarrow 5_{\text{Sp}(4)} \oplus 1_{\text{Sp}(4)}$$

$\text{Sp}(4)$ contains a $\text{SO}(4)$ subgroup:
identify with custodial symmetry!

Pions: $5_{\text{Sp}(4)} \rightarrow (2, 2) \oplus (1, 1)$

The minimal case

Galloway, Evans, Luty, Tacchi 1001.1361
G.C., F.Sannino 1402.0233

$$V_m = C_m f^4 \text{Tr}(\Sigma_B \cdot \Sigma) \\ \sim C_m \left(-4f^4 c_\theta + \sqrt{2} f^3 s_\theta h + \frac{1}{4} f^2 c_\theta (h^2 + \eta^2) + \dots \right)$$

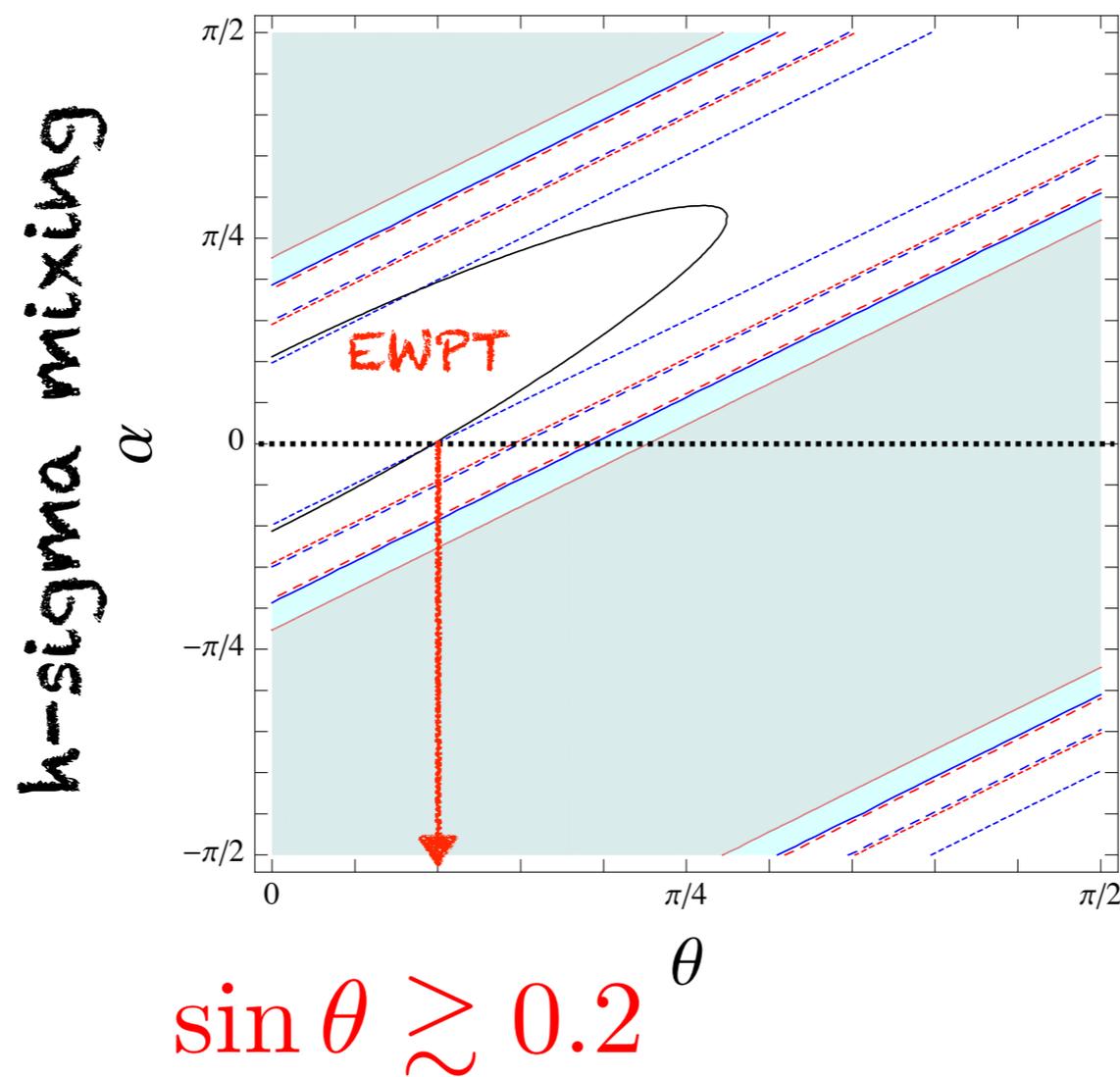
$$V_{top} = -C_t y_t'^2 f^4 \sum_{\alpha=1}^2 [\text{Tr}(P^\alpha \Sigma)]^2 \\ \sim -C_t y_t'^2 \left[f^4 s_\theta^2 + \frac{1}{\sqrt{2}} f^3 c_\theta s_\theta h + \frac{1}{8} f^2 (c_{2\theta} h^2 - s_\theta^2 \eta^2) + \dots \right]$$

$$m_h \sim \frac{\sqrt{C_t}}{2} y_t f \sin \theta \sim \frac{\sqrt{C_t}}{2} m_t \quad m_\eta \sim \frac{\sqrt{C_t}}{2} y_t f \sim \frac{m_h}{\sin \theta}$$

$$C_t \sim 2 \quad \text{reproduces data!}$$

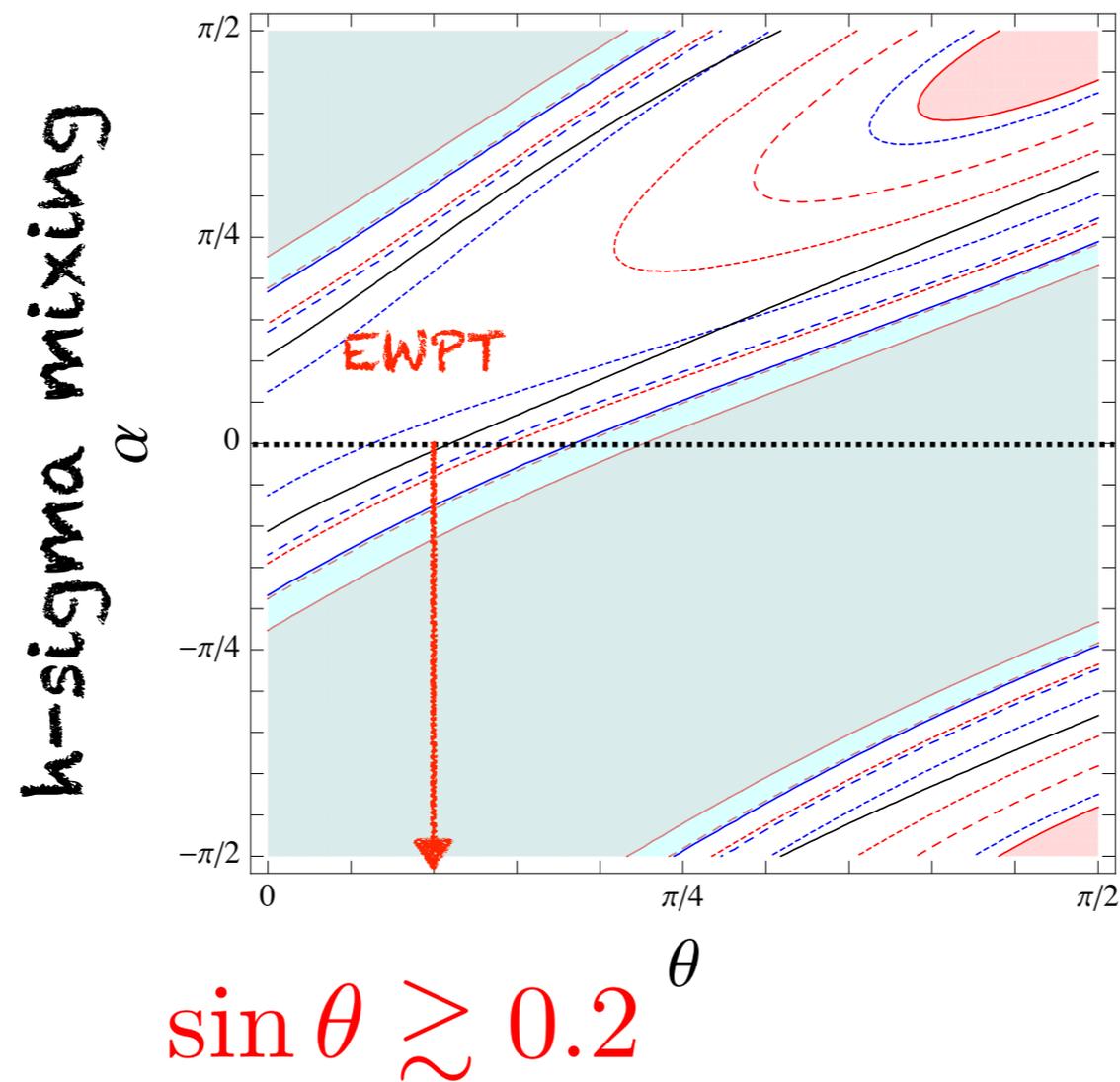
The minimal case

Arbey, G.C., Cai, Deandrea, Le Corre, Sannino
1502.04718



The minimal case

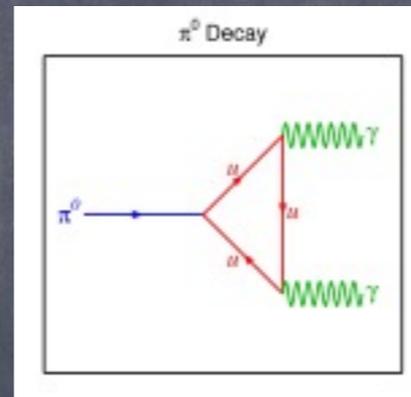
Arbey, G.C., Cai, Deandrea, Le Corre, Sannino
1502.04718



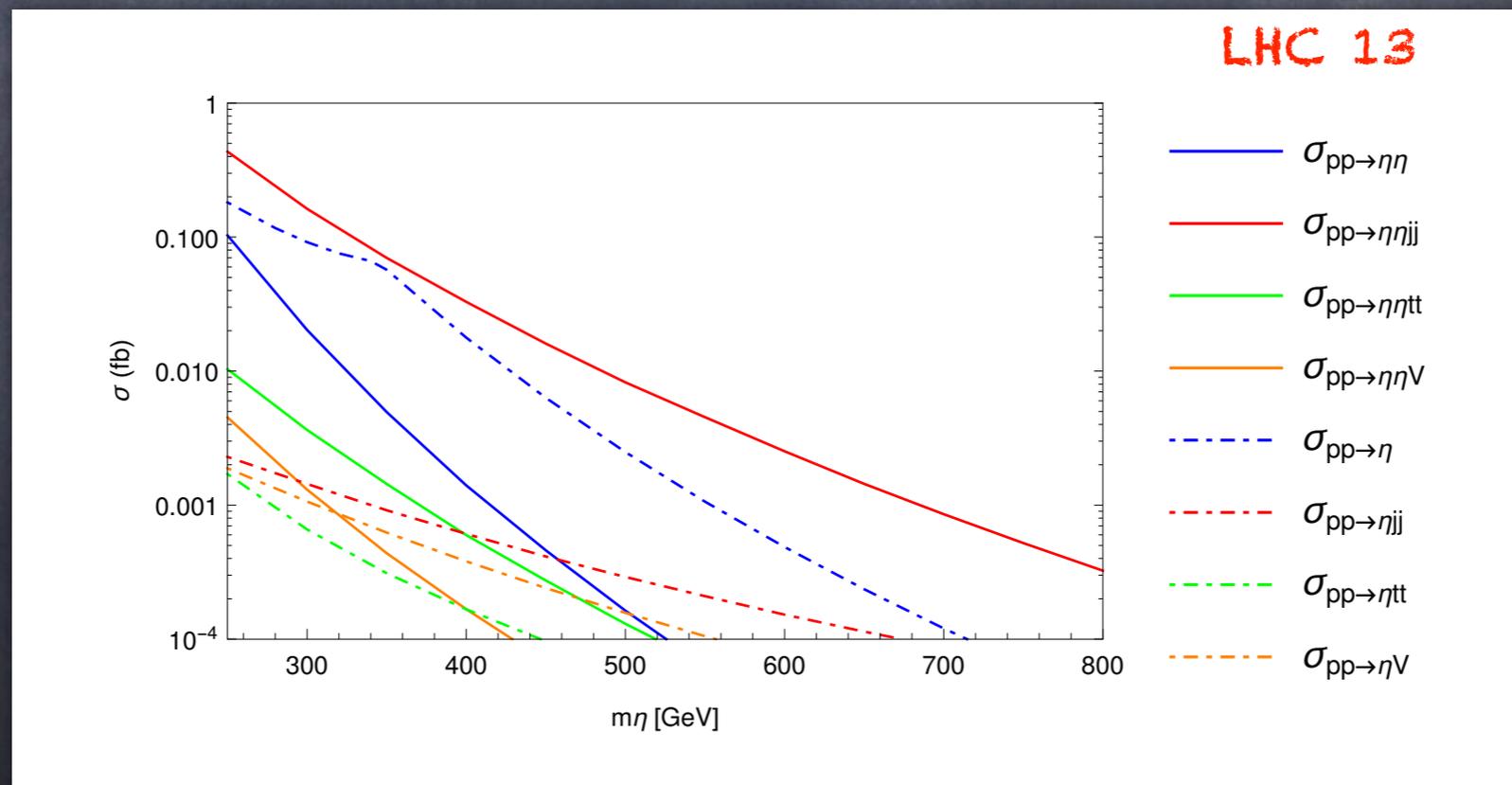
The minimal case

Arbey, G.C., Cai, Deandrea, Le Corre, Sannino
1502.04718

Singlet cannot be
Dark Matter!



Anomalous coupling to:
WW, ZZ, Z gamma



The minimal case

Lattice results:

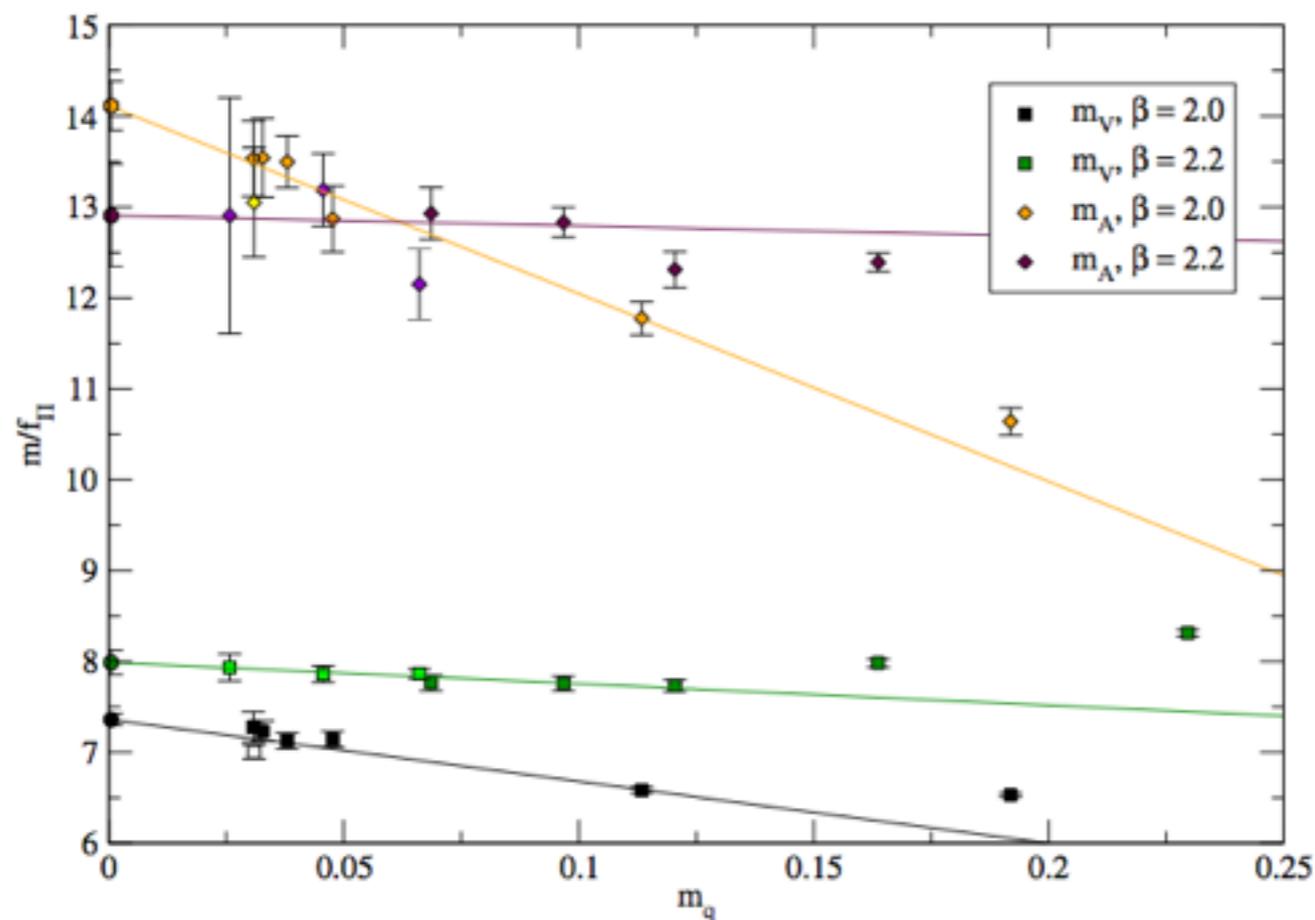


FIG. 6: The vector meson and axial vector meson masses in physical units. The chiral extrapolations have been performed using a linear fit to the points where $m_q < 0.12$.

$$m_a \sim \frac{3 \text{ TeV}}{\sin \theta} \gtrsim 15 \text{ TeV}$$

$$m_\rho \sim \frac{2.5 \text{ TeV}}{\sin \theta} \gtrsim 12 \text{ TeV}$$

$$m_\sigma \sim ???$$

$$m_\eta \sim \frac{m_h}{\sin \theta} \gtrsim 600 \text{ GeV}$$

$$m_h = 125 \text{ GeV}$$

A composite 2HDM

G.C., T.Ma
1508.07014

| | SU(N) | SU(2) _L | U(1) _Y |
|---|-------|----------------------|-------------------|
| $\psi_L = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$ | □ | 2 | 0 |
| $\psi_R = \begin{pmatrix} \psi_3 \\ \psi_4 \end{pmatrix}$ | □ | 1 1 | 1/2 -1/2 |

$$\mathcal{L}_{FCD} = i\bar{\psi}D_\mu\gamma^\mu\psi - \bar{\psi}M_Q\psi$$

$$\Pi = \frac{1}{2} \begin{pmatrix} \sigma_i\Delta^i + s/\sqrt{2} & -i\Phi_H \\ i\Phi_H^\dagger & \sigma_i N^i - s/\sqrt{2} \end{pmatrix}$$

A composite 2HDM

G.C., T.Ma
1508.07014

$$V_{\text{top}}(\theta) = -C_t f^4 \left[8|Y_t|^2 \sin^2 \theta + \leftarrow \text{Potential for theta} \right. \\ \left. 2\sqrt{2}|Y_t|^2 \sin(2\theta) \frac{h_1}{f} + \right. \\ \left. + 4\sqrt{2} \text{Im}(Y_D^* Y_t) \sin \theta \frac{h_2}{f} \right. \\ \left. + 2\sqrt{2} \text{Re}(Y_D^* Y_t) \sin(2\theta) \frac{A_0}{f} \right. \\ \left. + 4 \text{Im}(Y_T^* Y_t) \sin^2 \theta \frac{N_0 + \Delta_0}{f} + \dots \right]$$

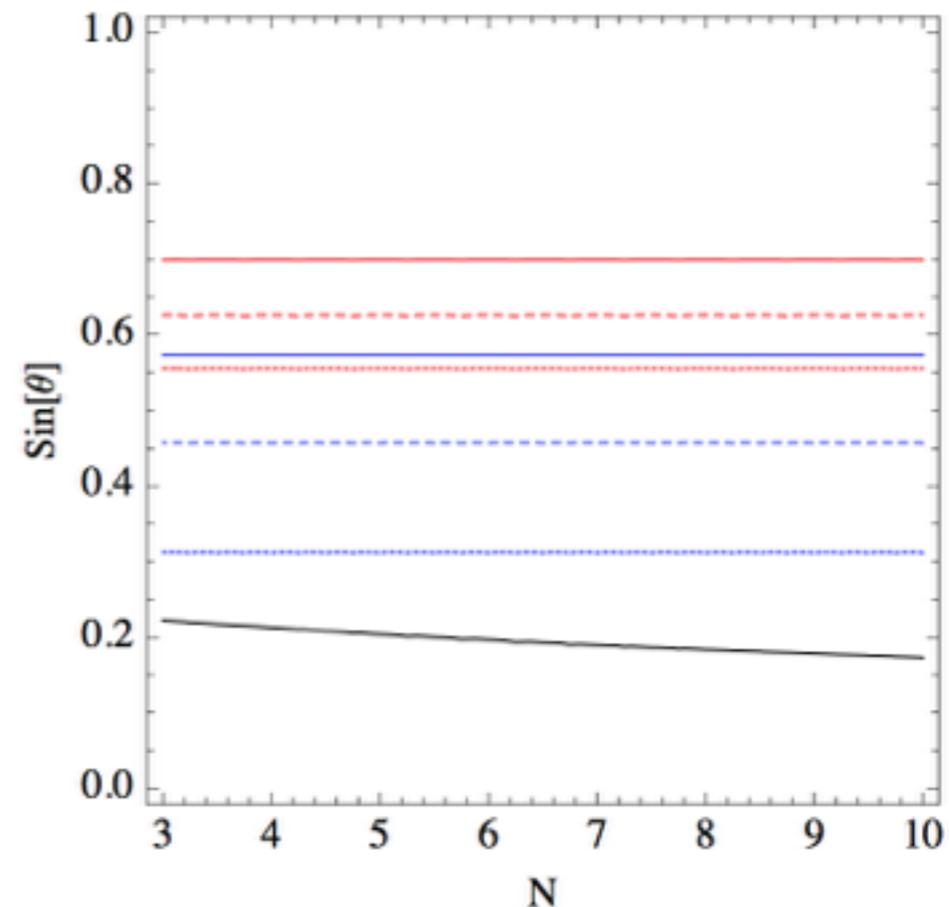
A composite 2HDM: EWPTs

$$\Delta S_{\text{Higgs}} = \frac{1 - \kappa_V^2}{6\pi} \ln \frac{\Lambda_{\text{FCD}}}{m_h}, \quad \Delta T_{\text{Higgs}} = -\frac{3(1 - \kappa_V^2)}{8\pi \cos^2 \theta_W} \ln \frac{\Lambda_{\text{FCD}}}{m_h},$$

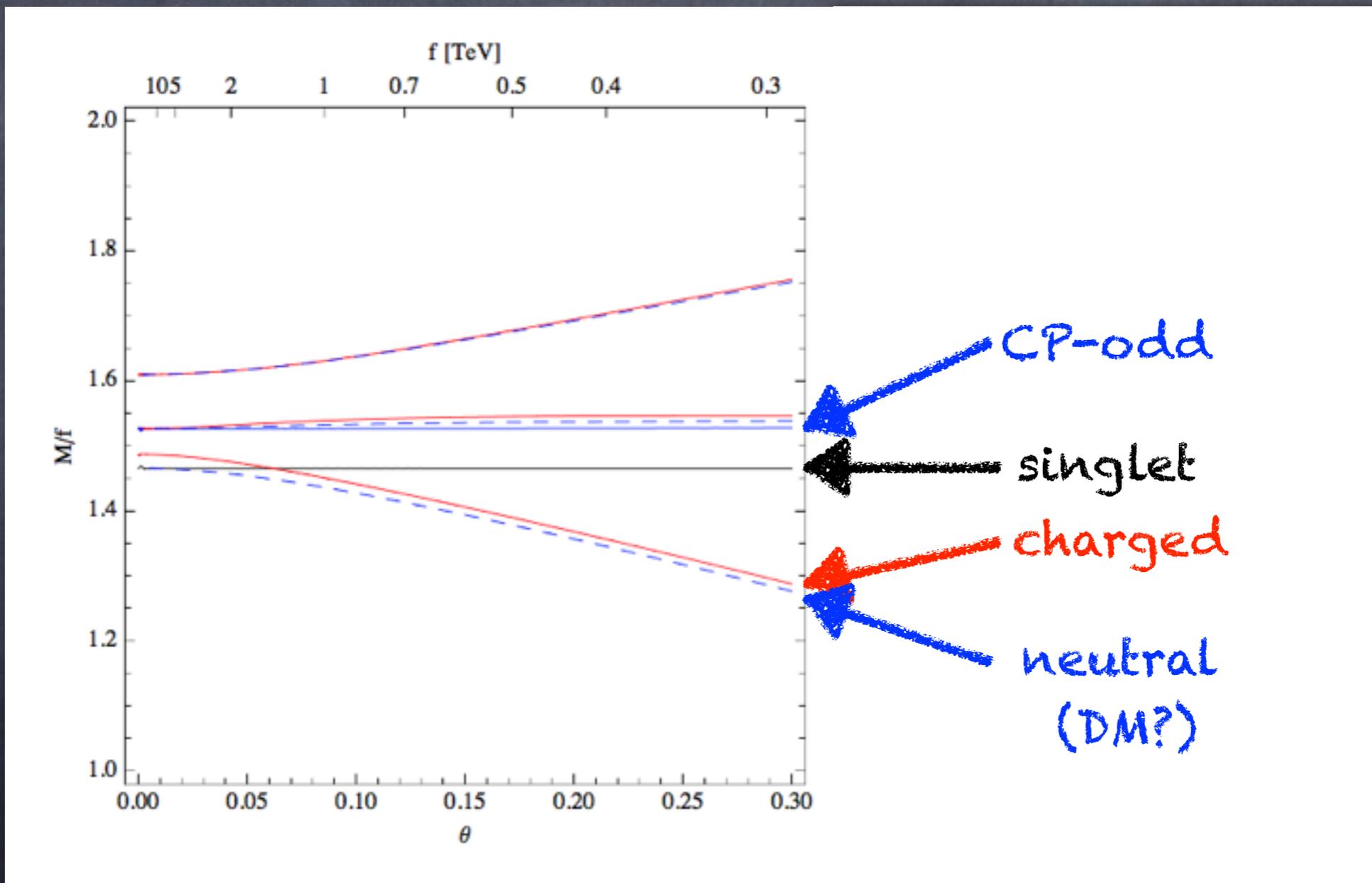
$$\Delta S_{\text{pNGB}} = -\frac{\sin^2 \theta}{4\pi}, \quad \Delta T_{\text{pNGB}} = \frac{\sin^2 \theta}{8\pi \sin^2 \theta_W} \frac{m_{H^\pm}^2 - m_{A_0}^2}{m_W^2} \ln \frac{\Lambda_{\text{FCD}}}{m_{\text{pNGB}}} \sim 0,$$

$$\Delta S_{\text{FCD}} = \frac{\sin^2 \theta}{3\pi} N, \quad \Delta T_{\text{FCD}} \sim 0,$$

Bounds similar
to minimal cases.



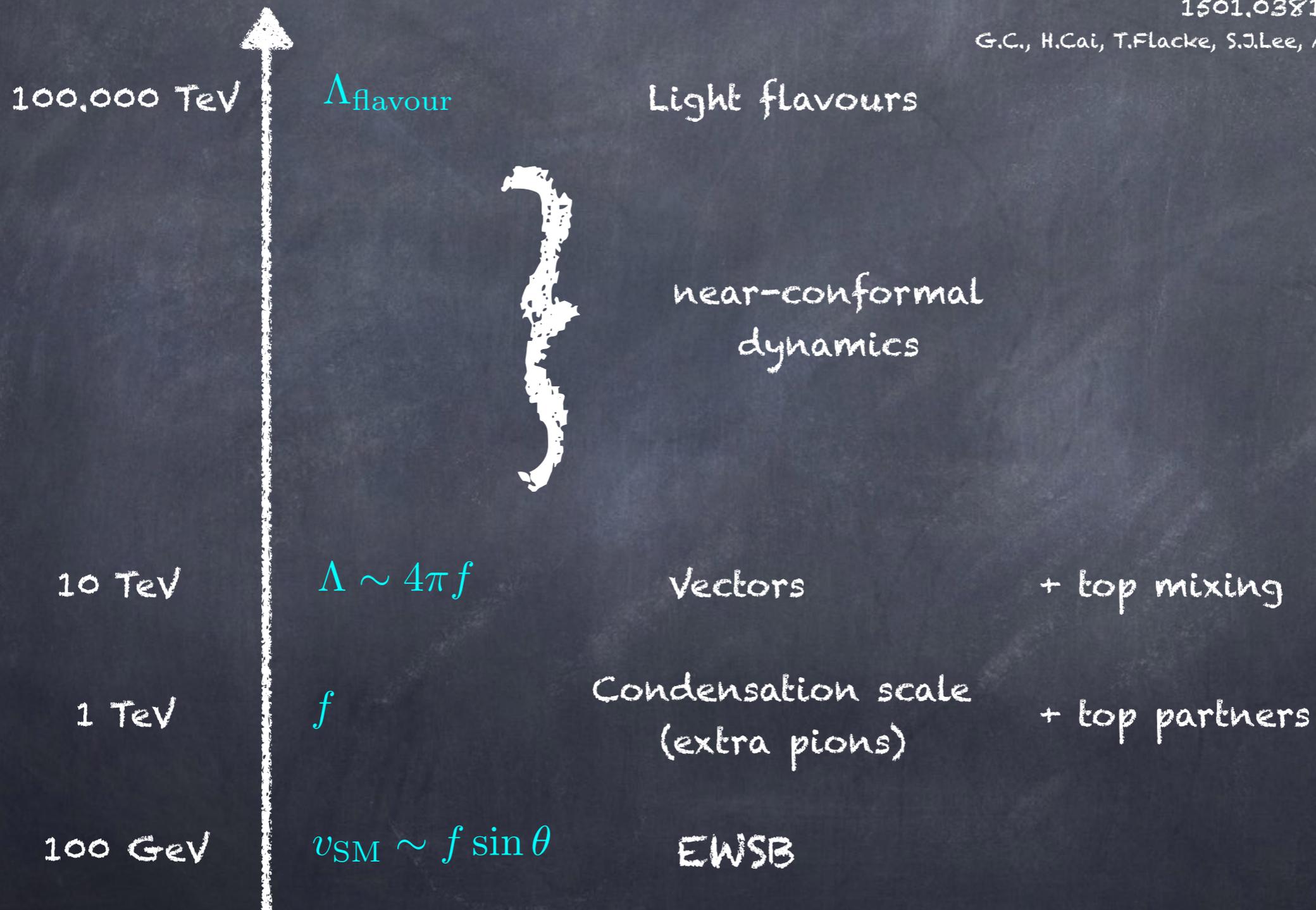
A composite 2HDM: spectrum



The hot potato: flavour!

1501.03818

G.C., H.Cai, T.Flacke, S.J.Lee, A.Parolini, H.Serodio



The hot potato: flavour!

- In FCD, the origin of the top partial compositeness can be probed!

$$\frac{c}{\Lambda_{\text{top}}^2} t\psi\psi \rightarrow y_t tT$$

- Why are the top partners lighter than other resonances?

Light t Hooft Top Partners

G.C., A. Parolini

1511.05163

| | gauged | global symmetries | | |
|--------|------------|-------------------|--------------|-------------------------------|
| | $Sp(2N_c)$ | $SU(N_Q)$ | $SU(N_\chi)$ | $U(1)$ |
| Q | \square | N_Q | 1 | 1 |
| χ | \square | 1 | N_χ | $-\frac{N_Q}{2N_\chi(N_c-1)}$ |

TABLE I: Fermionic field content of the first model.

Barnard, Gherghetta, Ray 1311.6562

Ferretti, Karateev 1312.5330

- $\langle QQ \rangle \neq 0$ breaks $SU(N_Q) \rightarrow Sp(N_Q)$, giving a pNGB Higgs.
- In the vacuum $\langle \chi\chi \rangle = 0$, $SU(N_\chi)$ is unbroken.

$$A_{SU(N_\chi)^2} \propto \dim(\chi) = (2N_c + 1)(N_c - 1)$$

Light t Hooft Top Partners

G.C., A. Parolini

1511.05163

global symmetries

| | $SU(N_Q) \times SU(N_\chi)$ | $Sp(N_Q) \times SU(N_\chi)$ | $d_{Sp(N_Q)}$ |
|------------------------|------------------------------------|----------------------------------|----------------------------|
| χQQ | (\mathbf{A}, \mathbf{F}) | $(\mathbf{1}, \mathbf{F})$ | 1 |
| $\chi \bar{Q} \bar{Q}$ | (\mathbf{S}, \mathbf{F}) | (\mathbf{A}, \mathbf{F}) | $\frac{N_Q(N_Q-1)}{2} - 1$ |
| $\bar{\chi} \bar{Q} Q$ | $(\mathbf{1}, \bar{\mathbf{F}})$ | $(\mathbf{S}, \bar{\mathbf{F}})$ | $\frac{N_Q(N_Q+1)}{2}$ |
| | $(\mathbf{Adj}, \bar{\mathbf{F}})$ | $(\mathbf{A}, \bar{\mathbf{F}})$ | 1 |
| | | $(\mathbf{S}, \bar{\mathbf{F}})$ | $\frac{N_Q(N_Q-1)}{2} - 1$ |
| | | | $\frac{N_Q(N_Q+1)}{2}$ |

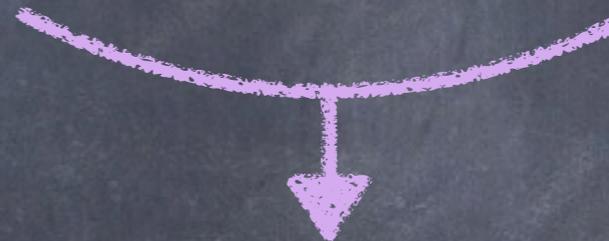
$$A_{SU(N_\chi)^2} \propto \dim(\chi) = (2N_c + 1)(N_c - 1)$$

- Matched if $N_Q = 2N_c$. For $N_Q = 4$, $A = \mathfrak{so}(5)$ of $Sp(4) \sim SO(5)$!!

...a few words on the diphoton excess

More precisely, the global symmetries are:

$$SU(N_Q) \times SU(N_\chi) \times U(1)_Q \times U(1)_\chi$$



$$\begin{aligned} \frac{\Gamma_{\sigma \rightarrow gg}}{\Gamma_{\sigma \rightarrow WW}} &= 5.1 \dots 3.3, \\ \frac{\Gamma_{\sigma \rightarrow ZZ}}{\Gamma_{\sigma \rightarrow WW}} &= 0.29 \dots 0.31, \\ \frac{\Gamma_{\sigma \rightarrow Z\gamma}}{\Gamma_{\sigma \rightarrow WW}} &= 0.19 \dots 0.17, \\ \frac{\Gamma_{\sigma \rightarrow \gamma\gamma}}{\Gamma_{\sigma \rightarrow WW}} &= 0.021 \dots 0.033, \end{aligned}$$

Anomalous $U(1) \rightarrow$ heavy σ'

Orthogonal $U(1) \rightarrow$ pNGB σ

Decays and production
only via WZW anomaly.

Summary and outlook

- FCD is a guide to build composite Higgs models!
- A very simple Higgs sector is possible ($SU(4)/Sp(4)$)
- Less minimal models feasible: DM? LHC signatures?
- FCD allows to probe fermion mass generations (UV completion, conformal windows...)
- FCD models can be tested on the Lattice!