

Dark Matter

I. Evidence

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São Paulo



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South American Institute for Fundamental Research

*School on DM and neutrinos
ICTP-SAIFR, July 23, 2018*

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- Let us have some practical examples of scenarios (model building : axions, SUSY WIMPs).
- How can we look for the very nature of these new particles? (Direct and indirect searches, colliders.)

Spiral galaxies

disk dynamical structure

$$F_{grav} = F_{cent}$$

$$G_N \frac{M}{R^2} = \frac{v_c^2}{R}$$

$$v_c \propto \sqrt{\frac{M}{R}}$$



Disk is Rotation supported:
observable velocity traces enclosed mass

Spiral galaxies disk dynamical structure

$$M = M(R)$$

$$v_c = v_c(R)$$

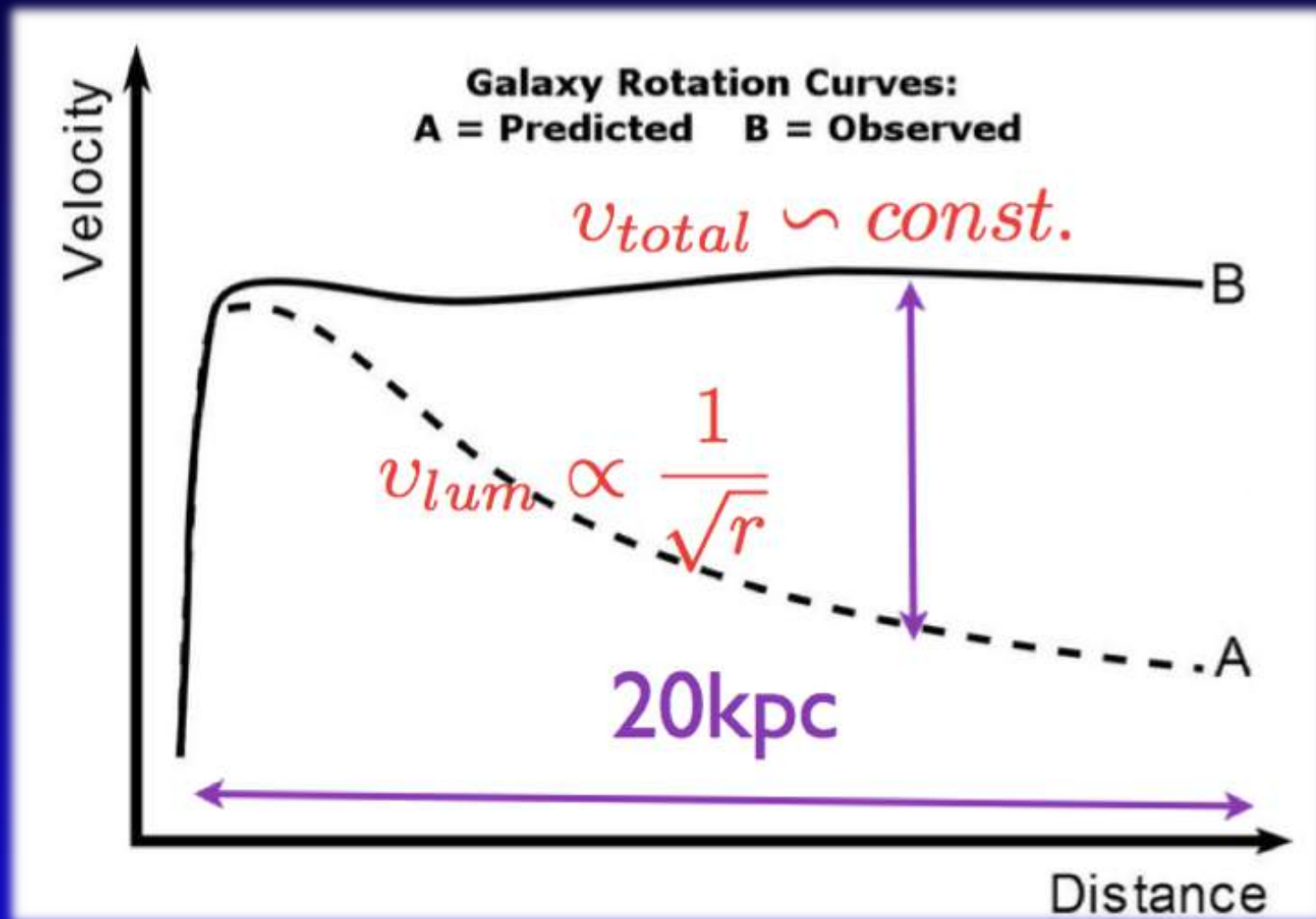
$$v_c(R) \propto \Phi(R)$$

$$v_c(R) \propto \sqrt{\frac{M(R)}{R}}$$



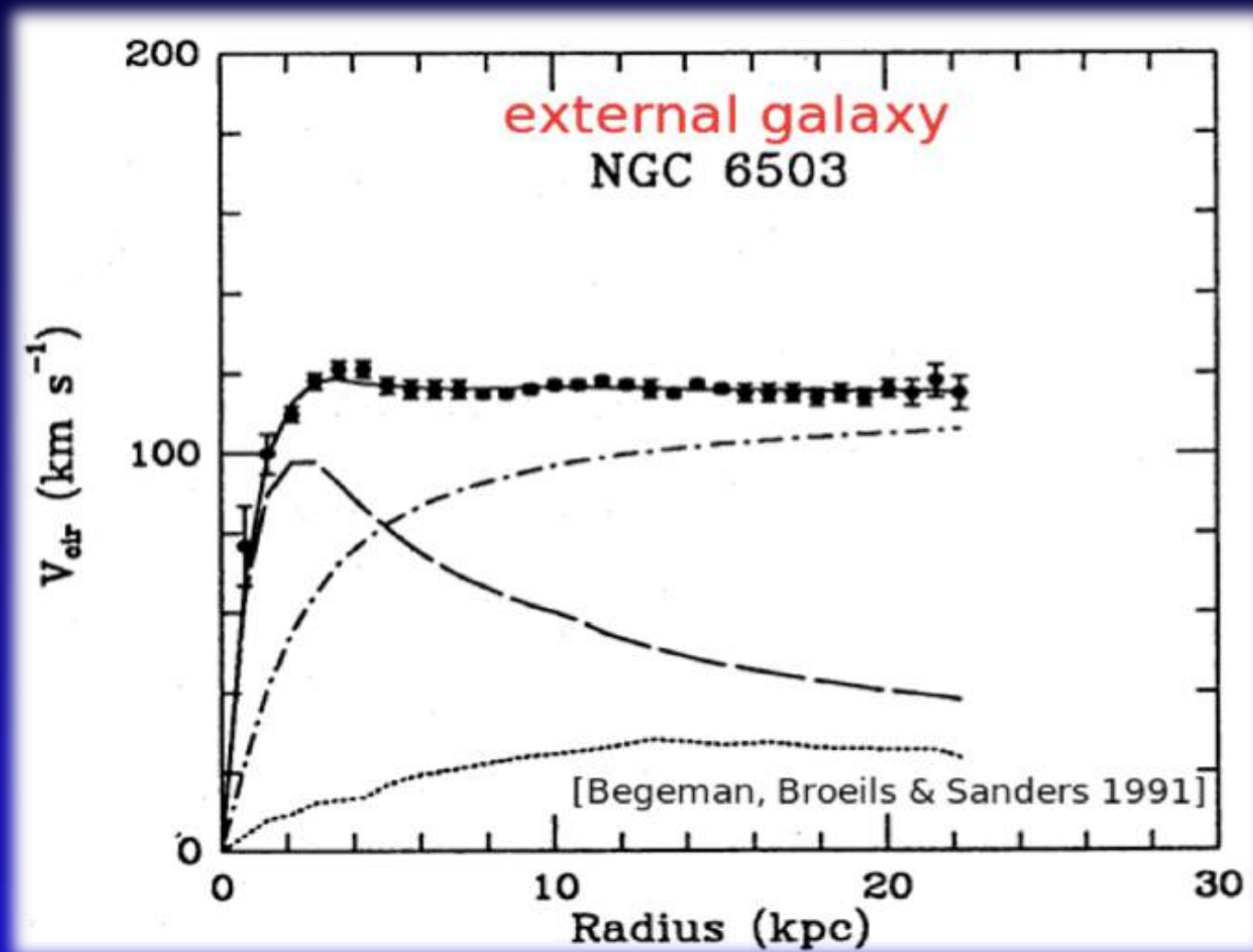
Using observed circular velocities to infer the potential
(total enclosed mass)

Rotation Curves in local galaxies: an evergreen classic (with interesting twists)



discrepancy between observed and predicted (from visible matter only)

Rotation Curves in local galaxies: an evergreen classic



discrepancy between observed and predicted (from visible matter only)

Not only the disk: Jeans analysis

Fluid continuity equation

$$\frac{\partial f}{\partial t} + \sum_{\alpha=1}^6 \frac{\partial}{\partial w_{\alpha}} (f \dot{w}_{\alpha}) = 0$$

Collisionless Boltzmann Equation

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \vec{\nabla} f - \vec{\nabla} \Phi \cdot \frac{\partial f}{\partial \vec{v}} = 0$$

(CBE)

Jeans analysis

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \vec{\nabla} f - \vec{\nabla} \Phi \cdot \frac{\partial f}{\partial \vec{v}} = 0$$

Collisionless Boltzmann Equation

$$\nu(\vec{x}) \equiv \int f(\vec{x}, \vec{v}) d^3 \vec{v}$$

$$\bar{v}_i(\vec{x}) \equiv \frac{1}{\nu(\vec{x})} \int v_i f(\vec{x}, \vec{v}) d^3 \vec{v}$$

$$\overline{v_i v_j}(\vec{x}) \equiv \frac{1}{\nu(\vec{x})} \int v_i v_j f(\vec{x}, \vec{v}) d^3 \vec{v} = \sigma_{ij}^2 + \bar{v}_i \bar{v}_j$$

$$\nu \frac{\partial \bar{v}_j}{\partial t} + \bar{v}_i \nu \frac{\partial \bar{v}_j}{\partial x_i} = -\nu \frac{\partial \Phi}{\partial x_j} - \frac{\partial}{\partial x_i} (\nu \sigma_{ij}^2) \quad (j = 1, 2, 3)$$

acceleration + kinematic viscosity = gravity + pressure

*Jeans analysis:
a practical example in a specific case
cylindrical symmetry*

$$\frac{1}{\nu} \frac{\partial}{\partial z} (\nu \sigma_{zz}^2) = -\frac{\partial \Phi}{\partial z} = -2\pi G \Sigma(z) \quad \Sigma(z) = \int_{-z}^z \rho(z') dz'$$

$$\Sigma(z) = -\frac{\sigma_{zz}^2}{2\pi G} \left(\frac{d}{dz} \ln \nu \right)$$

$$h^{-1} \equiv \frac{d}{dz} \ln \nu$$

From observations

$$\sigma_{zz}^2 = (20 \text{ km/s})^2$$

$$h = 300 \text{ pc}$$



$$\Sigma \approx 50 M_{\odot} \text{ pc}^{-2}$$

Some words about Gravitational lensing

Gravitational lensing geometry

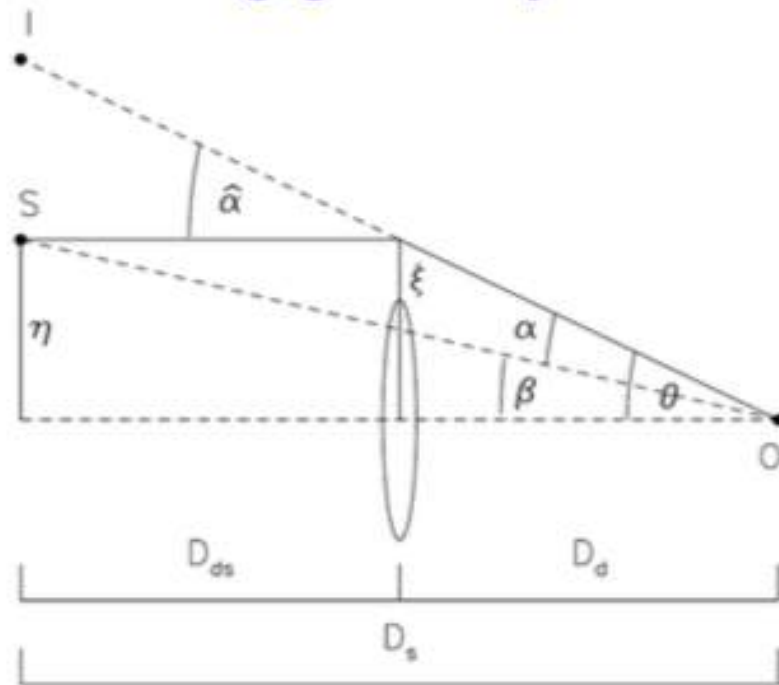


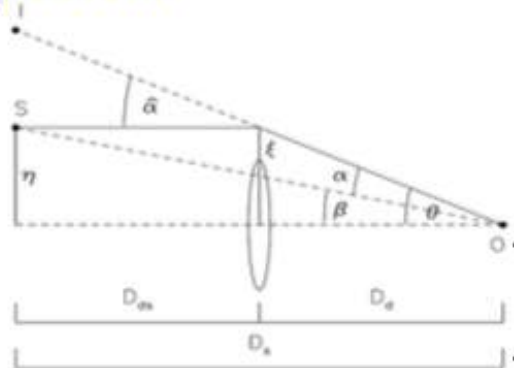
Figure from Narayan & Bartelmann (1996; arXiv:astro-ph/9606001)

► Parameters:

- Angular diameter distances: D_s , D_d and D_{ds}
- Source position: $\beta = \eta/D_s$ - *Impact parameter*
- Image position: $\theta = \xi/D_d$
- Deflexion angle: $\alpha = \hat{\alpha} D_{ds}/D_s$

Some words about Gravitational lensing

The Lensing Equation



$$\eta = \frac{D_s}{D_d} \xi - D_{ds} \hat{\alpha}(\xi)$$

- ▶ Using the angular quantities instead of the physical ones we get:

$$\beta = \theta - \frac{D_{ds}}{D_s} \hat{\alpha}(D_d \theta)$$

- ▶ Then by using the *physical* deflexion angle we arrive to the **Lens Equation**:

$$\beta = \theta - \alpha(\theta)$$

Some words about Gravitational lensing

The deflection angle

- ▶ From GR one can estimate the deflection angle due to a mass point:

$$\hat{\alpha} = \frac{4 G M}{c^2 \xi}$$

- ▶ The projected angle is then:

$$\alpha = |\alpha| = \frac{D_{ds}}{D_s} \frac{4 G M}{c^2 D_d |\theta|}$$

- ▶ Considering now a direction, as all the angles lies in the image-lens direction, we have:

$$\alpha = \frac{D_{ds}}{D_s} \frac{4 G M}{c^2 D_d |\theta|} \frac{\theta}{|\theta|}$$

The Einstein angle

- ▶ The lens equation can be easily solved in this case.
- ▶ Lets first define a 'natural scale', the *Einstein angle* for this lens:

$$\theta_E = \sqrt{\frac{4GM}{c^2} \frac{D_{ds}}{D_s D_d}}$$

- ▶ The lens equation then becomes:

$$\beta = \theta - \frac{\theta_E^2}{\theta}$$

- ▶ It has two solutions:

$$\theta_{\pm} = \frac{1}{2} \left(\beta \pm \sqrt{\beta^2 + 4\theta_E^2} \right)$$

The Einstein angle

- ▶ What is 'natural' about the *Einstein angle or radius* ?



- ▶ If $\beta = 0 \rightarrow \theta = \theta_E$

- ▶ The distance between two images (solutions) is $2\theta_E$
- ▶ The density inside the the Einstein radius is the critical lensing density

$$\Sigma_{cr} = \frac{c^2}{4\pi G} \frac{D_s}{D_d D_{ds}}$$

above which multiple imaging occur.

Lensing regimes

The lensing effect

- ▶ Gravitational lensing causes several effects on the images of the sources:
 1. Radial displacement
 2. Multiple imaging (given certain conditions)
 3. Magnification of the angular size
surface brightness conservation → and flux
 4. Distortion
 5. Time Delay
- ▶ The prevalence/interest of one or more of those effects over the others have to do with the *lensing regime*.

Lensing regimes

Lensing regimes

- ▶ Different lensing regimes differentiate themselves by:
 1. The lens mass distribution
 2. The distances involved
 3. The impact parameter
- ▶ The main lensing regimes are:
 1. Micro lensing $\theta \leq \theta_E$ *dist.* $\sim kpc$
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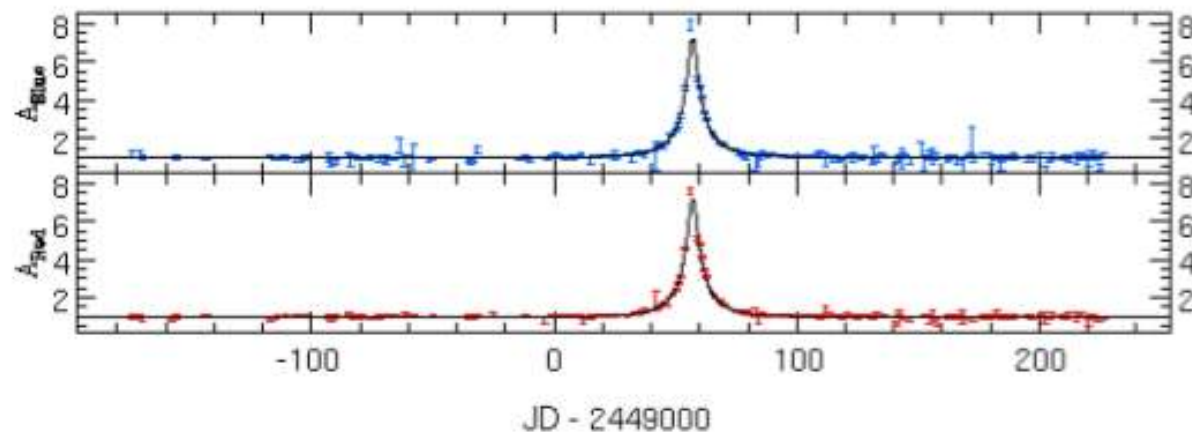
Micro Lenses

Micro lenses

- ▶ Typical *Einstein radius*

$$\theta_E = 0.902 \text{mas} \left(\frac{M}{M_\odot} \right)^{1/2} \left(\frac{D_d}{10 \text{kpc}} \right)^{-1/2} \left(1 - \frac{D_d}{D_s} \right)^{-1/2}$$

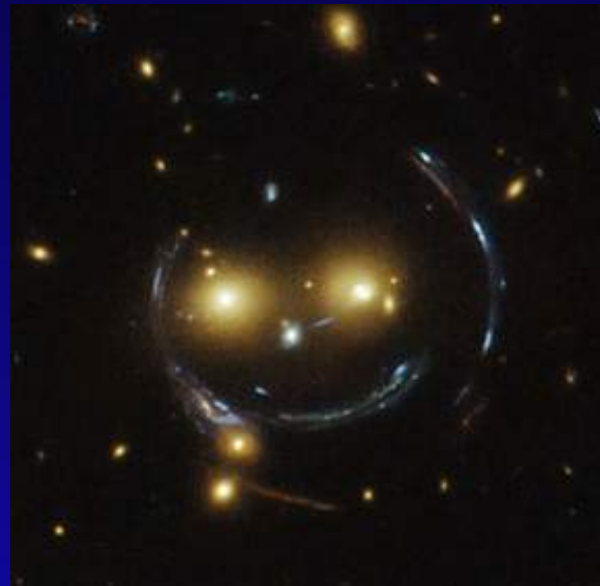
- ▶ We cannot observationally resolve the multiple imaging, but we can measure the variation of the flux due to the magnification



- ▶ Note that microlensing should produce symmetric and achromatic light curves as above.

Weak Lenses

Strong Lenses



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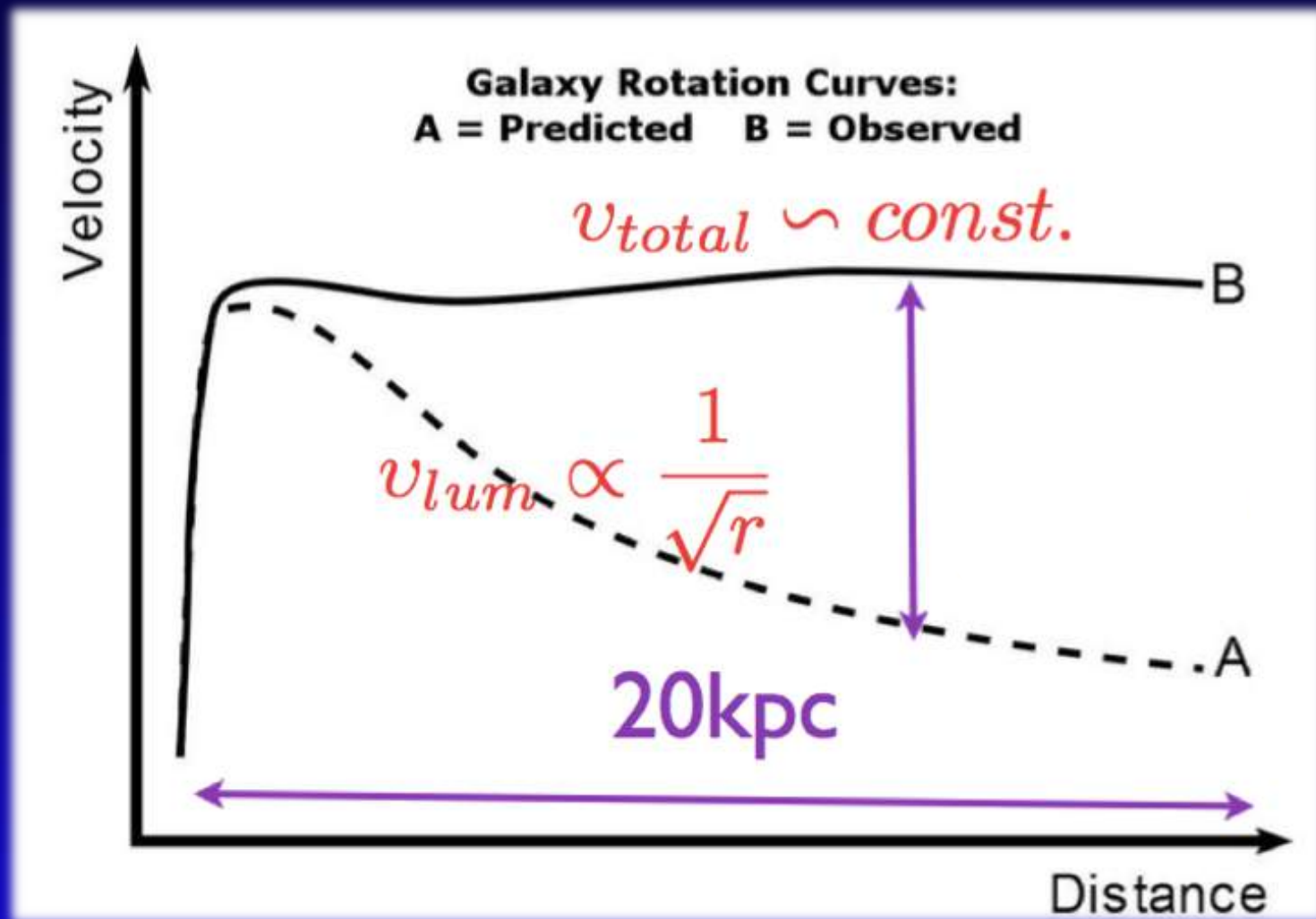
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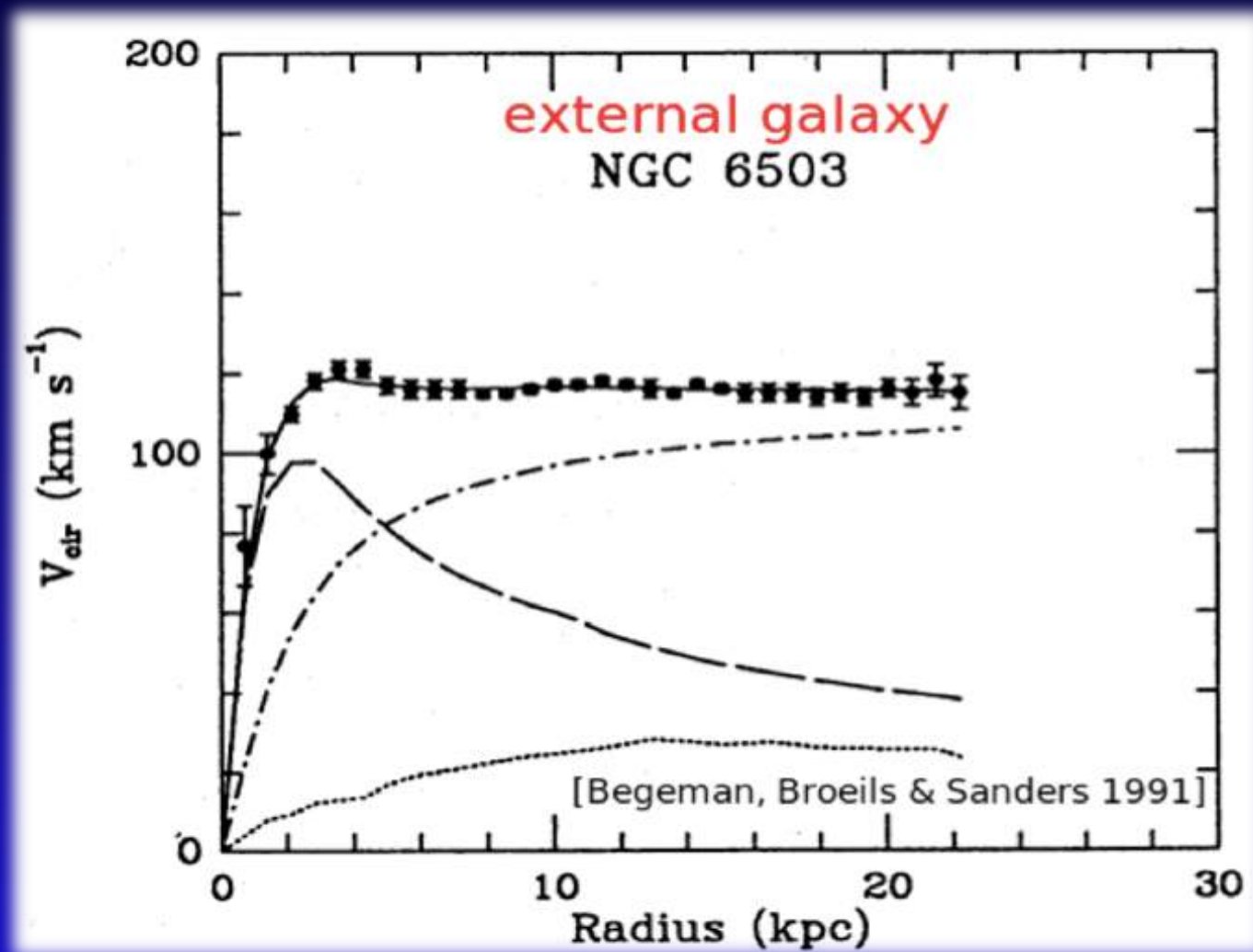
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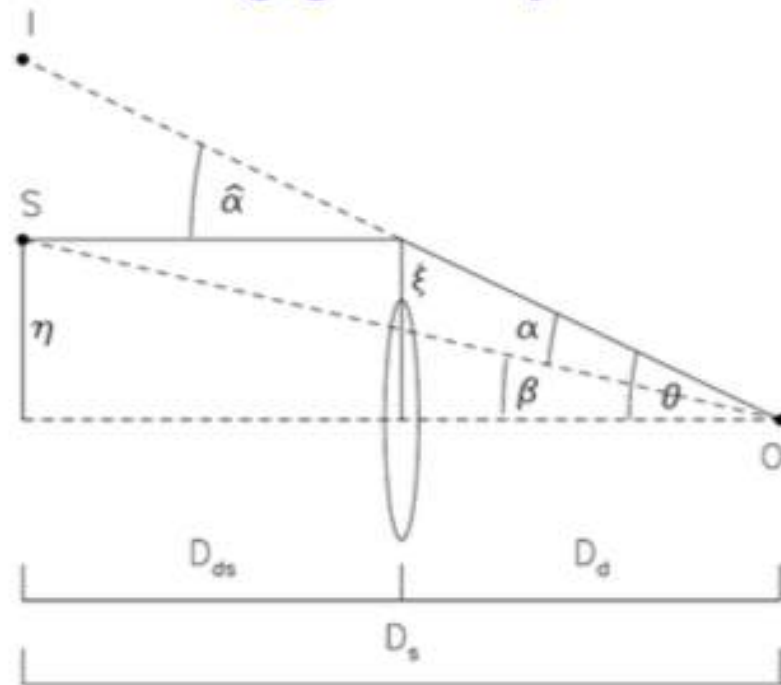


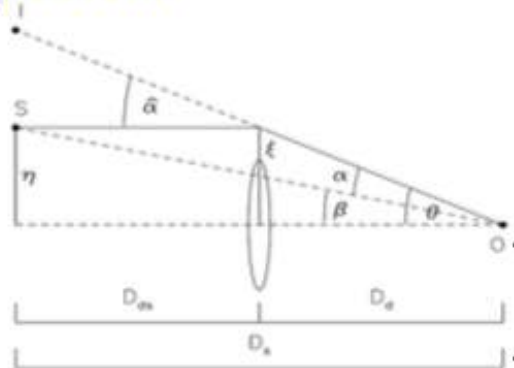
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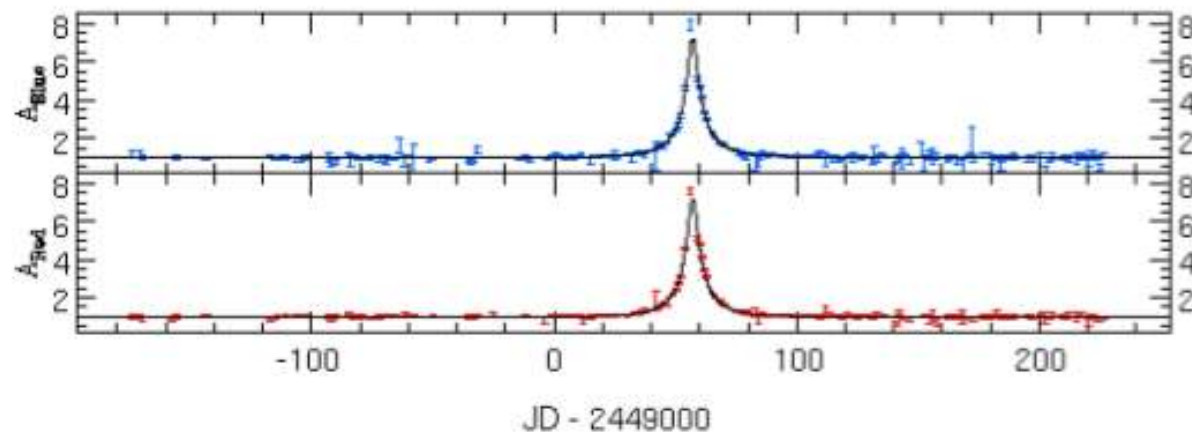
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