# Mesoscopic Physics of Photons

#### Eric Akkermans

**Photo Archive** 

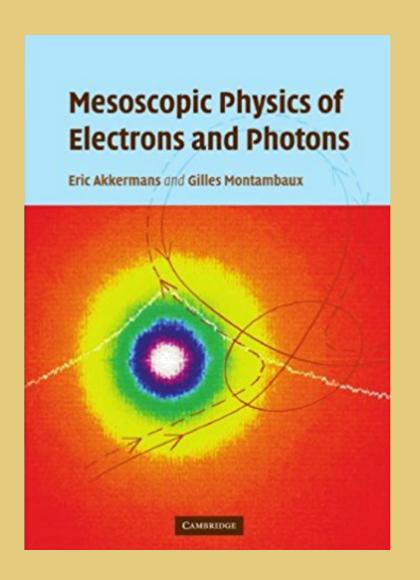






School on Interaction of Light with Cold Atoms, Sept. 16-27, 2019, Sao Paulo, Brazil, ICTP-SAIFR/IFT-UNESP.

Based on *Mesoscopic physics of electrons and photons*, by Eric Akkermans and Gilles Montambaux, Cambridge University Press, 2007



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## Introduction to mesoscopic physics

• The Aharonov-Bohm effect in disordered conductors.

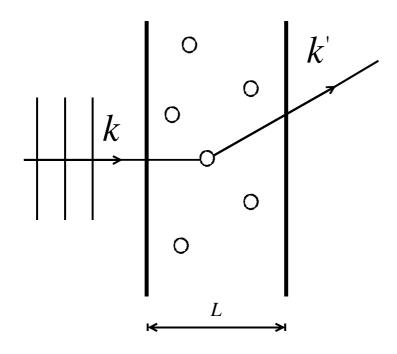
- The Aharonov-Bohm effect in disordered conductors.
- Phase coherence and effect of disorder.

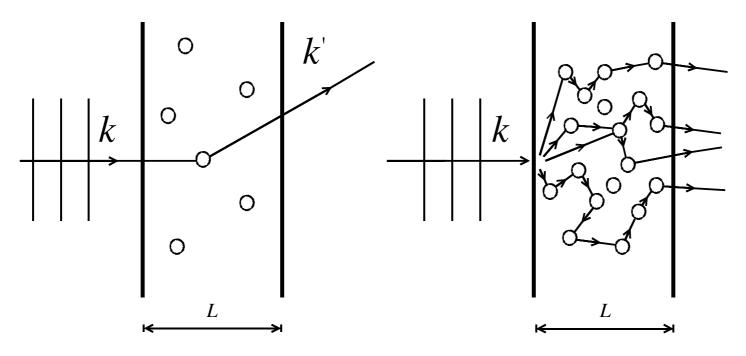
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- Average coherence: Sharvin<sup>2</sup> effect and coherent backscattering.

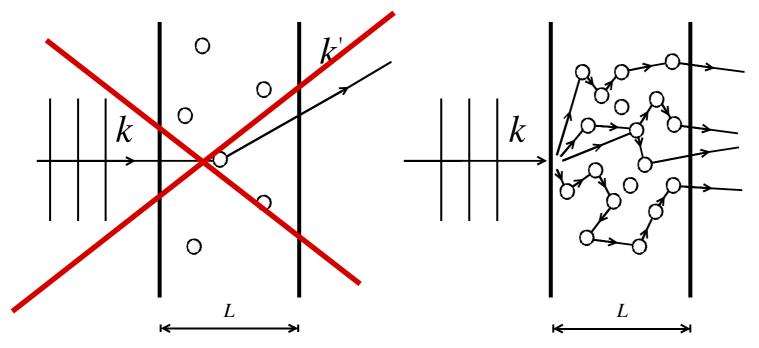
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- Phase coherence and self-averaging: universal fluctuations.

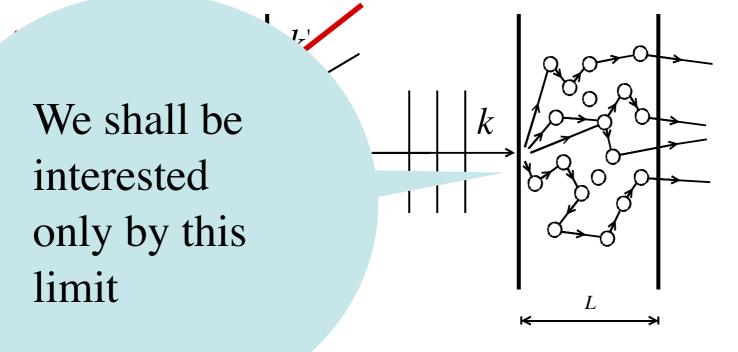
- The Aharonov-Bohm effect in disordered conductors.
- Phase coherence and effect of disorder.
- Average coherence: Sharvin<sup>2</sup> effect and coherent backscattering.
- Phase coherence and self-averaging: universal fluctuations.
- Classical probability and quantum crossings.

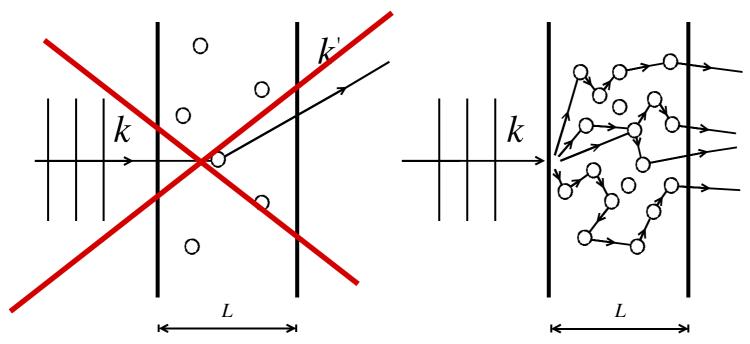
#### The framework:







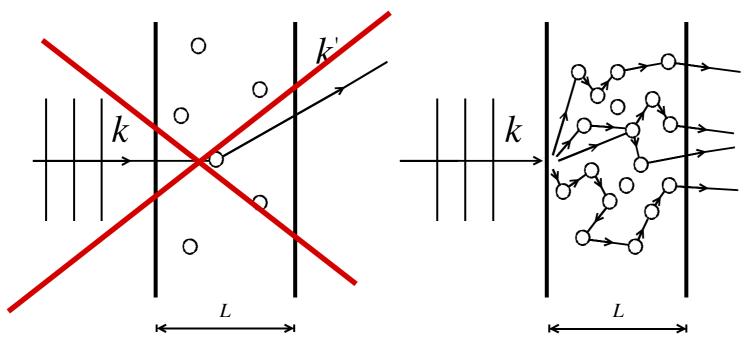




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Wavelength:  $\lambda_F = k_F^{-1}$ 

Elastic mean free path: *l* 



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Weak disorder  $\lambda_F \ll l$ : independent scattering events

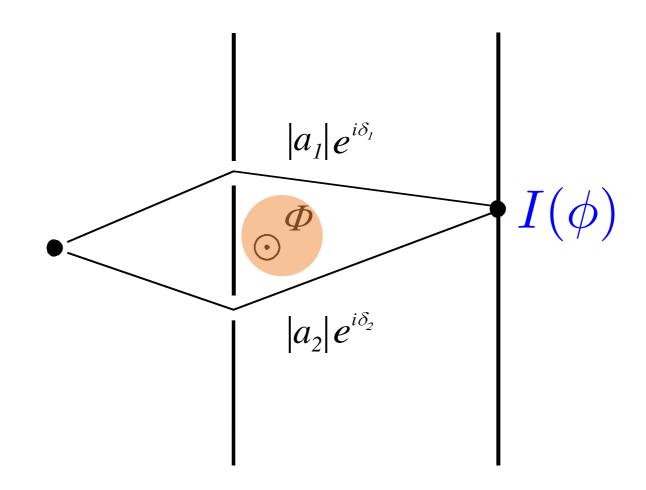
# A "canonical" mesoscopic effect

# The Aharonov-Bohm effect

Aharonov-Bohm (1959)

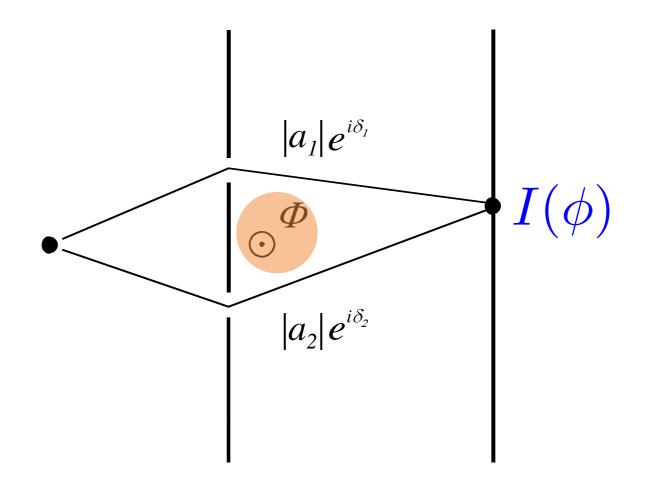
#### Aharonov-Bohm effect

No magnetic field on the electrons: no Lorentz force and no orbital motion.



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The quantum amplitudes  $a_{1,2} = |a_{1,2}|e^{i\delta_{1,2}}$  have phases:

$$\delta_1 = \delta_1^{(0)} - \frac{e}{\hbar} \int_1 \mathbf{A} \cdot d\mathbf{l}$$
 and  $\delta_2 = \delta_2^{(0)} - \frac{e}{\hbar} \int_2 \mathbf{A} \cdot d\mathbf{l}$ 

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where  $\phi_0 = h/e$  is the quantum of magnetic flux.

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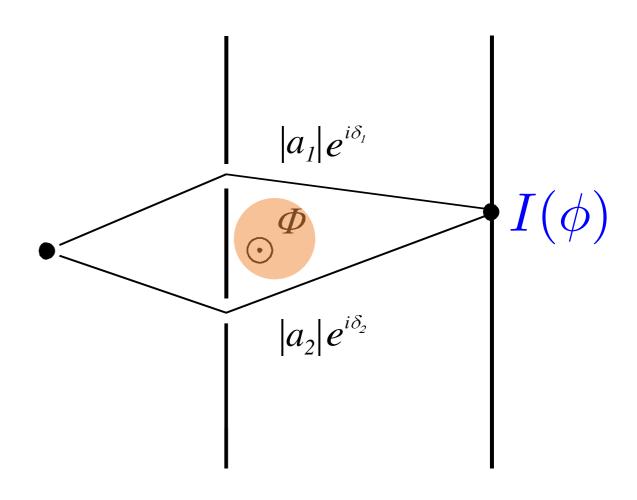
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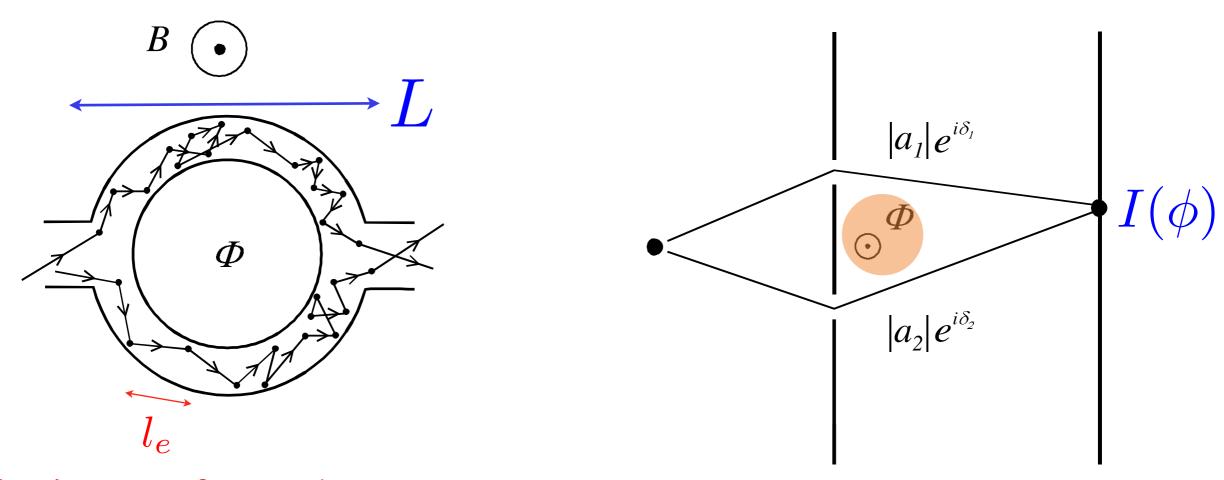
There is a continuous change of the state of interference:

Aharonov-Bohm effect (1959).

Implementation in metals : the conductance  $G(\phi)$  is the analog of the intensity.

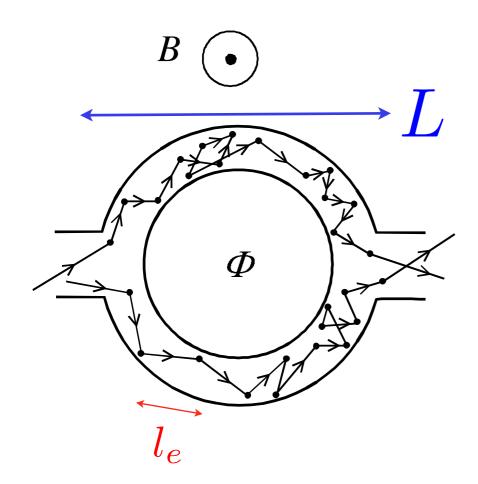


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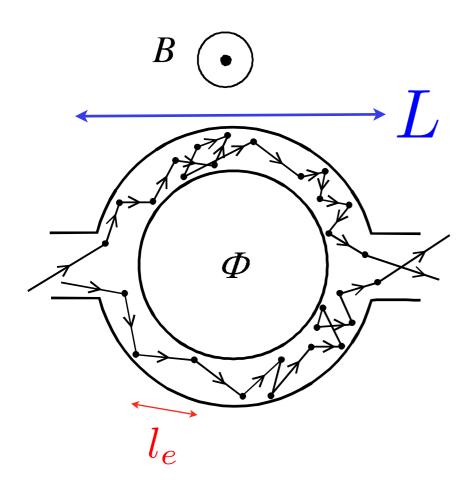
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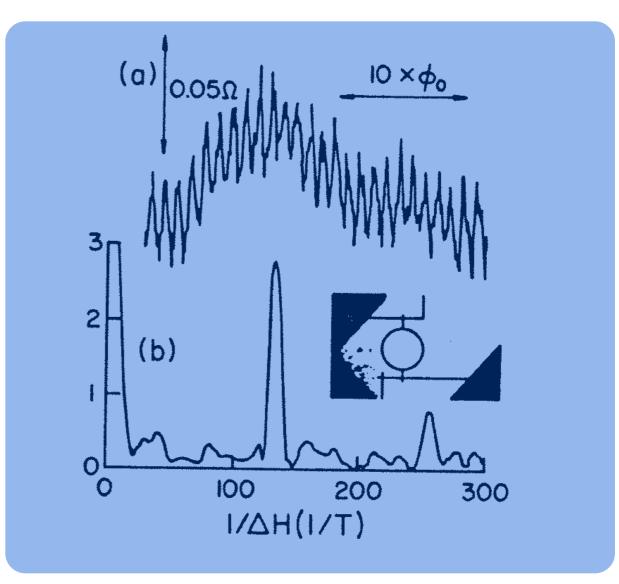


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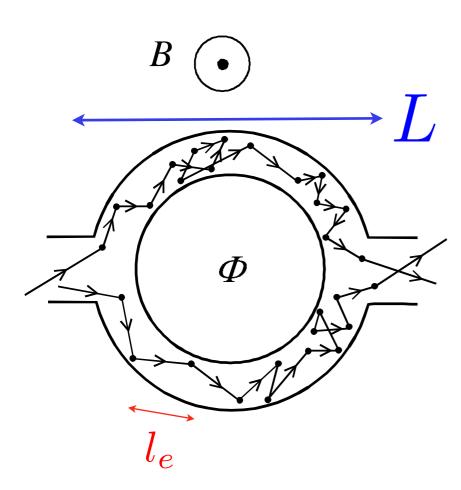
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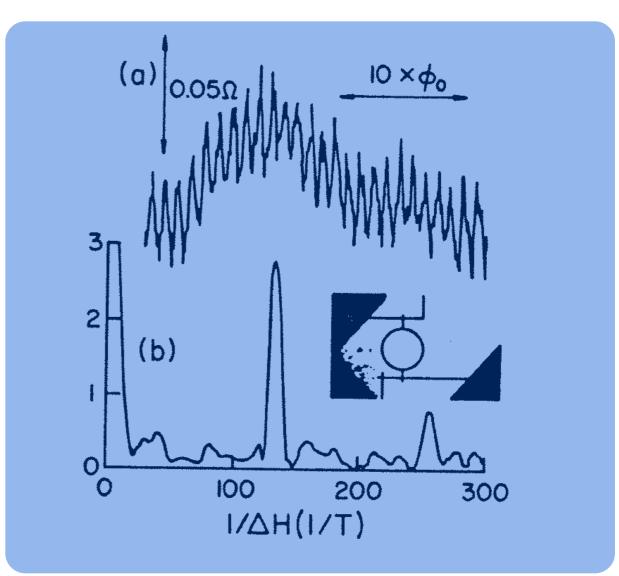
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elastic mean free path



$$G(\phi) = G_0 + \delta G \cos(\Delta \delta^{(0)}) + 2\pi \frac{\phi}{\phi_0}$$
 Webb et al. 1985

Phase coherent effects subsist in disordered metals.

Reconsider the Drude theory.

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It must exist a characteristic length  $L_\phi$  called phase coherence length beyond which all coherent effects disappear.

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The observation of coherent effects requires



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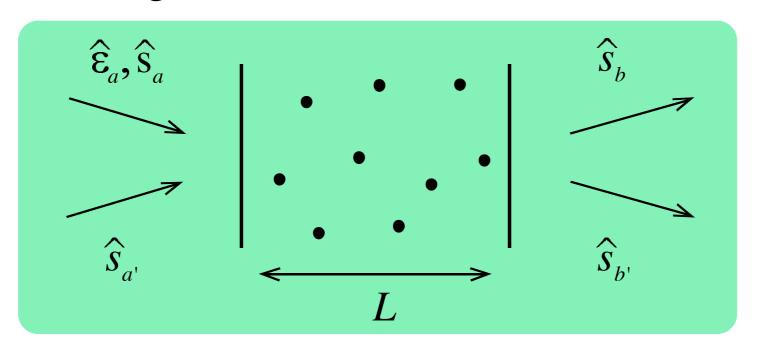
$$(G(\phi)) = G_0$$

Disorder seems to erase coherent effects....

# Formulate the same question: disorder vs. coherent effects in optics

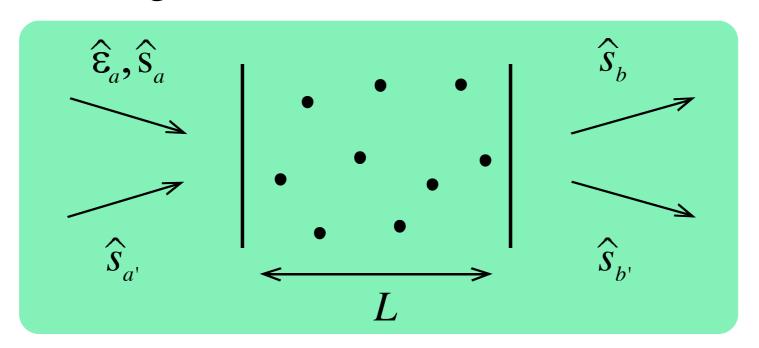
#### An analogous problem: Speckle patterns in optics

Consider the elastic multiple scattering of light transmitted through a fixed disorder configuration.

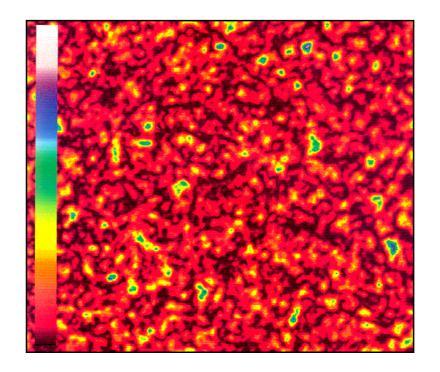


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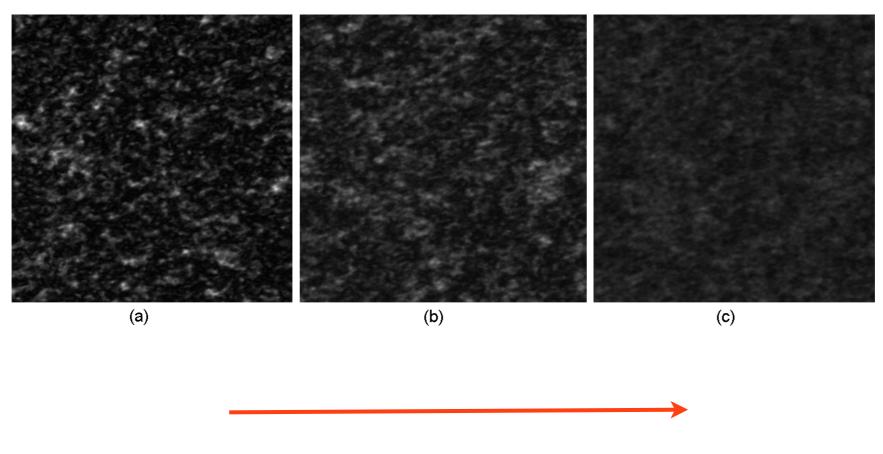


Outgoing light builds a speckle pattern i.e., an interference picture:



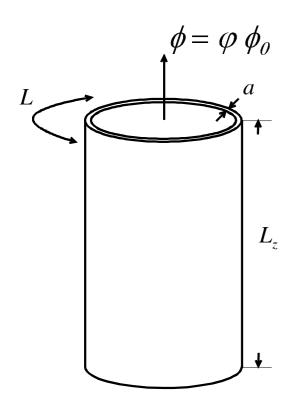
### Averaging over disorder erases the speckle pattern:

Integration over the motion of the scatterers leads to self-averaging

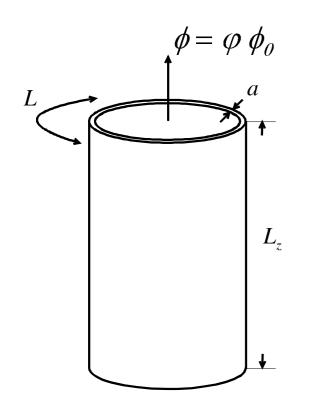


Time averaging

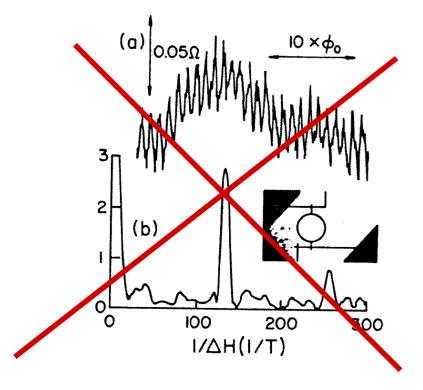
# There is an equivalent for the Aharonov-Bohm effect



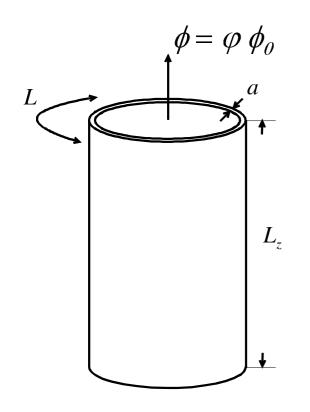
Experiment analogous to that of *Webb* but performed on a hollow cylinder of height larger than  $L_{\phi}$  pierced by a Aharonov-Bohm flux. Ensemble of rings identical to those of *Webb* but incoherent between themselves.



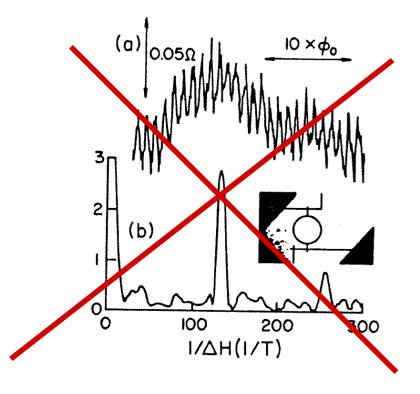
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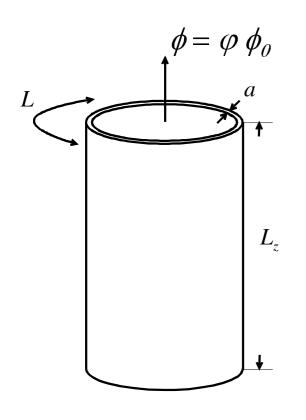
The signal modulated at  $\phi_0$  disappears



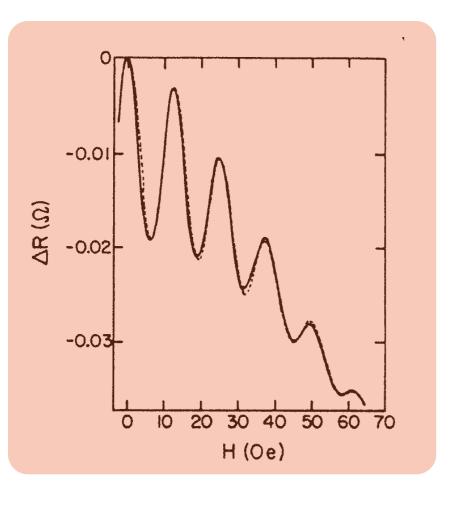
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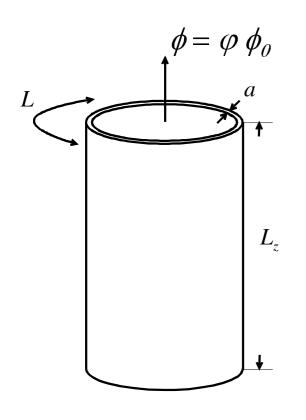
The signal modulated at  $\phi_0$  disappears but, instead, it appears a new contribution modulated at  $\phi_0/2$ 



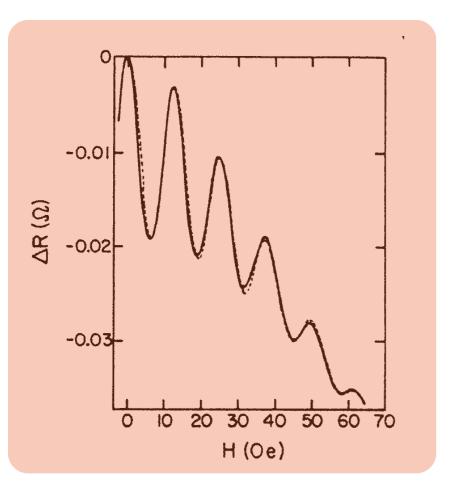
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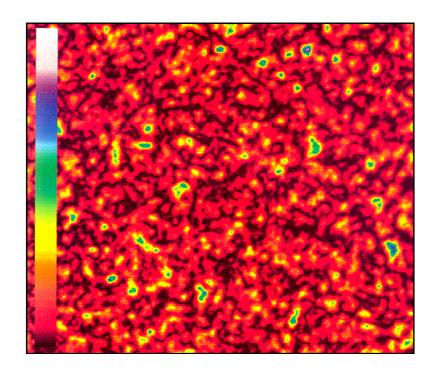
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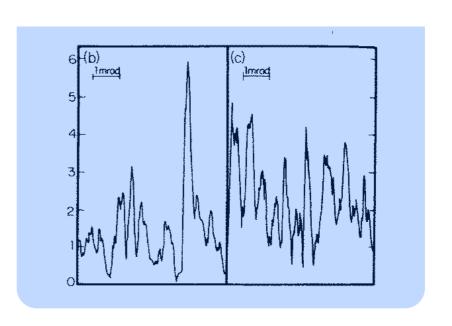


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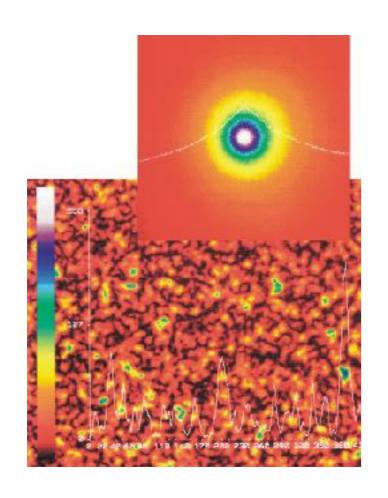
After all, disorder does not seem to erase coherent effects, but to modify them....

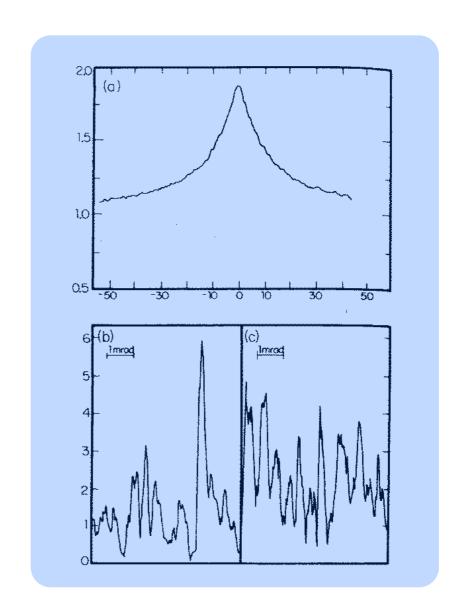
## What about speckle patterns?





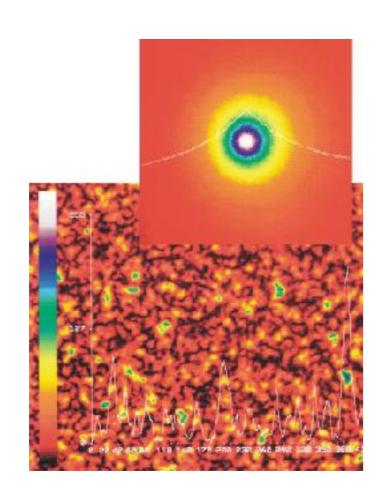
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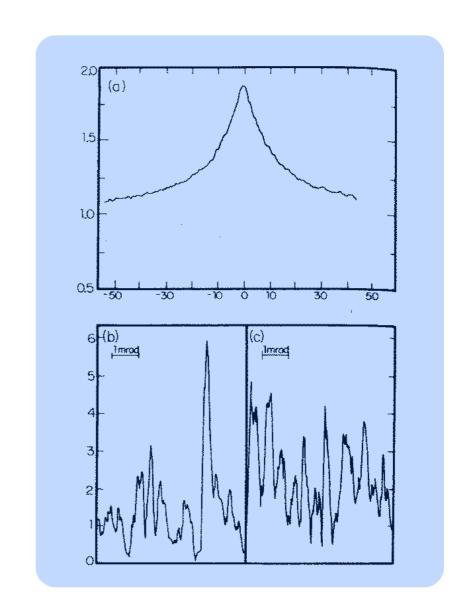




Averaging over disorder does not produce incoherent intensity only, but also an angular dependent part, the <u>coherent backscattering</u>, which is a coherence effect. We may conclude:

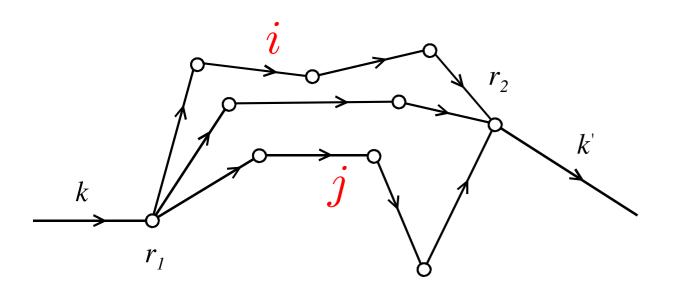
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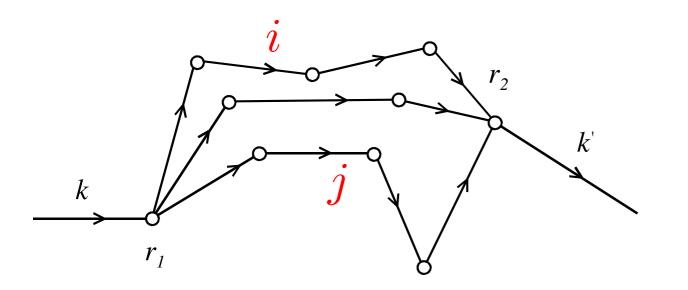


Averaging over disorder does not produce incoherent intensity only, but also an angular dependent part, the <u>coherent backscattering</u>, which is a coherence effect. We may conclude:

Elastic disorder is not related to decoherence: disorder does not destroy phase coherence and does not introduce irreversibility.

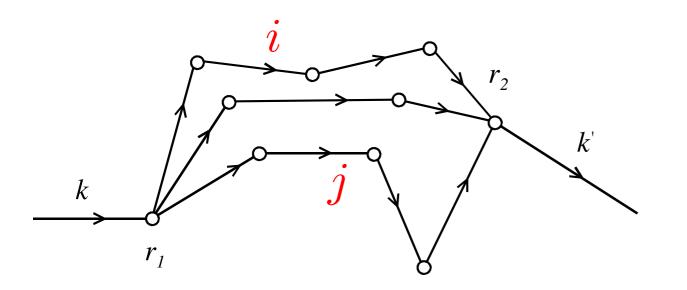


Complex amplitude  $A(\mathbf{k}, \mathbf{k}')$  associated to the multiple scattering of a wave (electron or photon) incident with a wave vector  $\mathbf{k}$  and outgoing with  $\mathbf{k}'$ 



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the complex amplitude  $f(\mathbf{r_1}, \mathbf{r_2}) = \sum_j |a_j| e^{i\delta_j}$  describes the propagation of the wave between  $\mathbf{r_1}$  and  $\mathbf{r_2}$ .

$$|A(\mathbf{k}, \mathbf{k}')|^2 = \sum_{\mathbf{r_1}, \mathbf{r_2}} \sum_{\mathbf{r_3}, \mathbf{r_4}} f(\mathbf{r_1}, \mathbf{r_2}) f^*(\mathbf{r_3}, \mathbf{r_4}) e^{i(\mathbf{k}.\mathbf{r_1} - \mathbf{k}'.\mathbf{r_2})} e^{-i(\mathbf{k}.\mathbf{r_3} - \mathbf{k}'.\mathbf{r_4})}$$

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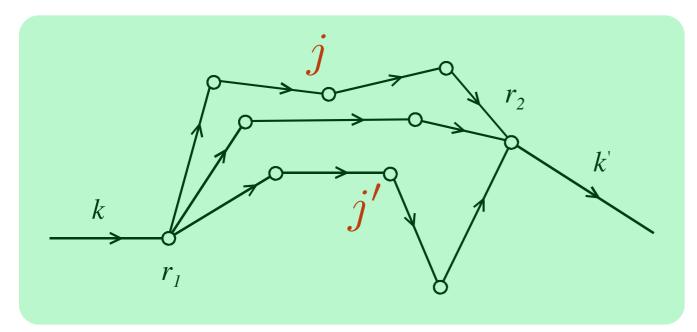
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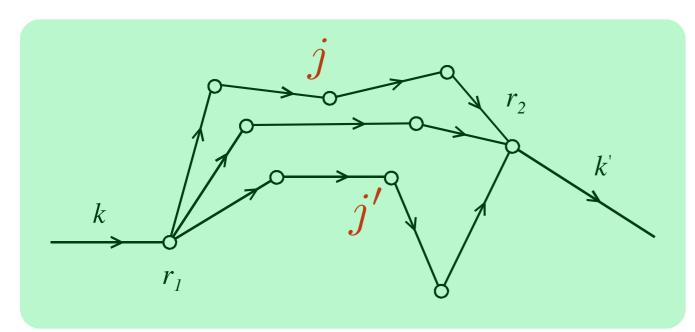


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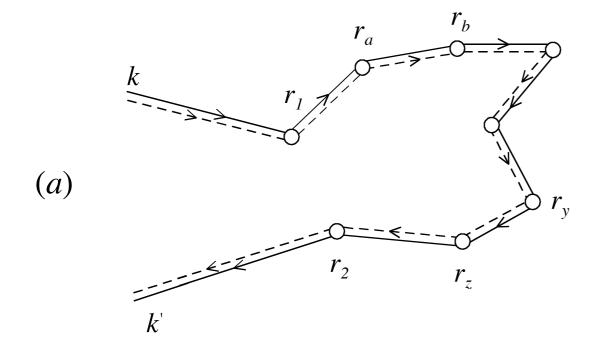
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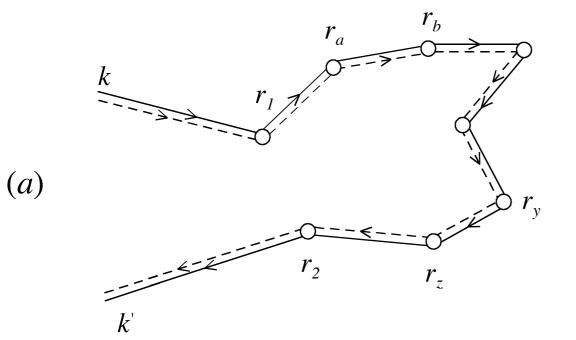


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The only remaining contributions to the intensity correspond to terms with zero dephasing, i.e., to identical trajectories.

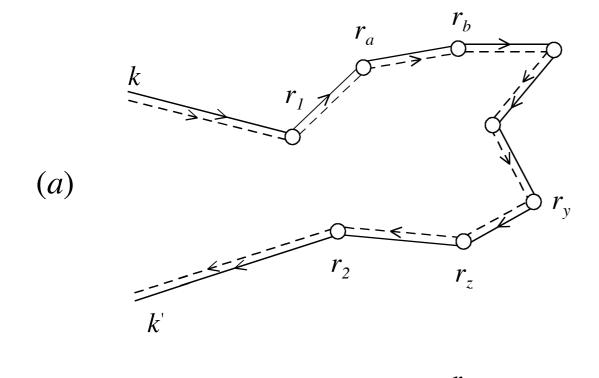


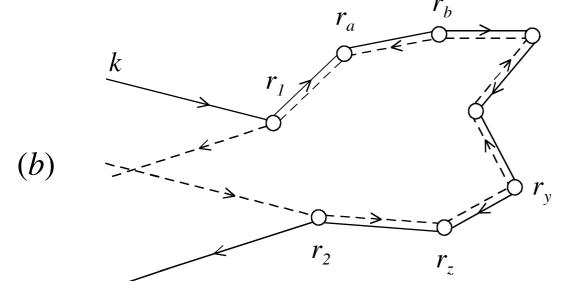
$$\mathbf{r_1} \rightarrow \mathbf{r_a} \rightarrow \mathbf{r_b} \cdots \rightarrow \mathbf{r_y} \rightarrow \mathbf{r_z} \rightarrow \mathbf{r_2}$$



$$r_1 \rightarrow r_a \rightarrow r_b \cdots \rightarrow r_y \rightarrow r_z \rightarrow r_2$$

If I see you, then you see me.

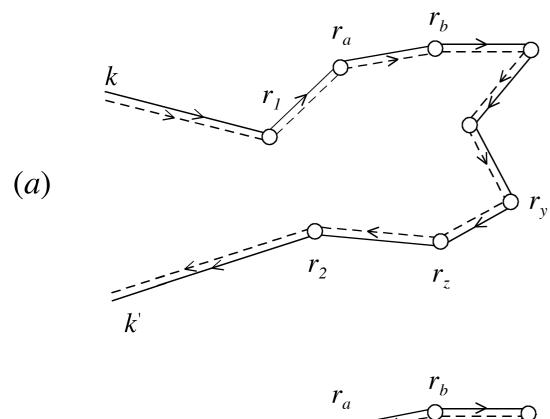


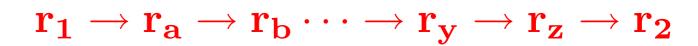


$$r_1 \rightarrow r_a \rightarrow r_b \cdots \rightarrow r_y \rightarrow r_z \rightarrow r_2$$

If I see you, then you see me.

$$\mathbf{r_2} \rightarrow \mathbf{r_z} \rightarrow \mathbf{r_y} \cdots \rightarrow \mathbf{r_b} \rightarrow \mathbf{r_a} \rightarrow \mathbf{r_1}$$





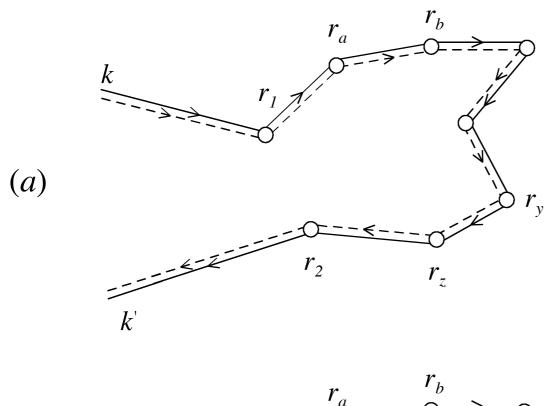
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$$(b) \qquad \qquad r_{1} \qquad \qquad r_{2} \qquad \qquad r_{3} \qquad \qquad r_{4} \qquad \qquad r_{5} \qquad \qquad r_{5}$$

The total average intensity is:

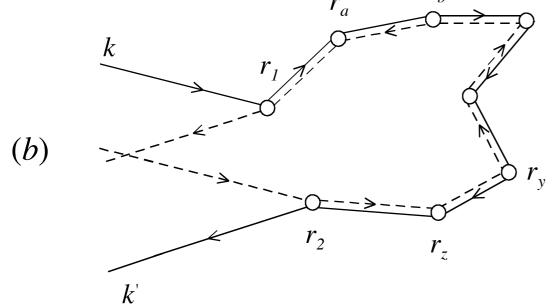
$$\overline{|A(\mathbf{k}, \mathbf{k}')|^2} = \sum_{\mathbf{r_1}, \mathbf{r_2}} |f(\mathbf{r_1}, \mathbf{r_2})|^2 \left[1 + e^{i(\mathbf{k} + \mathbf{k}') \cdot (\mathbf{r_1} - \mathbf{r_2})}\right]$$



$$\mathbf{r_1} 
ightarrow \mathbf{r_a} 
ightarrow \mathbf{r_b} \cdots 
ightarrow \mathbf{r_y} 
ightarrow \mathbf{r_z} 
ightarrow \mathbf{r_2}$$

If I see you, then you see me.

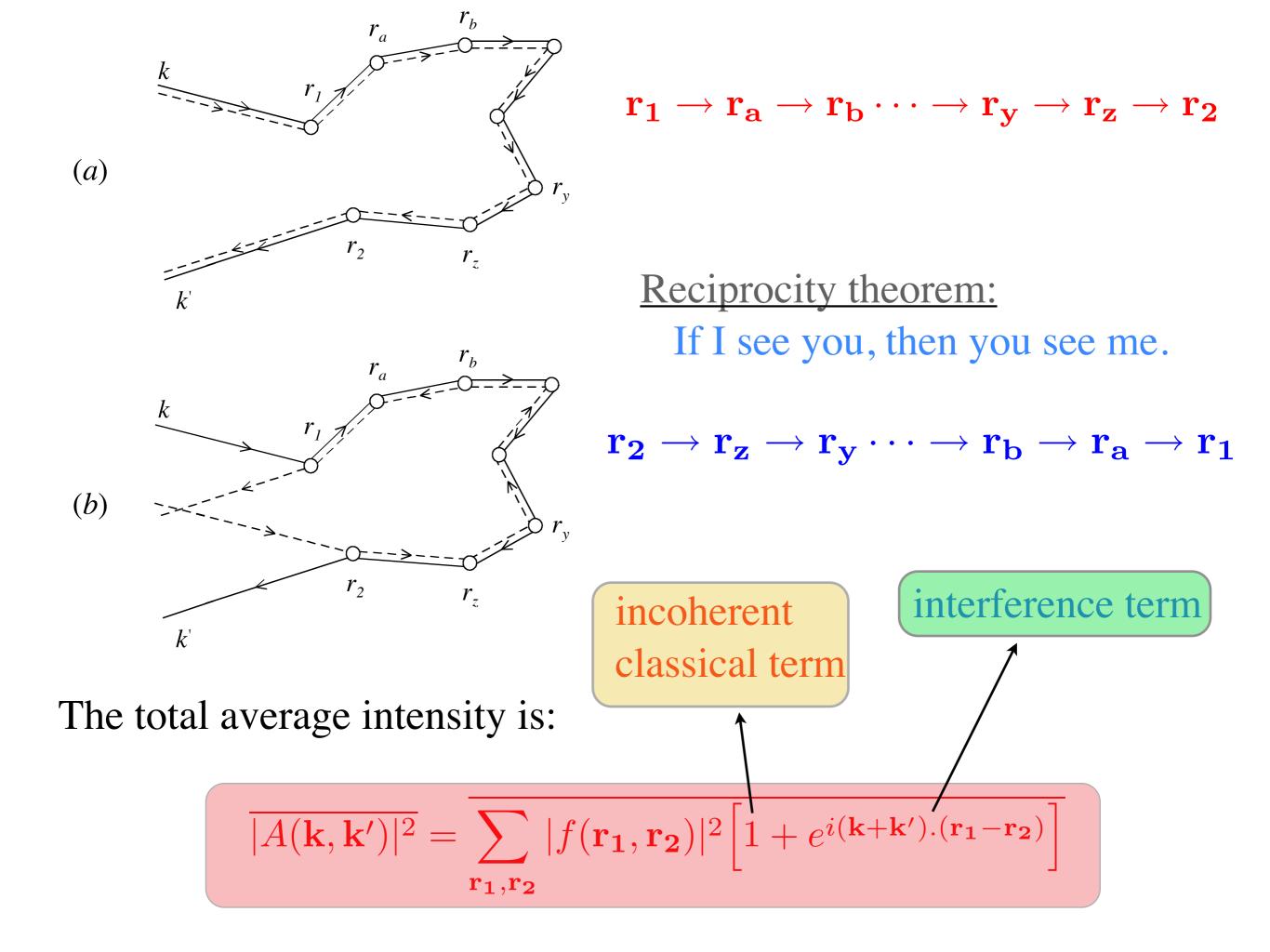
$$r_{2} \rightarrow r_{z} \rightarrow r_{y} \cdots \rightarrow r_{b} \rightarrow r_{a} \rightarrow r_{1}$$



incoherent classical term

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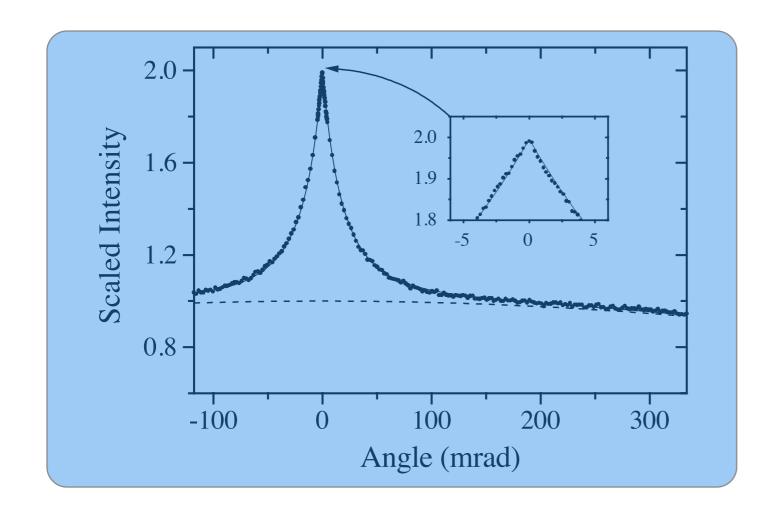
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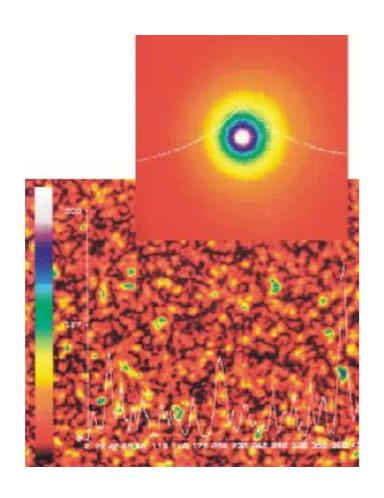
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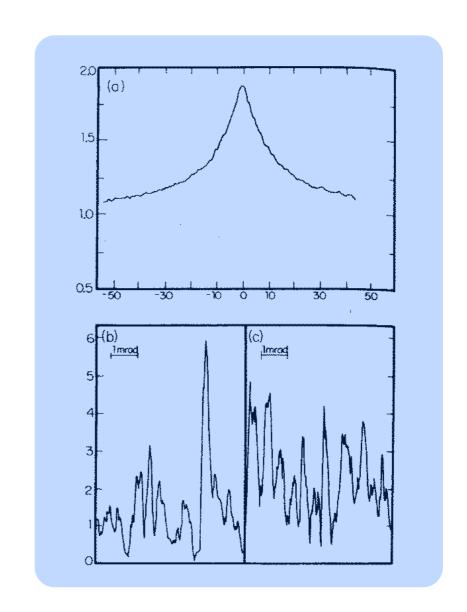
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Coherent backscattering

#### What about speckle patterns?





Averaging over disorder does not produce incoherent intensity only, but also an angular dependent part, the <u>coherent backscattering</u>, which is a coherence effect.

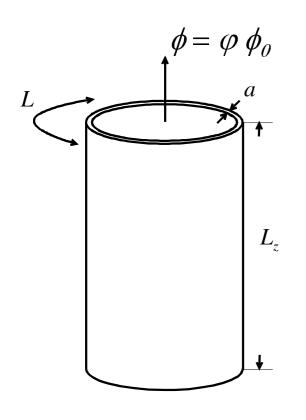
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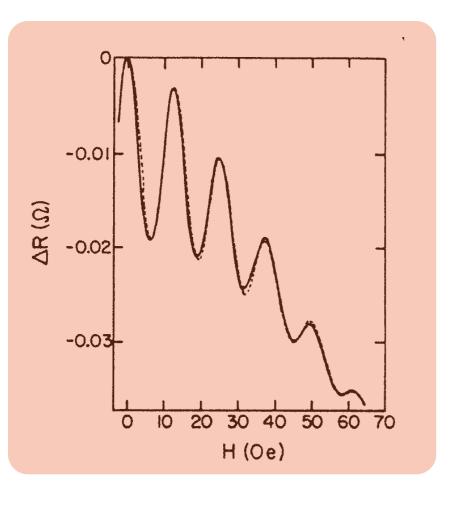
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 ${\bf r_1} - {\bf r_2} \simeq 0$ : closed loops, weak localization and  $\phi_0/2$  periodicity of the Sharvin effect.

### The Sharvin<sup>2</sup> experiment



Experiment analogous to that of *Webb* but performed on a hollow cylinder of height larger than  $L_{\phi}$  pierced by a Aharonov-Bohm flux. Ensemble of rings identical to those of *Webb* but incoherent between themselves.



The signal modulated at  $\phi_0$  disappears but, instead, it appears a new contribution modulated at  $\phi_0/2$ 

### Quantum complexity

# Random quantum systems (quantum complexity)

Disorder does not break phase coherence and it does not introduce irreversibility

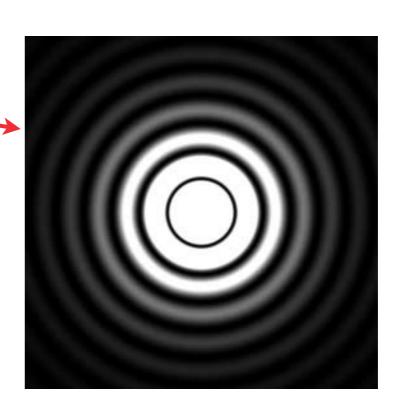
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Disorder does not break phase coherence and it does not introduce irreversibility

It introduces randomness and complexity: all symmetries are lost, there are no good quantum numbers.

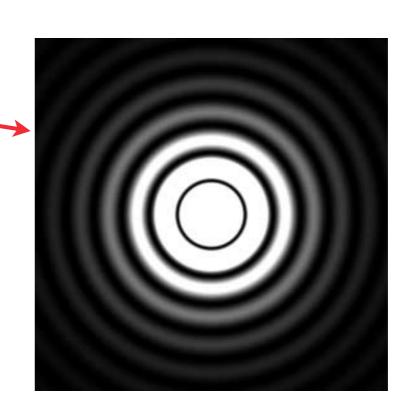
### Exemple: speckle patterns in optics

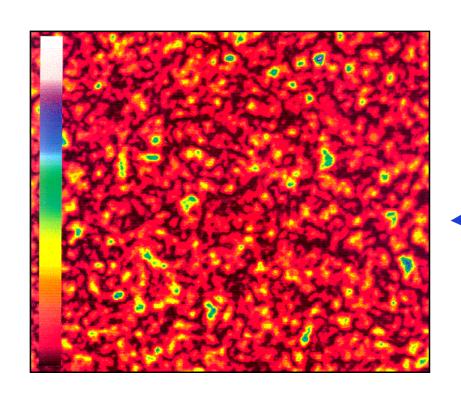
Diffraction — through a circular aperture: order in interference



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Transmission of light through a — disordered suspension: complex system

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- Decoherence: irreversible loss of quantum coherence  $L\gg L_{\varphi}$

A mesoscopic quantum system is a coherent complex quantum system with  $L \leq L_{\varphi}$ 

### An Exemple

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The system performs an average over realizations of the disorder.

For  $L \ll L_{\varphi}$  , we expect deviations from self-averaging which reflect the underlying quantum coherence.

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- a description of fluctuations and coherence in a quantum complex system.

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Due to disorder there is a finite conductance which is a quantum observable.

Classically, the conductance of a cubic sample of volume  $L^d$  is given by Ohm's law:  $G = \sigma L^{d-2}$  where  $\sigma$  is the conductivity.

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In the mesoscopic limit, the electrical conductance is not self-averaging.

