

Quantum Information & Quantum Thermodynamics

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Brazilian National Institute for Science and Technology of Quantum Information

QIS @ UFABC

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RESEARCH INTERESTS

Theoretical and experimental

- Quantum information Science
- Quantum thermodynamics
- Experimental tests-of-principles in QIS
- Entanglement and non-locality
- Open quantum systems and decoherence
- Quantum computation
- Quantum coherence aspects in biology
- Quantum metrology

Our quantum playground

Quantum Optics twin-photons & squeezed light

Trapped ions



NMR spin qubits





Quantum Thermodynamics

Computation



Quantum key distribution (cryptograph)

Quantum Simulation

Quantum Technology

Quantum

Quantum Info. Theory

Quantum Al Quantum Metrology



Logic Gates



Name			AND AB			NAND AB			OR $A+B$			$\overrightarrow{A+B}$			$\begin{array}{c} \mathbf{XOR} \\ A \oplus B \end{array}$			$\frac{\text{XNOR}}{\overline{A \oplus B}}$		
Alg. Expr.																				
Symbol																				
Truth Table	A 0 1	X 1 0	B 0 1 1	A 0 1 0 1	X 0 0 1	B 0 1 1	A 0 1 0 1	X 1 1 1 0	B 0 1 1	A 0 1 0 1	X 0 1 1 1	B 0 1 1	A 0 1 0 1	X 1 0 0 0	B 0 1 1	A 0 1 0 1	X 0 1 1 0	B 0 0 1 1	A 0 1 0 1	X 1 0 0 1





ENIAC (1946)



Frist quantum revolution microelectronics



vacuum valves

Frist transistor 1947



Single electron transistor 1987

Thermodynamics and information processing



M. P. Frank, The Physical Limits of Computing, Comput. Sci. Eng. 4 (3), 16 (2002)

Trend of minimum transistor switching energy Min transistor switching energy, kTs 1000000 -High 100000 l ow 10000 1000 trend 100 10 1995 2005 2015 2025 2035 Year of First Product Shipment

Figure 7: Trendline of minimum $\frac{1}{2}CV^2$ transistor switching energy. Values were calculated using the goals for high-performance and low-power processors in the 1999 International Technology Roadmap for Semiconductors [3]. Energy is expressed as a multiple of room-temperature kT, which also is the number of nats of information associated with that energy. If the trend is followed, thermal noise will begin to become significant in the 2030s, when transistor energies approach small multiples of kT. Shortly after 2035 (if not sooner) this trend will be *forced* to begin leveling off, since a bit of information *requires* at least 0.69 kT to carry it [37]. However, if reversible operations are used, the order-kT bit energies need not be *dissipated* [19,18], and so the dissipation per reversible bit manipulation might continue decreasing along this curve a while longer.

According Landauer the limit for the irreversible computational model is

 $-k_BT\ln 2$ $pprox 0.69 \ k_BT$ per erased bit



The ending of Moore Law

of transistors in chip doubles about ever two years (not even more!)

Single core performance





"....One never realizes experiments with a single electron or an atom or a small molecule. In thought experiments, one assumes that sometimes this is possible; invariably, this leads to ridiculous consequences....One may say that one does not realize experiments with single particles, more than one raises ichthyosaurs in the zoo."

E. Schrödinger (1952) Br. J. Philosophy Sci. **3** 109 (1952)





Single photons, single atoms, S. Haroche group (Nobel prize, 2012) Paris, France







QUANTUM COMPUTING

Devices based on quantum systems (employing subatomic physics, photons, and/or nano-technology) could make calculation far faster than conventional machines if nothing spoils the quantum weirdness.

Quantum bits (qubits) - superposition logic



Bit A classical computer encodes information in strings of bits which can take 0 or 1 value

 $|0\rangle$

 $|1\rangle$

 $|\Psi\rangle$

Qubit Quantum Bits can be encoded in two-level systems (i.e. 1/2-spin, up and down state) and can exist in a superposition of 0 and 1 simultaneously. It can be represented as a vector in the Bloch sphere.

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle$$

Quantum advantage massive parallel processing w/ quantum logic



Nature isn't classical, dammit, and if you want to make a simulation of nature, you'd better make it quantum mechanical, and by golly it's a wonderful problem, because it doesn't look so easy.

— Richard P. 7eynman —



Huge investment on quantum technology - second quantum revolution



"Technology Quarterly Here, There and Everywhere," The Economist (2017)



Brazilian National Institute for Science and Technology of Quantum Information

The

Economist

120 Principal investigators29 research institutions13 states e 25 citiesHundreds of PhD and MSc studentsDozens of Postdoctoral researchers

"Technology Quarterly Here, There and Everywhere," The Economist (2017)

Excited states

Patent applications to 2015, in:



Sources: UK Intellectual Property Office; European Commission

*By location of corporate headquarters



AIP American Institute of Physics

National Quantum Initiative Signed into Law



The enactment of the National Quantum Initiative Act on Dec. 21 creates a multiagency program spanning the National Institute of Standards and Technology, National Science Foundation, and Department of Energy. As part of the initiative, NSF and DOE will each establish between two and five competitively awarded research centers.



President Trump poses with the signed National Quantum Initiative Act. Behind him



European Commission

Quantum Technologies Flagship

The Quantum Technologies Flagship aims to place Europe at the forefront of the second quantum revolution, bringing transformative advances to science, industry and society.



THE QUANTUM LEAP

Quantum is a fundamentally new type of computer, based on the principles of quantum mechanics, that promises exponentially more power and speed than today's fastest supercomputers.





NewScientist Jan. 2019

DAILY NEWS 8 January 2019

IBM unveils its first commercial quantum computer





IBM quantum ide







Ground-to-satellite quantum teleportation & quantum cryptograph



NATURE | NEWS

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Science

China's quantum satellite clears major hurdle on way to ultrasecure communications

Probe sends entangled photons — which could underpin quantum-based data encryption — over unprecedented distance.

Davide Castelvecchi

15 June 2017

WIRED

SOPHIA CHEN SCIENCE 11.01.18 08:00 AM

QUANTUM PHYSICISTS FOUND A NEW, SAFER WAY TO NAVIGATE





Quantum magnetometer

After a series of high profile GPS hacks and failures, WIRED reported, the U.S. military and several national labs are working on new **quantum navigators** that could revolutionize global positioning systems by cutting out the need for satellite.

NV center in diamond



Let's Back to computation Easy and difficult computational tasks



Adding n digits numbers requires n elementary operations

11245 x 33456

Adding n digits numbers requires n² elementary operations

 $11245334 = p_1 \times p_2$



Factoring a n digit number requires

$$\approx \exp\left[\sqrt[3]{\frac{64}{9}\left(\ln n\right)\left(\ln\ln n\right)^2}\right] > n^k$$

elementary operations

Computational complexity

travelling salesman is a NP problem

3-sat logical clauses satisfability is a NP problem



Public-Key Cryptography





Quantum Computation

The most general state of a quantum bit (qubit)

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle$$

$$|0\rangle \rightarrow \begin{pmatrix} 1\\0 \end{pmatrix} \quad |1\rangle \rightarrow \begin{pmatrix} 0\\1 \end{pmatrix} \quad |\psi\rangle \rightarrow \begin{pmatrix} \cos\theta/2\\e^{i\varphi}\sin\theta/2 \end{pmatrix}$$
$$\vec{v}_{p} = \begin{pmatrix} x, y, z \end{pmatrix} \quad (0, 1)$$

$$x = \langle \psi | \hat{\sigma}_{x} | \psi \rangle,$$

$$y = \langle \psi | \hat{\sigma}_{y} | \psi \rangle,$$

$$z = \langle \psi | \hat{\sigma}_{z} | \psi \rangle.$$

$$\hat{\sigma}_{x} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix},$$
$$\hat{\sigma}_{y} = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix},$$
$$\hat{\sigma}_{z} = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix},$$
$$Base \left\{ |0\rangle, |1\rangle \right\}$$





Classical computational model

Logic Gates

Name			NOTAND \overline{A} AB			NAND \overline{AB}			OR $A+B$			$\frac{NOR}{\overline{A+B}}$			$\begin{array}{c} \mathbf{XOR} \\ A \oplus B \\ \hline \end{array} \\ \hline \end{array}$			XNOR		
Alg. Expr.																		$\overline{A \oplus B}$		
Symbol			A B X																	
Truth Table	A 0 1	X 1 0	B 0 0	A 0 1 0	X 0 0 0	B 0 0	A 0 1 0	X 1 1 1	B 0 0 1	A 0 1 0	X 0 1 1	B 0 0	A 0 1 0	X 1 0 0	B 0 0	A 0 1 0	X 0 1 1	B 0 0	A 0 1 0	X 1 0 0
			1	1	1	1	1	0	1	1	1	Ĩ	1	0	1	1	0	1	1	1





Qubit rotations

$$\hat{R}_{\vec{n}}\left(\theta\right) = \exp\left(-i\frac{\theta}{2}\vec{n}\cdot\hat{\vec{\sigma}}\right) = \cos\left(\frac{\theta}{2}\right)\hat{1} - i\sin\left(\frac{\theta}{2}\right)\vec{n}\cdot\hat{\vec{\sigma}}$$

$$R_{X}(\theta) \equiv \cos\left(\frac{\theta}{2}\right)\hat{1} - i\sin\left(\frac{\theta}{2}\right)\hat{X} = \begin{bmatrix} \cos\left(\frac{\theta}{2}\right) & -i\sin\left(\frac{\theta}{2}\right) \\ -i\sin\left(\frac{\theta}{2}\right) & \cos\left(\frac{\theta}{2}\right) \end{bmatrix}$$

$$R_{Y}(\theta) \equiv \cos\left(\frac{\theta}{2}\right)\hat{1} - i\sin\left(\frac{\theta}{2}\right)\hat{Y} = \begin{bmatrix}\cos\left(\frac{\theta}{2}\right) & -\sin\left(\frac{\theta}{2}\right)\\\\\sin\left(\frac{\theta}{2}\right) & \cos\left(\frac{\theta}{2}\right)\end{bmatrix}$$

$$R_{Z}(\theta) \equiv \cos\left(\frac{\theta}{2}\right)\hat{1} - i\sin\left(\frac{\theta}{2}\right)\hat{Z} = \begin{bmatrix} e^{-i\theta/2} & 0\\ 0 & e^{i\theta/2} \end{bmatrix}$$

Some identities

HXH = Z HYH = -Y HZH = X

$$H = \frac{1}{\sqrt{2}} \left(X + Z \right) = R_{Y} \left(\frac{\pi}{2} \right)$$

No-cloning theorem

It is impossible to copy unknown quantum information

Suppose that, we want to copy the state psi to a blank state of an auxiliary system

 $|\psi\rangle_{_S}\otimes|B\rangle_{_M}$.

A unitary operation can produce the copy as

$$|\psi\rangle_{S}\otimes|B\rangle_{M}\xrightarrow{U}\hat{U}(|\psi\rangle_{S}\otimes|B\rangle_{M})=|\psi\rangle_{S}\otimes|\psi\rangle_{M}.$$

A universal copy machine should work for a different state, let us say phi

$$\hat{U}(|\psi\rangle_{S}\otimes|B\rangle_{M}) = |\psi\rangle_{S}\otimes|\psi\rangle_{M},$$
$$\hat{U}(|\varphi\rangle_{S}\otimes|B\rangle_{M}) = |\varphi\rangle_{S}\otimes|\varphi\rangle_{M}.$$

What happens when we apply this unitary to a superposition state?

$$\hat{U}\Big[\Big(|\psi\rangle_{S} + |\varphi\rangle_{S}\Big) \otimes |B\rangle_{M}\Big] = \hat{U}\Big(|\psi\rangle_{S} \otimes |B\rangle_{M}\Big) + \hat{U}\Big(|\varphi\rangle_{S} \otimes |B\rangle_{M}\Big)$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\varphi\rangle_{S} \otimes |\varphi\rangle_{M}$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\varphi\rangle_{S} \otimes |\varphi\rangle_{M}$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\varphi\rangle_{S} \otimes |\varphi\rangle_{M}$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\varphi\rangle_{M} \otimes |\varphi\rangle_{M}$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\psi\rangle_{M} \otimes |\varphi\rangle_{M} = |\psi\rangle_{S} \otimes |\psi\rangle_{M}$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\psi\rangle_{M} \otimes |\psi\rangle_{M} + |\psi\rangle_{M} \otimes |\psi\rangle_{M} + |\psi\rangle_{M} \otimes |\psi\rangle_{M}$$

$$= |\psi\rangle_{S} \otimes |\psi\rangle_{M} + |\psi\rangle_{M} \otimes |\psi\rangle_{M} + |\psi\rangle_{$$

W. K. Wootters and W. H. Zurek, Nature 299, 802 (1982); D. Dieks, Phys. Lett. A 92,. 271 (1982).

Two qubit logic gates

The control not gate (equivalent to the XOR)



Computational basis
$$\{|00\rangle, |01\rangle, |10\rangle, |11\rangle$$



Example:



 $|0\rangle \langle 0| \otimes \mathbb{1} + |1\rangle \langle 1| \otimes \sigma_x$

Some quantum gates identities







One qubit operations and Cnot gates are universal



Deutsch Algorithm

quantum parallelism

D. Deutsch, Proc. R. Soc. London A 400, 97 (1985).

Let us consider the 1-bit function $f(x) : \{0,1\} \rightarrow \{0,1\}$

The function f(x) is computed by a oracle that you have no access



$$\begin{array}{l} (0) \quad |\psi_{0}\rangle_{ab} = |0\rangle_{a} |1\rangle_{b} \\ (1) \quad |\psi_{1}\rangle_{ab} = H^{\otimes 2} |0\rangle_{a} |1\rangle_{b} = \frac{1}{2} (|0\rangle_{a} + |1\rangle_{a}) \otimes (|0\rangle_{b} - |1\rangle_{b}) \\ (2) \quad |\psi_{2}\rangle_{ab} = \frac{1}{2} \hat{U}_{f(x)}^{ab} \Big[|0\rangle_{a} (_{a} |0\rangle_{b} - |1\rangle_{b}) + |1\rangle_{a} (|0\rangle_{b} - |1\rangle_{b}) \Big] \\ = \frac{1}{2} \Big[(-1)^{f(0)} |0\rangle_{a} (_{a} |0\rangle_{b} - |1\rangle_{b}) + (-1)^{f(1)} |1\rangle_{a} (|0\rangle_{b} - |1\rangle_{b}) \Big] \\ |\psi_{2}\rangle_{ab} = \begin{cases} \frac{1}{2} (-1)^{f(0)} (|0\rangle_{a} + |1\rangle_{a}) (_{a} |0\rangle_{b} - |1\rangle_{b}), & \text{if} \quad f(0) = f(1) \\ \frac{1}{2} (-1)^{f(0)} (|0\rangle_{a} - |1\rangle_{a}) (_{a} |0\rangle_{b} - |1\rangle_{b}), & \text{if} \quad f(0) \neq f(1) \end{cases}$$

$$(3) \quad |\psi_{3}\rangle_{ab} = H^{(a)} |\psi_{2}\rangle_{ab} = \begin{cases} \overline{\sqrt{2}} (-1) & |0\rangle_{a} (a|0\rangle_{b} - |1\rangle_{b}), & y = f(0) = f(1) \\ \frac{1}{\sqrt{2}} (-1)^{f(0)} |1\rangle_{a} (a|0\rangle_{b} - |1\rangle_{b}), & if = f(0) \neq f(1) \end{cases}$$

$$\left|\psi_{3}\right\rangle_{ab} = \left(-1\right)^{f(0)} \left|f(0) \oplus f(1)\right\rangle_{a} \frac{1}{\sqrt{2}} \left(\left|0\right\rangle_{b} - \left|1\right\rangle_{b}\right)$$

$$f(0) \oplus f(1) = \begin{cases} 0, & \text{if} \quad f(0) = f(1) \\ 1, & \text{if} \quad f(0) \neq f(1) \end{cases}$$

Shor's factoring algorithm

Exponential complexity reduction

Quantum communication

Superdense Coding

One-time pad cryptograph

Η

 \mathbf{V}^{j}

0

0

In the Bob's lab

$$\begin{split} \left| B_{x,y} \right\rangle_{ab} &= \frac{1}{\sqrt{2}} \left(\left| 0 \right\rangle_{a} \left| y \right\rangle_{b} + \left(-1 \right)^{x} \left| 1 \right\rangle_{a} \left| \overline{y} \right\rangle_{b} \right) \\ \left| B_{0,0} \right\rangle_{ab} &= \frac{1}{\sqrt{2}} \left(\left| 0 \right\rangle_{a} \left| 0 \right\rangle_{b} + \left| 1 \right\rangle_{a} \left| 1 \right\rangle_{b} \right) = \left| \Phi^{+} \right\rangle_{ab} \\ \left| B_{1,0} \right\rangle_{ab} &= \frac{1}{\sqrt{2}} \left(\left| 0 \right\rangle_{a} \left| 0 \right\rangle_{b} - \left| 1 \right\rangle_{a} \left| 1 \right\rangle_{b} \right) = \left| \Phi^{-} \right\rangle_{ab} \\ \left| B_{0,1} \right\rangle_{ab} &= \frac{1}{\sqrt{2}} \left(\left| 0 \right\rangle_{a} \left| 1 \right\rangle_{b} + \left| 1 \right\rangle_{a} \left| 0 \right\rangle_{b} \right) = \left| \Psi^{+} \right\rangle_{ab} \\ \left| B_{1,1} \right\rangle_{ab} &= \frac{1}{\sqrt{2}} \left(\left| 0 \right\rangle_{a} \left| 1 \right\rangle_{b} - \left| 1 \right\rangle_{a} \left| 0 \right\rangle_{b} \right) = \left| \Psi^{-} \right\rangle_{ab} \end{split}$$

$$|input\rangle_{ab} = |B_{x,y}\rangle_{ab}$$

$$= \frac{1}{\sqrt{2}} \left(|0\rangle_a |y\rangle_b + (-1)^x |1\rangle_a |\overline{y}\rangle_b \right)$$

$$\stackrel{CNOT_{ab}}{\rightarrow} \frac{1}{\sqrt{2}} \left(|0\rangle_a |y\rangle_b + (-1)^x |1\rangle_a |1 \oplus \overline{y}\rangle_b \right)$$

$$\stackrel{H_a}{\rightarrow} \frac{1}{\sqrt{2}} \left[\left(\frac{1 + (-1)^x}{\sqrt{2}} \right) |0\rangle_a |y\rangle_b + \left(\frac{1 - (-1)^x}{\sqrt{2}} \right) |1\rangle_a |1 \oplus \overline{y}\rangle_b \right]$$

$$\boxed{1 \oplus \overline{y} = y}$$

$$\begin{aligned} x &= 0 \\ |output\rangle &= |0\rangle_a |y\rangle_b = \begin{cases} |0\rangle_a |0\rangle_b, \ y &= 0 \rightarrow |input\rangle = |\Phi^+\rangle \\ |0\rangle_a |1\rangle_b, \ y &= 1 \rightarrow |input\rangle = |\Psi^+\rangle \end{cases} \\ x &= 1 \\ |output\rangle &= |1\rangle_a |y\rangle = \begin{cases} |1\rangle_a |0\rangle_b, \ y &= 0 \rightarrow |input\rangle = |\Phi^-\rangle \\ |1\rangle_a |1\rangle_b, \ y &= 1 \rightarrow |input\rangle = |\Psi^-\rangle \end{cases} \end{aligned}$$

How to implement qubits in physical systems

Traped lons

Photonic systems

Nuclear spins - NMR

Superconducting Circuit electrodynamics

Many others candidates

Open Quantum System - Master Equation two level system interacting with a Markovian heat reservoir

System environment interaction

UFABC

Initially uncorrelated

 $\rho_{SR}(0) = \rho_S(0) \otimes \rho_R(0)$

Open Quantum System - Master Equation two level system interacting with a Markovian heat reservoir

$$\begin{split} \frac{1}{t}\rho\left(t\right) &= -\frac{\imath}{\hbar}\left[H\left(t\right),\rho\left(t\right)\right] \\ &+ \gamma_{\downarrow}\left[\Gamma_{\downarrow}\rho\left(t\right)\Gamma_{\uparrow}^{\dagger} - \frac{1}{2}\left\{\rho\left(t\right),\Gamma_{\uparrow}^{\dagger}\Gamma_{\downarrow}\right\}\right] \\ &+ \gamma_{\uparrow}\left[\Gamma_{\uparrow}\rho\left(t\right)\Gamma_{\downarrow}^{\dagger} - \frac{1}{2}\left\{\rho\left(t\right),\Gamma_{\downarrow}^{\dagger}\Gamma_{\uparrow}\right\}\right] \end{split}$$

where

 $rac{a}{d}$

$$\gamma_{\downarrow} = \gamma_0 \left(N_{
m BE} + 1
ight)$$
 $\gamma_{\uparrow} =$

$$\gamma_{\uparrow} = \gamma_0 N_{
m BE}$$

$$N_{\rm BE} = \left(e^{\beta\hbar\omega} - 1\right)^{-1}$$

 γ_0

is the Bose-Einstein distribution

is the vacuum decay rate

Quantum information w/ Nuclear Spins

Liquid state NMR experimental setup

Thank you for your attention

MSc, PhD, and Postdoctoral positions in QIS available @ UFABC

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