# Active Matter Julia Yeomans University of Oxford

- 1. What is active matter and why is it interesting?
- 2. Background 1: low Reynolds number hydrodynamics
- 3. Background 2: nematic liquid crystals
- 4. Active nematics and active turbulence
- 5. Self-propelled topological defects
- 6. Confining active turbulence
- 7. Bacteria: the hare and the tortoise
- 8. Eukaryotic cells as an active system

Active matter: takes energy from its surroundings on a single particle level and uses it to do work.





cells



active colloids



microswimmers



Kinesin walking, from Inner Life of a Cell







Bacterial flagellar motor





### Active propulsion:





Di Leonardo, Sokolov,





Galajda et al, J Modern Optics 2011



heamatite

hydrogen peroxide



Pine group, New York

## Active turbulence: bacteria





Dense suspension of microswimmers

## Active turbulence: eukaryotic cells







Why study active matter:

- 1. to understand biological systems: biomechanics and self-assembly
- 2. To create new types of micro-engines
- 3. As examples of non-equilibrium statistical physics

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 $\rho \left\{ \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} \right\} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$ 



$$\rho \left\{ \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} \right\} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$$
  
inertial terms viscous terms

$$\tilde{v} = \frac{v}{V_0} \qquad \qquad \tilde{x} = \frac{x}{L_0}$$

$$\tilde{\nabla} = L_0 \nabla$$
  $\tilde{t} = \frac{V_0}{L_0} t$   $\frac{\partial}{\partial \tilde{t}} = \frac{L_0}{V_0} \frac{\partial}{\partial t}$ 

$$\frac{V_0^2}{L_0}\rho\left\{\frac{\partial\tilde{\mathbf{v}}}{\partial\tilde{t}} + (\tilde{\mathbf{v}}\cdot\tilde{\nabla})\tilde{\mathbf{v}}\right\} = -\nabla p + \frac{V_0}{L_0^2}\mu\tilde{\nabla}^2\tilde{\mathbf{v}} + \mathbf{f}$$

$$\left\{\frac{\partial \tilde{\mathbf{v}}}{\partial \tilde{t}} + (\tilde{\mathbf{v}} \cdot \tilde{\nabla})\tilde{\mathbf{v}}\right\} = -\frac{L_0}{V_0^2 \rho} \nabla p + \frac{\mu}{L_0 V_0 \rho} \tilde{\nabla}^2 \tilde{\mathbf{v}} + \frac{L_0}{V_0^2 \rho} \mathbf{f}$$

$$Re = \frac{\text{inertial response}}{\text{viscous response}} \sim \frac{\rho L_0 V_0}{\mu}$$



Low Re

## High Re





$$\frac{V_0^2}{L_0}\rho\left\{\frac{\partial\tilde{\mathbf{v}}}{\partial\tilde{t}} + (\tilde{\mathbf{v}}\cdot\tilde{\nabla})\tilde{\mathbf{v}}\right\} = -\nabla p + \frac{V_0}{L_0^2}\mu\tilde{\nabla}^2\tilde{\mathbf{v}} + \mathbf{f}$$

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Re ~ 0

## **Stokes equations**

 $\nabla p = \mu \nabla^2 \mathbf{v} + \mathbf{f}$  $\nabla \cdot \mathbf{v} = 0$ 



## Purcell's Scallop Theorem

#### A swimmer strokes must be non-invariant under time reversal







## Green function of the Stokes equation (Stokeslet)





# Far flow field of a swimmer

Swimmers have dipolar far flow fields because they have no net force acting on them



## Dipolar far flow field



# Far flow field of a swimmer

# Swimmers have dipolar far flow fields because they have no net force acting on them

$$v_r = \frac{f}{4\pi\mu} \frac{L}{r^2} (3\cos^2\theta - 1)$$

## Swimmer and colloidal flow fields



 $v \sim \frac{1}{r^2}$ 

 $v \sim \frac{1}{r}$ 





Goldstein group, Cambridge

E-coli





Extensile Pusher Contractile Puller





Extensile Pusher Contractile Puller

Flow field has nematic symmetry

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## nematic symmetry



### nematic symmetry

## nematic phase



 $Q_{ij} = \langle n_i n_j - \frac{\delta_{ij}}{3} \rangle$ 







# Tensor order parameter, Q

$$Q_{ij} = \langle n_i n_j - \frac{\delta_{ij}}{3} \rangle$$

Landau-de Gennes free energy

$$F = \frac{K}{2} (\partial_k Q_{ij})^2 + \frac{A}{2} Q_{ij} Q_{ji} + \frac{B}{3} Q_{ij} Q_{jk} Q_{ki} + \frac{C}{4} (Q_{ij} Q_{ji})^2$$



An 'elastic liquid'

## Topological defects in nematic liquid crystals



topological charge  
m = 
$$\frac{1}{2\pi} \int_{dS} d\theta$$



### liquid crystals





### crystal dislocations

magnetic monopoles in spin ice

topological insulators

quantum vortex in a superfluid

### cosmic strings in the early universe