COSMOLOGICAL BACKGROUNDS of GRAVITATIONAL WAVES 2nd Lecture



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A primer on Inflation

INFLATIONARY COSMOLOGY



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Inflation (basics)



Needed for Consistency of the Big Bang theory

$$a \sim e^{H_* t} \gtrsim e^{60}$$



$$\frac{d^2a}{dt^2} > 0 \Leftrightarrow \frac{d}{dt}\mathcal{H}^{-1} < 0$$





* **Pefinition:**
$$\frac{d^2a}{dt^2} > 0 \Leftrightarrow \frac{d}{dt} \mathcal{H}^{-1} < 0$$
INF
* **Consequences:** If N $\gtrsim 60$ Horizon Problem Solved !
Bonus: Null Curvature
$$(|\frac{K}{\mathcal{H}^2}| \sim |K/H_i^2|e^{-2N} = |K/H_i^2|e^{-120} \ll 1)$$

* **Definition:**

$$\frac{d^{2}a}{dt^{2}} > 0 \Leftrightarrow \frac{d}{dt}\mathcal{H}^{-1} < 0$$
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$$(|\frac{K}{\mathcal{H}^{2}}| \sim |K/\mathcal{H}_{i}^{2}|e^{-2N} = |K/\mathcal{H}_{i}^{2}|e^{-120} \ll 1))$$
* **Implementation:**

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^{2} + V(\phi) \quad (\phi \text{ Inflaton})$$
F $V(\phi) \gg \frac{1}{2}\dot{\phi}^{2}, \frac{1}{2}(\nabla\phi)^{2}$
 $\psi \equiv \frac{p_{\phi}}{\rho_{\phi}} = \frac{\frac{1}{2}\dot{\phi}^{2} - \frac{1}{6\pi^{2}}(\nabla\phi)^{2} - V(\phi)}{\frac{1}{2}\dot{\phi}^{2} + \frac{1}{2\pi^{2}}(\nabla\phi)^{2} + V(\phi)} \simeq \frac{-V(\phi)}{V(\phi)} \simeq -1$

* Implementation:
$$\mathcal{L} = \frac{1}{2}(\partial \phi)^2 + V(\phi)$$
 (\$\phi\$ Inflaton)

Slow Roll (SR) Regime:

$$\epsilon = \frac{\dot{\phi}^2}{2m_p^2 H^2} \ll 1 \ ; \ \eta = -\frac{\ddot{\phi}}{H\dot{\phi}} \ll 1$$

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$$\epsilon_V \equiv \frac{m_p^2}{2} \left(\frac{V'}{V}\right)^2$$

$$\eta_V \equiv m_p^2 \left(\frac{V''}{V}\right)$$

$$(\epsilon \simeq \epsilon_V, \quad \eta \simeq \eta_V - \epsilon_V)$$

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$$(\epsilon \simeq \epsilon_V, \quad \eta \simeq \eta_V - \epsilon_V)$$

If
$$\epsilon_V, \eta_V \ll 1 \Rightarrow \epsilon, \eta \ll 1$$

SR \Rightarrow **quasidS** for $\Delta N = 60$
 $||||$
 $a(t) \simeq a_i e^{\int_t H(\phi) dt'}$

* Implementation:
$$\mathcal{L} = \frac{1}{2}(\partial \phi)^2 + V(\phi)$$
 (\$\phi\$ Inflaton)

Case of Study: $V(\phi) = \frac{1}{2}m_{\phi}^{2}\phi^{2}$ $\epsilon_{V}(\phi) = \eta_{V}(\phi) = 2(m_{p}/\phi)^{2}$ $N(\phi) = (\phi/2m_{p})^{2} - 1/2$

* Implementation:
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Case of Study:
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$$N(\phi) = (\phi/2m_p)^2 - 1/2$$

$$V(\phi) \quad \epsilon_V(\phi) \quad e_V(\phi) = (\phi/2m_p)^2 - 1/2$$

$$V(\phi) \quad e_V(\phi) \quad e_V(\phi) = 1$$

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* Is that ALL? NO!

$$\phi(t)$$

$$g_{\mu\nu}(t)$$
(Background)



* Is that ALL? NO!









$$\hat{\phi}(\vec{x},t) = \phi(t) + \hat{\delta\phi}(\vec{x},t) \quad \rightarrow \quad \langle \hat{\delta\phi}^2(\vec{x},t) \rangle \neq 0$$



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$$ds^2 = g^{\text{tot}}_{\mu\nu} dx^{\mu} dx^{\nu} = (g_{\mu\nu}(t) + \delta g_{\mu\nu}(\vec{x}, t)) dx^{\mu} dx^{\nu}$$

$$\hat{\phi}(\vec{x},t) = \phi(t) + \hat{\delta\phi}(\vec{x},t) \quad \rightarrow \quad \langle \hat{\delta\phi}^2(\vec{x},t) \rangle \neq 0$$

$$= -(1+2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1-2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$

$$ds^2 = -(1+2\Phi)dt^2 + 2B_i dx^i dt + a^2[(1-2\Psi)\delta_{ij} + E_{ij}]dx^i dx^j$$
$$\phi(\vec{x},t) = \phi(t) + \delta\phi(\vec{x},t)$$

$$ds^{2} = -(1+2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1-2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$

$$\phi(\vec{x},t) = \phi(t) + \delta\phi(\vec{x},t)$$

$$E_{ij} = 2\partial_{ij}E + 2\partial_{(i}F_{j)} + h_{ij}$$

Inflation: A generator of Primordial Fluctuations



Expanding U. \longrightarrow Vector Perturbations $S_i, F_i \propto \frac{1}{a}$

$$ds^{2} = -(1+2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1-2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$

$$\phi(\vec{x},t) = \phi(t) + \delta\phi(\vec{x},t)$$

$$E_{ij} = 2\partial_{ij}E + 2\partial_{(i}F_{j)} + h_{ij}$$

$$\partial_{i}h_{ij} = h_{ii} = 0$$
(tensors = GWs)

$$ds^{2} = -(1+2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1-2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$
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$$\phi(\vec{x},t) = \phi(t) + \delta\phi(\vec{x},t)$$

$$\zeta = -[\Psi + (H/\dot{\rho})\delta\rho_{\phi}] \xrightarrow{\text{Piff.}} \zeta$$

$$R = [\Psi + (H/\dot{\phi})\delta\phi] \xrightarrow{\text{Piff.}} R$$

$$Q = [\delta\phi + (\dot{\phi}/H)\Psi] \xrightarrow{\text{Piff.}} Q$$

$$AII$$

$$Gauge$$

$$Inv. I$$

$$ds^{2} = -(1 + 2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1 - 2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$

$$\phi(\vec{x}, t) = \phi(t) + \delta\phi(\vec{x}, t)$$

$$\zeta \equiv -[\Psi + (H/\dot{\rho})\delta\rho_{\phi}] \xrightarrow{\text{Piff.}} \zeta$$

$$\mathcal{R} \equiv [\Psi + (H/\dot{\phi})\delta\phi] \xrightarrow{\text{Piff.}} Q$$

$$Q \equiv [\delta\phi + (\dot{\phi}/H)\Psi] \xrightarrow{\text{Piff.}} Q$$

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$$Curvature \qquad \text{Tensor}$$

$$Pert. (GW)$$

$$Fixing Gauge: e.g. \quad E, \delta\phi = 0 \Rightarrow g_{ij} = a^{2}[(1 - 2R)\delta_{ij} + h_{ij}]$$

$$ds^{2} = -(1+2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1-2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$
$$\phi(\vec{x},t) = \phi(t) + \delta\phi(\vec{x},t)$$

$$g_{ij} = a^2 [(1 - 2\mathcal{R})\delta_{ij} + h_{ij}]$$
Inflation: A generator of Primordial Fluctuations

$$ds^{2} = -(1+2\Phi)dt^{2} + 2B_{i}dx^{i}dt + a^{2}[(1-2\Psi)\delta_{ij} + E_{ij}]dx^{i}dx^{j}$$
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$$g_{ij} = a^{2}[(1-2R)\delta_{ij} + h_{ij}] \quad S = \frac{m_{p}^{2}}{2}\int d^{4}x\sqrt{-g}\{R - (\partial\phi)^{2} - 2V(\phi)\}$$

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$$S = S_{(0)} + S_{(2)}^{(s)} + S_{(2)}^{(t)}$$

$$S_{(2)}^{(s)} = \frac{1}{2} \int d^{4}x \ a^{3} \frac{\dot{\phi}^{2}}{H^{2}} \left[\dot{R}^{2} - a^{-2}(\partial_{i}R)^{2}\right]$$
Background
$$S_{(2)}^{(t)} = \frac{m_{p}^{2}}{8} \int dt dx^{3} a^{3} \left[(\dot{h}_{ij})^{2} - a^{-2}(\partial_{l}h_{ij})^{2}\right]$$

Inflation: A generator of Primordial Fluctuations

Scalar Fluctuations:

 $S_{(2)}^{(s)} = \frac{1}{2} \int d^4x \ a^3 \frac{\dot{\phi}^2}{H^2} \left[\dot{\mathcal{R}}^2 - a^{-2} (\partial_i \mathcal{R})^2 \right]$

Inflation: A generator of Primordial Fluctuations

$$\begin{aligned} & d\tau \equiv dt/a(t) \\ S_{(2)}^{(s)} &= \frac{1}{2} \int d^4x \ a^3 \frac{\dot{\phi}^2}{H^2} \left[\dot{\mathcal{R}}^2 - a^{-2} (\partial_i \mathcal{R})^2 \right] \\ & \left[\begin{array}{c} & t \\ \frac{1}{2} \int d\tau dx^3 \left[(v')^2 - (\nabla v)^2 + \frac{z''}{z} v^2 \right] \right] \\ & \left[v \equiv z \mathcal{R}, \ z \equiv a \frac{\dot{\phi}}{H} \right] \text{ (Mukhanov variable)} \end{aligned} \end{aligned}$$

Inflation: A generator of Primordial Fluctuations

$$S_{(2)}^{(s)} = \frac{1}{2} \int d^4x \ a^3 \frac{\dot{\phi}^2}{H^2} \left[\dot{\mathcal{R}}^2 - a^{-2} (\partial_i \mathcal{R})^2 \right] = \left[\frac{1}{2} \int d\tau dx^3 \left[(v')^2 - (\nabla v)^2 + \frac{z''}{z} v^2 \right] \right]$$

(F.T.:
$$v(\mathbf{x},t) = \int d\mathbf{k} \, e^{-i\mathbf{k}\mathbf{x}} \, v_{\mathbf{k}}(t)$$
)
 $v_{\vec{k}}'' + (k^2 - z''/z)v_{\vec{k}} = 0$ with $\left(\frac{z''}{z} = \frac{1}{\tau^2} \left(\nu^2 - \frac{1}{4}\right), \quad \nu \equiv \frac{3}{2} + 2\epsilon - \eta$

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Quantization: $v_{\vec{k}}(t) \to v_k(t)\hat{a}_{\vec{k}} + v_k^*(t)\hat{a}_{-\vec{k}}^{\dagger}$, $[a_{\vec{k}}, a_{\vec{k}'}^{\dagger}] = (2\pi)^3 \delta(\vec{k} - \vec{k}')$

Inflation: A generator of Primordial Fluctuations

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$$v_k(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau)$$

(we keep only one, $\hat{H}v_k = +kv_k$)

Inflation: A generator of Primordial Fluctuations

$$S_{(2)}^{(s)} = \frac{1}{2} \int d^4x \ a^3 \frac{\dot{\phi}^2}{H^2} \left[\dot{\mathcal{R}}^2 - a^{-2} (\partial_i \mathcal{R})^2 \right] = \left[\frac{1}{2} \int d\tau dx^3 \left[(v')^2 - (\nabla v)^2 + \frac{z''}{z} v^2 \right] \right]$$

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 with $\frac{z''}{z} = \frac{1}{\tau^2} \left(\nu^2 - \frac{1}{4}\right), \quad \nu \equiv \frac{3}{2} + 2\epsilon - \eta$

Quantization: $v_{\vec{k}}(t)$

$$v_{\vec{k}}(t) \to v_k(t)\hat{a}_{\vec{k}} + v_k^*(t)\hat{a}_{-\vec{k}}^{\dagger} , \quad [a_{\vec{k}}, a_{\vec{k}'}^{\dagger}] = (2\pi)^3 \delta(\vec{k} - \vec{k}')$$

$$v_{k}(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau) \xrightarrow{-k\tau \gg 1} \frac{1}{\sqrt{2k}} e^{-ik\tau}$$
(sub-Hubble)
(we keep only one, $\hat{H}v_{k} = +kv_{k}$)
Positive define freq

Inflation: A generator of Primordial Fluctuations

$$S_{(2)}^{(s)} = \frac{1}{2} \int d^4x \ a^3 \frac{\dot{\phi}^2}{H^2} \left[\dot{\mathcal{R}}^2 - a^{-2} (\partial_i \mathcal{R})^2 \right] = \left[\frac{1}{2} \int d\tau dx^3 \left[(v')^2 - (\nabla v)^2 + \frac{z''}{z} v^2 \right] \right]$$

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Quantization: $v_{\vec{k}}(t) \to v_k(t)\hat{a}_{\vec{k}} + v_k^*(t)\hat{a}_{-\vec{k}}^{\dagger}$, $[a_{\vec{k}}, a_{\vec{k}'}^{\dagger}] = (2\pi)^3 \delta(\vec{k} - \vec{k}')$

$$v_k(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau)$$

(Bunch-Davies) Vacuum Fluct.

Inflation: A generator of Primordial Fluctuations

$$v_k(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau)$$
$$\hat{v}_{\vec{k}}(t) \to v_k(t) \hat{a}_{\vec{k}} + v_k^*(t) \hat{a}_{-\vec{k}}^{\dagger}$$

Inflation: A generator of Primordial Fluctuations

$$v_k(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau)$$
$$\hat{v}_{\vec{k}}(t) \to v_k(t) \hat{a}_{\vec{k}} + v_k^*(t) \hat{a}_{-\vec{k}}^{\dagger}$$

$$\left[v \equiv z\mathcal{R} \,, \quad z \equiv a\frac{\dot{\phi}}{H} \right] \longrightarrow \left\langle \hat{\mathcal{R}}_{\vec{k}}\hat{\mathcal{R}}_{\vec{k}'} \right\rangle \equiv \frac{1}{z^2} \left\langle \hat{v}_{\vec{k}}\hat{v}_{\vec{k}'} \right\rangle \equiv (2\pi)^3 \frac{H^2}{a^2 \dot{\phi}^2} \left| v_k(\eta) \right|^2 \delta(\vec{k} + \vec{k}')$$

Inflation: A generator of Primordial Fluctuations

$$v_k(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau)$$
$$\hat{v}_{\vec{k}}(t) \to v_k(t) \hat{a}_{\vec{k}} + v_k^*(t) \hat{a}_{-\vec{k}}^{\dagger}$$

$$\begin{bmatrix} v \equiv z\mathcal{R}, & z \equiv a\frac{\dot{\phi}}{H} \end{bmatrix} \longrightarrow \begin{cases} \left\langle \hat{\mathcal{R}}_{\vec{k}}\hat{\mathcal{R}}_{\vec{k}'} \right\rangle \equiv \frac{1}{z^2} \left\langle \hat{v}_{\vec{k}}\hat{v}_{\vec{k}'} \right\rangle \equiv (2\pi)^3 \frac{H^2}{a^2 \dot{\phi}^2} |v_k(\eta)|^2 \delta(\vec{k} + \vec{k}') \\ \equiv P_{\mathcal{R}}(k,\eta) & \text{Scalar} \\ \text{Fower Spectrum} \end{cases}$$

Inflation: A generator of Primordial Fluctuations

$$v_k(\tau) = \frac{\sqrt{\pi}}{2} e^{i\pi(\nu+1/2)/2} (-\tau)^{1/2} H_{\nu}^{(1)}(-k\tau)$$
$$\hat{v}_{\vec{k}}(t) \to v_k(t) \hat{a}_{\vec{k}} + v_k^*(t) \hat{a}_{-\vec{k}}^{\dagger}$$

Inflation: A generator of Primordial Fluctuations

Tensor Fluctuations: $S_{(2)}^{(t)} = \frac{m_p^2}{8} \int dt dx^3 a^3 \left[(\dot{h}_{ij})^2 - a^{-2} (\partial_l h_{ij})^2 \right]$

Inflation: A generator of Primordial Fluctuations

$$\begin{aligned} \frac{1}{S_{(2)}^{(t)} = \frac{m_p^2}{8} \int dt dx^3 a^3 \left[(\dot{h}_{ij})^2 - a^{-2} (\partial_l h_{ij})^2 \right]} & = \frac{1}{2} \int d\tau d^3 \mathbf{k} \left[(v_{\mathbf{k}}^{s'})^2 - \left(k^2 - \frac{a''}{a} \right) (v_{\mathbf{k}}^{s})^2 \right]} \\ h_{ij}(\vec{k}, \tau) &= \epsilon_{ij}^{(s)} h_{\vec{k}}^{(s)} \longrightarrow v^{(s)} \equiv \frac{a}{2} m_p h_{\vec{k}}^{(s)} \end{aligned}$$

Inflation: A generator of Primordial Fluctuations

$$S_{(2)}^{(t)} = \frac{m_p^2}{8} \int dt dx^3 a^3 \left[(\dot{h}_{ij})^2 - a^{-2} (\partial_l h_{ij})^2 \right] = \left[\sum_s \frac{1}{2} \int d\tau d^3 \mathbf{k} \left[(v_{\mathbf{k}}^{s\prime})^2 - \left(k^2 - \frac{a''}{a} \right) (v_{\mathbf{k}}^{s})^2 \right] \right]$$

Same Procedure as with Scalar Pert. Quantize+Bunch-Davies+Power Spectrum Quantize+Bunch-Davies+Power Spectrum

Inflation: A generator of Primordial Fluctuations

$$S_{(2)}^{(t)} = \frac{m_p^2}{8} \int dt dx^3 a^3 \left[(\dot{h}_{ij})^2 - a^{-2} (\partial_l h_{ij})^2 \right] = \left[\sum_s \frac{1}{2} \int d\tau d^3 \mathbf{k} \left[(v_{\mathbf{k}}^{s\prime})^2 - \left(k^2 - \frac{a''}{a} \right) (v_{\mathbf{k}}^{s})^2 \right] \right]$$

Same Procedure as with Scalar Pert. Quantize+Bunch-Davies+Power Spectrum

Quantization of Gravity dof!





Quantum fluctuations !









End primer on Inflation

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$$g_{\mu\nu} = g^{(B)}_{\mu\nu} + \delta g_{\mu\nu} \quad ; \quad [\delta g_{\mu\nu}]^{TT} = h_{ij} , \begin{cases} h_{ii} = 0\\ \partial_i h_{ij} = 0 \end{cases}$$



$$\left\langle h_{ij}(\vec{k},t)\right\rangle = 0$$

$$\begin{cases} \mathsf{Quantum}\\\mathsf{Fluctuations} \end{cases}$$

$$\left\langle h_{ij}(\vec{k},t)h_{ij}^*(\vec{k}',t)\right\rangle \equiv (2\pi)^3 \frac{2\pi^2}{k^3} \Delta_h^2(k)\delta(\vec{k}-\vec{k}') \end{cases}$$

$$g_{\mu\nu} = g^{(B)}_{\mu\nu} + \delta g_{\mu\nu} \quad ; \quad [\delta g_{\mu\nu}]^{TT} = h_{ij} , \begin{cases} h_{ii} = 0\\ \partial_i h_{ij} = 0 \end{cases}$$



$$\left\langle h_{ij}(\vec{k},t) \right\rangle = 0$$

$$\begin{cases} \mathsf{Quantum}\\ \mathsf{Fluctuations} \end{cases}$$

$$\left\langle h_{ij}(\vec{k},t)h_{ij}^*(\vec{k}',t) \right\rangle \equiv (2\pi)^3 \frac{2\pi^2}{k^3} \Delta_h^2(k)\delta(\vec{k}-\vec{k}') \end{cases}$$

$$\Delta_{h}^{2}(k) = \frac{2}{\pi^{2}} \left(\frac{H}{m_{p}} \right)^{2} \left(\frac{k}{aH} \right)^{n_{t}}$$

$$n_{t} \equiv -2\epsilon$$
energy scale

$$\Delta_{h}^{2}(k) = \frac{2}{\pi^{2}} \left(\frac{H}{m_{p}}\right)^{2} \left(\frac{k}{aH}\right)^{n_{t}}$$

$$n_{t} \equiv -2\epsilon$$
energy scale

$$\Omega_{\rm GW}^{(o)}(f) \equiv \frac{1}{\rho_c^{(o)}} \left(\frac{d\log\rho_{\rm GW}}{d\log k} \right)_o = \underbrace{\frac{\Omega_{\rm Rad}^{(o)}}{24} \Delta_{h_*}^2(k)}_{\text{Transfer Funct.:} T(k) \propto k^0 (\text{RD})} \Delta_h^2(k) = \frac{2}{\pi^2} \left(\underbrace{\frac{H}{m_p}}_{m_p} \right)^2 \left(\frac{k}{aH} \right)^{n_t}_{n_t \equiv -2\epsilon}$$









Shall we coffee break ?



Let's continue ...

Inflating !
Inflation
$$\rightarrow \begin{pmatrix} initial \\ cond. \end{pmatrix} = \begin{pmatrix} Primordial \\ perturbations \end{pmatrix} \begin{pmatrix} Scalar \\ \\ Tensor \end{pmatrix} \qquad Irreducible GW \bigvee \\ Background \end{pmatrix}$$







that it can play an important rale also





that it can play an important role alco

INFLATIONARY MODELS Axion-Inflation



Gauge fields source a blue tilted Non-Gaussian, & Chiral GW Background

Bartolo et al '16, 1610.06481

INFLATIONARY MODELS Axion-Inflation



Gauge fields source a blue tilted Non-Gaussian, & Chiral GW Background

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What if there are arbitrary fields coupled to the inflaton ? (i.e. no need of extra symmetry)

large excitation of fields !? will they create GWs?

$$\begin{aligned} \text{inflaton } \phi &\longrightarrow V(\phi) \\ -\mathcal{L}_{\chi} &= (\partial \chi)^2 / 2 + g^2 (\phi - \phi_0)^2 \chi^2 / 2 \quad \text{Scalar Fld} \\ -\mathcal{L}_{\psi} &= \bar{\psi} \gamma^{\mu} \partial_{\mu} \psi + g(\phi - \phi_0) \bar{\psi} \psi \quad \text{Fermion Fld} \\ \mathcal{L} &= -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - |(\partial_{\mu} - gA_{\mu})\Phi)|^2 - V(\Phi^{\dagger}\Phi) \quad \text{Gauge Fld } (\Phi = \phi e^{i\theta}) \end{aligned}$$

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Let us suppose
$$\left| \Delta_{\mathcal{R}}^2 \gg \Delta_{\mathcal{R}}^2 \right|_{\text{CMB}} \sim 3 \cdot 10^{-9}$$
, @ small scales

$$ds^{2} = a^{2}(\eta) \left[-(1+2\Phi)d\eta^{2} + \left[(1-2\Psi)\delta_{ij} + 2F_{(i,j)} + h_{ij} \right] dx^{i} dx^{j} \right]$$

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 $h_{ij}'' + 2\mathcal{H}h_{ij}' + k^2h_{ij} = S_{ij}^{TT} \sim \Phi * \Phi$ (2nd Order Pert.)

$$\begin{split} \overbrace{S_{ij}} &= 2\Phi\partial_i\partial_j\Phi - 2\Psi\partial_i\partial_j\Phi + 4\Psi\partial_i\partial_j\Psi + \partial_i\Phi\partial_j\Phi - \partial^i\Phi\partial_j\Psi - \partial^i\Psi\partial_j\Phi + 3\partial^i\Psi\partial_j\Psi \\ &- \frac{4}{3(1+w)\mathcal{H}^2}\partial_i(\Psi' + \mathcal{H}\Phi)\partial_j(\Psi' + \mathcal{H}\Phi) \\ &- \frac{2c_s^2}{3w\mathcal{H}}\left[3\mathcal{H}(\mathcal{H}\Phi - \Psi') + \nabla^2\Psi\right]\partial_i\partial_j(\Phi - \Psi) \end{split}$$
 D. Wands et al, 2006-2010
 Baumann et al, 2007
 Peloso et al, 2018

 $\begin{array}{c|c} \text{INFLATION} & \longrightarrow & \text{IF} \left\{ \begin{array}{c} \text{non-monotonic} \\ \text{multi-field} \end{array} \right\} \xrightarrow{} & \text{possible to} \\ \text{enhance } \Delta^2_{\mathcal{R}} \end{array}$ (at small scales) **BBN** $\Omega_{gw,0} < 1.5 \times 10^{-5} \longrightarrow \Delta_{\mathcal{R}}^2 < 0.1$ **LIGO** $\Omega_{gw,0} < 6.9 \times 10^{-6}$ _____ $\triangle_{\mathcal{R}}^2 < 0.07$ **PTA** $\Omega_{gw,0} < 4 \times 10^{-8}$ \longrightarrow $\Delta_{\mathcal{R}}^2 < 5 \times 10^{-3}$ **LISA** $\Omega_{gw,0} < 10^{-13}$ \longrightarrow $\Delta_{\mathcal{R}}^2 < 1 \times 10^{-5}$ **BBO** $\Omega_{gw,0} < 10^{-17}$ \longrightarrow $\Delta_R^2 < 3 \times 10^{-7}$

Phys.Rev. D81 (2010) 023527

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Has LIGO detected PBH's ?



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Primordial Black Holes (PBH) may be produced!

Has LIGO detected PBH's ?



'We will know soon, determining mass/spin distributions' (M. Fishbach (LIGO), Moriond'19)

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PBH candidate for Dark Matter ?

Clesse & Garcia-Bellido, 2015-2017 Ali-Haimoud et al 2016-2017 **Window is very narrow**

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* If PBH are the DM, what is the GW from 2nd $O(\Phi)$? Bartolo et al, '18

* If GW from from 2nd O(Φ) PBH, then Non-Gaussianity? Bartolo et al, '19

* If GW from from 2nd O(Φ) PBH, then Anisotropies? Bartolo et al, '19



GWs from Reheating





1) $V(\phi, \chi) = \frac{1}{4}\lambda\phi^4 + \frac{1}{2}m_{\chi}^2\chi^2 + \frac{1}{2}g^2\phi^2\chi^2$ (Chaotic) 2) $V(\phi, \chi) = \frac{1}{2}\mu^2\phi^2 + \frac{\lambda}{4}(\chi^2 - v^2)^2 + \frac{1}{2}g^2\phi^2\chi^2$ (Hybrid) INFLATON MATTER COUPLING

 $\begin{cases} \ddot{\phi}(t) + 3H\dot{\phi} + V'(\phi) = 0 \quad \text{(Inflaton Zero-Mode : Damped Oscillator)} \\ \Box \phi_k + F(\int dq \phi_q \chi_{|k-q|})\phi_k + \dots = 0 \quad \text{(Inflaton Fluctuations)} \\ \Box \chi_k + F(\int dq \chi_q, \phi_{|k-q|})\chi_k + \dots = 0 \quad \text{(Matter Fluctuations)} \end{cases}$

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DYNAMICS:

Non-Linear, Non-Perturbative & Far-From-Equilibrium $\mathbf{k}_i \pm \Delta \mathbf{k}_i \rightarrow \varphi_k(t), n_k(t) \sim exp\{\mu_k t\}$

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1) Chaotic Scenarios: PARAMETRIC RESONANCE

 $V(\phi,\chi) = V(\phi) + \frac{1}{2}m_{\chi}^2\chi^2 + \frac{1}{2}g^2\phi^2\chi^2$ (Chaotic Models)



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INFLATIONARY PREHEATING

Physics of (p)REHEATING: $\ddot{\varphi}_k + \omega^2(k,t)\varphi_k = 0$

 $\begin{cases} \text{Hybrid Preheating}: \quad \omega^2 = k^2 + m^2(1 - Vt) < 0 \quad \text{(Tachyonic)} \\ \text{Chaotic Preheating}: \quad \omega^2 = k^2 + \Phi^2(t) \sin^2 \mu t \quad \text{(Periodic)} \end{cases} \end{cases}$

At \mathbf{k}_i : $\varphi_{k_i}, n_{k_i} \sim e^{\mu(k,t)t} \Rightarrow$ Inhomogeneities: $\begin{cases} L_i \sim 1/k_i \\ \delta \rho / \rho \gtrsim 1 \\ v \approx c \end{cases}$

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Non - linear dynamics

Lattice Simulations

Let's look at an example ...

Lattice Simulations: Dynamics non-linear out-Eq

Lattice Simulations: Dynamics **non-linear**

• Scalars $(n_k \gg 1)$: $\Box \phi + V_{,\phi} = 0, \ \Box \chi_a + V_{,\chi_a} = 0$

Semi-classical regime $\pi_k \approx \kappa \phi_k + \dots$ (Squeezed States)

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• FLRW:
$$H^2 = \frac{8\pi G}{3}\rho$$
, $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p)$, $\begin{cases} \rho = \langle \rho_{\phi} + \rho_{\chi} + ... \rangle \\ p = \langle p_{\phi} + p_{\chi} + ... \rangle \end{cases}$

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• GW: $h_{ij}'' + 2\mathcal{H}h_{ij}' - \nabla^2 h_{ij} = 16\pi G \Pi_{ij}^{\mathrm{TT}}, \quad \Pi_{ij}^{TT} = \{\partial_i \chi^a \partial_j \chi^a\}^{TT}$ $ds^2 = a^2 (-d\eta^2 + (\delta_{ij} + h_{ij}) dx^i dx^j), \quad \mathrm{TT} : \begin{cases} h_{ii} = 0\\ h_{ij,j} = 0 \end{cases}$

Lattice Simulations: Dynamics out-Eq

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Hybrid Preheating

$$V(\phi, \chi) = \frac{\lambda}{4} (|\chi|^2 - v^2)^2 + \frac{1}{2} |\chi|^2 \phi^2 + V(\phi)$$



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3 stages: **Exp. Instabilities** \rightarrow **Non-linearities** \rightarrow **Relaxation**



Parameter Dependence (Peak amplitude)

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 $A^2 - \omega^6$

 $\omega^2 \equiv V''(\Phi_I)$

1/2

q



 $\Omega^{(o)}_{
m GW}$

Chag

Modele



 $g^2\Phi_{
m i}^2$

 $q \equiv$

Parameter Dependence (Peak amplitude)

Chaotic Models:
$$\Omega_{\rm GW}^{(o)} \sim 10^{-11}$$
@ $f_o \sim 10^8 - 10^9 \ {\rm Hz}$ Large amplitude !... at high Frequency !

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$$\Omega_{\rm GW} \propto q^{-1/2} \longrightarrow$$
 Spectroscopy of particle couplings?



different couplings ... different peaks?

Parameter Dependence (Peak amplitude)

Chaotic Models:
$$\Omega_{\rm GW}^{(o)} \sim 10^{-11}$$
,@ $f_o \sim 10^8 - 10^9 \ {\rm Hz}$ Large amplitude !... at high Frequency !

Very unfortunate... no detectors there !



Parameter Dependence (Peak amplitude)

Hybrid Models:
$$\Omega_{
m GW}^{(o)} \propto \left(rac{v}{m_p}
ight)^2 imes f(\lambda, g^2)$$
 , $f_o \sim \lambda^{1/4} imes 10^9 ~{
m Hz}$

$$\begin{array}{ll} \Omega_{\rm GW}^{(o)} \sim 10^{-11} \,, & @ & \begin{cases} f_o \sim 10^8 - 10^9 \,\, {\rm Hz} \,\, 0.1 \\ & & & \\ f_o \sim 10^2 \,\, {\rm Hz} \,\, (natural) \\ & & & \\ \hline \lambda \sim 10^{-28} \,\, (natural) \\ & & & \\ \hline \lambda \sim 10^{-28} \,\, (natural) \\ & & & \\ \hline \lambda \sim 10^{-28} \,\, (natural) \\ & & & \\ \hline \end{array} \right)$$

realistically speaking ...



EARLY UNIVERSE





BACK SLIDES

PROPAGATION OF TENSORS

Irreducible GW background from Inflation

Tensors = GWs




$$\begin{split} \hat{h}_{ij}(\mathbf{x},\eta) &= \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \begin{pmatrix} h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r}^+ \end{pmatrix} e_{ij}^r(\hat{\mathbf{k}}) \\ \mathbf{quantum fields} \\ \mathbf{Polarizations: +, x} \\ \rho_{\rm GW}(t) &= \frac{1}{32\pi Ga^2(t)} \left\langle \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \right\rangle_V \\ &\equiv \frac{1}{32\pi Ga^2(t)} \frac{1}{V} \int_V d\mathbf{x} \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \end{split}$$

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$$\hat{h}_{ij}(\mathbf{x},\eta) = \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \left(h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r}^+ \right) e_{ij}^r(\hat{\mathbf{k}})$$
quantum fields
$$p_{\text{GW}}(t) = \frac{1}{32\pi G a^2(t)} \left\langle \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \right\rangle_V$$

$$\equiv \frac{1}{32\pi G a^2(t)} \frac{1}{V} \int_V d\mathbf{x} \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t)$$

$$= \frac{1}{32\pi G a^2(t)} \int \frac{d\mathbf{k}}{(2\pi)^3} \frac{d\mathbf{k}'}{(2\pi)^3} \dot{h}_{ij}(\mathbf{k},t) \dot{h}_{ij}^*(\mathbf{k}',t)$$

$$\times \frac{1}{V} \frac{c}{(2\pi)^3 \delta^{(3)}(\mathbf{k}-\mathbf{k}')}$$

$$\begin{split} \hat{h}_{ij}(\mathbf{x},\eta) &= \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \begin{pmatrix} h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r}^+ \end{pmatrix} e_{ij}^r(\hat{\mathbf{k}}) \\ \mathbf{quantum fields} \end{split}$$
Polarizations: +, x
$$\begin{split} \rho_{\rm GW}(t) &= \frac{1}{32\pi Ga^2(t)} \left\langle \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \right\rangle_V \\ &\equiv \frac{1}{32\pi Ga^2(t)} \frac{1}{V} \int_V d\mathbf{x} \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \\ &= \frac{1}{32\pi Ga^2(t)V} \int \frac{d\mathbf{k}}{(2\pi)^3} \dot{h}_{ij}(\mathbf{k},t) \dot{h}_{ij}^*(\mathbf{k},t) \end{split}$$

$$\hat{h}_{ij}(\mathbf{x},\eta) = \sum_{r=+,\times} \int \frac{d^3\mathbf{k}}{(2\pi)^{3/2}} \left(h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r}^+ \right) e_{ij}^r(\hat{\mathbf{k}})$$
quantum fields
Polarizations: +, X

$$\rho_{\rm GW}(t) = \frac{1}{32\pi Ga^2(t)} \left\langle \dot{h}_{ij}(\mathbf{x},t)\dot{h}_{ij}(\mathbf{x},t) \right\rangle_V \longrightarrow \text{Volume/Time Average}$$

$$\hat{h}_{ij}(\mathbf{x},\eta) = \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \left(h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r}^+ \right) e_{ij}^r(\hat{\mathbf{k}})$$
quantum fields
Polarizations: +, x

$$\rho_{\rm GW}(t) = \frac{1}{32\pi Ga^2(t)} \left\langle h_{ij}(\mathbf{x},t)h_{ij}(\mathbf{x},t) \right\rangle_{\rm QM} \longrightarrow \text{ensemble average}$$

$$\begin{split} \hat{h}_{ij}(\mathbf{x},\eta) &= \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \begin{pmatrix} h_k(\eta) \, e^{i\mathbf{k}\mathbf{x}} \, \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \, \hat{a}_{\mathbf{k}r}^+ \end{pmatrix} \, e_{ij}^r(\hat{\mathbf{k}}) \\ \mathbf{quantum fields} & \mathbf{Polarizations:} +, \mathbf{x} \\ \rho_{\rm Gw}(t) &= \frac{1}{32\pi G a^2(t)} \left\langle \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \right\rangle_{\rm QM} \longrightarrow \text{ensemble average} \\ &= \frac{1}{32\pi G a^2(t)} \int \frac{d\mathbf{k}}{(2\pi)^3} \frac{d\mathbf{k}'}{(2\pi)^3} \, e^{i\mathbf{x}(\mathbf{k}-\mathbf{k}')} \left\langle \dot{h}_{ij}\left(\mathbf{k},t\right) \dot{h}_{ij}^*\left(\mathbf{k}',t\right) \right\rangle \end{split}$$

$$\begin{split} \hat{h}_{ij}(\mathbf{x},\eta) &= \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \begin{pmatrix} h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \, \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \, \hat{a}_{\mathbf{k}r}^+ \end{pmatrix} e_{ij}^r(\hat{\mathbf{k}}) \\ \mathbf{quantum fields} & \text{Polarizations: +, } \mathbf{x} \\ \rho_{\text{GW}}(t) &= \frac{1}{32\pi G a^2(t)} \left\langle \dot{h}_{ij}(\mathbf{x},t) \dot{h}_{ij}(\mathbf{x},t) \right\rangle_{\text{GM}} \rightarrow \text{ensemble average} \\ &= \frac{1}{32\pi G a^2(t)} \int \frac{d\mathbf{k}}{(2\pi)^3} \frac{d\mathbf{k}'}{(2\pi)^3} e^{i\mathbf{x}(\mathbf{k}-\mathbf{k}')} \left\langle \dot{h}_{ij}(\mathbf{k},t) \, \dot{h}_{ij}^*(\mathbf{k}',t) \right\rangle \\ & \left\langle \dot{h}_{ij}\left(\mathbf{k},t\right) \dot{h}_{ij}^*\left(\mathbf{k}',t\right) \right\rangle \equiv (2\pi)^3 \mathcal{P}_{\dot{h}}(k,t) \delta^{(3)}(\mathbf{k}-\mathbf{k}') \end{split}$$

$$\begin{split} \hat{h}_{ij}(\mathbf{x},\eta) &= \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \begin{pmatrix} h_k(\eta) \, e^{i\mathbf{k}\mathbf{x}} \, \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \, \hat{a}_{\mathbf{k}r}^+ \end{pmatrix} e_{ij}^r(\hat{\mathbf{k}}) \\ \mathbf{quantum fields} & \mathbf{Polarizations: +, x} \\ \rho_{\rm GW}(t) &= \frac{1}{(4\pi)^3 Ga^2(t)} \int \frac{dk}{k} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \\ &= \int \frac{d\rho_{\rm GW}}{d\log k} \, d\log k \end{split}$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

$$\hat{h}_{ij}(\mathbf{x},\eta) = \sum_{r=+,\times} \int \frac{d^3 \mathbf{k}}{(2\pi)^{3/2}} \left(h_k(\eta) e^{i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r} + h_k^*(\eta) e^{-i\mathbf{k}\mathbf{x}} \hat{a}_{\mathbf{k}r}^+ \right) e_{ij}^r(\hat{\mathbf{k}})$$
quantum fields
Polarizations: +, ×
$$\rho_{\rm GW}(t) = \frac{1}{(4\pi)^3 G a^2(t)} \int \frac{dk}{k} k^3 \mathcal{P}_{\dot{h}}(k,t) = \int \frac{d\rho_{\rm GW}}{d\log k} d\log k$$

$$\frac{d\rho_{\rm GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t)$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

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$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)
ight
angle \equiv (2\pi)^{3}\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

Horizon Re-entry tensors propagate
Rad Dom:
$$h_r(\mathbf{k}, \eta) = \frac{A_r(\mathbf{k})}{a(\eta)} e^{ik\eta} + \frac{B_r(\mathbf{k})}{a(\eta)} e^{-ik\eta}$$

$$\frac{d\rho_{\scriptscriptstyle \rm GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

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$$\left\{\begin{array}{l} @ \text{Horizon}: \left\{\begin{array}{l} h = h_* \\ \dot{h}_* = 0 \end{array}\right\} \\ A = B = \frac{1}{2}a_*h_* \end{array}\right\}$$

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

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$$\left\langle \dot{h}\dot{h}\right\rangle = k^2 \langle hh\rangle = \left(\frac{a_*}{a}\right)^2 \frac{k^2}{2} \langle |h_*|^2 \rangle = \left(\frac{a_o}{a}\right)^2 \frac{k^2}{2(1+z_*)^2} \frac{2\pi^2}{k^3} \Delta_{h_*}^2$$

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

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$$\begin{array}{c} \text{Inflationary} \\ \text{Tensor Spectrum} \end{array}$$

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

$$\mathcal{P}_{\dot{h}} = \left(\frac{a_o}{a}\right)^2 \frac{k^2}{2(1+z_*)^2} \frac{2\pi^2}{k^3} \Delta_{h_*}^2 \rightleftharpoons \frac{d\log\rho_{\rm GW}}{d\log k} = \frac{1}{8} \frac{a_o^2}{a^4} \frac{m_p^2 k^2}{(1+z_*)^2} \Delta_{h_*}^2$$

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

$$\mathcal{P}_{h} = \left(\frac{a_{o}}{a}\right)^{2} \frac{k^{2}}{2(1+z_{*})^{2}} \frac{2\pi^{2}}{k^{3}} \Delta_{h_{*}}^{2} \Longrightarrow \frac{d \log \rho_{\rm GW}}{d \log k} = \frac{1}{8} \frac{a_{o}^{2}}{a^{4}} \frac{m_{p}^{2}k^{2}}{(1+z_{*})^{2}} \Delta_{h_{*}}^{2}$$

$$(1+z_{*})_{\rm RD}^{-2} = \Omega_{\rm Rad}^{(o)} \frac{a_{o}^{2}H_{o}^{2}}{k^{2}} \Longrightarrow \frac{d \log \rho_{\rm GW}}{d \log k} = \frac{\Omega_{\rm Rad}^{(o)}}{24} \left(\frac{a_{o}}{a}\right)^{4} 3m_{p}^{2}H_{o}^{2} \Delta_{h}^{2}$$

*

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

$$\mathcal{P}_{\dot{h}} = \left(\frac{a_o}{a}\right)^2 \frac{k^2}{2(1+z_*)^2} \frac{2\pi^2}{k^3} \Delta_{h_*}^2 \Longrightarrow \frac{d\log\rho_{\rm GW}}{d\log k} = \frac{1}{8} \frac{a_o^2}{a^4} \frac{m_p^2 k^2}{(1+z_*)^2} \Delta_{h_*}^2$$

$$(1+z_*)_{\rm RD}^{-2} = \Omega_{\rm Rad}^{(o)} \frac{a_o^2 H_o^2}{k^2} \quad \Longrightarrow \quad \frac{d\log\rho_{\rm GW}}{d\log k} = \frac{\Omega_{\rm Rad}^{(o)}}{24} \left(\frac{a_o}{a}\right)^4 3m_p^2 H_o^2 \Delta_{h_*}^2$$

$$\Omega_{\rm GW}^{(o)} \equiv \frac{1}{\rho_c^{(o)}} \left(\frac{d \log \rho_{\rm GW}}{d \log k} \right)_o = \frac{\Omega_{\rm Rad}^{(o)}}{24} \Delta_{h_*}^2$$

$$\frac{d\rho_{\rm\scriptscriptstyle GW}}{d\log k}(k,t) = \frac{1}{(4\pi)^3 G \, a^2(t)} \, k^3 \, \mathcal{P}_{\dot{h}}(k,t) \label{eq:GW}$$

$$\left\langle \dot{h}_{ij}\left(\mathbf{k},t\right)\dot{h}_{ij}^{*}\left(\mathbf{k}',t\right)\right\rangle \equiv (2\pi)^{3}\,\mathcal{P}_{\dot{h}}(k,t)\delta^{(3)}(\mathbf{k}-\mathbf{k}')$$

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$$\Omega_{\rm GW}^{(o)}(f) \equiv \frac{1}{\rho_c^{(o)}} \left(\frac{d \log \rho_{\rm GW}}{d \log k} \right)_o = \frac{\Omega_{\rm Rad}^{(o)}}{24} \Delta_{h_*}^2(k) \qquad (k = 2\pi f)$$

GW normalized Inflationary tensor spectrum (today)





Inflationary Hubble Rate

$$\Delta_h^2(k) = \frac{2}{\pi^2} \left(\frac{H}{m_p}\right)^2 \left(\frac{k}{aH}\right)^{n_t}$$
$$n_t \equiv -2\epsilon$$

Small red-tilt, i.e. (almost-) scale-invariant



STIFF EQUATION of STATE & INFLATIONARY GW HF-TAIL





Figure 1. Λ CDM+inflation expansion history with a stiff epoch.





$$\Omega_{\rm GW}(f) \propto H_{\rm inf}^2 \left(\frac{f}{f_{\rm RD}}\right)^{\frac{2(w-1/3)}{(w+1/3)}}$$

Not Scale Invariant !



$$\Omega_{\rm GW}(f) \propto H_{\rm inf}^2 \left(\frac{f}{f_{\rm RD}}\right)^{\frac{2(w-1/3)}{(w+1/3)}}$$

DGF & Tanin (preliminar)



$$\Omega_{\rm GW}(f) \propto H_{\rm inf}^2 \left(\frac{f}{f_{\rm RD}}\right)^{\frac{2(w-1/3)}{(w+1/3)}}$$

DGF & Tanin (preliminar) CMB B-Modes from INFLATION

$$\langle \mathcal{E}^{2} \rangle, \ \langle \mathcal{B}^{2} \rangle \rightarrow \langle |e_{lm}|^{2} \rangle \equiv C_{l}^{E}, \ \langle |b_{lm}|^{2} \rangle \equiv C_{l}^{B}$$
Polarization Angular
Power Spectrum
Pepends on Scalar
(also tensor) Perturbations

$$Pepends only on
Tensor Perturbations!
$$= \underbrace{|}_{E < 0}$$

$$= \underbrace{|}_{E > 0}$$

$$= \underbrace{|}_{B < 0}$$

$$= \underbrace{|}_{B > 0}$$$$



$$\langle \mathcal{E}^2 \rangle, \langle \mathcal{B}^2 \rangle \rightarrow \langle |e_{lm}|^2 \rangle \equiv C_l^E, \langle |b_{lm}|^2 \rangle \equiv C_l^B$$
Polarization Angular
Power Spectrum
Repends on Scalar
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Perturbations
Perturbations !
Perturbations !
















