



G_2 3 NC bosons

φ 1 real scalar singlet under $SU(3)_C \times U(1)_{em}$

CCWZ

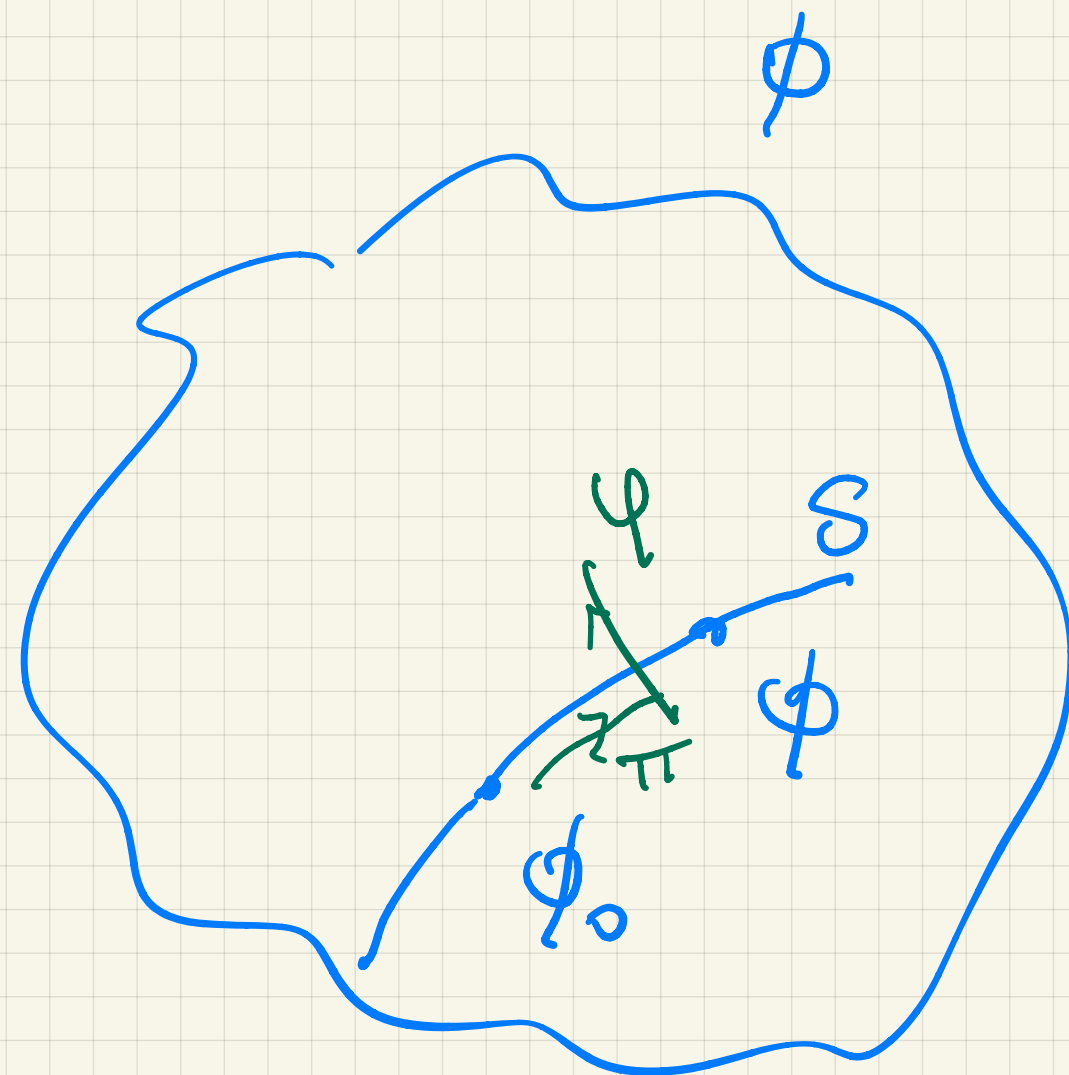
ϕ scalars

$$g \in G: \phi \rightarrow g\phi$$

ϕ_0 ground state

$$H = \{h \in G: h\phi_0 = \phi_0\} \text{ unbroken}$$

$$H \neq G$$



Along δ : $\phi = \exp(\pi_0 \hat{t}_0) \phi_0$ $\hat{t}_0$ broken generators

transf. of π under G

$$\phi = \exp(\pi_0 \hat{t}_0) \quad g\phi = \exp(\pi'_0 \hat{t}_0) \phi, \quad g: \pi \rightarrow \pi'$$

$$g \exp(\pi_0 \hat{t}_0) = \exp(\pi'_0 \hat{t}_0) h \quad h = h(g, \pi) \in H$$

$$\begin{aligned} g(\phi) &= g \exp(\pi_0 \hat{t}_0) \phi_0 = \exp(\pi'_0 \hat{t}_0) h \phi_0 \\ &= \exp(\pi'_0 \hat{t}_0) \phi_0 \end{aligned}$$

CCWZ:

$\exists \varphi$ "Standard Coordinates"

$$\phi = (\pi_\phi, \varphi)$$

$$\phi_0 \subset \rightarrow \varphi = \sim$$

φ transform linearly under H

$$\varphi \rightarrow D(h)\varphi$$

$D(h)$
linear repr.
of H

$$\left\{ \begin{array}{l} \pi \rightarrow \pi' \\ \varphi \rightarrow D(h(g, \pi))\varphi \end{array} \right\} g \in G$$

$$G = SU(2)_L \times U(1)_Y$$

$$H = U(1)_{em}$$

$$g \in G \quad g = u_L e^{i\alpha y}$$

$$u_L \in SU(2)_L$$

$$h \in H \quad h = e^{i\alpha q}$$

q unbroken

t_a^L gen of $SU(2)_L$ broken

$$\Sigma = e^{i g_a t_a^L}$$

$$g \Sigma = \Sigma' h \quad h(g, \Sigma)$$

$$g\Sigma = u_L e^{i\alpha y} \Sigma = u_L \Sigma e^{i\alpha(q-t_3)} =$$

$$\text{SU}(2)_L = \underbrace{u_L \Sigma e^{i\alpha t_3}}_{\text{SU}(2)_L} \underbrace{e^{i\alpha q}}_{\text{U}(1)_{\text{em}}} = \Sigma' h$$

$$\boxed{\begin{aligned} \Sigma &\rightarrow \Sigma' = u_L \Sigma e^{i\alpha t_3} \\ h &= e^{i\alpha q} \end{aligned}}$$

$$\mathcal{L}_\Sigma = \frac{v^2}{2} \text{Tr}[(D_\mu \Sigma)^\dagger (D^\mu \Sigma)] + \textcircled{+} \frac{v^2}{2} \text{Tr}(\Sigma^\dagger D_\mu \Sigma \sigma_3)^2$$

$$D_\mu \Sigma = \partial_\mu \Sigma + ig W_\mu^a \frac{\sigma_a}{2} \Sigma - ig' \Sigma \frac{\sigma_3}{2} B_\mu$$

$$W_\mu = W_\mu^a \frac{\sigma_a}{2} \quad \tilde{W}_\mu = \Sigma^\dagger W_\mu \Sigma - \frac{i}{g} \Sigma^\dagger \partial_\mu \Sigma$$

SU(2)_L gauge transh
by Σ or W_μ

$$\mathcal{L}_\Sigma = M_{W,0}^2 \left(1 + \frac{\partial_T}{2}\right) \tilde{W}_\mu^+ \tilde{W}^{\mu-} + \frac{M_{Z,0}^2}{2} \left(1 - \frac{\partial_T}{2}\right) \tilde{Z}_\mu \tilde{Z}^\mu$$

$$M_{W,0}^2 = \frac{g'^2 v^2}{2}$$

$$\tilde{W}_\mu^\pm \rightarrow \dots$$

$$\tilde{Z}_\mu \rightarrow \dots$$

$$M_{Z,0}^2 = \frac{g'^2 + g^2}{2} v^2$$

Unitary gauge $\Sigma = 1$

$$\mathcal{L}_\Sigma = M_W^2 W_\mu^+ W^{\mu-} + \frac{M_Z^2}{2} Z_\mu Z^\mu$$

$$M_W^2 = M_{W,0}^2 \left(1 + \frac{\partial T}{2}\right)$$

$$M_Z^2 = M_{Z,0}^2 \left(1 - \frac{\partial T}{2}\right)$$

$$\rho = \frac{M_W^2}{C_W M_Z^2} = \frac{1 + \partial T/2}{1 - \partial T/2}$$

$$\rho_{exp} = 1 + O\left(\frac{1}{200}\right) \quad |\partial T| \lesssim \frac{1}{200}$$

$\longleftrightarrow \mathcal{L}(Z)$ is approximately $SU(2)_L \times SU(2)_R$ symm.

$$\Sigma \rightarrow U_L \Sigma U_R^\dagger \quad U_L \in SU(2)_L \quad U_R \in SU(2)_R$$

$$\Sigma = 1 \quad SU(2)_L \times SU(2)_R \rightarrow SU(2)_C \quad \text{diagonal} \\ \text{(custodial } SU(2))$$

$$\varphi \rightarrow \mathcal{D}(h(g, \bar{z}))\varphi$$

φ singlet of $SU(3)_c \times U(1)_{em}$.

$$\varphi \rightarrow \varphi$$

$$\mathcal{D} = 11$$

Singlet under the full $SU(2)_L \times U(1)_Y$

$$\mathcal{L} = \frac{v^2}{2} \uparrow_{\partial_T=0} \left[\partial_\mu \bar{z} + (\partial^\mu \bar{z}) \right] \left(1 + \sqrt{2} \partial \frac{\varphi}{v} + b \frac{\varphi^2}{2v^2} + \dots \right)$$

$$+ \frac{1}{2} (\partial_\mu \varphi)^2 - V(\varphi)$$

$$W_L W_L \rightarrow W_L W_L$$

Equivalence theorem:

$$= \times \left(1 + \mathcal{O}\left(\frac{v^2}{E^2}\right)\right)$$

$$A(G^+ G^- \rightarrow G^+ G^-)_{\text{tree}} \approx \frac{s+t}{2v^2} (1 - \partial^2) + \mathcal{O}\left(\frac{v^2}{E^2}\right)$$

"Unitarisation" for $\partial = 1$

$$A(G^+ G^- \rightarrow \varphi\varphi)_{\text{tree}} = \frac{s}{2v^2} (1 - \partial^2) + \mathcal{O}(\dots)$$

"Unit." for $1 = \partial^2 = 1$

$$a \approx 1$$

$$SM: a=b=c \approx 1 \quad (\text{no high orders in } \phi^n)$$

$$H = \begin{pmatrix} 0 \\ v + \frac{\phi}{\sqrt{2}} \end{pmatrix}$$

recover $\mathcal{L}_{SM}$

$$V_{SM} \rightarrow$$

$$V(\sqrt{1+H^2} - v^2)$$

$$S = \frac{1}{2}$$

$$H = \langle H \rangle + H'$$

$$\int_{\text{SM}} \mathcal{L}_{\text{Higgs}}$$

$$= \lambda_{ij}^e \bar{e}_{Ri} L_j H^* + \lambda_{ij}^D \bar{Q}_{Ri} Q_j H^* + \lambda_{ij}^U \bar{U}_{Ri} Q_j H + \text{h.c.}$$

$$= \dots \dots \dots \langle H \rangle^* \dots \dots \dots \langle H \rangle^* \dots \dots \dots \langle H \rangle \dots \dots$$

$$+ H'^* \quad H'^* \quad H'$$

$$\mathcal{L}_{\text{free fermions}} = \overline{u_{iL}} i\gamma^\mu \partial_\mu u_{iL} + \overline{d_{iL}} i\gamma^\mu \partial_\mu d_{iL} + \overline{e_{iL}} -$$

$$\overline{u_{iR}} i\gamma^\mu \partial_\mu u_{iR} + \overline{d_{iR}} i\gamma^\mu \partial_\mu d_{iR} + \dots$$

$$- m_{ij}^U \overline{u_{iR}} u_{jL} - m_{ij}^D \overline{d_{iR}} d_{jL} - m_{ij}^E \overline{e_{iR}} e_{jL} + \text{h.c.}$$

$$u'_{iR} = U_{ij}^{uR} u_{jR}$$

$$u'_{iL} = U_{ij}^{uL} d_{jL}$$

$$d'_{iR} = U_{ij}^{dR} d_{jR}$$

$$d'_{iL} = U_{ij}^{dL} d_{iL}$$

1) kin. terms preserved

2) mass terms $\rightarrow$

$$- m_{u_i} \overline{u'_{iR}} u'_{iL} - m_{d_i} \overline{d'_{iR}} d'_{iL} + \text{h.c.}$$

$$- m_{e_i} \overline{e'_{iR}} e'_{iL}$$

1) U unitary

2) $m_U = U_{UR}^{\dagger} m_U^{\text{diag}} U_{UL}$

↑
diagonal > 0

$$m_D = U_{DR}^{\dagger} m_D^{\text{diag}} U_{DL}$$

$$m_E = U_{ER}^{\dagger} m_E^{\text{diag}} U_{EL}$$

Th. In generic $n \times n$ complex matrix

$\exists U_R, U_L$ $n \times n$ unitary s.t. $m = U_R^{\dagger} m^{\text{diag}} U_L$

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≥ 0