Scattering in chiral strong backgrounds

Tim Adamo University of Edinburgh

based on work w/ L. Mason & A. Sharma

12 June 2020

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

see 2003.13501, 2007.xxxxx

Motivation

Tree-level gauge theory/gravity in Minkowski space:

We know everything

[Parke-Taylor, Witten, Roiban-Spradlin-Volovich, Hodges, Cachazo-Skinner, Cachazo-He-Yuan,...]

Motivation

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ● ●

Tree-level gauge theory/gravity in Minkowski space:

We know everything

[Parke-Taylor, Witten, Roiban-Spradlin-Volovich, Hodges, Cachazo-Skinner, Cachazo-He-Yuan,...]

- All multiplicity formulae for tree-level S-matrix
- Important data for massless g = 0 sector of string theory

Motivation

Tree-level gauge theory/gravity in Minkowski space:

We know everything

[Parke-Taylor, Witten, Roiban-Spradlin-Volovich, Hodges, Cachazo-Skinner, Cachazo-He-Yuan,...]

- All multiplicity formulae for tree-level S-matrix
- Important data for massless g = 0 sector of string theory

But what about gauge theory/gravity/string theory in **non-trivial/strong** background fields?

Much less known – even in simplest strong backgrounds

• Tree-level 4-points (AdS, plane waves) $D^{Hoker et al., Raju}$,

TA-Casali-Mason-Nekovar]

 Unclear how 'novel' methods (unitarity, double copy, scattering equations) extend to strong backgrounds

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Much less known – even in simplest strong backgrounds

- Tree-level 4-points (AdS, plane waves) [D'Hoker et al., Raju, TA-Casali-Mason-Nekovar]
- Unclear how 'novel' methods (unitarity, double copy, scattering equations) extend to strong backgrounds

Today: Can we make all-multiplicity statements on (any) strong backgrounds?

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ● ●

Much less known – even in simplest strong backgrounds

- Tree-level 4-points (AdS, plane waves) [D'Hoker et al., Raju, TA-Casali-Mason-Nekovar]
- Unclear how 'novel' methods (unitarity, double copy, scattering equations) extend to strong backgrounds

Today: Can we make all-multiplicity statements on (any) strong backgrounds?

Yes!

Class of chiral, asymp. flat strong backgrounds in 4d gauge theory and gravity

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Radiative backgrounds

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

Radiative field is completely characterized by free data at \mathscr{I}^{\pm} Asymp. flat gauge/gravitational fields:

$$A = \mathcal{A}^{0}(u, z, \bar{z}) \, \mathrm{d}z + \bar{\mathcal{A}}^{0}(u, z, \bar{z}) \, \mathrm{d}\bar{z} + O(r^{-1})$$

$$\mathrm{d}s^2 = \mathrm{d}s_{\mathbb{M}}^2 - 2\frac{m_B(u, z, \bar{z})}{r} \,\mathrm{d}u^2 + \frac{\sigma^0(u, z, \bar{z})}{r} \,\mathrm{d}z^2 + \frac{\bar{\sigma}^0(u, z, \bar{z})}{r} \,\mathrm{d}\bar{z}^2 + \frac{\bar{\sigma}^0(u, z, \bar{z})}{r} \,\mathrm{d}\bar{z}^2$$

Radiative backgrounds

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

Radiative field is completely characterized by free data at \mathscr{I}^{\pm} Asymp. flat gauge/gravitational fields:

$$A = \mathcal{A}^{0}(u, z, \bar{z}) \, \mathrm{d}z + \bar{\mathcal{A}}^{0}(u, z, \bar{z}) \, \mathrm{d}\bar{z} + O(r^{-1})$$

$$\mathrm{d}s^2 = \mathrm{d}s_{\mathbb{M}}^2 - 2\frac{m_B(u, z, \bar{z})}{r} \,\mathrm{d}u^2 + \frac{\sigma^0(u, z, \bar{z})}{r} \,\mathrm{d}z^2 + \frac{\bar{\sigma}^0(u, z, \bar{z})}{r} \,\mathrm{d}\bar{z}^2 + \frac{\bar{\sigma}^0(u, z, \bar{z})}{r} \,\mathrm{d}\bar{z}^2$$

Free data are spin-weighted functions on \mathscr{I}^+ :

$$\mathcal{A}^{\mathsf{0}}(u,z,ar{z}) \leftrightarrow s = 1\,, \qquad \sigma^{\mathsf{0}}(u,z,ar{z}) \leftrightarrow s = 2$$

SD Radiative fields

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Lorentzian-real fields ${\cal A}^0,\,\bar{{\cal A}}^0$ and $\sigma^0,\,\bar{\sigma}^0$ related by complex conjugation

Complexify $\Rightarrow A^0$, \tilde{A}^0 and σ^0 , $\tilde{\sigma}^0$ independent

SD Radiative fields

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

Lorentzian-real fields \mathcal{A}^0 , $\bar{\mathcal{A}}^0$ and σ^0 , $\bar{\sigma}^0$ related by complex conjugation

Complexify $\Rightarrow A^0$, \tilde{A}^0 and σ^0 , $\tilde{\sigma}^0$ independent

A SD radiative field has $\tilde{\mathcal{A}}^0 = 0/\tilde{\sigma}^0 = 0$

SD Radiative fields

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

Lorentzian-real fields \mathcal{A}^0 , $\bar{\mathcal{A}}^0$ and σ^0 , $\bar{\sigma}^0$ related by complex conjugation

Complexify $\Rightarrow A^0$, \tilde{A}^0 and σ^0 , $\tilde{\sigma}^0$ independent

A SD radiative field has
$$ilde{\mathcal{A}}^0 = 0/ ilde{\sigma}^0 = 0$$

Idea:

• Exploit integrability of SD sector to compute tree-level S-matrix

Twistor theory

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Twistor space: $Z^{A}=(\mu^{\dot{lpha}},\lambda_{lpha})$ homog. coords. on \mathbb{CP}^{3}

$$\mathbb{PT} = \mathbb{CP}^3 \setminus \{\lambda_\alpha = 0\}$$

 $x\in\mathbb{C}^4$ given by $X\cong\mathbb{CP}^1\subset\mathbb{PT}$ via $\mu^{\dotlpha}=x^{lpha\dotlpha}\lambda_lpha$

Twistor theory

Twistor space: $Z^{A}=(\mu^{\dot{lpha}},\lambda_{lpha})$ homog. coords. on \mathbb{CP}^{3}

$$\mathbb{PT} = \mathbb{CP}^3 \setminus \{\lambda_\alpha = 0\}$$

 $x\in\mathbb{C}^4$ given by $X\cong\mathbb{CP}^1\subset\mathbb{PT}$ via $\mu^{\dotlpha}=x^{lpha\dotlpha}\lambda_lpha$

Many applications in flat background:

- Massless free fields \leftrightarrow cohomology on \mathbb{PT} [Penrose, Sparling, Eastwood-Penrose-Wells]
- Representation for on-shell scattering kinematics [Hodges]
- Full tree-level S-matrix of $\mathcal{N} = 4$ SYM [Witten, Berkovits, Roiban-Spradlin-Volovich]
- Full tree-level S-matrix of $\mathcal{N}=8$ SUGRA [Cachazo-Skinner]

Trivializing SD sector

Theorem [Ward 1977]

There is a 1:1 correspondence between:

- SD SU(N) Yang-Mills fields on \mathbb{C}^4 , and
- rank N holomorphic vector bundles E → PT trivial on every X ⊂ PT (+ technical conditions)

Trivializing SD sector

Theorem [Ward 1977]

There is a 1:1 correspondence between:

- SD SU(N) Yang-Mills fields on \mathbb{C}^4 , and
- rank N holomorphic vector bundles E → PT trivial on every X ⊂ PT (+ technical conditions)

Theorem [Penrose 1976]

There is a 1:1 correspondence between:

- SD 4-manifolds *M*, and
- $\mathbb{P}\mathscr{T}$ a complex deformation of $\mathbb{P}\mathbb{T}$ with rational curve $X \subset \mathbb{P}\mathscr{T}$ with normal bundle $N_X \cong \mathcal{O}(1) \oplus \mathcal{O}(1)$

Upshot

SD sector encoded by integrable $\mathbb C\text{-structures}$ on twistor space

- Gauge theory: $E \to \mathbb{PT}$ partial connection $\overline{D} = \overline{\partial} + A$, $A \in \Omega^{0,1}(\operatorname{End} E)$, $\overline{D}^2 = 0$
- Gravity: $\mathcal{T}_{\mathbb{P}\mathscr{T}}$ complex structure $\overline{\nabla} = \overline{\partial} + V$, $V \in \Omega^{0,1}(\mathcal{T}_{\mathbb{P}\mathscr{T}}), \ \overline{\nabla}^2 = 0$

Upshot

SD sector encoded by integrable $\mathbb C\text{-structures}$ on twistor space

- Gauge theory: $E \to \mathbb{PT}$ partial connection $\overline{D} = \overline{\partial} + A$, $A \in \Omega^{0,1}(\operatorname{End} E)$, $\overline{D}^2 = 0$
- Gravity: $\mathcal{T}_{\mathbb{P}\mathscr{T}}$ complex structure $\overline{\nabla} = \overline{\partial} + V$, $V \in \Omega^{0,1}(\mathcal{T}_{\mathbb{P}\mathscr{T}}), \ \overline{\nabla}^2 = 0$

For SD radiative fields [Sparling, Newman, Eastwood-Tod]

- $A(Z) = \partial_u \mathcal{A}^0(\mu^{\dot{\alpha}}\bar{\lambda}_{\dot{\alpha}},\lambda,\bar{\lambda}) D\bar{\lambda}$
- $V(Z) = \sigma^{0}(\mu^{\dot{\alpha}}\bar{\lambda}_{\dot{\alpha}},\lambda,\bar{\lambda}) \,\mathrm{D}\bar{\lambda}\,\bar{\lambda}^{\dot{\alpha}}\,\frac{\partial}{\partial\mu^{\dot{\alpha}}}$

Gauge theory

▲□▶ ▲□▶ ▲ □▶ ▲ □▶ □ のへぐ

Restrict to SD rad. gauge fields valued in Cartan $E|_X$ trivialized by holomorphic frame

$$\begin{split} \bar{\partial}|_{X}\mathsf{H}(x,\lambda) &= -\mathsf{A}|_{X}\,\mathsf{H}(x,\lambda)\,,\\ \mathsf{A}|_{X} &= \bar{\partial}g(x,\lambda)\,, \quad \text{as } H^{0,1}(\mathbb{P}^{1},\mathcal{O}) = \emptyset\,,\\ &\Rightarrow \quad \mathsf{H}(x,\lambda) = \exp\left[-g(x,\lambda)\right] \end{split}$$

Gauge theory

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Restrict to SD rad. gauge fields valued in Cartan $E|_X$ trivialized by holomorphic frame

$$\begin{split} \bar{\partial}|_{X}\mathsf{H}(x,\lambda) &= -\mathsf{A}|_{X}\,\mathsf{H}(x,\lambda)\,,\\ \mathsf{A}|_{X} &= \bar{\partial}g(x,\lambda)\,, \quad \text{as } H^{0,1}(\mathbb{P}^{1},\mathcal{O}) = \emptyset\,,\\ &\Rightarrow \quad \mathsf{H}(x,\lambda) = \exp\left[-g(x,\lambda)\right] \end{split}$$

Recover space-time gauge field from

$$\partial_{lpha\dot{lpha}}g(x,\lambda) = A_{lpha\dot{lpha}}(x) + \lambda_{lpha} g_{\dot{lpha}}(x,\lambda)$$

Gluon perturbations

Gluon perturbations in background encoded by cohomology:

+ helicity gluon $\leftrightarrow a \in H^{0,1}_{\overline{D}}(\mathbb{PT}, \mathcal{O} \otimes \operatorname{End} E)$ - helicity gluon $\leftrightarrow b \in H^{0,1}_{\overline{D}}(\mathbb{PT}, \mathcal{O}(-4) \otimes \operatorname{End} E)$

Gluon perturbations

A D N A 目 N A E N A E N A B N A C N

Gluon perturbations in background encoded by cohomology: + helicity gluon $\leftrightarrow a \in H^{0,1}_{\bar{D}}(\mathbb{PT}, \mathcal{O} \otimes \operatorname{End} E)$ - helicity gluon $\leftrightarrow b \in H^{0,1}_{\bar{D}}(\mathbb{PT}, \mathcal{O}(-4) \otimes \operatorname{End} E)$

E.g. + helicity with asymp. momentum $k_{lpha\dot{lpha}}=\kappa_lpha ilde{\kappa}_{\dot{lpha}}$

$$\mathsf{a}(Z) = \mathsf{T} \, \frac{\langle \xi \, \lambda \rangle}{\langle \xi \, \kappa \rangle} \bar{\partial} \left(\frac{1}{\langle \lambda \, \kappa \rangle} \right) \, \mathrm{e}^{\mathrm{i} \frac{\langle \xi \, \kappa \rangle}{\langle \xi \, \lambda \rangle} \, [\mu \, \tilde{\kappa}]}$$

Leads to

$$a_{lpha\dot{lpha}}^{(+)}(x) = \mathsf{T}\,rac{\xi_lphaig(ilde{\kappa}_{\dot{lpha}} + e\,g_{\dot{lpha}}(x,\kappa)ig)}{\langle\xi\,\kappa
angle}\,\exp\left[\mathrm{i}\,k\cdot x + e\,g(x,\kappa)
ight]$$

How do we use all this to actually compute amplitudes?

<□ > < □ > < □ > < Ξ > < Ξ > < Ξ > Ξ · のQ@

How do we use all this to actually compute amplitudes?

Lesson from flat background: [Nair, Witten, Berkovits,

Roiban-Spradlin-Volovich]

- Tree-level YM amplitudes live on rational holomorphic curves in $\mathbb{P}\mathbb{T}$
- N^kMHV degree related to degree of curve d = k + 1

$$\mu^{\dot{\alpha}}(\sigma) = u^{\dot{\alpha}}_{\mathbf{a}(d)} \,\sigma^{\mathbf{a}(d)} \,, \qquad \lambda_{\alpha}(\sigma) = \lambda_{\alpha \,\mathbf{a}(d)} \,\sigma^{\mathbf{a}(d)}$$

How do we use all this to actually compute amplitudes?

Lesson from flat background: [Nair, Witten, Berkovits,

Roiban-Spradlin-Volovich]

- Tree-level YM amplitudes live on rational holomorphic curves in $\mathbb{P}\mathbb{T}$
- N^kMHV degree related to degree of curve d = k + 1

$$\mu^{\dot{\alpha}}(\sigma) = u^{\dot{\alpha}}_{\mathbf{a}(d)} \,\sigma^{\mathbf{a}(d)} \,, \qquad \lambda_{\alpha}(\sigma) = \lambda_{\alpha \,\mathbf{a}(d)} \,\sigma^{\mathbf{a}(d)}$$

Colour-ordered N^{d-1}MHV amplitude of $\mathcal{N} = 4$ SYM:

$$\int \frac{\mathrm{d}^{2d+2}\lambda}{\operatorname{vol}\operatorname{GL}(2,\mathbb{C})} \,\delta^{2|4(d+1)} \left(\sum_{i=1}^{n} s_{i}\,\tilde{\kappa}_{i}\,\sigma_{i}^{\mathbf{a}(d)}\right) \\ \times \prod_{j=1}^{n} \frac{\mathrm{d}s_{j}\operatorname{D}\sigma_{j}}{s_{j}\,(j\,j+1)}\,\overline{\delta}^{2}(\kappa_{j}-s_{j}\,\lambda(\sigma_{j}))$$

Central conjecture

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Twistor string theory coupled to background field A (trivial extension to $\mathcal{N}=4$ SYM)

Worldsheet action:

$$S = rac{1}{2\pi} \int_{\Sigma} Y_I \, \bar{\partial} Z' + \mathrm{tr}(j \, \mathsf{A}) + S_{\mathfrak{g}} + \mathrm{ghosts} \, ,$$

Vertex operators:

$$V = \int_{\Sigma} \operatorname{tr}(j \, a(Z))$$

Central conjecture

Twistor string theory coupled to background field A (trivial extension to $\mathcal{N}=4$ SYM)

Worldsheet action:

$$S = \frac{1}{2\pi} \int_{\Sigma} Y_I \,\overline{\partial} Z' + \operatorname{tr}(j \operatorname{A}) + S_{\mathfrak{g}} + \operatorname{ghosts},$$

Vertex operators:

$$V = \int_{\Sigma} \operatorname{tr}(j \, a(Z))$$

Conjecture

N^kMHV amplitude on SD rad. background given by worldsheet correlator at degree d = k + 1

Amplitude formula

Let
$$Z'(\sigma) = U'_{a(d)} \sigma^{a(d)}$$
, and
A $(Z(\sigma)) = \bar{\partial} g(U, \sigma)$

Result for N^kMHV color-ordered tree amplitude:

$$\int \frac{\mathrm{d}^{4|4(d+1)}U}{\mathrm{vol}\,\mathrm{GL}(2,\mathbb{C})} \prod_{i=1}^{n} \frac{\mathrm{d}s_{i}\,\mathrm{D}\sigma_{i}}{s_{j}\left(i\,i+1\right)} \,\overline{\delta}^{2}(\kappa_{i}-s_{i}\,\lambda(\sigma_{i})) \\ \times \exp\left(\left[\mu(\sigma_{i})\,\tilde{\kappa}_{i}\right]+e_{i}\,g(U,\sigma_{i})\right)$$

<ロト < 団ト < 団ト < 団ト < 団ト 三 のへで</p>

Properties

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQで

Formula passes many checks

- Explicit 3- and 4-point checks
- Flat & perturbative limits
- Proved for MHV (d = 1)

Important differences from RSVW

- No momentum conservation
- 4*d* residual integrals (obstructed by background functional freedom)

Properties

Formula passes many checks

- Explicit 3- and 4-point checks
- Flat & perturbative limits
- Proved for MHV (d = 1)

Important differences from RSVW

- No momentum conservation
- 4*d* residual integrals (obstructed by background functional freedom)

```
Similar - more complicated - story for gravity!
```

All-multiplicity formulae on some strong backgrounds!

 \Rightarrow interesting probe of non-linearities in QFT

Any use/interest for string theory?

- SD *plane waves* \subset SD radiative backgrounds
- Vacuum plane waves are string backgrounds [Amati-Klimcik, Horowitz-Steif]
- So these formulae provide data for the $\alpha' \to 0$ limit of strings in plane wave backgrounds

Further directions

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Many interesting open questions:

- Prove N^kMHV formulae for k > 1 (analyticity constraints [Ilderton-MacLeod] ?)
- Other chiral backgrounds (with sources, not asymptotically flat)
- Non-chiral backgrounds, scattering equations, ambitwistor strings

Further directions

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Many interesting open questions:

- Prove N^kMHV formulae for k > 1 (analyticity constraints [Ilderton-MacLeod] ?)
- Other chiral backgrounds (with sources, not asymptotically flat)
- Non-chiral backgrounds, scattering equations, ambitwistor strings

Thanks!