

BLACK HOLE THERMODYNAMICS

Lecture 4

Euclidean Methods

BH Mechanics

K : surface gravity
constant on \mathcal{H}^+

M, J, Q conserved
quantities.

$$dM = \frac{k}{8\pi} dA_{\mathcal{H}^+} + \Omega_{\mathcal{H}^+} dJ + \Phi_{\mathcal{H}^+} dQ$$

$\delta A_{\mathcal{H}^+} \geq 0$ area increase

BH Thermodynamics

$k_{\mathcal{H}^+} \sim$ temp
canonical ensemble
(temp constant in thermal eq.)

$E \sim M$ energy system

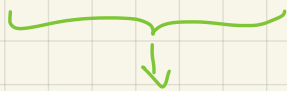
$$dE = T dS + \sum_i \mu_i dN_i$$

$dS \geq 0$ entropy increase

We will find that

$$T_{\mathcal{H}^+} = \frac{\hbar}{2\pi} K \rightarrow \text{Hawking Temperature}$$

$$S_{\text{BH}} = \frac{A_{\mathcal{H}^+}}{4 l_p^2} \rightarrow \text{Bekenstein-Hawking entropy formula.}$$



Einstein-Hilbert $(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G T_{\mu\nu})$

Euc Schw BH:

Schw:

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \frac{dr^2}{\left(1 - \frac{2M}{r}\right)} + r^2 (d\theta^2 + \sin^2\theta d\phi^2)$$

↓ Euc Schw: $t = -it_E$ Wick Rotation

$$ds_E^2 = \left(1 - \frac{2M}{r}\right) dt_E^2 + \frac{dr^2}{\left(1 - \frac{2M}{r}\right)} + r^2 (d\theta^2 + \sin^2\theta d\phi^2)$$

Zoom near $r = 2M$, define new coords

$$r = \frac{\rho^2}{8M} + 2M \quad dr = \frac{\rho}{4M} d\rho$$

look at behaviour as $\rho \rightarrow 0$ ($r \rightarrow 2M$)

$$1 - \frac{2M}{r} = 1 - \frac{1}{1 + \frac{\rho^2}{16M^2}} = \frac{\rho^2}{16M^2} + O(\rho^4)$$

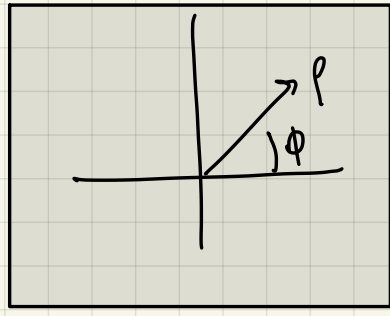
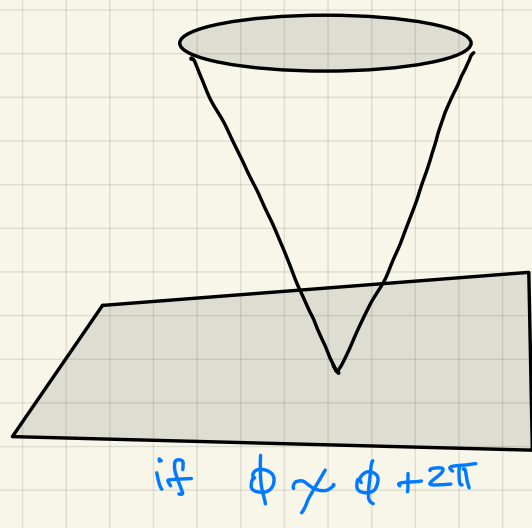
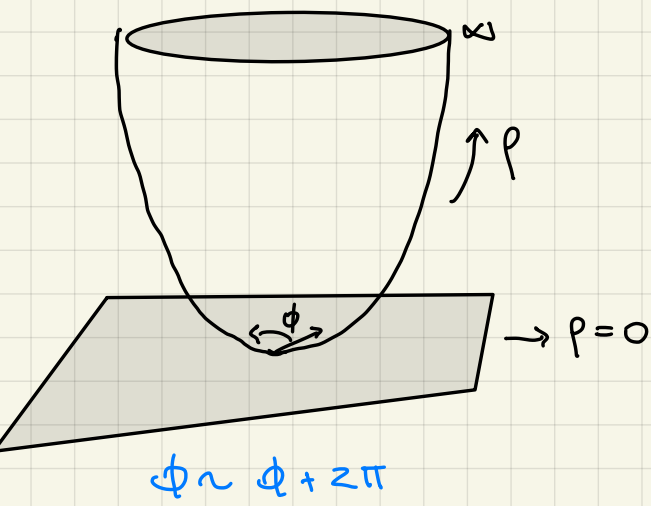
$$ds_E^2 \underset{\rho \rightarrow 0}{=} \frac{\rho^2}{16M^2} dt_E^2 + \frac{1}{\cancel{\rho^2/16M^2}} \frac{\cancel{\rho^2}}{16M^2} d\rho^2 + (2M)^2 (d\theta^2 + \sin^2\theta d\phi^2) + \dots$$

$$= \underbrace{\frac{\rho^2}{16M^2} dt_E^2 + d\rho^2}_{\mathbb{R}^2 \text{ in polar coords}} + \underbrace{(2M)^2 (d\theta^2 + \sin^2\theta d\phi^2)}_{S^2} + \dots$$

\mathbb{R}^2 in polar coords

$$\rho^2 d\phi^2 + d\rho^2 \Rightarrow \phi = \frac{t_E}{4M}$$

$$\rho \in [0, \infty) \quad \phi \in [0, 2\pi)$$



If we want a regular Euclidean geometry for Schw
 $\phi \sim \phi + 2\pi$

$$\Rightarrow t_E \sim t_E + 2\pi (4M)$$

QFT interpretation

What does it mean for Euc time to be periodic?

$$t_E \sim t_E + \beta$$

From a path integral we would have

$$Z[\Phi_i, \Phi_f] = \int_{\Phi(0)=\Phi_i}^{\Phi(t_E)=\Phi_f} \mathcal{D}\Phi e^{\frac{iS[\Phi]}{\hbar}} = \langle \Phi_i | e^{-\frac{it_E H}{\hbar}} | \Phi_f \rangle$$

$$iS[\Phi] = i \int_{t=it_E} dt d^3x L = -S_E = - \int dt_E d^3x L$$

$$e^{-itH/\hbar} \rightarrow e^{-tE/\hbar}$$

$$\Rightarrow t_E \sim t_E + \beta$$

$$Z[\beta] = \text{Tr}(e^{-\beta H}) \quad : \quad \text{partition function @ finite temp}$$

for any momentum

$$Z[\beta, \Omega] = \text{Tr}(e^{-\beta(H - \Omega J)}) \quad T = 1/\beta$$

Interpretation implemented for Euc Schw

$$t_E \sim t_E + 2\pi(4M)$$

$$\Rightarrow \beta = 2\pi \cdot \underbrace{4M}_{k^{-1}} \Rightarrow \frac{1}{4} = \frac{1}{8\pi M}$$

$$\frac{1}{4} = \frac{k}{2\pi}$$

Rindler spacetime: back to Lorentzian

$$ds_E = \underbrace{\frac{p^2}{16M^2} dt_E^2 + dp^2}_{R^2 \text{ in polar coords}} + \underbrace{(2M)^2 (d\theta^2 + \sin^2\theta d\phi^2)}_{S^2} + \dots$$

↓ back to Lorentzian

$$= - \underbrace{\frac{p^2}{16M^2} dt^2 + dp^2}_{\text{Rindler}} + (2M)^2 (d\theta^2 + \sin^2\theta d\phi^2)$$

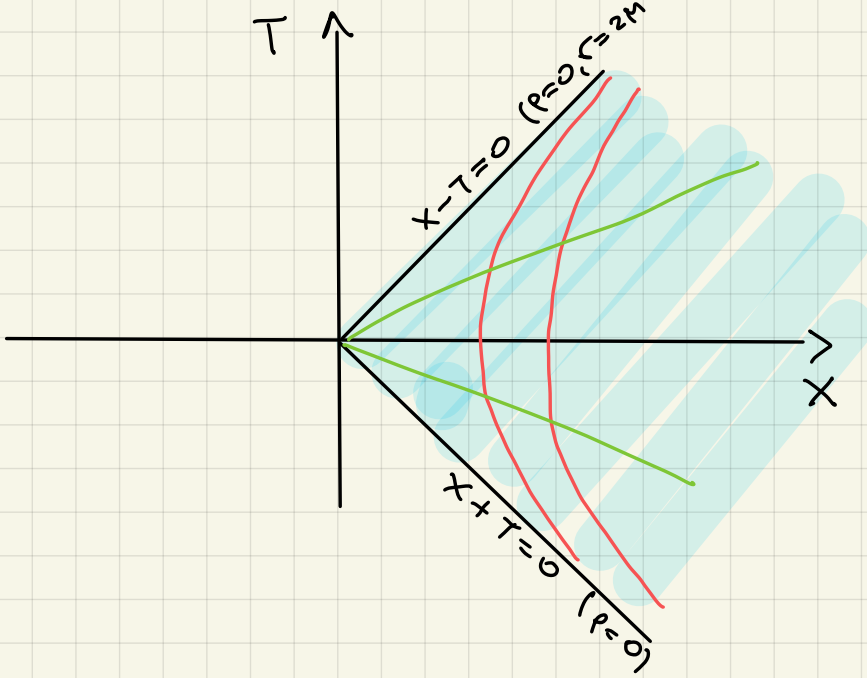
$$T = p \sinh\left(\frac{t}{4M}\right)$$

$$X = p \cosh\left(\frac{t}{4M}\right)$$

$$= -dT^2 + dX^2 + \dots$$

$$X^2 - T^2 = p^2 \quad \text{---}$$

$$\frac{T}{X} = \tanh\left(\frac{t}{4M}\right) \quad \text{---}$$



Rindler: acc obs
in Minkowski

$$\text{acc} : \frac{1}{4M} = k$$