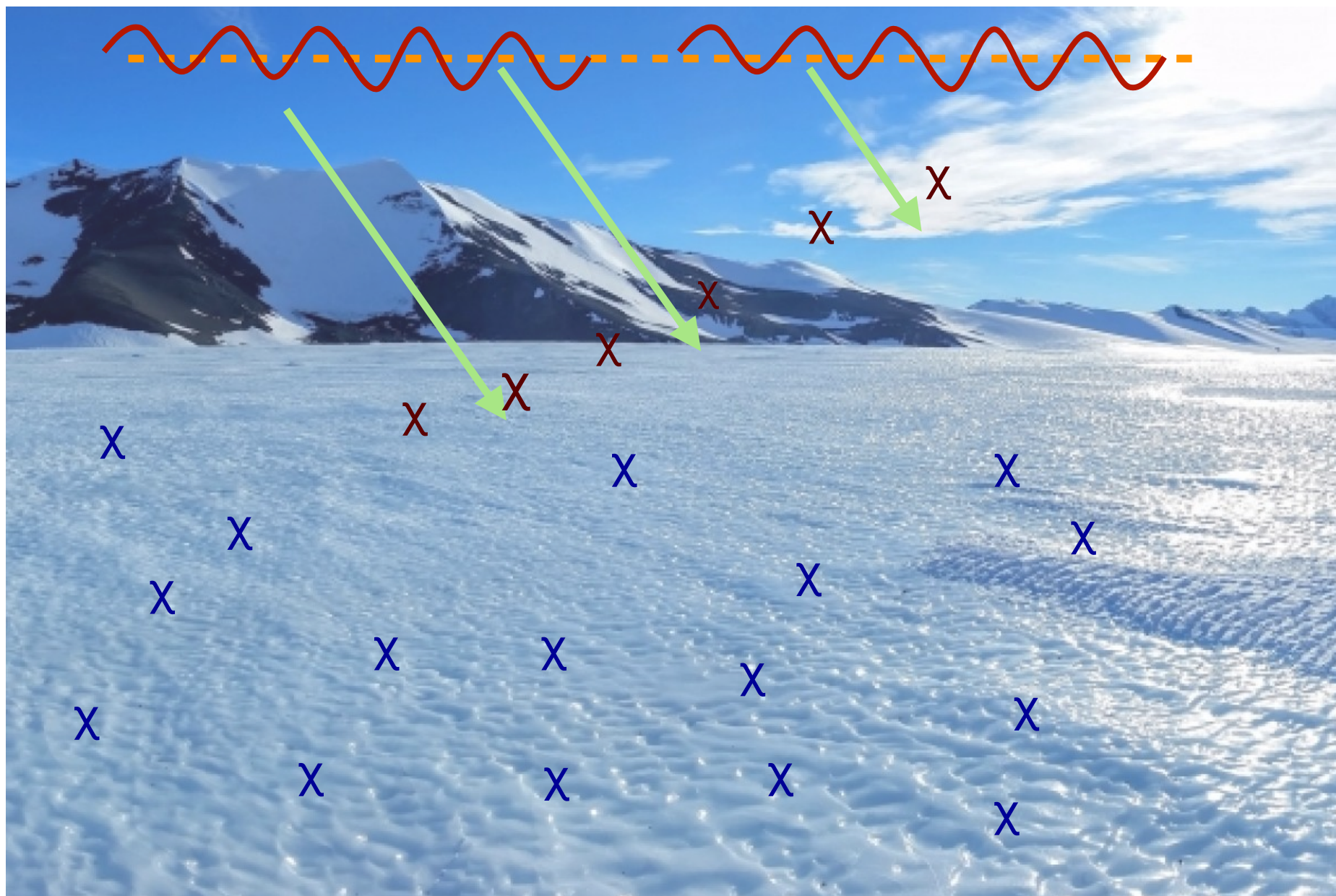


Freeze-in versus Glaciation: Freezing into a thermalized hidden sector

Nicolas Fernandez

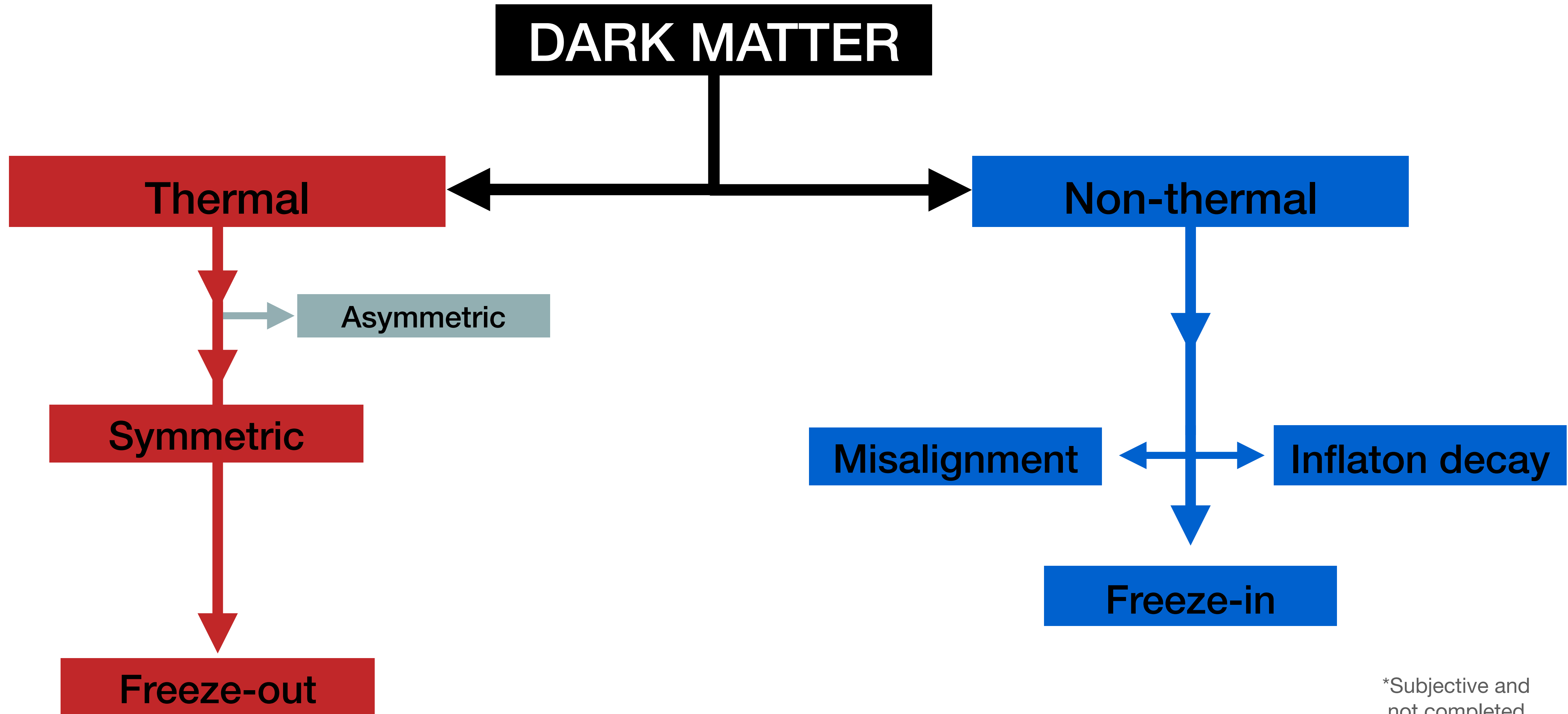
UIUC

*Based on **NF**, Kahn and Shelton in prep.*



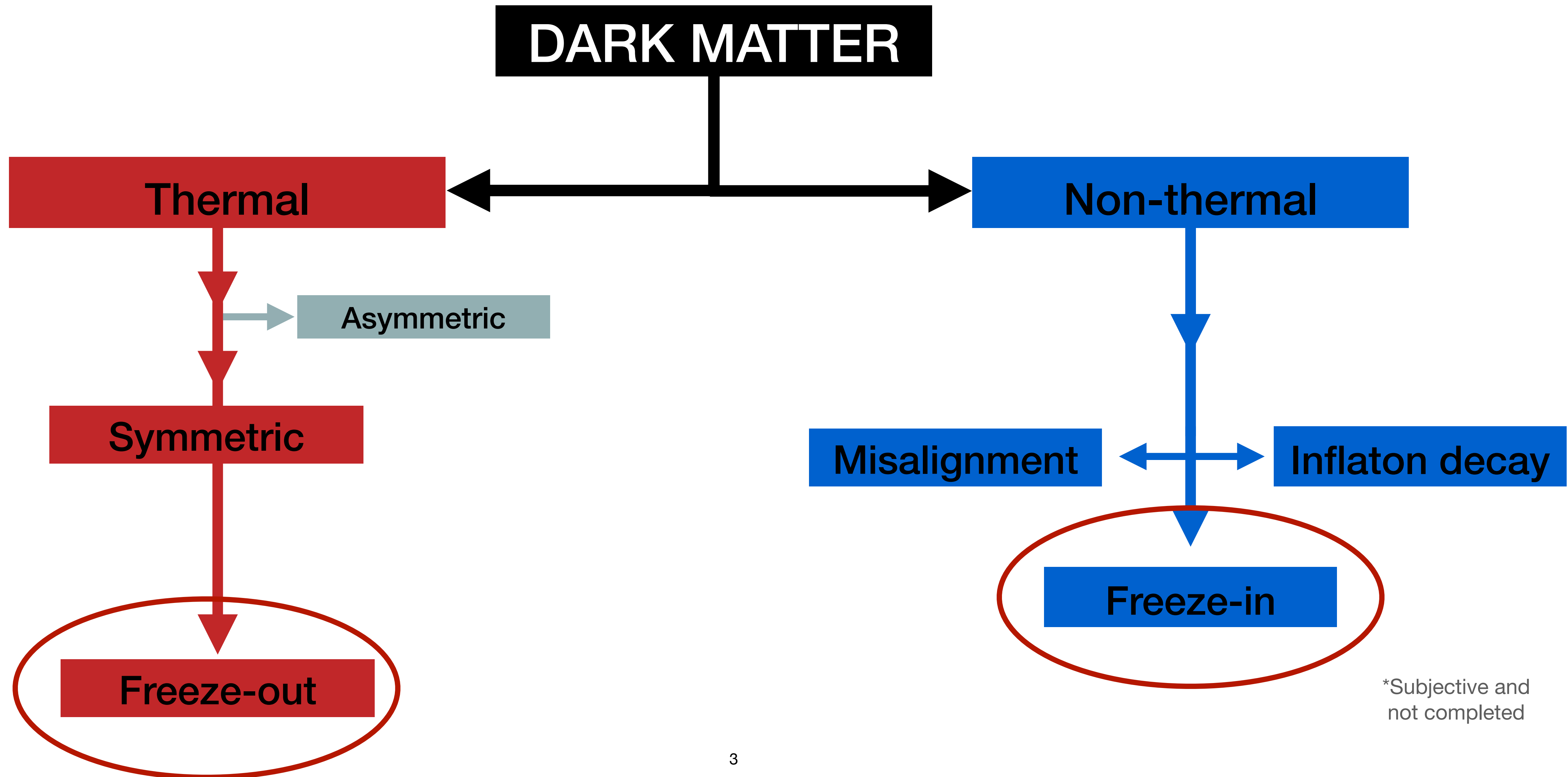
I ILLINOIS

Dark Matter Production Flow*



*Subjective and not completed

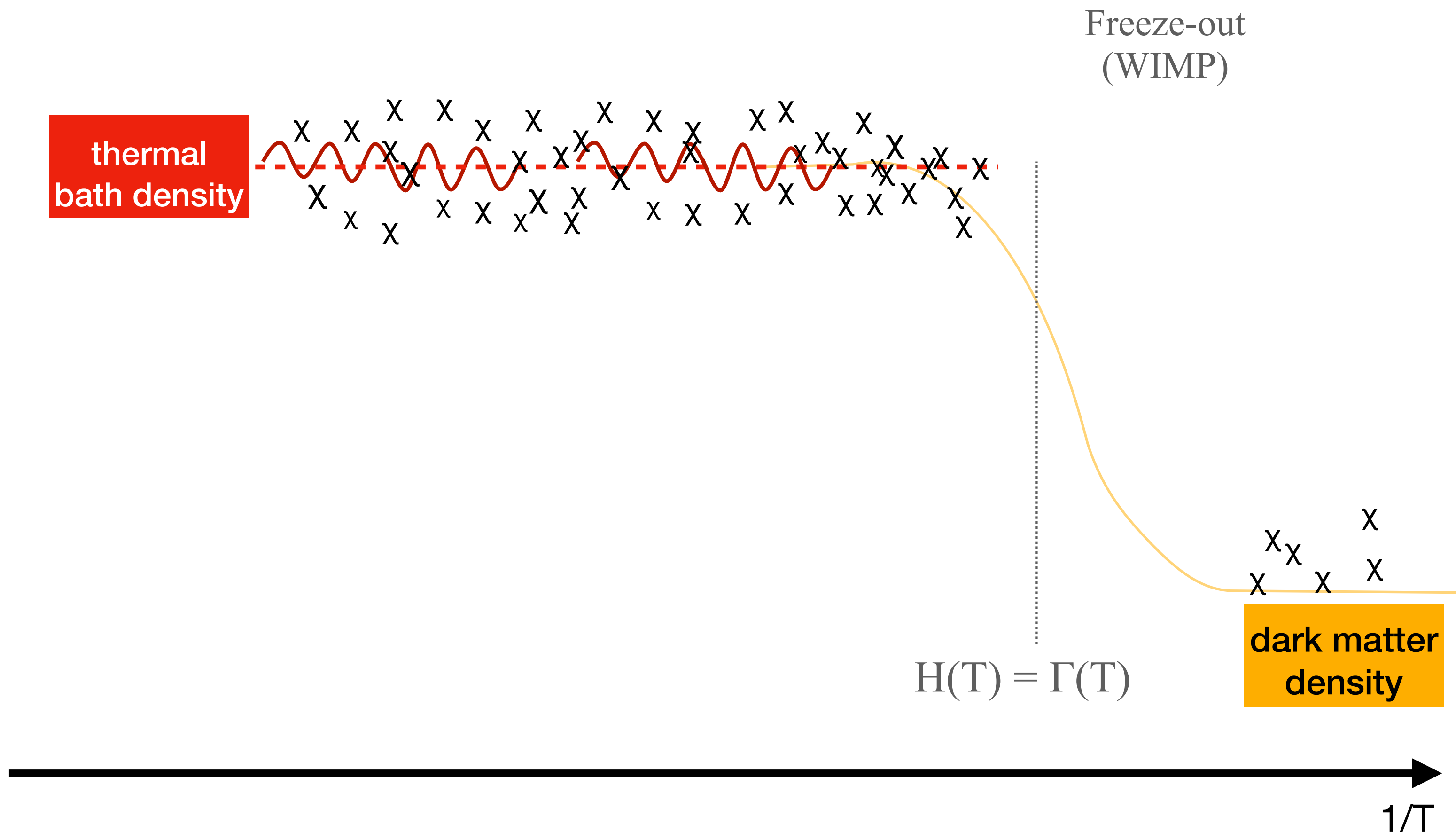
Dark Matter Production Flow*



*Subjective and not completed

Freeze-out

Or the art of getting rid of stuff

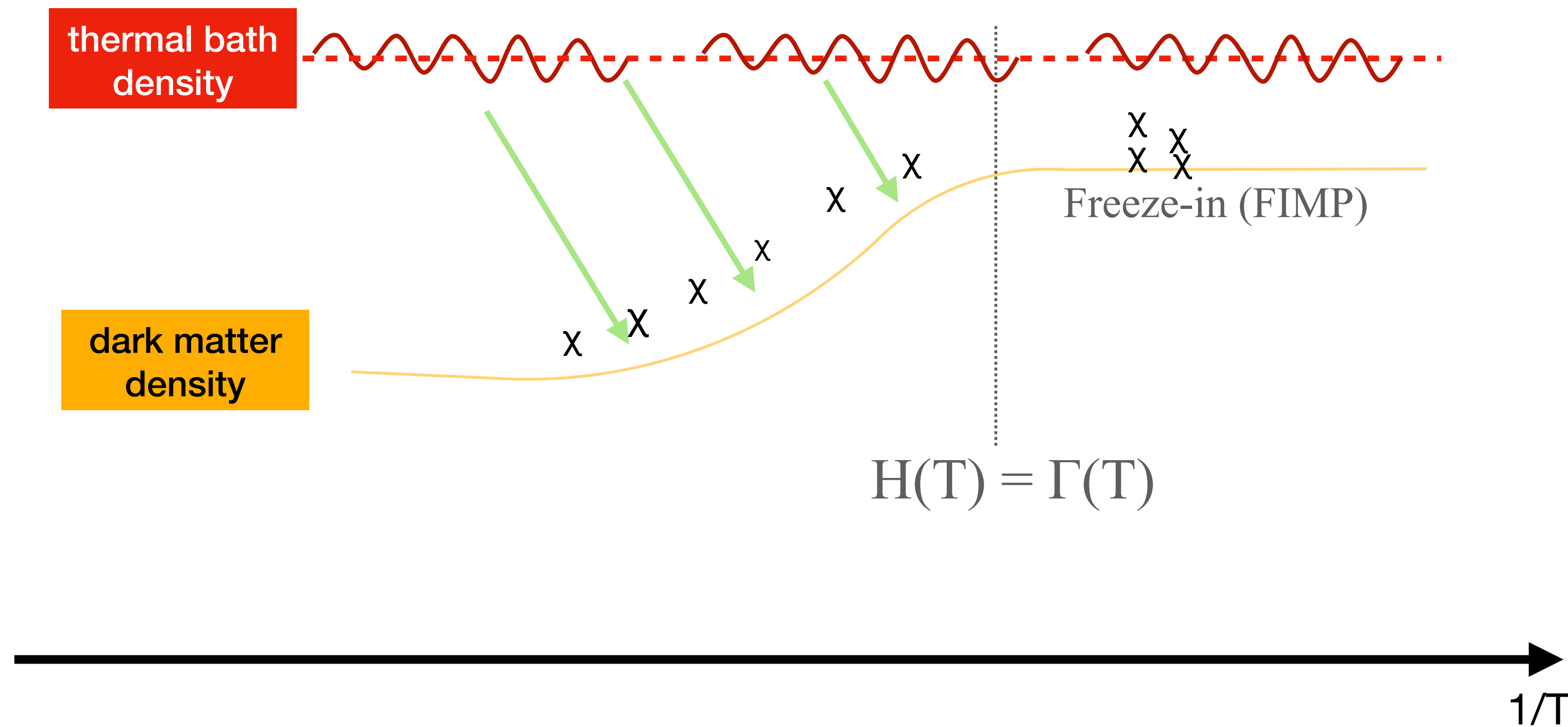


- Relic abundance is independent of initial conditions
- Fine with BBN (masses $>$ few MeV)
- Experimentally testable

Freeze-in

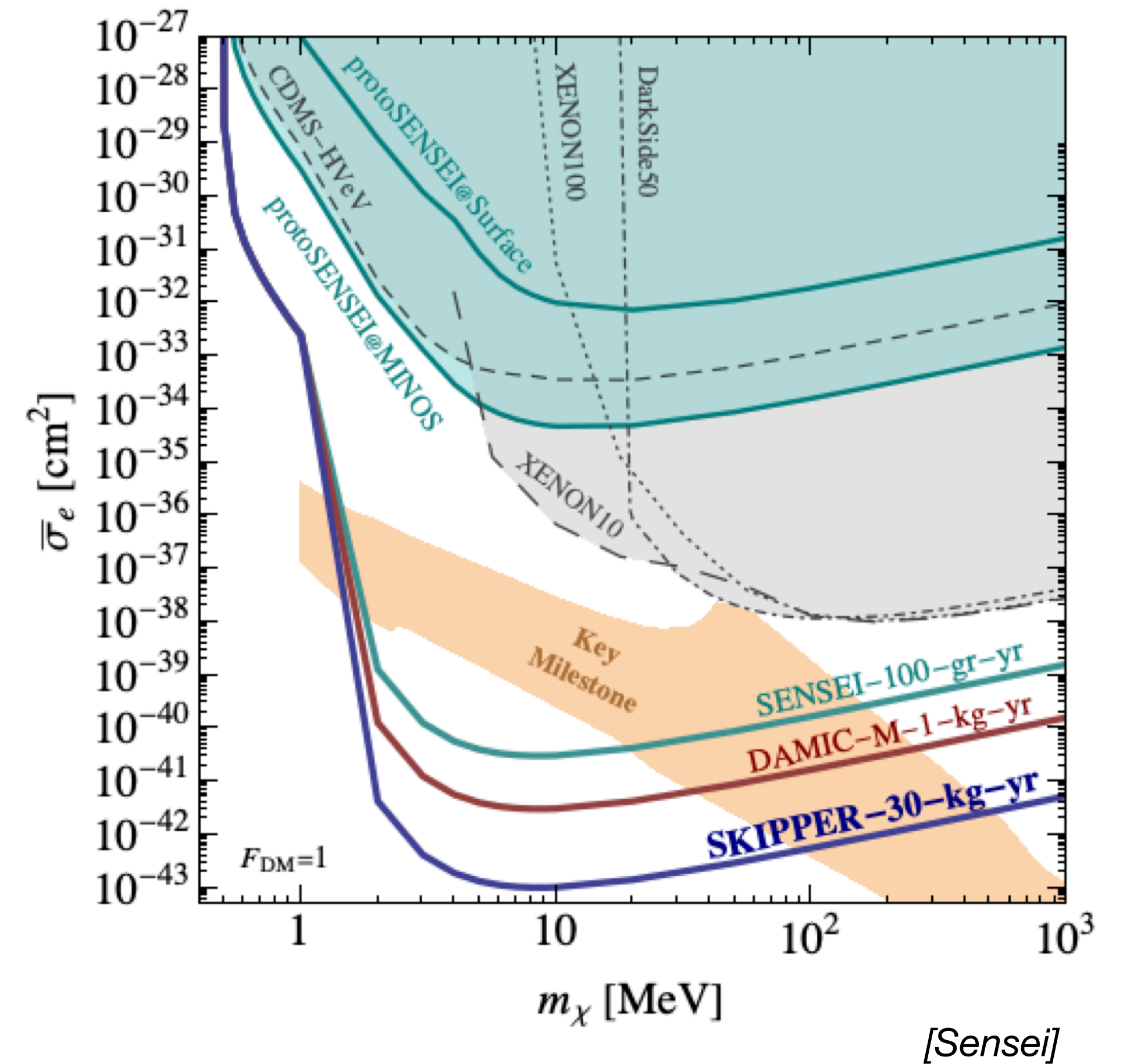
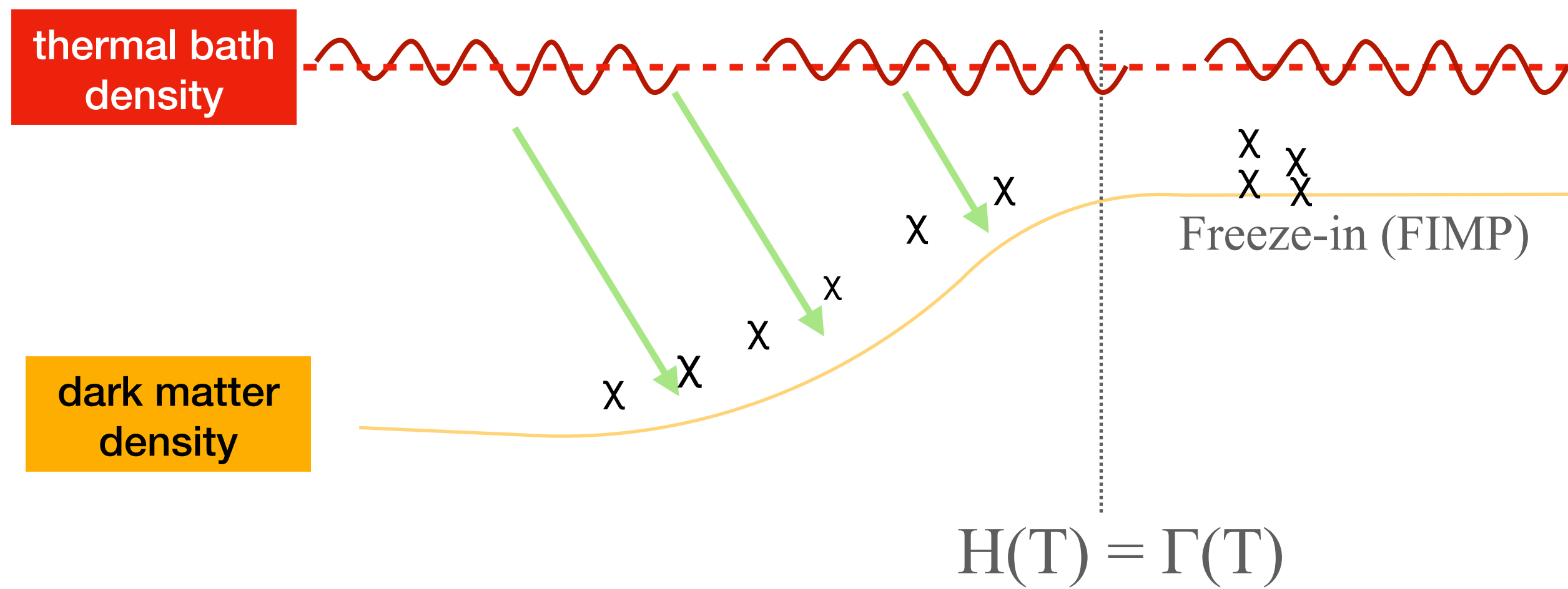
Or the art of getting less but enough

- Relic abundance is independent of initial conditions*
- Fine with BBN and Neff (masses $> \text{keV}$)
- Experimentally testable soon! Very exciting!



Freeze-in

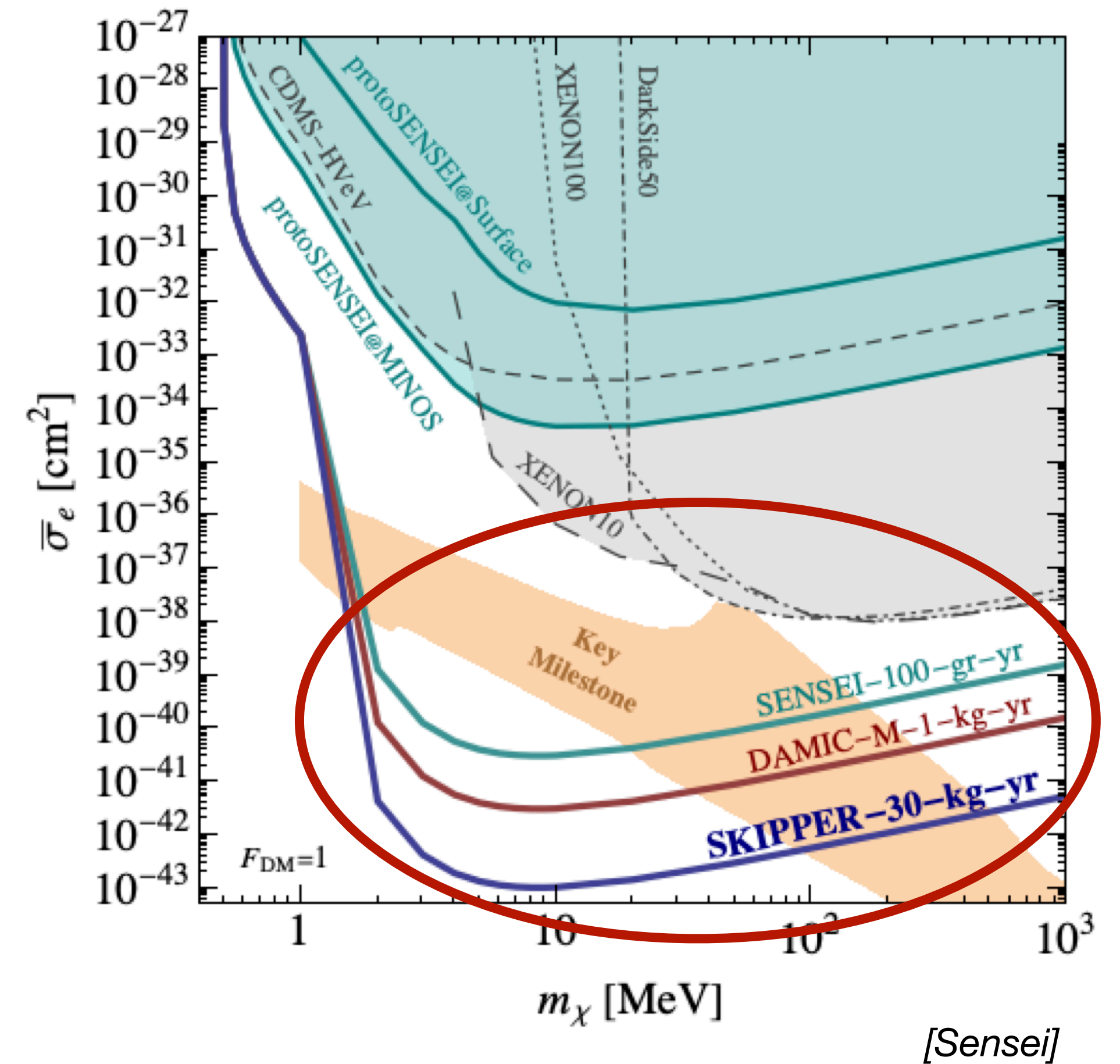
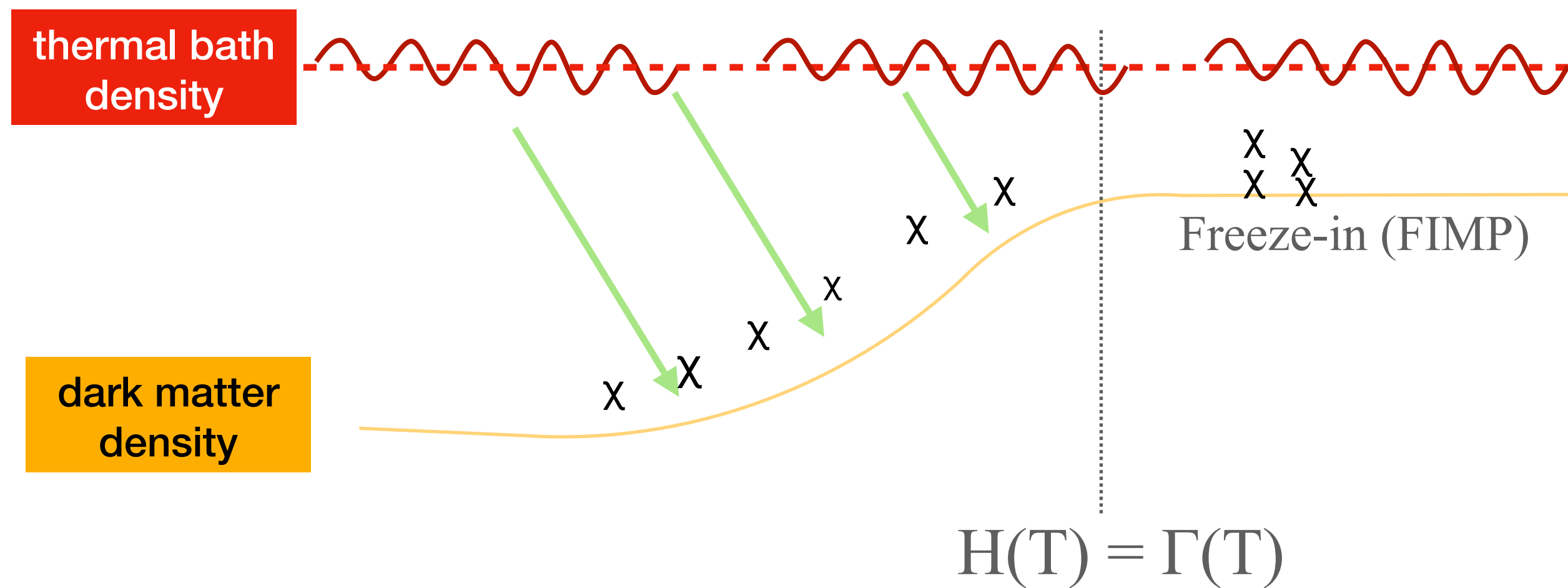
Or the art of getting less and just enough



Freeze-in

Or the art of getting less and just enough

Tien-Tien Yu Talk

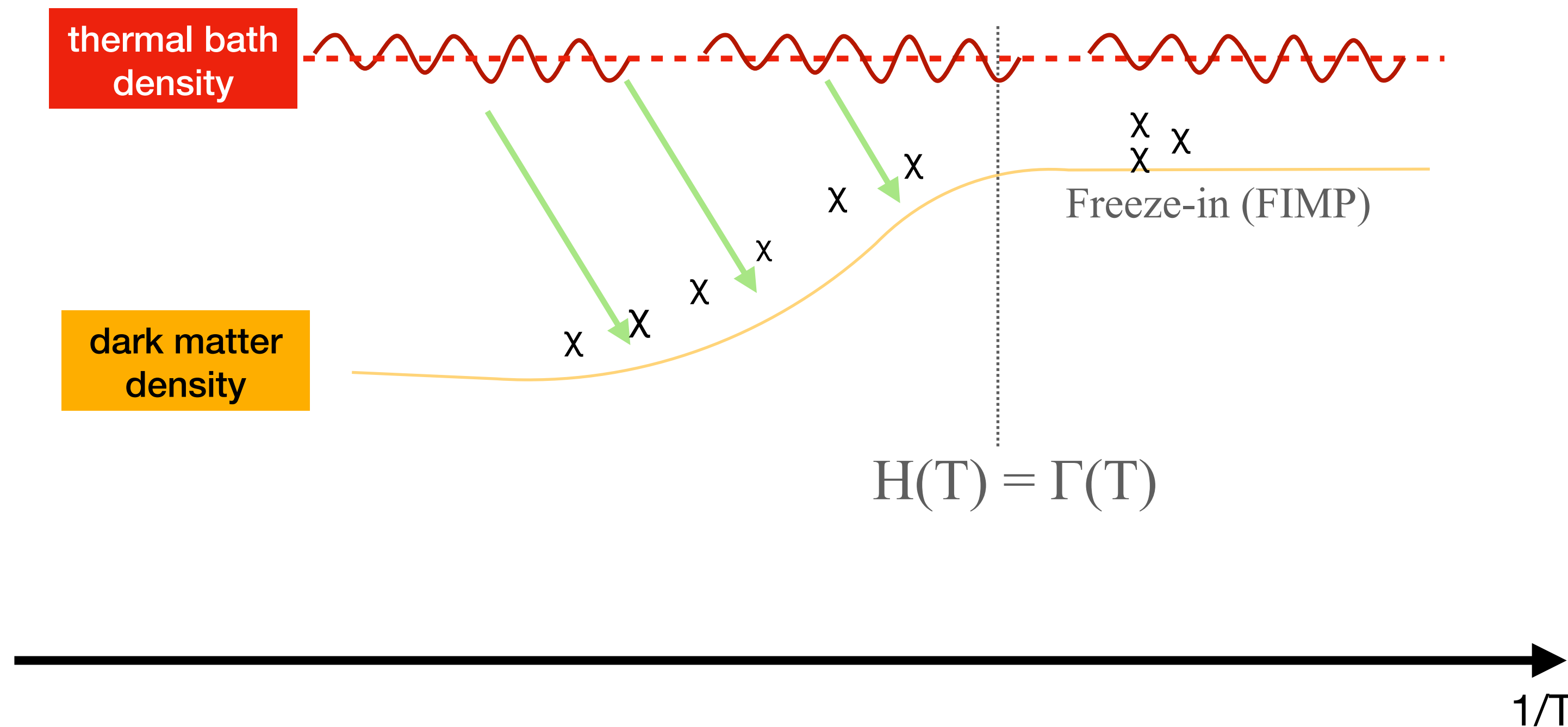


Freeze-in

Or the art of getting less and just enough

- Relic abundance is independent of initial conditions

- Fine with BBN and Neff (masses $> \text{keV}$)
- Experimentally testable soon! Very exciting!



Freeze-in

Or the art of getting less and just enough

The standard freeze-in paradigm has a hidden UV sensitivity in that the initial DM population is assumed to be exactly zero.

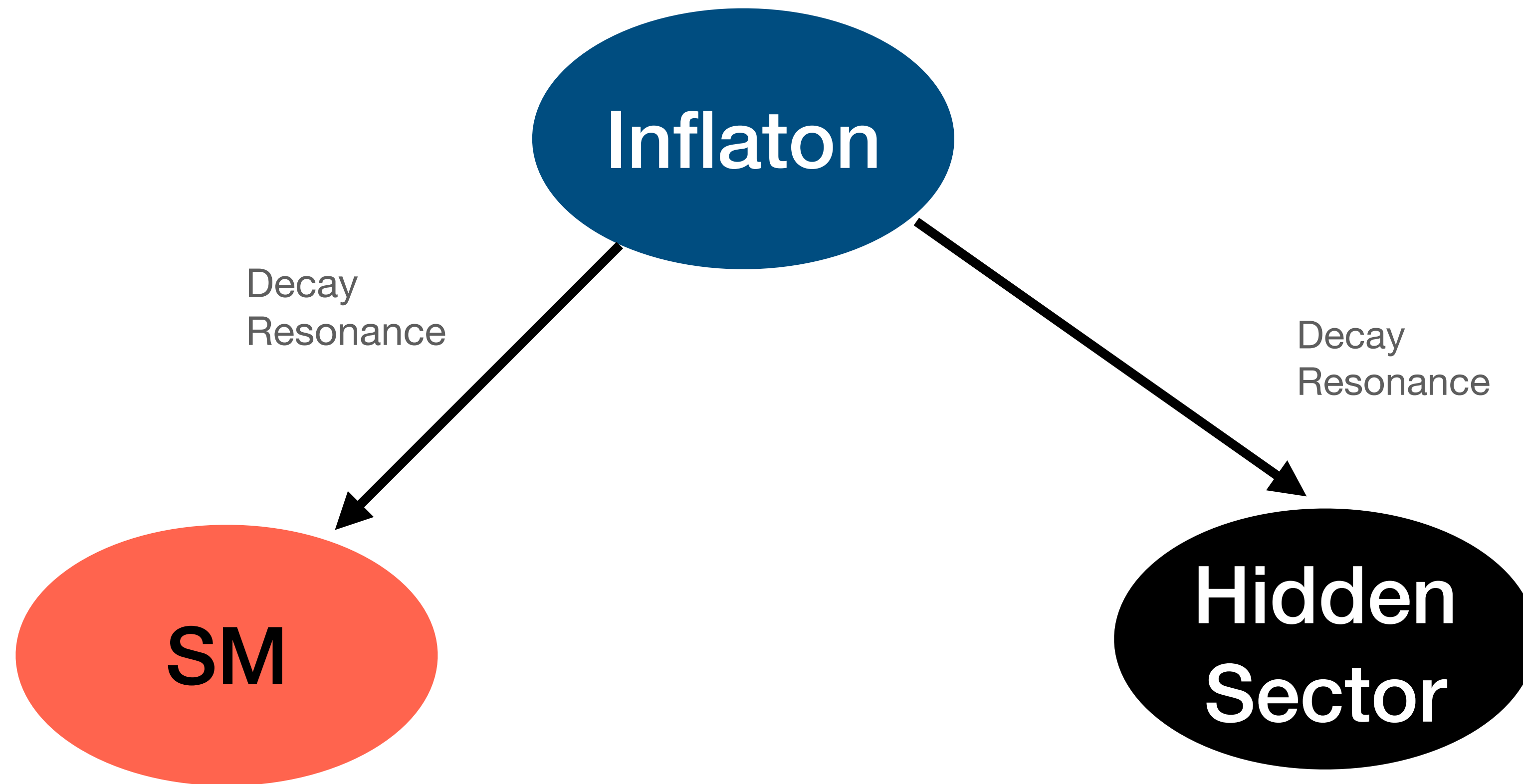
thermal
density

dark matter
density

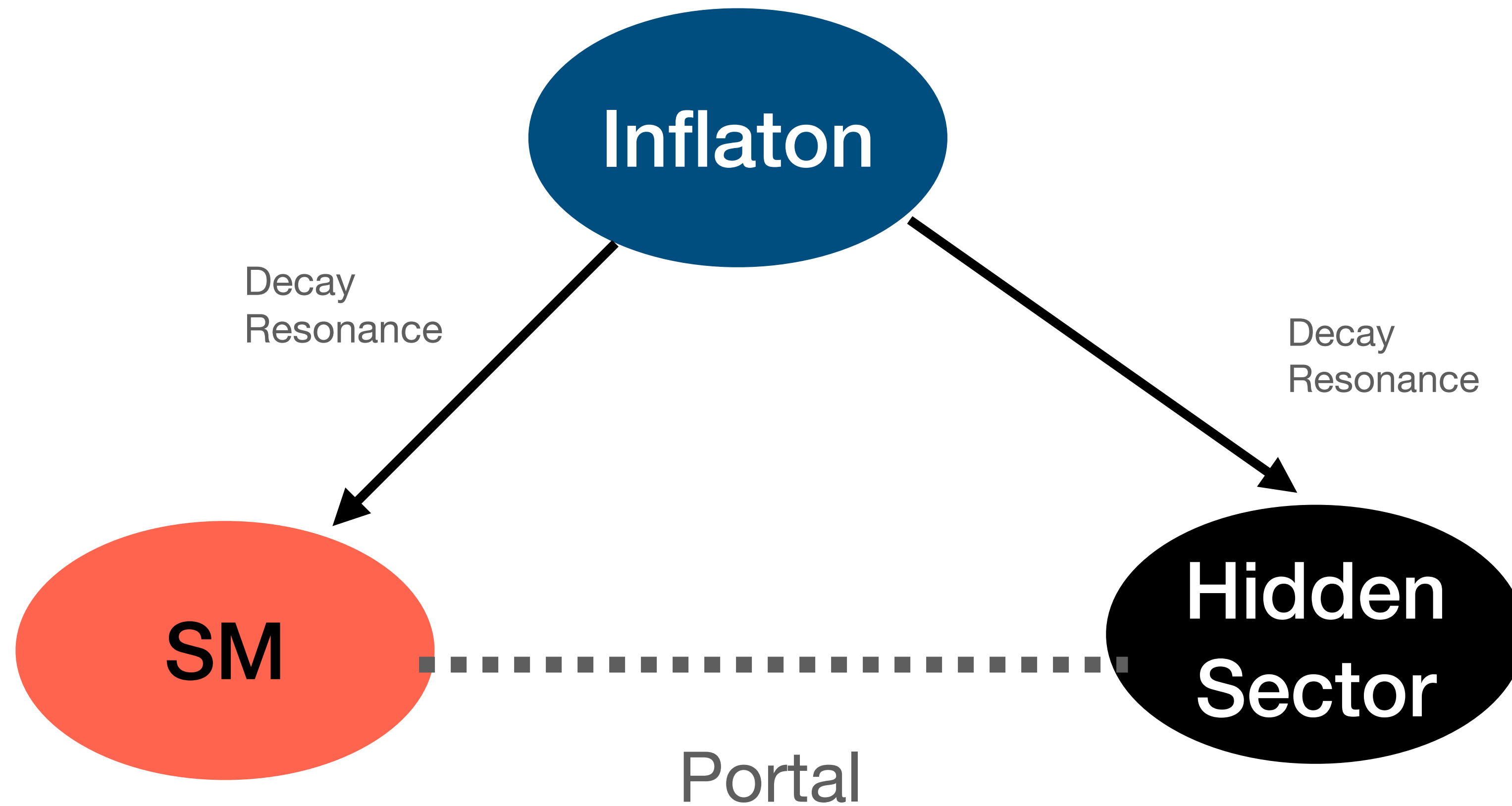
$$\Omega(1) - \Omega(1)$$

1/T

What if...



What if...



Explicit Model: Kinetic mixing portal

$$\mathcal{L} = -\frac{1}{4}\tilde{Z}'_{\mu\nu}\tilde{Z}'^{\mu\nu} - \frac{\epsilon}{2\cos\theta_W}\tilde{Z}'_{\mu\nu}\tilde{B}^{\mu\nu} - \frac{1}{2}m_{Z_D}^2\tilde{Z}_{D\mu}\tilde{Z}_D^\mu + g_\chi J_D^\mu\tilde{Z}_{D\mu} + \bar{\chi}(i\gamma^\mu\partial_\mu - m_\chi)\chi,$$

$$m_{Z_D} \ll m_Z \quad (\text{ultra light mediator})$$

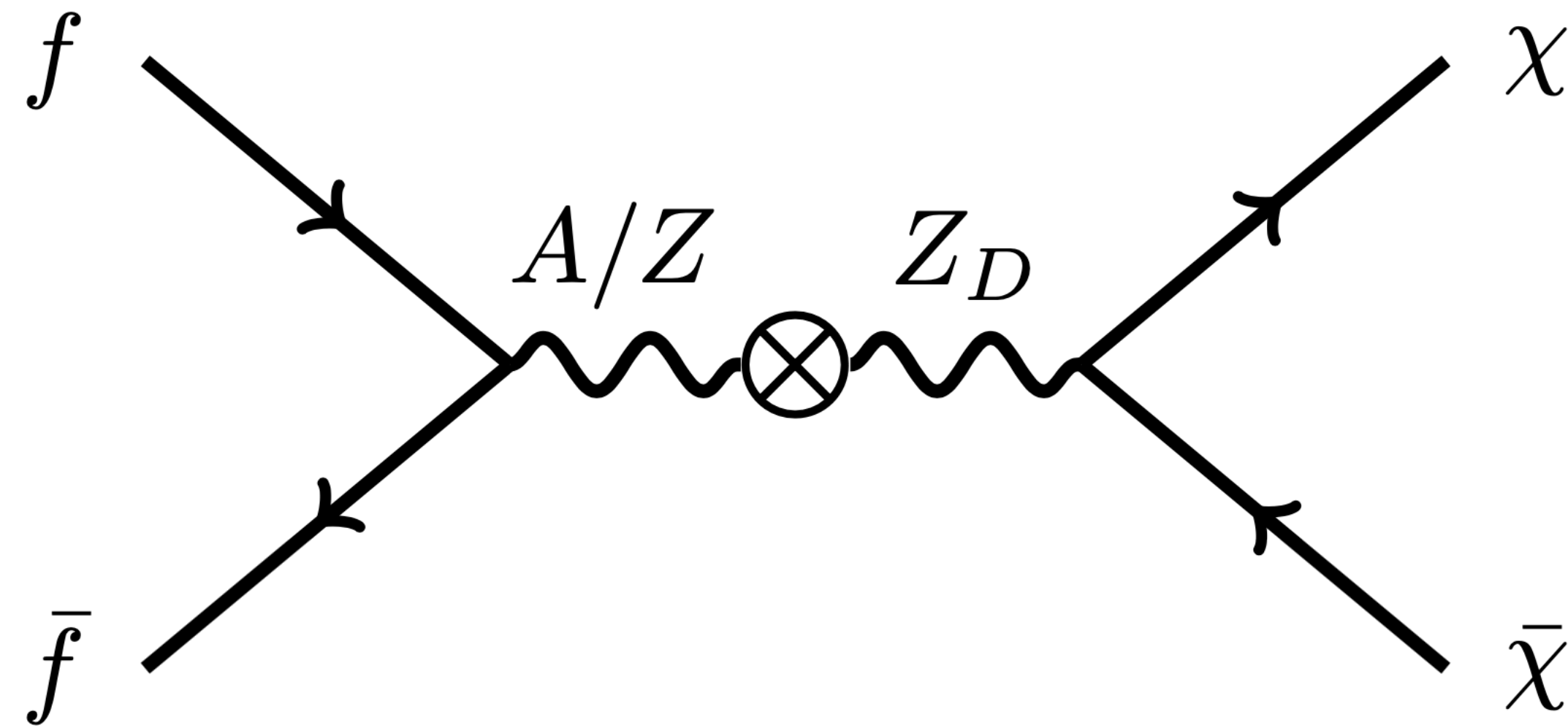
$$\mathcal{L} \supset -\epsilon e J_{\text{EM}}^\mu Z_{D\mu} + \epsilon g_\chi \tan\theta_W J_D^\mu Z_\mu + g_\chi J_D^\mu Z_{D\mu},$$

[X. Chu, T.H., M. Tytgat '11]

Explicit Model: Kinetic mixing portal

$$\mathcal{L} \supset -\epsilon e J_{\text{EM}}^\mu Z_{D\mu} + \epsilon g_\chi \tan \theta_W J_D^\mu Z_\mu + g_\chi J_D^\mu Z_{D\mu},$$

$$1\text{MeV} < m_\chi < 1\text{GeV}$$

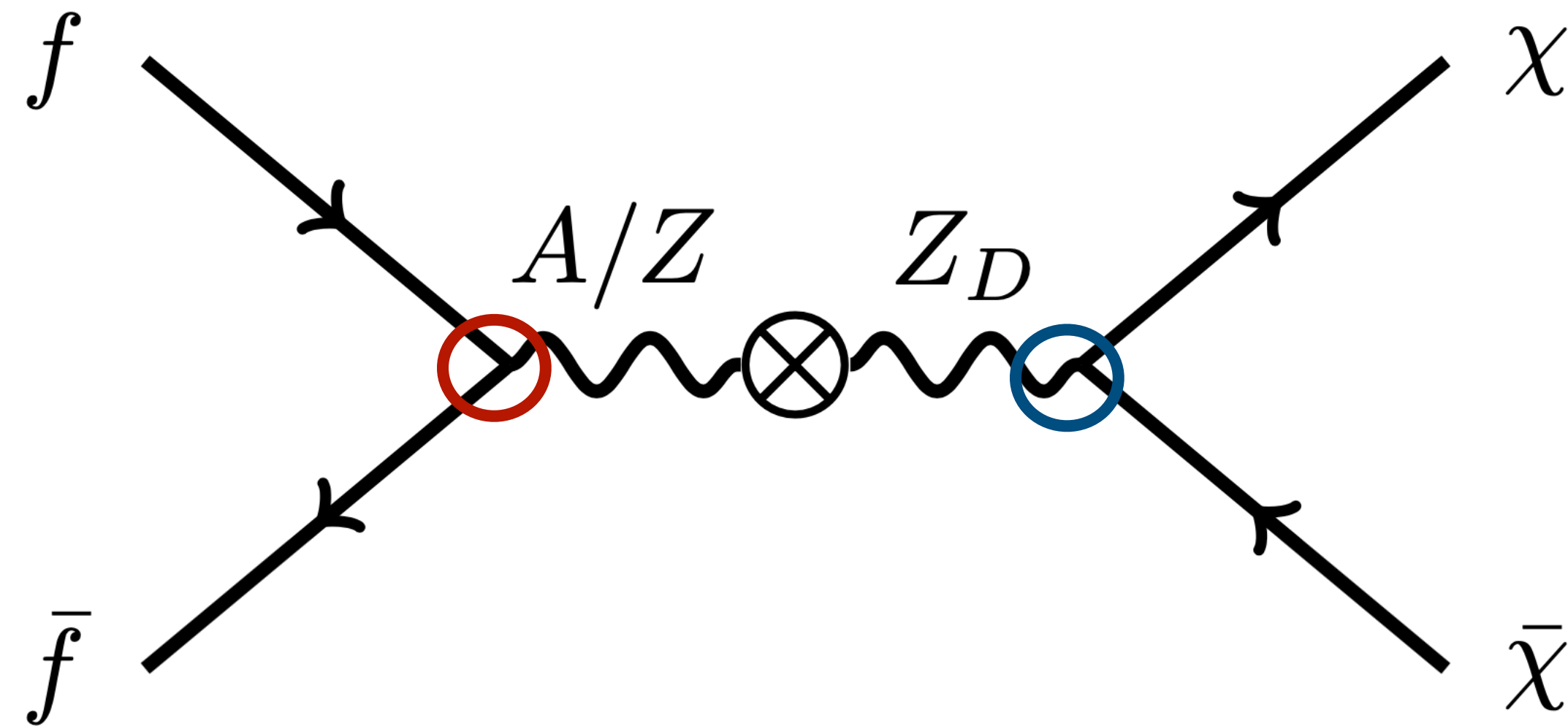


Explicit Model: Kinetic mixing portal

$$\mathcal{L} \supset -\epsilon e J_{\text{EM}}^\mu Z_{D\mu} + \epsilon g_\chi \tan \theta_W J_D^\mu Z_\mu + g_\chi J_D^\mu Z_{D\mu},$$

$$\alpha_D = \frac{g_\chi^2}{4\pi}$$

$$1\text{MeV} < m_\chi < 1\text{GeV}$$



$$\langle \sigma v \rangle \propto \epsilon^2 \alpha_D$$

Scenario

- We explore how a pre-existing population of DM, either alone or as part of a thermalized dark sector, affects the dynamics of freeze-in.
- For a kinetically mixed dark photon, the dominant source of energy injection into the hidden sector is through DM pair production (thermal corrections).
- Elastic processes are fast enough for instantaneous kinetic equilibrium.

Boltzmann Equation

Number density of DM:

$$\dot{n}_\chi + 3Hn_\chi = -\langle\sigma v\rangle_{fo}^{\tilde{T}}(n_\chi^2 - n_{eq}^2(\tilde{T})) + \langle\sigma v\rangle_{fi}^T n_{eq}^2$$

Energy density of the HS:

$$\dot{\rho}_{HS} + 3H(\rho_{HS} + P_{HS}) = \langle\sigma v E\rangle_{fi}^T n_{eq}^2(T)$$



Boltzmann Equation

$$\chi\chi \leftrightarrow Z_D Z_D$$

Number density of DM:

$$\dot{n}_\chi + 3Hn_\chi = -\langle\sigma v\rangle_{f_o}^{\tilde{T}}(n_\chi^2 - n_{eq}^2(\tilde{T})) + \langle\sigma v\rangle_{f_i}^T n_{eq}^2$$

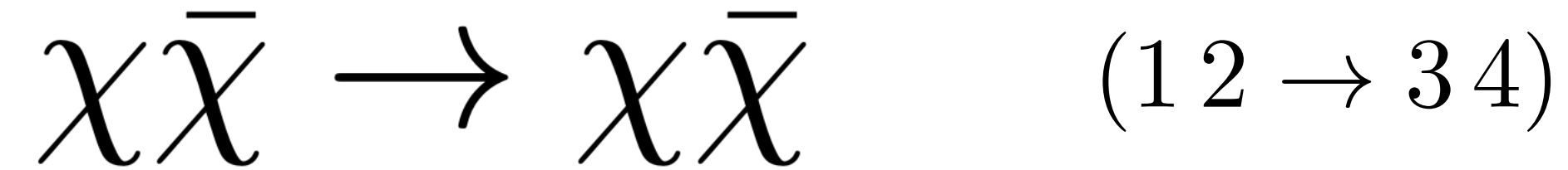
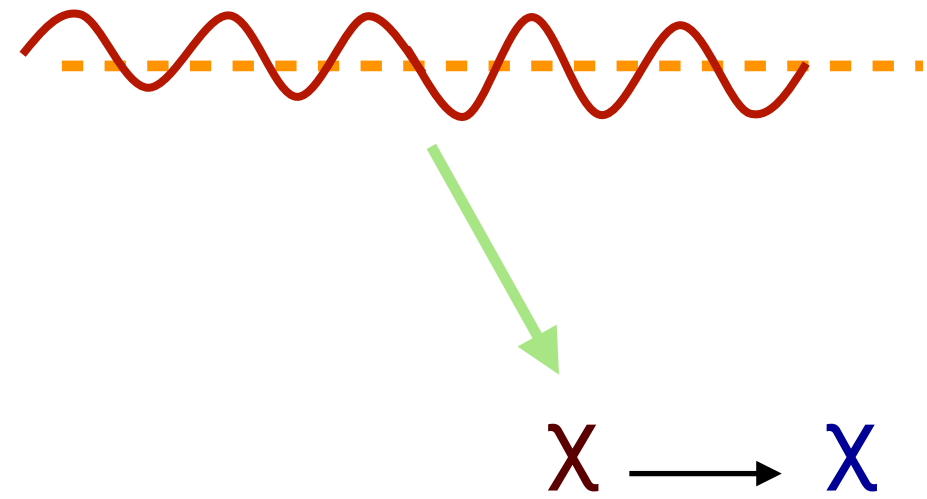
Energy density of the HS:

$$\dot{\rho}_{HS} + 3H(\rho_{HS} + P_{HS}) = \langle\sigma v E\rangle_{f_i}^T n_{eq}^2(T)$$

$$f\bar{f} \rightarrow \chi\chi$$



Instantaneous kinetic equilibration of DM



$$T_1 \neq T_2$$

$$x_i = \frac{m_i}{T_i}$$

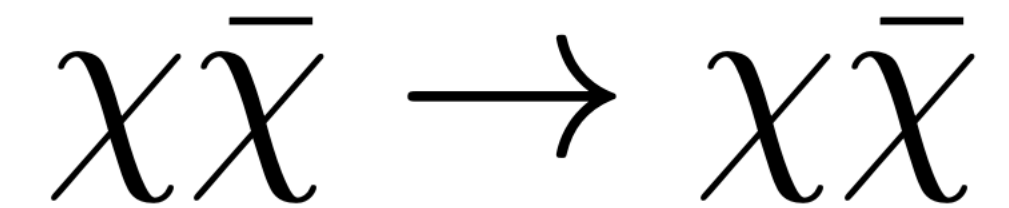
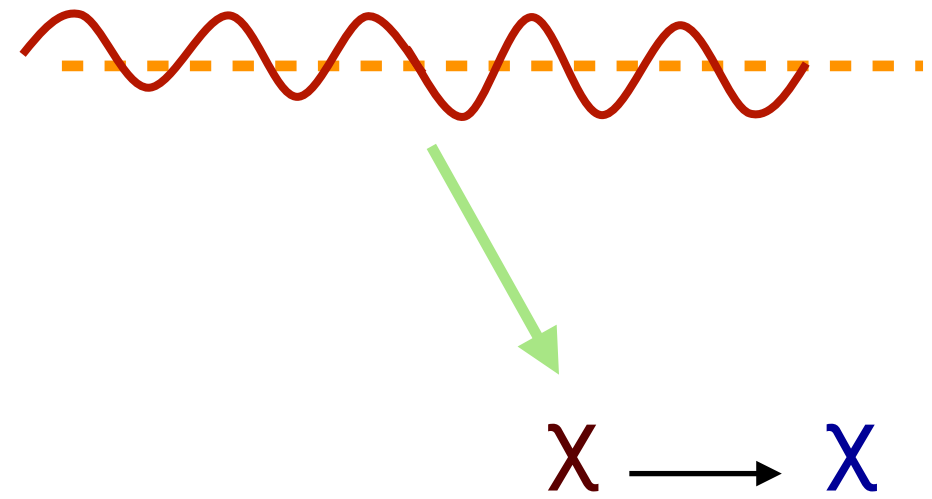
Momentum transferred:

$$\begin{aligned} C_{12 \rightarrow 34}^p(T, \tilde{T}) &= n_1^{\text{eq}}(T) n_2^{\text{eq}}(\tilde{T}) \langle \sigma v p \rangle \\ &= -\frac{g_1 g_2 T^4 \tilde{T}^3}{32\pi^4} \int_{\tilde{s}_{\min}}^{\infty} d\tilde{s} \frac{\lambda^{\frac{1}{2}}(\tilde{s}^2, x_1, x_2)}{\tilde{s}} \sigma(s) \left(\lambda(\tilde{s}^2, x_1, x_2) K_2(\tilde{s}) + 4\tilde{s} x_1^2 K_1(\tilde{s}) \right) \\ &\quad s = \tilde{s}^2 T \tilde{T} + (T - \tilde{T})(T x_1^2 - \tilde{T} x_2^2), \quad \tilde{s}_{\min} = x_1 + x_2 \end{aligned}$$

Turning the tables

Have you seen this formula?

Instantaneous kinetic equilibration of DM



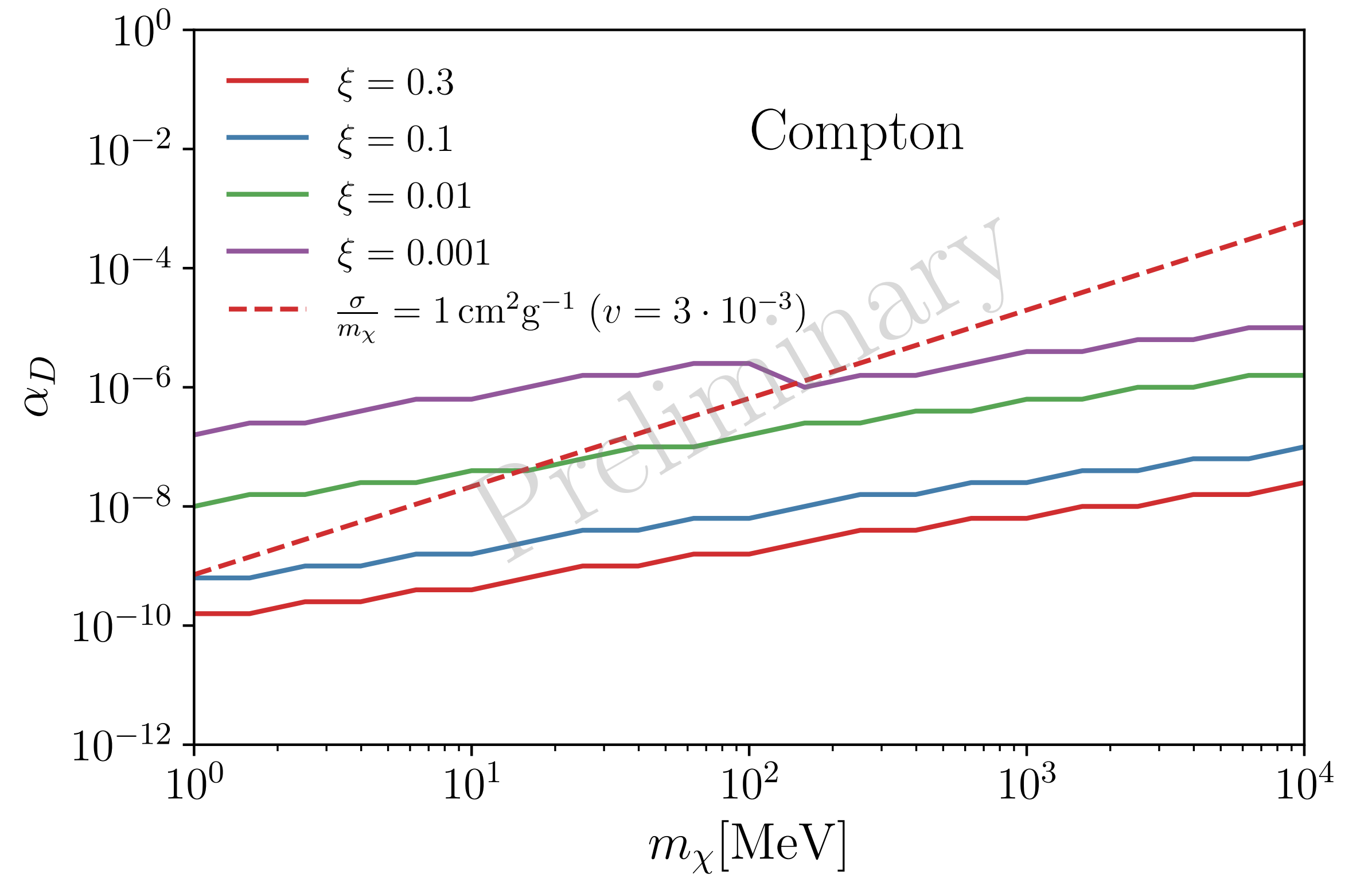
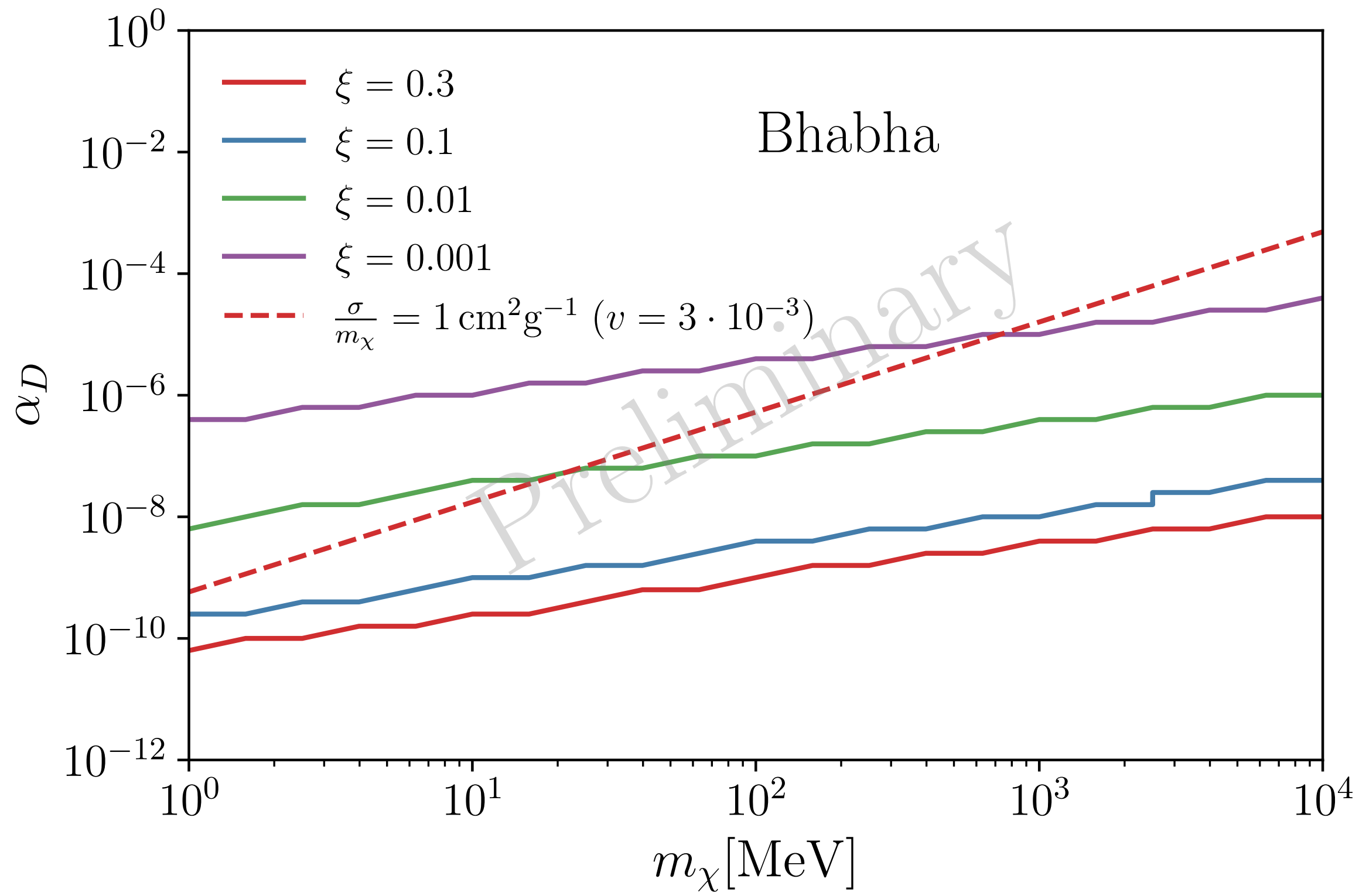
(1 2 \rightarrow 3 4)

$T_1 \neq T_2$

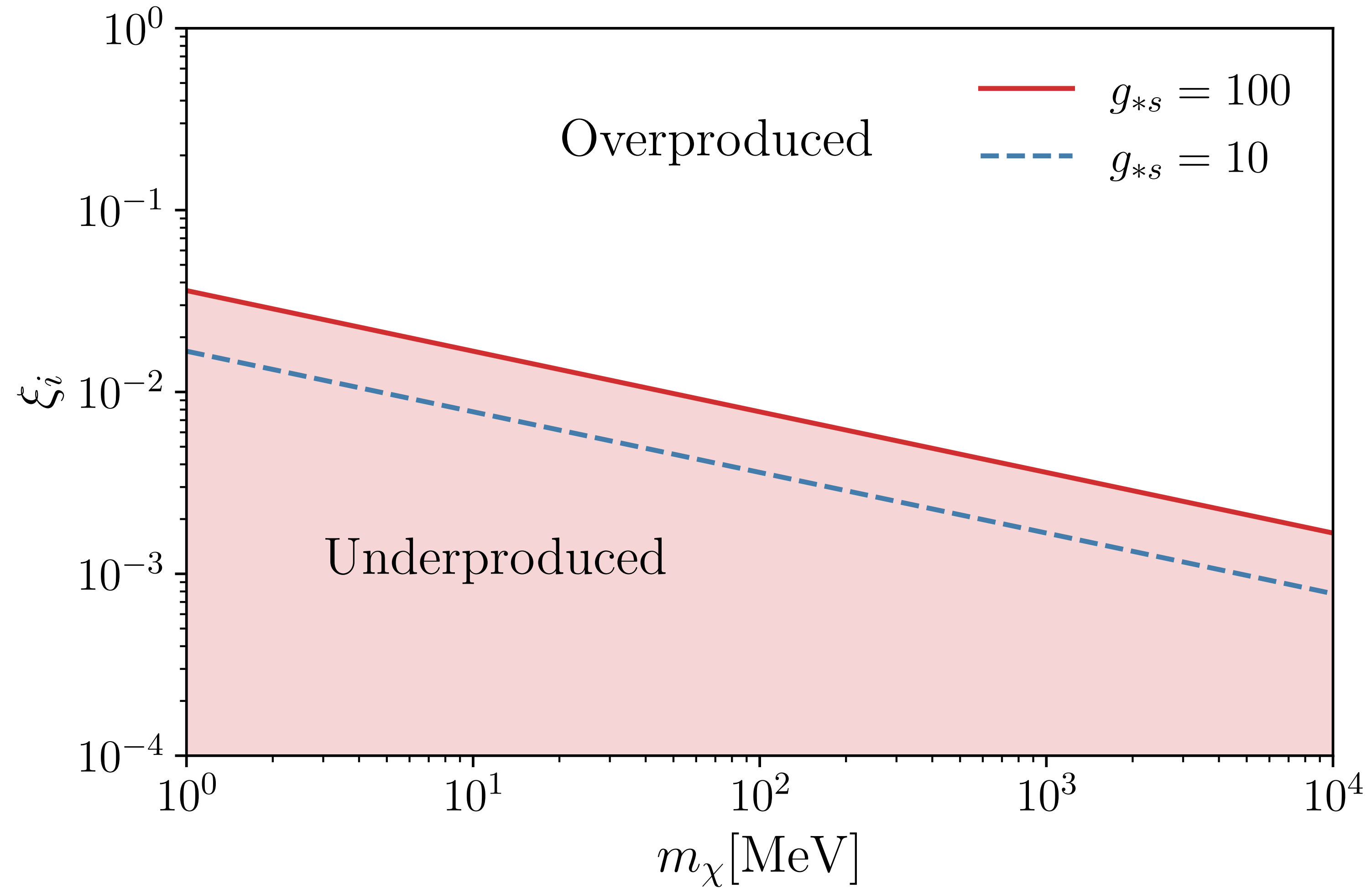
Thermally averaged momentum loss:

$$\Gamma_{p \text{ loss}} \approx \left\langle \frac{dp}{dt} \right\rangle \frac{1}{\langle p \rangle} = \frac{n_{2\text{eq}}(\tilde{T}) \langle \sigma_T v p \rangle}{\langle p \rangle}$$

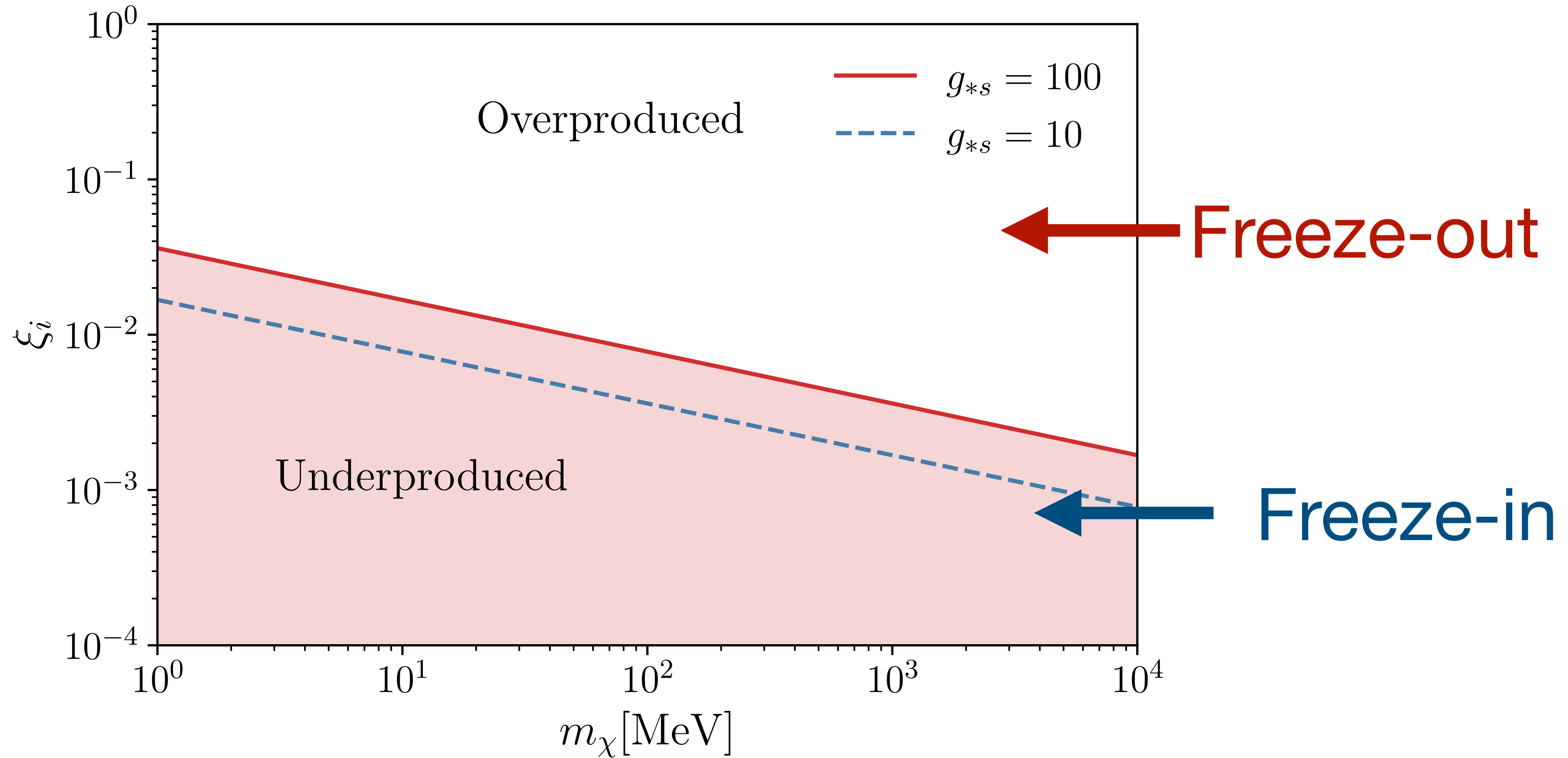
Instantaneous kinetic equilibration of DM



Initial DM

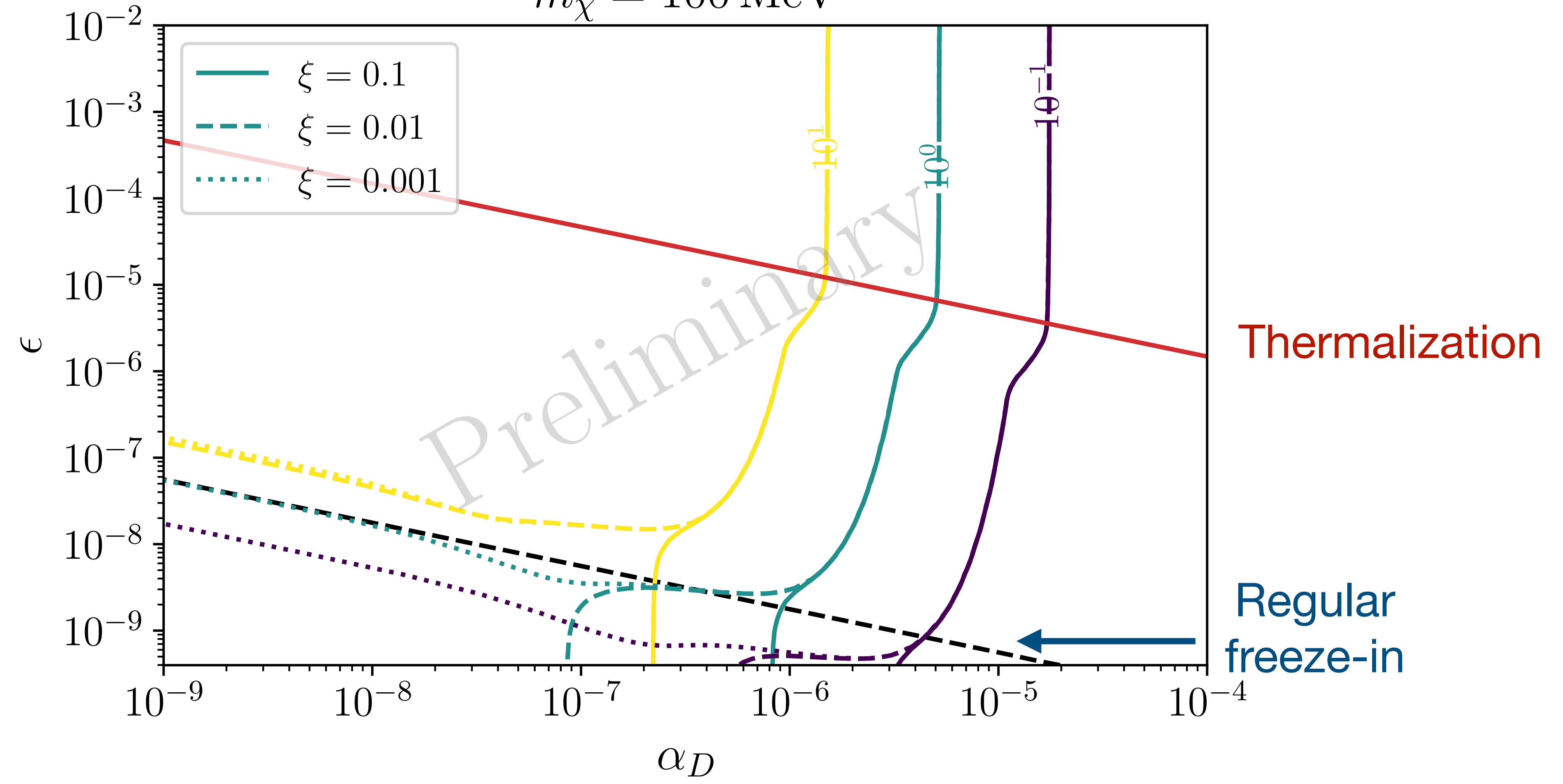


Freeze-in or freeze-out?

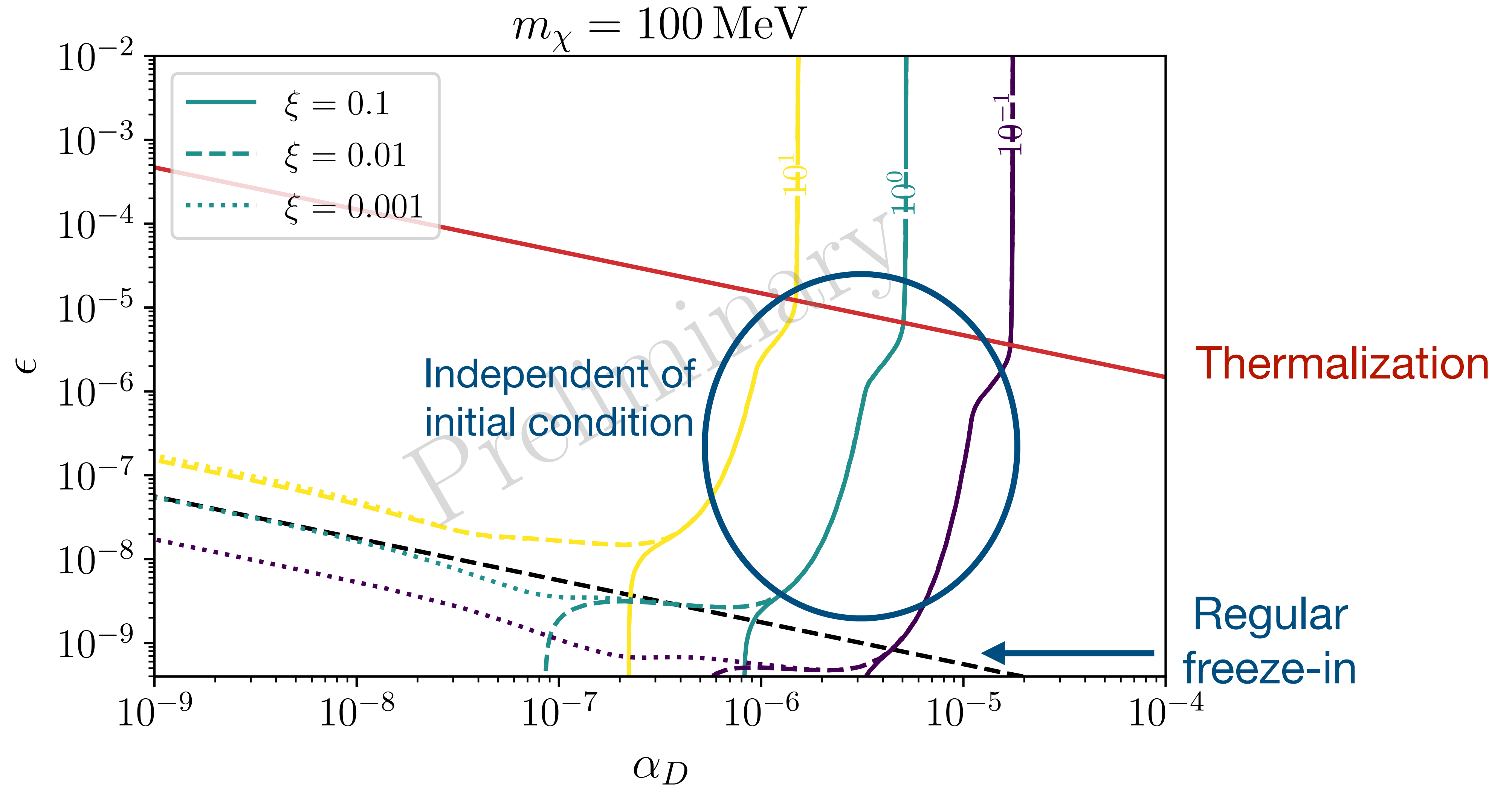


Instantaneous kinetic equilibration of DM

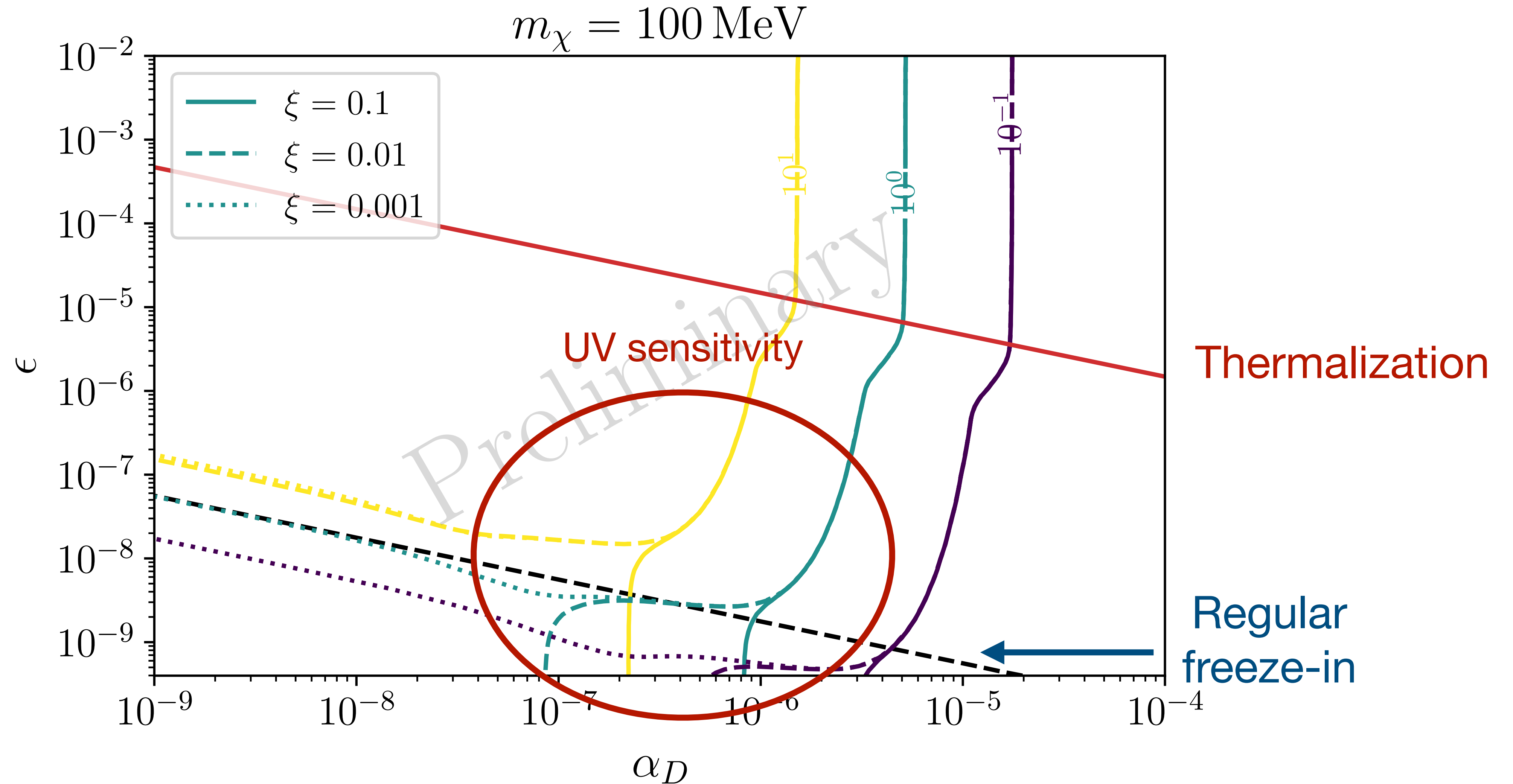
$m_\chi = 100 \text{ MeV}$



Instantaneous kinetic equilibration of DM



Instantaneous kinetic equilibration of DM



Conclusion

- The standard freeze-in paradigm, this same combination of couplings appears in the annihilation cross section, leading to a 1-to-1 relation between thermal history parameter space and direct detection parameter space. As soon as one allows for an initial thermalized population in the dark sector, this “freeze-in line” expands to a “glaciation band” because there are multiple points in the $\varepsilon - \alpha_D$ plane which achieve the correct relic abundance.

Obrigado!