

Distinguishing Random States and Black Holes

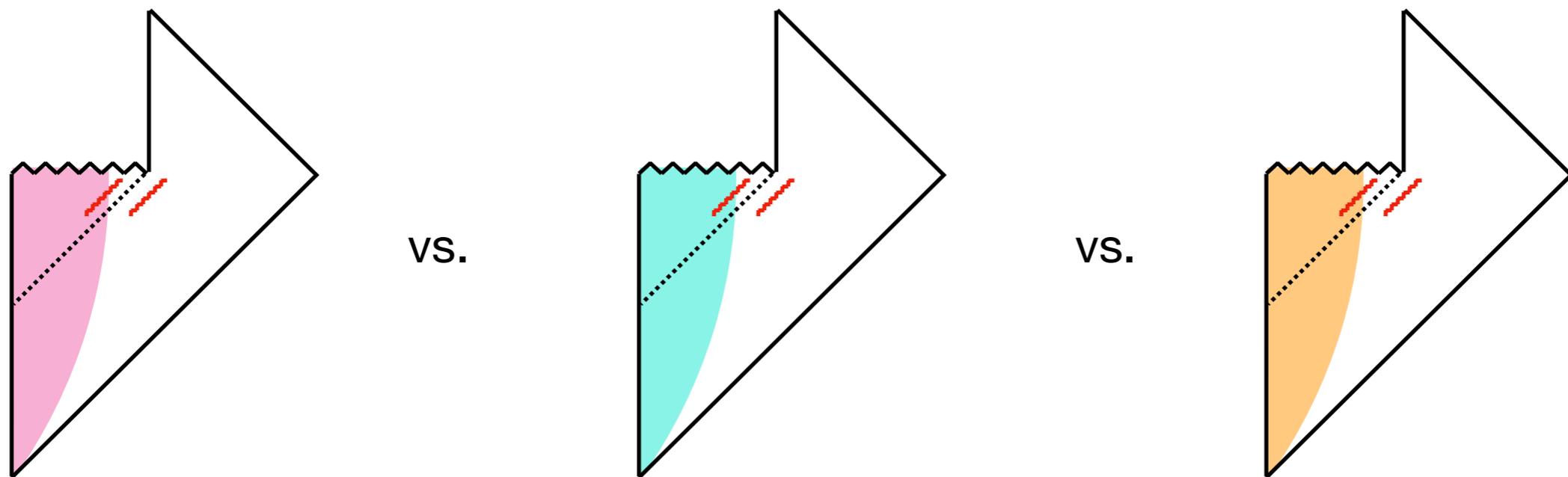
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June 23, 2021
Strings 2021

Based on:

1. arXiv:2102.05053 [hep-th] Phys. Rev. Lett. 126, 171603 (2021)
2. arXiv:210X.xxxxx (w/ V. Narovlansky & S. Ryu)

A Black Hole Information Problem

- Hawking told us that black holes will radiate **thermal** radiation
- This is a breakdown of quantum mechanics (unitarity) — **different** black hole microstates (with same M, Q, J) evolve to the same final state
- How can we **distinguish** different black hole microstates from the radiation?
- Must probe physics **beyond** the **Page curve** for von Neumann entropy



Quantum Hypothesis Testing

- In quantum information theory, this problem is referred to as **quantum hypothesis testing**
- What is the best we can do at distinguishing states ρ and σ using a **quantum measurement** (POVM)?
- For POVM $\{\hat{A}, 1 - \hat{A}\}$, we conclude we have ρ if we get measurement outcome corresponding to \hat{A} and σ otherwise
- **Type I error**: probability of concluding we have σ when we really have ρ , $\alpha(\hat{A}) := \text{Tr}[(1 - \hat{A})\rho]$. **Type II error**: probability of concluding we have ρ when we really have σ , $\beta(\hat{A}) := \text{Tr}[\hat{A}\sigma]$.
- We optimize our POVM to minimize the error probability:

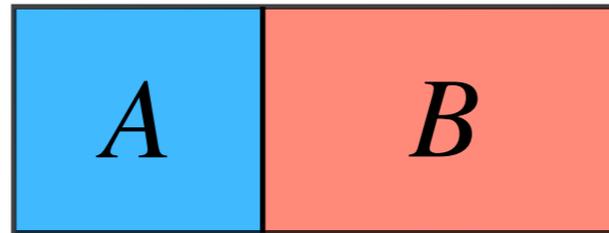
$$\min_{\hat{A}} \left[\alpha(\hat{A}) + \beta(\hat{A}) \right] := 1 - T(\rho || \sigma)$$

Trace distance

[Helstrom (1969)]

Replica tricks for distinguishability measures

$$\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$$



Petz Renyi Relative Entropy measures the rate at which states can be distinguished as one is given more copies

$$D_\alpha(\rho_A || \sigma_A) = \lim_{m \rightarrow 1-\alpha} \frac{1}{\alpha - 1} \log \left[\text{Tr} \left[\rho_A^\alpha \sigma_A^m \right] \right]$$

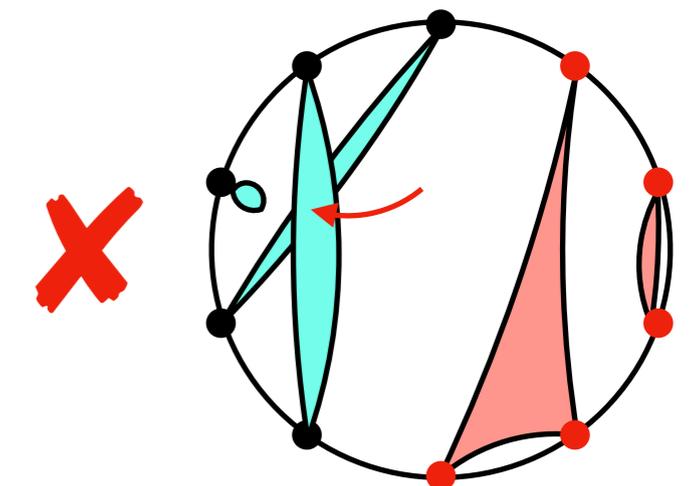
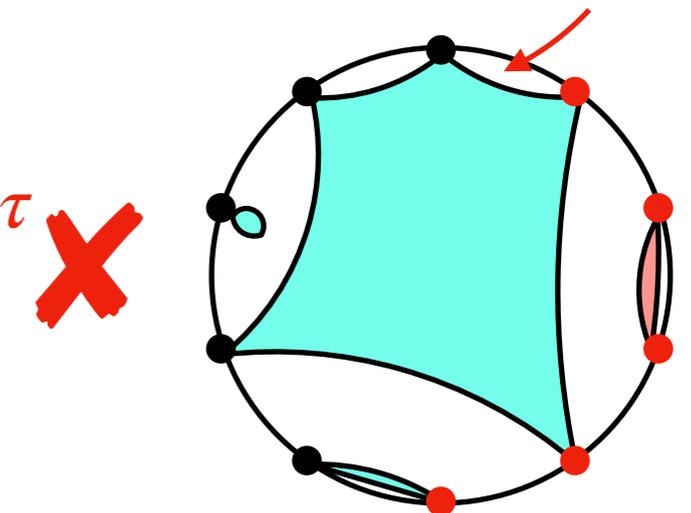
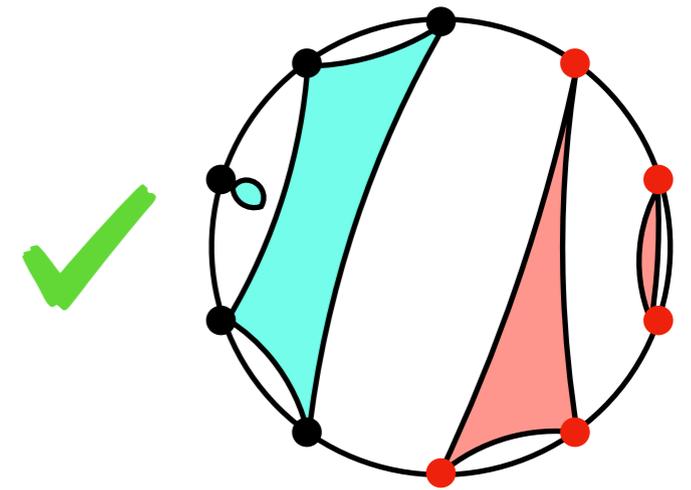
ρ_A and σ_A are independently Haar random

$$\overline{\text{Tr} \left[\rho_A^\alpha \sigma_A^m \right]} = \frac{1}{(d_A d_B)^{\alpha+m}} \sum_{\tau \in \mathcal{S}_\alpha \times \mathcal{S}_m} d_A^{C(\eta^{-1} \circ \tau)} d_B^{C(\tau)}$$

of cycles in τ

Can be exactly evaluated using the combinatorics of non-crossing partitions

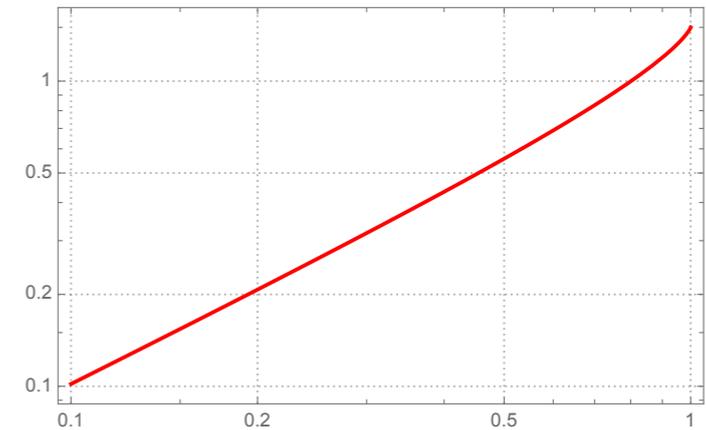
$$D_\alpha(\rho_A || \sigma_A) = \frac{1}{\alpha - 1} \begin{cases} \log \left[{}_2F_1 \left(1 - \alpha, -\alpha; 2; \frac{d_A}{d_B} \right) {}_2F_1 \left(\alpha - 1, \alpha; 2; \frac{d_A}{d_B} \right) \right], & d_A < d_B \\ \log \left[\frac{d_B {}_2F_1 \left(1 - \alpha, -\alpha; 2; \frac{d_B}{d_A} \right) {}_2F_1 \left(\alpha - 1, \alpha; 2; \frac{d_B}{d_A} \right)}{d_A} \right], & d_A > d_B \end{cases}$$



Page Curves for Relative Entropy

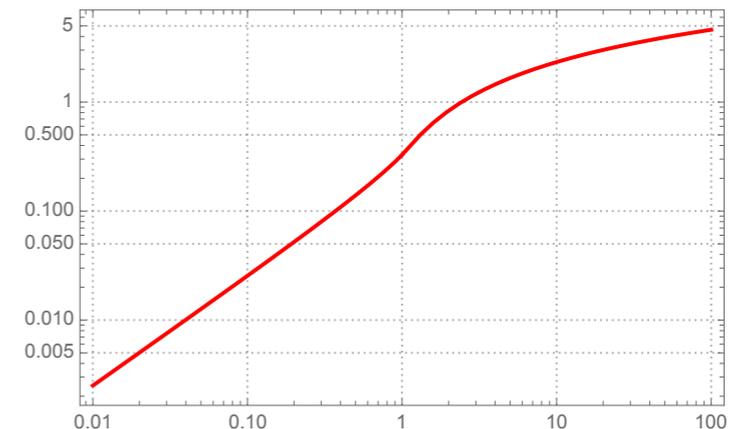
Relative Entropy:

$$D(\rho_A || \sigma_A) = \begin{cases} 1 + \frac{d_A}{2d_B} + \left(\frac{d_B}{d_A} - 1\right) \log\left(1 - \frac{d_A}{d_B}\right), & d_A < d_B \\ \infty, & d_A > d_B \end{cases}$$



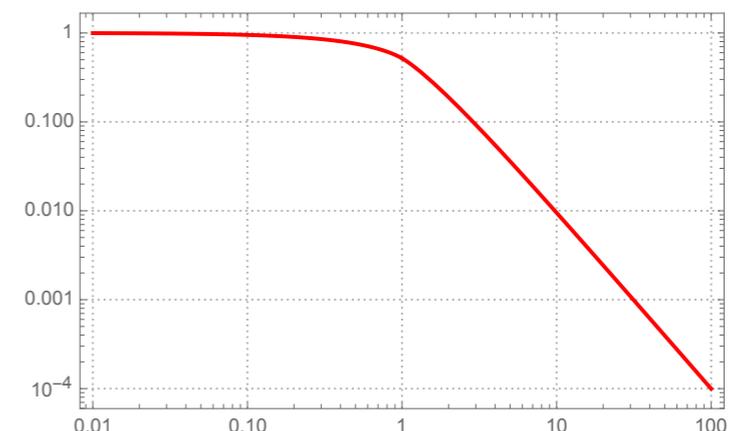
Chernoff Distance:

$$\xi(\rho_A || \sigma_A) = \begin{cases} -2 \log \left[{}_2F_1\left(\frac{1}{2}, -\frac{1}{2}; 2; \frac{d_A}{d_B}\right) \right], & d_A < d_B \\ -\log \left[\frac{d_B}{d_A} {}_2F_1\left(\frac{1}{2}, -\frac{1}{2}; 2; \frac{d_B}{d_A}\right)^2 \right], & d_A > d_B \end{cases}$$



Fidelity:

$$F_H(\rho_A || \sigma_A) = \begin{cases} {}_2F_1\left(\frac{1}{2}, -\frac{1}{2}; 2; \frac{d_A}{d_B}\right)^4, & d_A < d_B \\ \frac{d_B^2}{d_A^2} {}_2F_1\left(\frac{1}{2}, -\frac{1}{2}; 2; \frac{d_B}{d_A}\right)^4, & d_A > d_B \end{cases}$$



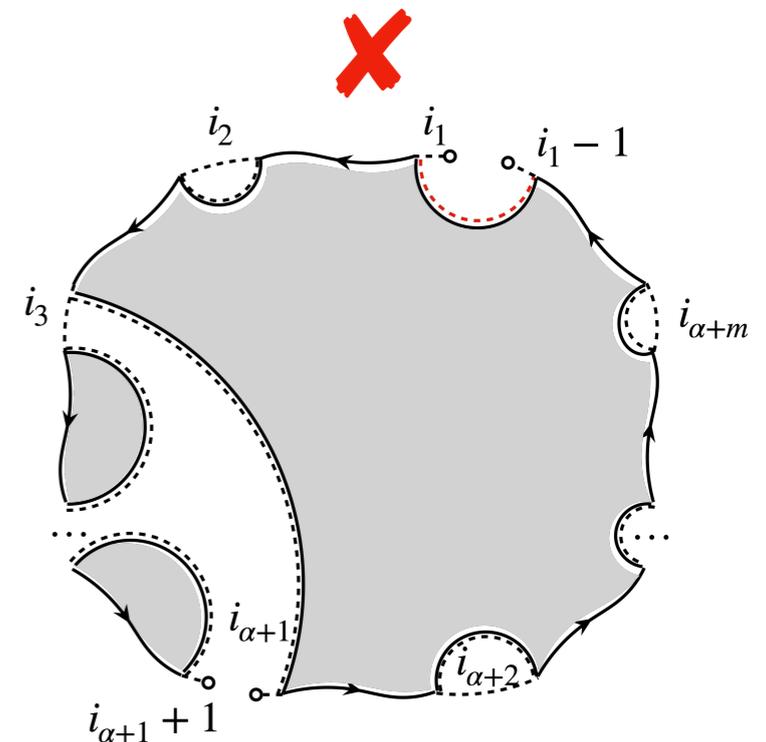
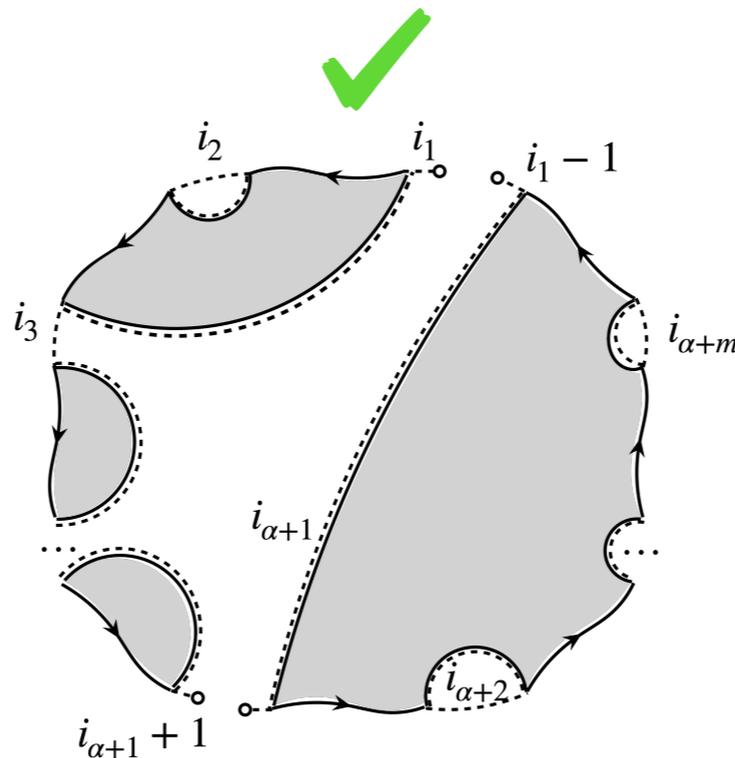
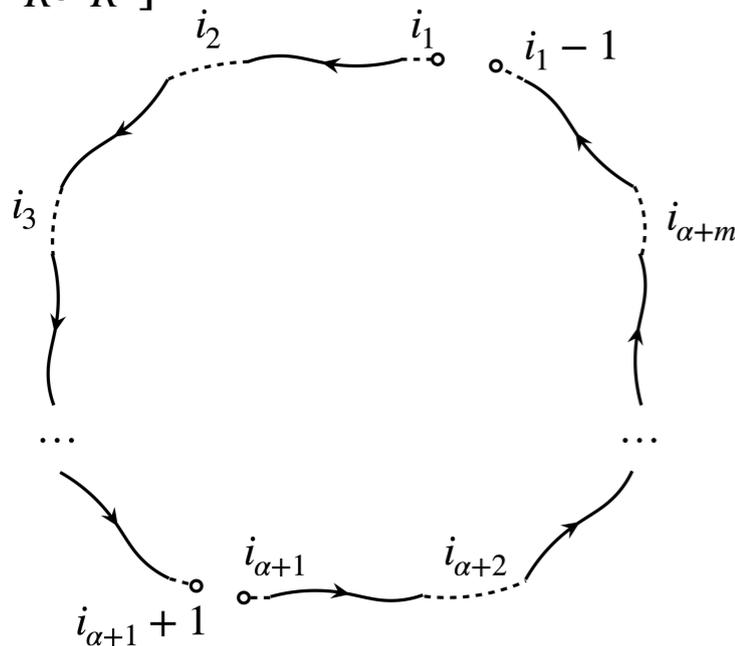
West Coast Model

- We use the PSSY (west coast) model as a toy model of black hole evaporation – JT gravity decorated with non dynamical end-of-world branes with k flavors

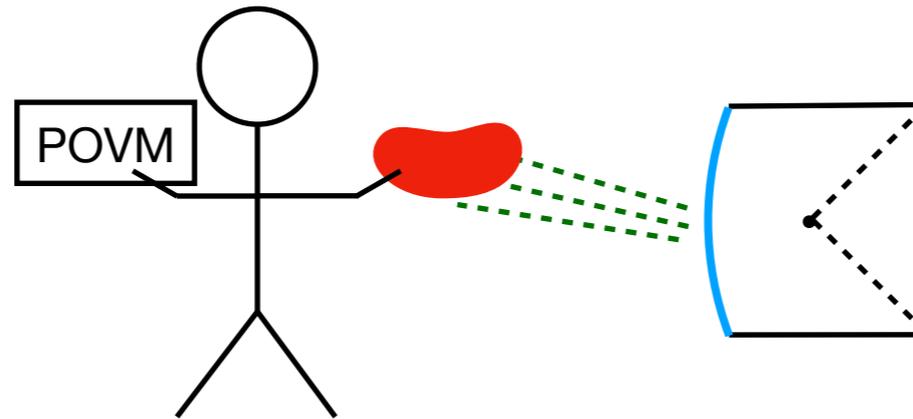
$$|\Psi\rangle = \frac{1}{\sqrt{k}} \sum_{i=1}^k |\psi_i\rangle_B |i\rangle_R \quad |\Psi'\rangle = \frac{1}{\sqrt{k}} \sum_{i=1}^k |\psi_i\rangle_B |i+1\rangle_R$$

$$\langle \psi_i | \psi_j \rangle = \begin{array}{c} i \text{-----} j \\ \curvearrowright \end{array}$$

$$\text{Tr} [\rho_R^\alpha \rho_R^{\prime m}] =$$



Implications



- Before the Page time, we need $O\left(e^{S_{BH}-S_R} \log[\epsilon^{-1}]\right)$ copies of the radiation to distinguish the microstates with error ϵ
- After the Page time, we only need a **single copy** of the radiation to distinguish the microstates with exponentially small error. However, this will be a “complex” measurement
- If we did not include the replica wormholes in the gravitational path integral, all relative entropies would be **zero** at all times i.e. **information loss**

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- “Page curves” for relative entropy, Petz Renyi Relative entropy, Sandwiched Renyi Relative entropy, trace distance, fidelity — agreement with small matrices
- Characterization of **distinguishability of black hole microstates** in AdS/CFT from subsets of boundary data (with and without fixed areas) — correction to JLMS
- Solution to distinguishing **generic random tensor network states** using flow network techniques
- **Derivation of subsystem ETH** (up until $f = 1/2$) for generic chaotic systems and holographic CFTs in all dimensions

Not in this talk but in:

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