EXTRACTING FUNDAMENTAL PHYSICS OUT OF THE LARGE-SCALE STRUCTURE OF THE UNIVERSE

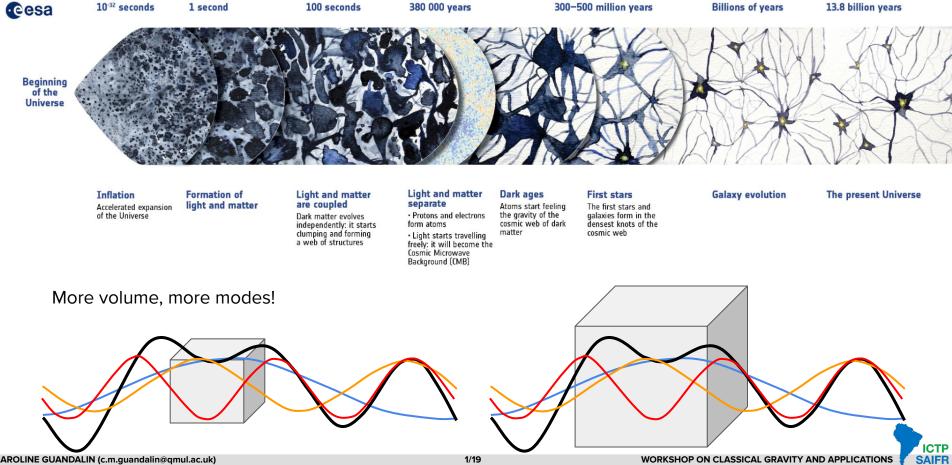
CAROLINE GUANDALIN (QMUL)

WORKSHOP ON CLASSICAL GRAVITY AND APPLICATIONS

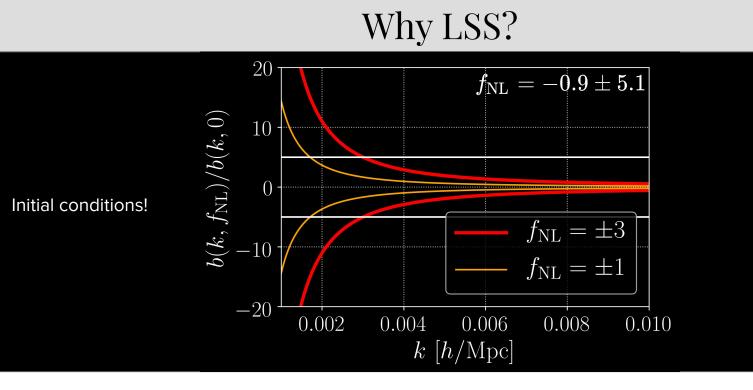


|International Centre for Theoretical Physics |South American Institute for Fundamental Research

Why LSS?



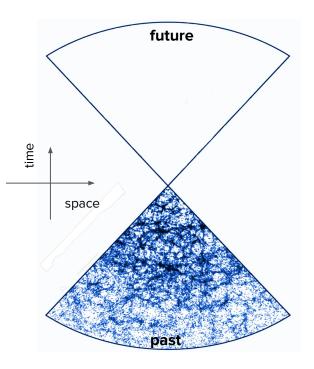
CAROLINE GUANDALIN (c.m.guandalin@qmul.ac.uk)



More volume, more modes!

Larger scales!

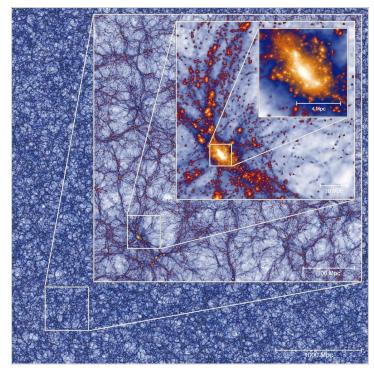
Why LSS?



More volume, more modes!

Larger scales!

Why LSS?



More volume, more modes!

Larger scales!

Credit: Millennium-XXL

Nonlinear physics!



General considerations

Understanding the large-scale distribution of dark matter

$$\delta(oldsymbol{x},t) = rac{
ho(oldsymbol{x},t)-ar{
ho}(t)}{ar{
ho}(t)} \Rightarrow \Delta(oldsymbol{n},z) = rac{N(oldsymbol{n},z)-ar{N}(z)}{ar{N}(z)}$$

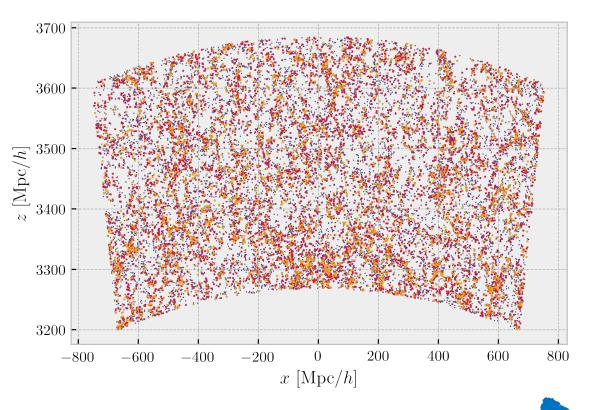
Different objects will trace the dark matter distribution differently.

$$P_lpha(k,z)=b_lpha^2(z)\,P_m(k,z)$$

Understanding the large-scale distribution of dark matter

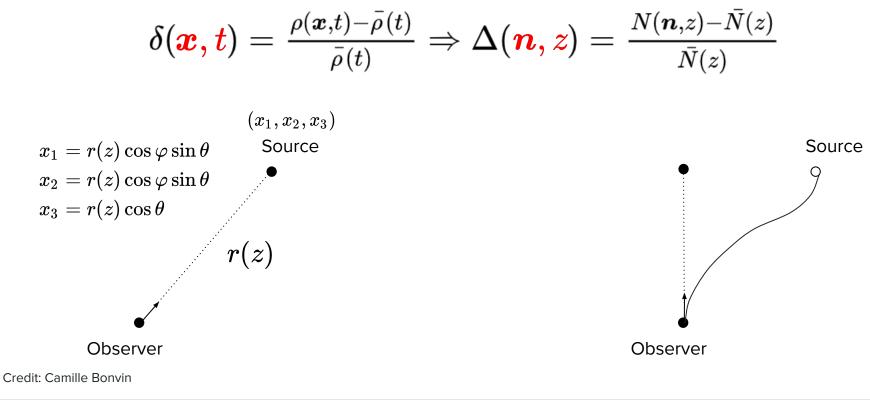
 $1.7 \leq z \leq 2.1$

	# halos	Mean mass $[M_{\odot}/h]$	Bias (fit)	$ar{n}(ar{z})$ $[Mpc/h]^{-3}$
		$\bar{z} = 1.89$)	
All	480643	4.41×10^{12}	2.927	7.053×10^{-4}
H_0	160081	1.86×10^{12}	2.551	2.349×10^{-4}
H_1	160547	2.83×10^{12}	2.758	2.356×10^{-4}
H ₂	160015	8.54×10^{12}	3.477	2.348×10^{-4}
		$\bar{z} = 2.29$)	
All	326899	3.85×10^{12}	3.469	4.666×10^{-4}
H_0	109003	1.83×10^{12}	3.020	1.556×10^{-4}
H_1	108809	2.66×10^{12}	3.270	1.553×10^{-4}
H_2	109087	7.05×10^{12}	4.154	1.557×10^{-4}
		$\bar{z} = 2.69$)	
All	205678	3.44×10^{12}	4.214	2.947×10^{-4}
H_0	68501	1.80×10^{12}	3.735	9.815×10^{-5}
H_1	68550	2.52×10^{12}	4.006	9.822×10^{-5}
H_2	68627	5.98×10^{12}	4.932	9.833×10^{-5}



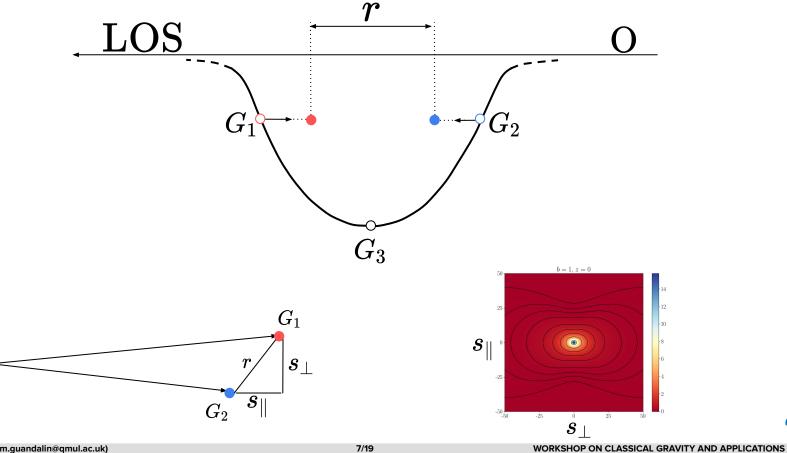
FAPESP 2018/10396-2

Understanding the large-scale distribution of dark matter



ICTP

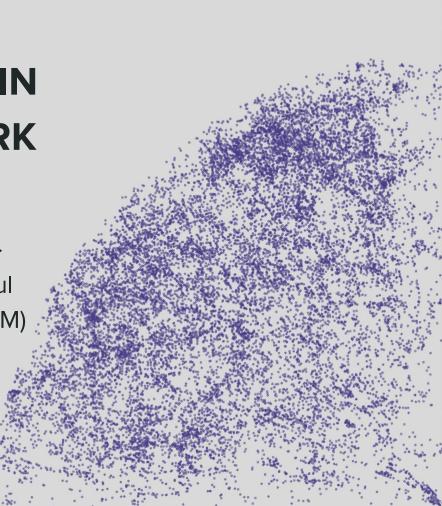
Redshift-space distortions



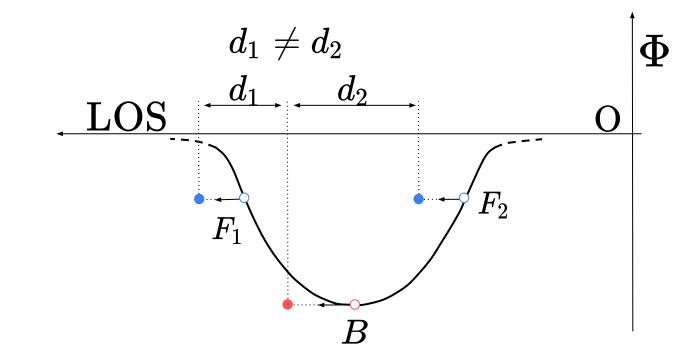
RELATIVISTIC FEATURES IN THE CLUSTERING OF DARK MATTER HALOS

Julian Adamek (UZH), Phil Bull (QMUL → Manchester), Chris Clarkson (QMUL), Raul Abramo (IFUSP) & Louis Coates (PhD at QM)

> 2009.02284 [MNRAS, 501(2), 2021, 2547-2561]



Relativistic effects: breaking the symmetry of the correlation function



Bonvin [1409.2224]

The theoretically observed number counts

e.g. Durrer 2021, Bonvin 2014

$$\Delta(\boldsymbol{n},z) = b_{\alpha}D - \frac{1}{\mathcal{H}}\partial_{r}(\boldsymbol{v}\cdot\hat{\boldsymbol{n}}) + \frac{1}{\mathcal{H}}\dot{\boldsymbol{v}}\cdot\hat{\boldsymbol{n}} + \left(1 - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - \frac{2-5s}{r_{s}\mathcal{H}} - 5s + b_{e}\right)\boldsymbol{v}\cdot\hat{\boldsymbol{n}} + \frac{1}{\mathcal{H}}\partial_{r}\Phi$$
$$+ \frac{5s-2}{2r_{s}}\int_{0}^{r_{s}}\mathrm{d}r \,\frac{r_{s}-r}{r}\Delta_{\Omega}(\Phi+\Psi) + \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} + \frac{2-5s}{r_{s}\mathcal{H}} + 5s - b_{e}\right)\left[\Phi + \int_{0}^{r}\mathrm{d}r'\,(\dot{\Psi}+\dot{\Phi})\right]$$
$$+ \frac{1}{\mathcal{H}}\dot{\Psi} + \Phi - (2-5s)\Psi + \frac{2-5s}{2r_{s}}\int_{0}^{r_{s}}\mathrm{d}r\,\left[2 - \frac{r_{s}-r}{r}\Delta_{\Omega}\right](\Psi+\Phi)$$

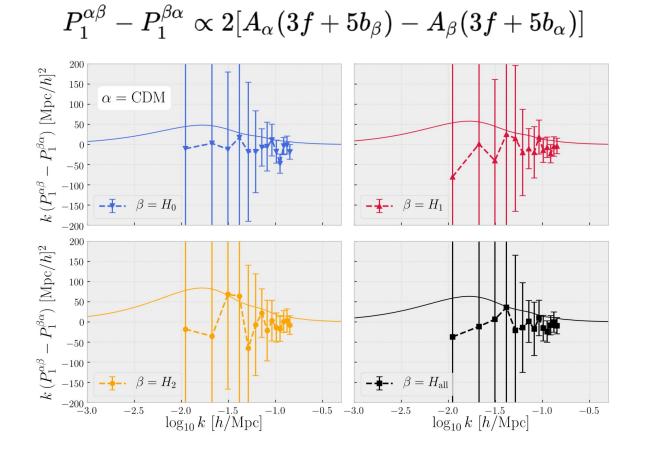
$$\delta_{\alpha}^{(s)}(\boldsymbol{n},z) = b_{\alpha}D - \frac{1}{\mathcal{H}}\partial_{r}(\boldsymbol{v}\cdot\hat{\boldsymbol{n}}) + \left(b_{e}^{\alpha} - 5s_{\alpha} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - \frac{2 - 5s_{\alpha}}{r_{s}\mathcal{H}}\right)\boldsymbol{v}\cdot\hat{\boldsymbol{n}}$$

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Dipole in the cross-correlation

$$\delta_{\alpha}^{(s)}(\boldsymbol{n},z) = D_{l}^{(g)} - \frac{1}{\mathcal{H}}\partial_{r}(\boldsymbol{v}\cdot\hat{\boldsymbol{n}}) + \left(b_{e}^{\alpha} - 5s_{\alpha} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - \frac{2 - 5s_{\alpha}}{r_{s}\mathcal{H}}\right)\boldsymbol{v}\cdot\hat{\boldsymbol{n}}$$
$$\delta_{\alpha}^{(s)}(\boldsymbol{k}) = \delta^{(r)}(\boldsymbol{k})\left[b_{\alpha} + f\mu_{\boldsymbol{k}}^{2} + if(\mathcal{H}k^{-1})A_{\alpha}\mu_{\boldsymbol{k}}\right]$$
$$P_{\alpha\beta}^{(s)}(\boldsymbol{k}) = P^{(r)}(k)\left[(b_{\alpha} + f\mu^{2})(b_{\beta} + f\mu^{2}) + A_{\alpha}A_{\beta}f^{2}\mu^{2}\frac{\mathcal{H}^{2}}{k^{2}} + if\mu\left[(b_{\beta} + f\mu^{2})A_{\alpha} - (b_{\alpha} + f\mu^{2})A_{\beta}\right]\frac{\mathcal{H}}{k}\right]$$
$$P_{1}^{(s)}(k) = iP^{(r)}(k)\frac{f}{5}\left[A_{\alpha}(3f + 5b_{\beta}) - A_{\beta}(3f + 5b_{\alpha})\right]\frac{\mathcal{H}}{k}$$

Relativistic effects: attempted measurement



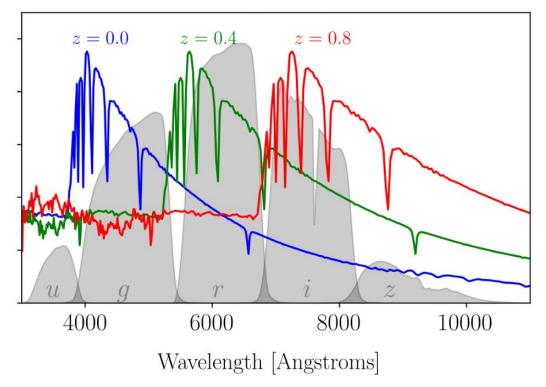
CLUSTERING REDSHIFTS WITH THE GALAXY-HI CROSS-BISPECTRUM

Isabella P. Carucci (UniTO), David Alonso (Oxford) & Kavilan Moodley (UKZN)

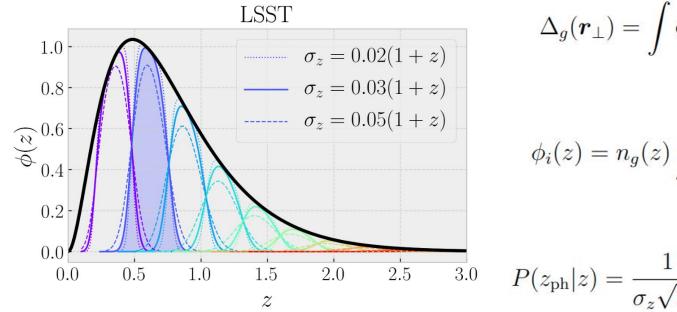
> 2112.05034 (published in MNRAS)

Photometric redshift surveys

Broad-band photometry



Photometric redshifts



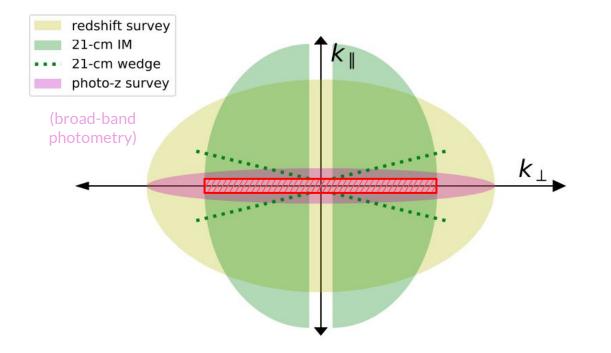
$$\Delta_g(\boldsymbol{r}_{\perp}) = \int \mathrm{d}r_{\parallel} \, \phi(r_{\parallel}) \, \delta_g(r_{\parallel}, \boldsymbol{r}_{\perp})$$

$$\phi_i(z) = n_g(z) \int_{z_{\rm ph}^i}^{z_{\rm ph}^{i+1}} \mathrm{d}z_{\rm ph} P(z_{\rm ph}|z)$$

$$P(z_{\rm ph}|z) = \frac{1}{\sigma_z \sqrt{2\pi}} \exp\left[-\frac{1}{2} \frac{(z_{\rm ph} - z)}{\sigma_z^2}\right]$$

CAROLINE GUANDALIN (c.m.guandalin@qmul.ac.uk)

Problem of cross-correlating photometric galaxies with HI IM

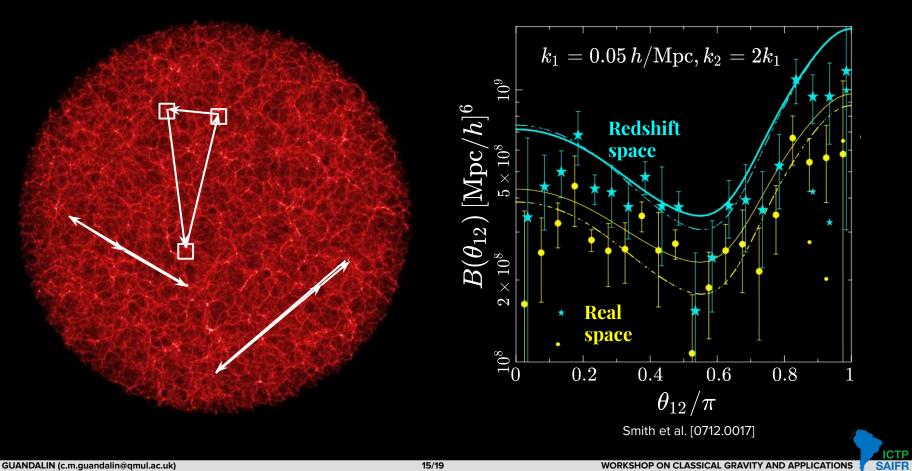


Modi et al. [2102.08116]



Jessica Muir's talk (photo-z)

Bispectrum



2- and 3- point statistics

Power spectrum

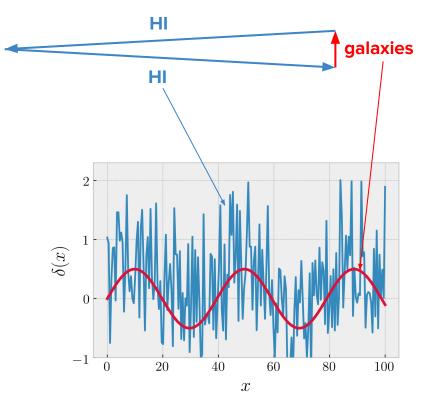
$$P_{xy}(k) = \mathcal{K}_{xy} P_{\rm nl}(k) + N_{xy} \,\delta_{xy}$$

Bispectrum

 $B_{xyz}(\boldsymbol{k}_1, \boldsymbol{k}_2, \boldsymbol{k}_3) = \mathcal{K}_{xyz}B(\boldsymbol{k}_1, \boldsymbol{k}_2, \boldsymbol{k}_3) + \mathcal{J}_{xyz}^{112}P(k_1)P(k_2) + \mathcal{J}_{xyz}^{121}P(k_1)P(k_3) + \mathcal{J}_{xyz}^{211}P(k_2)P(k_3)$

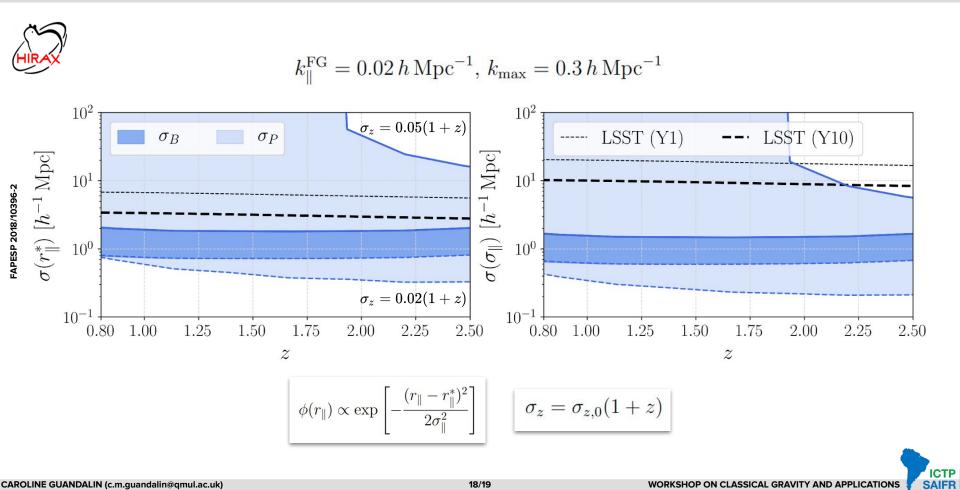


Attempt to reconstruct the long-wavelength modes





Impact of photo-*z* width



Conclusions

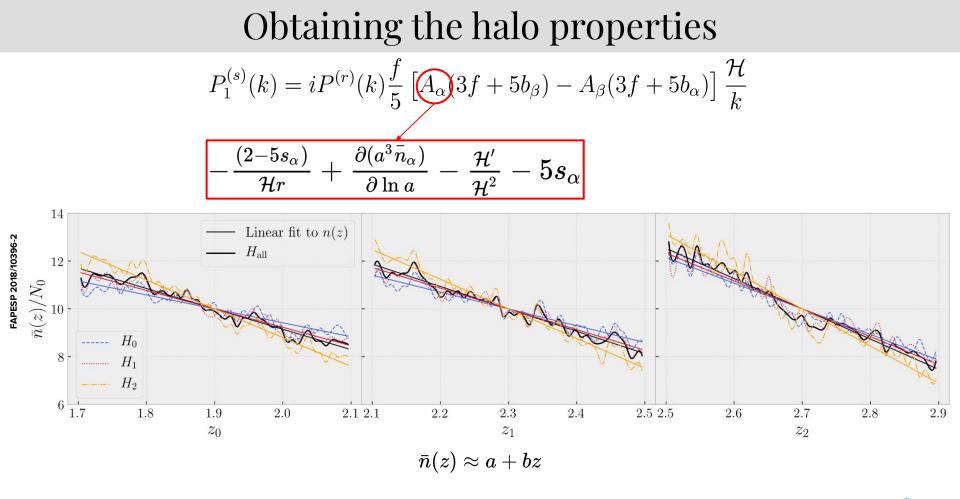
Relativistic effects:

- A dipole signal emerges from relativistic corrections in the cross-power spectrum of dark matter tracers;
- We developed a pipeline to model, estimate and interpret the dipolar modulation induced by relativistic effects in the 3d clustering of halos obtained from a fully relativistic simulation;
- Integrated effects were neglected;
- No dipole detection could be claimed with our simulations.

Clustering redshifts:

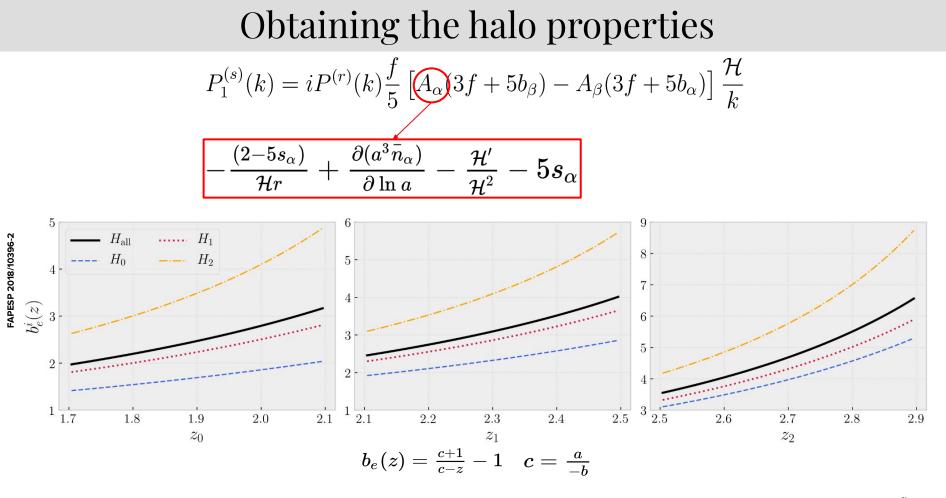
- We employed the multiple-tracers technique to combine photometric galaxies with 21-cm intensity mapping to recover the redshift distribution of a photometric galaxy sample;
- The galaxy-HI bispectra helps to recover the large radial scales apparently lost to foreground contamination;
- The bispectrum is capable of calibrating the redshift distribution in situations where the two-point function is not.

Relativistic effects (extra)



following Beutler & Di Dio [2004.08014]

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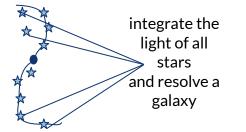


following Beutler & Di Dio [2004.08014]

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SAIFR

Galaxy-HI complementarity (extra)



Photometric surveys

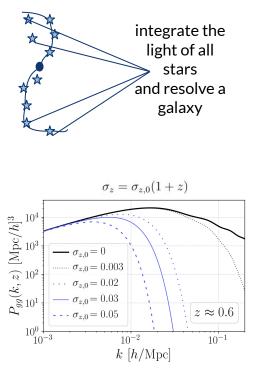
Individually resolved galaxies

High number density

Poorer redshifts

Low-pass k-filter





Photometric surveys

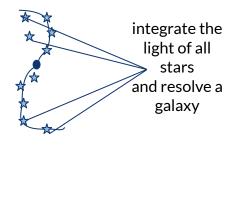
Individually resolved galaxies

High number density

Poorer redshifts

Low-pass k-filter





integrate the

HI emission of many galaxies

to resolve large-scale fluctuations

Photometric surveys

Individually resolved galaxies

Intensity Mapping

Unresolved map (intensity)

High number density

Large volumes

Poorer redshifts

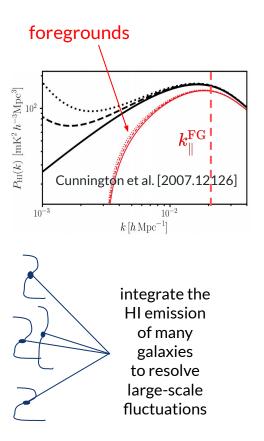
Low-pass k-filter

Excellent redshifts

High-pass k-filter



CAROLINE GUANDALIN (c.m.guandalin@qmul.ac.uk)



Photometric surveys

Individually resolved galaxies

High number density

Large volumes

Intensity Mapping

Unresolved map

(intensity)

Poorer redshifts

Low-pass k-filter

Excellent redshifts

High-pass k-filter



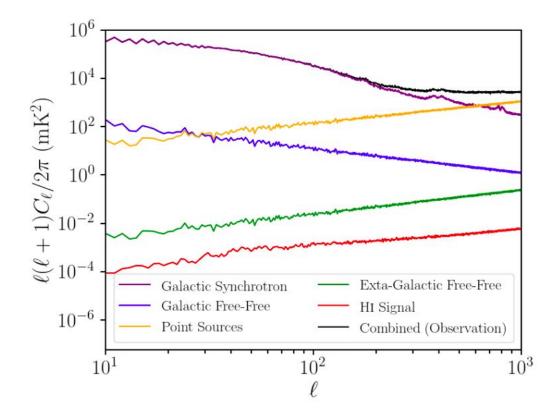
Clustering–z (extra)



 r_{\parallel}



Foregrounds



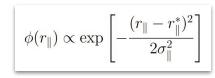
Cunnington et al. [1904.01479]

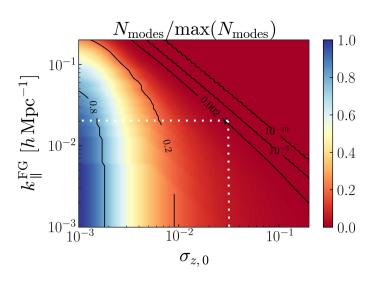
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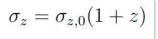
Clustering redshifts with 2-point statistics

$$F^{P}_{\alpha\beta} = \sum_{\boldsymbol{X}\boldsymbol{X}'} \frac{A}{4\pi} \int_{0}^{\infty} \mathrm{d}k_{\perp} \, k_{\perp} \, \partial_{\alpha} \mathcal{P}^{XY}(k_{\perp}) \, \partial_{\beta} \mathcal{P}^{X'Y'}(k_{\perp}) \, \mathcal{I}^{XX'}(k_{\perp}) \, \mathcal{I}^{YY'}(k_{\perp})$$

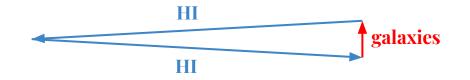
$$\sigma_P^{-1}(\theta) \propto \int \frac{\mathrm{d}k_{\parallel}}{2\pi} \, \mathcal{K}^2_{\theta}(k_{\parallel}, \sigma_{\parallel}) \,\mathrm{e}^{-k_{\parallel}^2 \sigma_{\parallel}^2} \frac{P_{gh}^2(\mathbf{k}_{\parallel}, k_{\perp})}{P_{hh}(k_{\parallel}, k_{\perp})}$$







3-point statistics



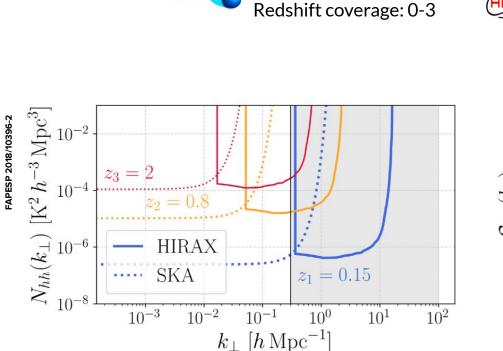
$$\begin{split} F_{\alpha\beta}^{B} &= \sum_{\boldsymbol{X}\boldsymbol{X}'} \frac{A}{4\pi} \frac{1}{6} \int_{0}^{\infty} \mathrm{d}\boldsymbol{k}_{\perp} \, \mathrm{d}\boldsymbol{q}_{\perp} \, \mathrm{d}\boldsymbol{p}_{\perp} \, \frac{\boldsymbol{k}_{\perp} \boldsymbol{q}_{\perp} \boldsymbol{p}_{\perp}}{\pi^{2} A_{T}} \, \partial_{\alpha} \mathcal{B}^{XYZ}(\boldsymbol{k}_{\perp}, \boldsymbol{p}_{\perp}, \boldsymbol{q}_{\perp}) \, \partial_{\beta} \mathcal{B}^{X'Y'Z'}(\boldsymbol{k}_{\perp}, \boldsymbol{p}_{\perp}, \boldsymbol{q}_{\perp}) \, \mathcal{I}^{XX'}(\boldsymbol{k}_{\perp}) \, \mathcal{I}^{YY'}(\boldsymbol{q}_{\perp}) \, \mathcal{I}^{ZZ'}(\boldsymbol{p}_{\perp}) \\ & \sigma_{B}^{-1}(\theta) \propto \int \frac{\mathrm{d}\boldsymbol{p}_{\parallel}}{2\pi} \frac{\mathrm{d}\boldsymbol{k}_{\parallel}}{2\pi} \, \mathcal{K}_{\theta}^{2}(\boldsymbol{k}_{\parallel}, \sigma_{\parallel}) \, \mathrm{e}^{-\boldsymbol{k}_{\parallel}^{2} \sigma_{\parallel}^{2}} \frac{B_{ghh}^{2}(\boldsymbol{k}_{\parallel}, \boldsymbol{p}_{\parallel}, -\boldsymbol{k}_{\parallel} - \boldsymbol{p}_{\parallel}; \boldsymbol{k}_{\perp}, \boldsymbol{p}_{\perp}, \boldsymbol{q}_{\perp})}{P_{hh}(\boldsymbol{p}_{\parallel}, \boldsymbol{q}_{\perp}) P_{hh}(-\boldsymbol{k}_{\parallel} - \boldsymbol{p}_{\parallel}, \boldsymbol{p}_{\perp})} \end{split}$$

We can avoid the foreground-dominated region!



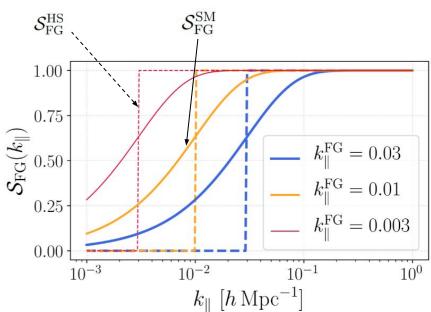
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Beam and noise



Single-dish mode

Area: 20000 deg²



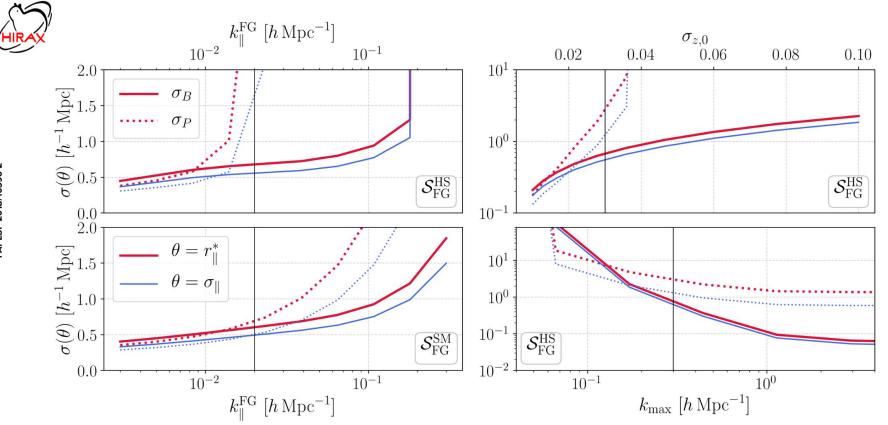
Interferometer mode

Redshift coverage: 0.8-2.5

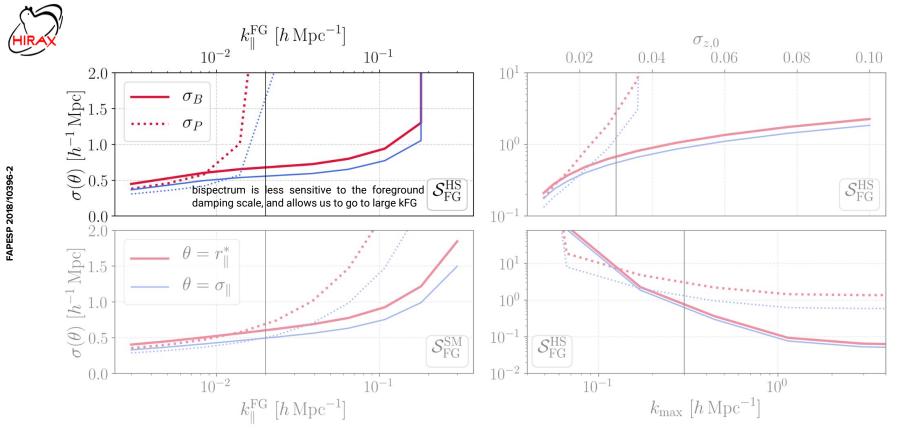
Area: 2000 deg²

 $^{-3}\,\mathrm{Mpc}^{3}$

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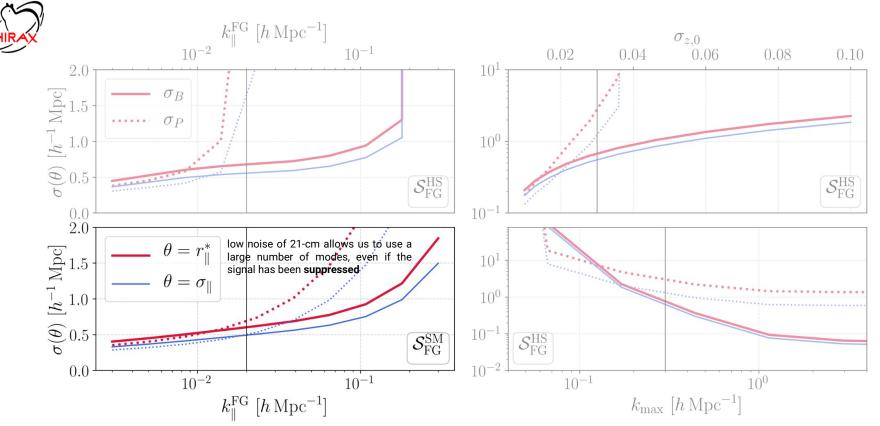


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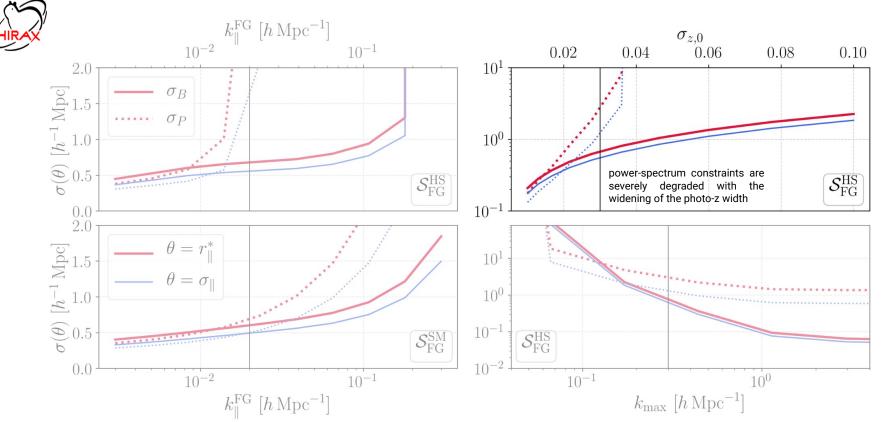


ICTP SAIFR

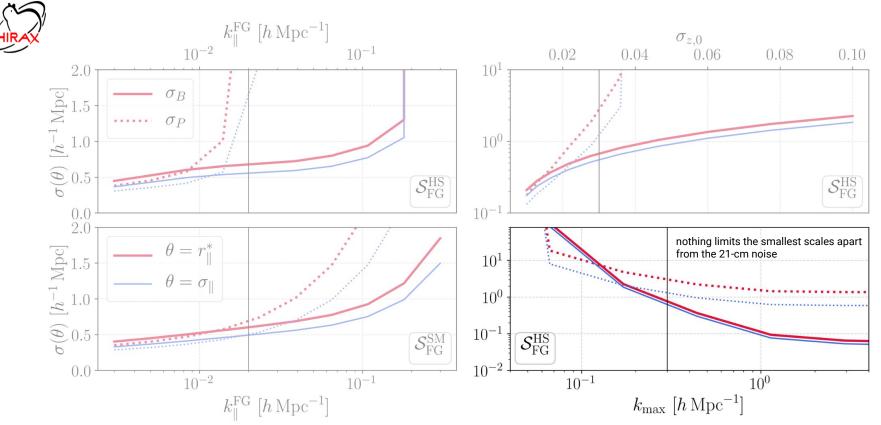
WORKSHOP ON CLASSICAL GRAVITY AND APPLICATIONS



WORKSHOP ON CLASSICAL GRAVITY AND APPLICATIONS



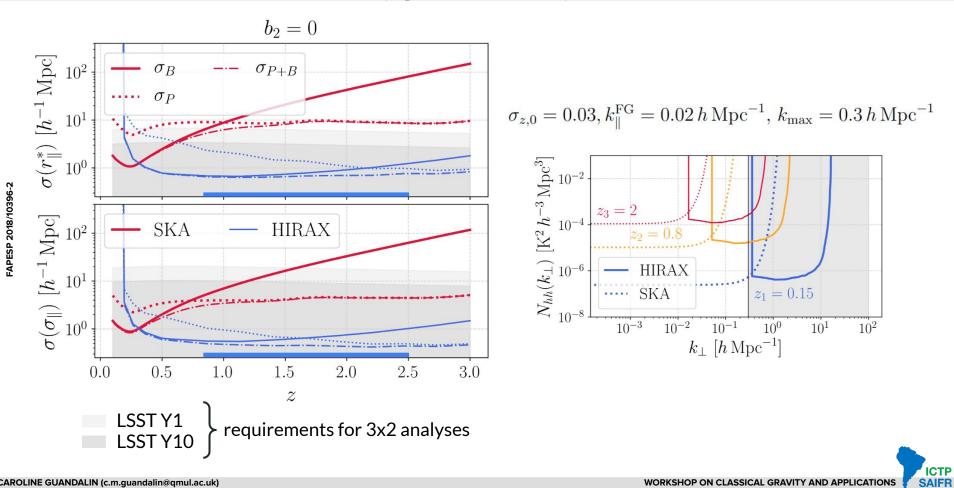
WORKSHOP ON CLASSICAL GRAVITY AND APPLICATIONS 💡



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Forecast for Stage-IV surveys

(single-dish vs. interferometer)



 $P_{xy}(k) = \mathcal{K}_{xy} P_{\mathrm{nl}}(k) + N_{xy} \,\delta_{xy}$

 $B_{ghh}(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) = \bar{T}_{h}^{2}(z)\mathcal{S}_{b}(p_{\perp}, q_{\perp})\mathcal{S}_{FG}(p_{\parallel}, q_{\parallel}) \times$ $\left[b_{g,1}b_{h,1}^{2}B(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) + b_{g,1}b_{h,1}b_{h,2}P(k)P(p) + b_{g,1}b_{h,2}b_{h,1}P(k)P(q) + b_{g,2}b_{h,1}^{2}P(p)P(q) \right]$

21-cm signal

$$\delta T_h(\boldsymbol{x}) \equiv \bar{T}_h(z)\delta_h(\boldsymbol{x}) \equiv \bar{T}_h(z)\left[b_{h,1}(z)\delta(\boldsymbol{x}) + \frac{b_{h,2}(z)}{2}\delta^2(\boldsymbol{x})\right]$$

 $b_h(z) = 1.307 \left(0.66655 + 0.17765 \, z + 0.050223 \, z^2 \right)$

 $\bar{T}_h(z) = (0.055919 + 0.23242 \, z - 0.024136 \, z^2) \, [\text{mK}]$

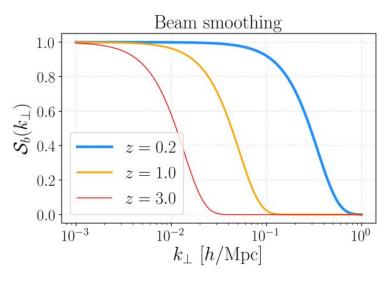
ICTP

 $P_{xy}(k) = \mathcal{K}_{xy} P_{\mathrm{nl}}(k) + N_{xy} \,\delta_{xy}$

 $B_{ghh}(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) = \bar{T}_{h}^{2}(z)\mathcal{S}_{b}(p_{\perp}, q_{\perp})\mathcal{S}_{\mathrm{FG}}(p_{\parallel}, q_{\parallel}) \times$ $\left[b_{g,1}b_{h,1}^{2}B(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) + b_{g,1}b_{h,1}b_{h,2}P(k)P(p) + b_{g,1}b_{h,2}b_{h,1}P(k)P(q) + b_{g,2}b_{h,1}^{2}P(p)P(q)\right]$

Single dish (SKA-like survey)

 $\delta T_h^{\text{obs}}(\boldsymbol{k}) = \mathcal{S}_{\text{b}}(k_{\perp}) \mathcal{S}_{\text{FG}}(k_{\parallel}) \, \delta T_h(\boldsymbol{k})$





 $P_{xy}(k) = \mathcal{K}_{xy} P_{\rm nl}(k) + N_{xy} \,\delta_{xy}$

 $B_{ghh}(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) = \bar{T}_{h}^{2}(z)\mathcal{S}_{b}(p_{\perp}, q_{\perp})\mathcal{S}_{FG}(p_{\parallel}, q_{\parallel}) \times$ $\left[b_{g,1}b_{h,1}^{2}B(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) + b_{g,1}b_{h,1}b_{h,2}P(k)P(p) + b_{g,1}b_{h,2}b_{h,1}P(k)P(q) + b_{g,2}b_{h,1}^{2}P(p)P(q)\right]$

Single dish (SKA-like survey) $\delta T_{h}^{\text{obs}}(\boldsymbol{k}) = S_{\text{b}}(k_{\perp}) S_{\text{FG}}(k_{\parallel}) \delta T_{h}(\boldsymbol{k})$ $N_{hh} = \sigma_{\text{pix}}^{2} V_{\text{pix}}$ $\sigma_{\text{pix}}^{2} = T_{\text{sys}}^{2} \frac{1}{\Delta \nu t_{\text{tot}}} \frac{\Omega_{\text{tot}}}{N_{\text{dishes}} N_{\text{beams}}}{1}$ $T_{\text{rx}} = 15 \text{ K} + 30 \text{ K} \left(\frac{\nu}{\text{GHz}} - 0.75\right)^{2}$

Bacon et al. 2020

Pourtsidou, Bacon & Crittenden [1610.04189]

ICTP

 $P_{xy}(k) = \mathcal{K}_{xy} P_{\mathrm{nl}}(k) + N_{xy} \,\delta_{xy}$

$$B_{ghh}(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) = \bar{T}_{h}^{2}(z) \mathcal{S}_{b}(p_{\perp}, q_{\perp}) \mathcal{S}_{\mathrm{FG}}(p_{\parallel}, q_{\parallel}) \times \left[b_{g,1} b_{h,1}^{2} B(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) + b_{g,1} b_{h,1} b_{h,2} P(k) P(p) + b_{g,1} b_{h,2} b_{h,1} P(k) P(q) + b_{g,2} b_{h,1}^{2} P(p) P(q) \right]$$

Single dish (SKA-like survey)

$$\delta T_h^{\mathrm{obs}}(\boldsymbol{k}) = \mathcal{S}_{\mathrm{b}}(k_{\perp}) \, \mathcal{S}_{\mathrm{FG}}(k_{\parallel}) \, \delta T_h(\boldsymbol{k})$$

$$N_{hh} = \sigma_{\rm pix}^2 V_{\rm pix}$$

$$\sigma_{\rm pix}^2 = T_{\rm sys}^2 \frac{1}{\Delta \nu \, t_{\rm tot}} \frac{\Omega_{\rm tot}}{\Omega_{\rm pix}} \frac{1}{N_{\rm dishes} N_{\rm beams}}$$

Pourtsidou, Bacon & Crittenden [1610.04189]

SKA1-MID: **133 dishes** with **15 m** in diameter. While it is required Bands 1 and 2 to fully observe the redshift range 0 < z < 3 [Santos et al. 1501.03989], for simplicity we fixed **Band 1 specifications**: sky coverage of **20000 deg**², frequency resolution of **15.2 kHz**, and integration time of **10000 hours** [SKASWG: Bacon et al. 2019].



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 $P_{xy}(k) = \mathcal{K}_{xy} P_{\mathrm{nl}}(k) + N_{xy} \,\delta_{xy}$

 $B_{ghh}(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) = \bar{T}_{h}^{2}(z)\mathcal{S}_{b}(p_{\perp}, q_{\perp})\mathcal{S}_{FG}(p_{\parallel}, q_{\parallel}) \times$ $\left[b_{g,1}b_{h,1}^{2}B(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) + b_{g,1}b_{h,1}b_{h,2}P(k)P(p) + b_{g,1}b_{h,2}b_{h,1}P(k)P(q) + b_{g,2}b_{h,1}^{2}P(p)P(q)\right]$

Interferometer (HIRAX-like survey)

 $\delta T_h^{\mathrm{obs}}(\boldsymbol{k}) = \mathcal{S}_{\mathrm{FG}}(k_{\parallel}) \, \delta T_h(\boldsymbol{k})$

 $N_{hh}(k_{\perp}) = \frac{4\pi f_{\rm sky} \,\chi^2(z) \,(1+z) \,T_{\rm sys}^2 \,\theta_{\rm FWHM}^2}{H(z) \,t_{\rm tot} \,\lambda_{21}(z) \,N_d(\boldsymbol{d} = \boldsymbol{k}_{\perp} \,\chi(z) \,\lambda_{21}(z)/2\pi)}$ $T_{\rm sys} = 50 \,\mathrm{K} + 60 \,\left[\frac{\nu_{21}(z)}{300 \,\mathrm{MHz}}\right]^{-2.5} \,\mathrm{K}$

e.g. Alonso et al. [1704.01941]

ICTP

 $P_{xy}(k) = \mathcal{K}_{xy} P_{\rm nl}(k) + N_{xy} \,\delta_{xy}$

$$B_{ghh}(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) = \bar{T}_{h}^{2}(z)\mathcal{S}_{b}(p_{\perp}, q_{\perp})\mathcal{S}_{\mathrm{FG}}(p_{\parallel}, q_{\parallel}) \times \left[b_{g,1}b_{h,1}^{2}B(\boldsymbol{k}, \boldsymbol{p}, \boldsymbol{q}) + b_{g,1}b_{h,1}b_{h,2}P(k)P(p) + b_{g,1}b_{h,2}b_{h,1}P(k)P(q) + b_{g,2}b_{h,1}^{2}P(p)P(q)\right]$$

Galaxy survey (LSST-like)

 $b_{q,1}(z) = 0.95 D^{-1}(z)$

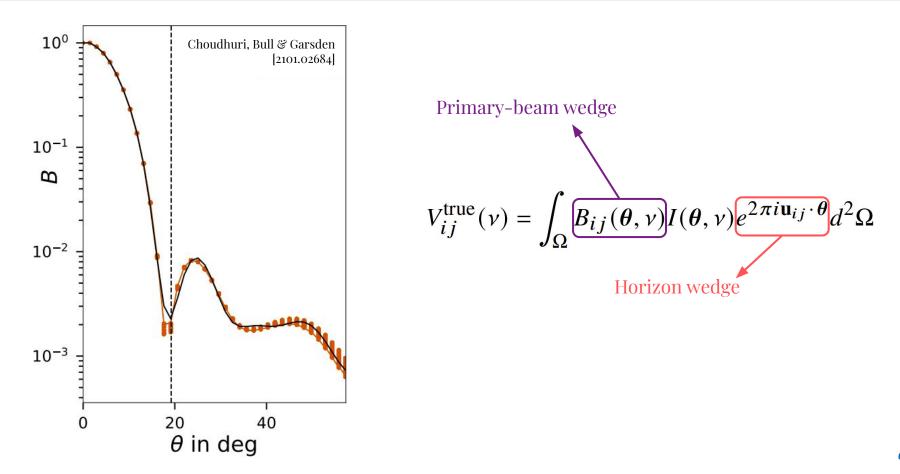
$$N_{gg} = \frac{1}{\bar{n}_g}$$
$$n_g(z) = A \, z^\alpha \, \exp\left[-\left(\frac{z}{z_0}\right)^\beta\right]$$

$\Omega_{\rm sky}$ (deg ²)	nsources	Zmin	Zmax	z_0	α	β
13800	48	0.2	3	2	0.9	0.28

e.g. Ballardini, Matthewson & Maartens [1906.04730]

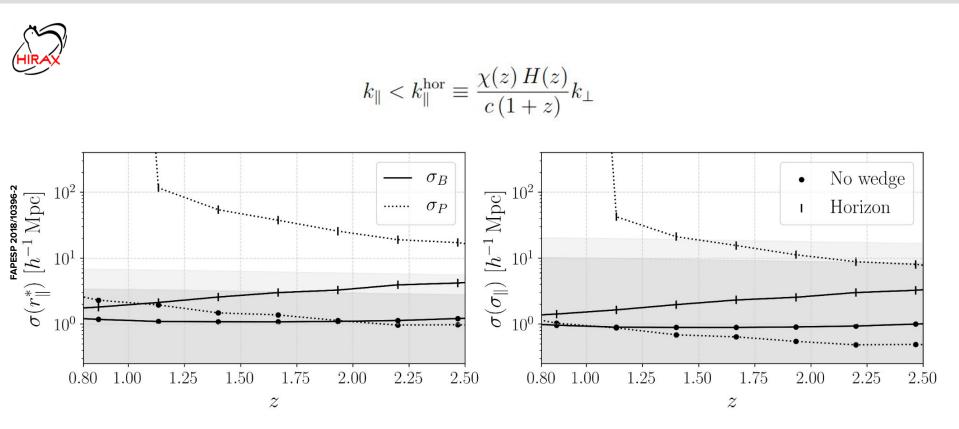


Horizon and primary-beam wedges



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Impact of horizon wedge



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