SIGNATURES OF THE YANG-MILLS DECONFINEMENT TRANSITION FROM THE GLUON TWO-POINT CORRELATOR

Duifje Maria van Egmond

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In collaboration with Urko Reinosa

MOTIVATION

- The deconfinement transition was shown to be altered by the presence of a magnetic field[1].
- One looks at the interplay between the Polyakov loop and a magnetic background field:

$$D_\mu = \partial_\mu - iA_\mu - iQa_\mu .$$

- We look at the interplay between gauge-dependent derivatives of the Polyakov loop and an arbitrary background gauge field

$$\bar{D}_\mu (A_\mu - \bar{A}_\mu) = 0, \text{ with } \bar{D}_\mu \equiv \partial_\mu - \left[\bar{A}_\mu, \cdot\right].$$

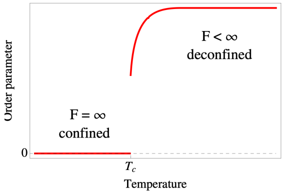
- Important difference: magnetic field is a **physical** field that alters the **physical** Polyakov loop, our background field is a **gauge parameter** that alters **gauge-fixed** quantities.
- Our work emphasizes the importance of an appropriate background (gauge) choice when considering certain symmetries.

[1] A.J. Mizher, M.N. Chernodub, E.S. Fraga, Phys.Rev.D 82 (2010) 105016

INTRODUCTION

- At some very high temperature T_c, hadrons become free quarks and gluons → quark-gluon plasma.
- An order parameter for this transition is the Polyakov loop:

$$\ell \sim \langle P e^{i \int_0^\beta d\tau A_0(\tau, x)} \rangle \sim e^{-\beta F}$$



- Under center symmetry, $\ell \rightarrow Z_N \ell$, so deconfinement is signaled by a broken center symmetry (in pure Yang-Mills).
- Confirmed by lattice data: second order transition for $SU(2)$, first order for $SU(3)$.

ENCODING OF THE TRANSITION

Polyakov loop:

$$\ell \sim \langle P e^{i\int_0^\beta d\tau A_0(\tau,x)} \rangle \sim e^{-\beta F}.$$

Because the Polyakov loop is related to A_0, it is expected that the transition is encoded in (the tower of)

$$\langle A_0 \rangle, \langle A_0 A_0 \rangle, \dots, \langle A_0^n \rangle.$$

For the appropriate choice of gauge, can the transition be reflected in the lowest order correlators?

LANDAU GAUGE CORRELATOR

In principle:

- $\langle A \rangle$ is found by minimizing the effective action $\Gamma[A]$. It represents the state of the system . $\langle A_0 \rangle \rightarrow$ order parameter.
- The two-point correlator derives from the effective action

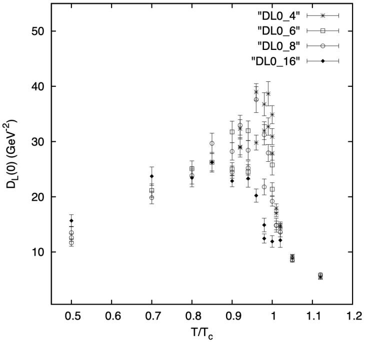
$$1 \Big/ \left. \frac{\partial^2 \Gamma}{\partial A^2} \right|_{A=\langle A \rangle} = \langle AA \rangle_c,$$

so for $SU(2)$, $\langle A_0 A_0 \rangle$ should diverge at T_c.

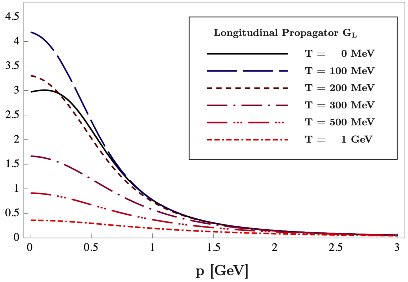
In practice:

- In the Landau gauge, $\partial_\mu A_\mu = 0$, then $\langle A_0 \rangle = 0$. $\rightarrow$ no order parameter.
- No evidence of divergence of $\langle A_0 A_0 \rangle$ was found on the (gauge-fixed) lattice and in the continuum.

$SU(2)$ LANDAU GAUGE CORRELATORS



Electric susceptibility (zero momentum longitudinal propagator)



Longitudinal gluon propagator

*T. Mendes and A. Cucchieri, PoS LATTICE2014, 183 (2015).

*L. Fister and J. M. Pawlowski, [arXiv:1112.5440 [hep- ph]](2012).

BACKGROUND FIELD GAUGES

- In the Landau gauge the effective action is not center-symmetric $\Gamma[A] \neq \Gamma[A^U]$.
- Introducing a background field $\bar{A}$, the effective action is gauge-invariant $\Gamma_{\bar{A}}[A] = \Gamma_{\bar{A}^U}[A^U]$. Landau-deWitt gauge:

$$\bar{D}_\mu(A_\mu - \bar{A}_\mu) = 0, \text{ with } \bar{D}_\mu \equiv \partial_\mu - \left[\bar{A}_\mu, \cdot\right].$$

- In the Background field effective action one looks at $\tilde{\Gamma}[\bar{A}]$ by taking $\bar{A} = \langle A \rangle$.
  - The two-point function $\langle AA \rangle_c$ is not directly accessible from $\tilde{\Gamma}[\bar{A}]$,
  - Relies on the strict background independence of $\tilde{\Gamma}[\bar{A}]$, which is not easy to maintain in the presence of truncations.
- We propose the Center-symmetric Landau gauge, which fixes $\bar{A} = \bar{A}_c$. Then $\Gamma_{\bar{A}}[A] = \Gamma_{\bar{A}^{U_c}}[A^{U_c}]$.

OUR SETUP

- We work in the Landau-deWitt gauge with a background field $\bar{A}_\mu$:

$$\bar{D}_\mu(A_\mu - \bar{A}_\mu) = 0, \text{ with } \bar{D}_\mu \equiv \partial_\mu - \left[\bar{A}_\mu, \cdot\right].$$

- We take $\bar{A}$ and $\langle A \rangle$ in the temporal direction, $\propto \delta_{\mu 0}$, and along the diagonal color directions (σ^3 for $SU(2)$, (λ^3, λ^8) for $SU(3)$), so that $\Gamma[A, \bar{A}] \propto V(A, \bar{A})$.
- The center-symmetric values for $\bar{A}$ and $\langle A \rangle$ are found by Weyl chambers. For example in $SU(2)$, $A_{c,\mu} = \bar{A}_{c,\mu} = \delta_{\mu 0} \frac{T}{g} \pi \frac{\sigma^3}{2}$.
- We fix $\bar{A} = \bar{A}_c$: Center-symmetric Landau gauge.
  Center-symmetric phase when $\langle A \rangle = A_c$, $\rightarrow$ order parameter.

CURCI-FERRARI MODEL

We have computed $\langle A \rangle$ and $\langle A(0,p)A(0,-p) \rangle$ up to first loop order in the finite temperature Curci-Ferrari model:

$$S = S_{YM} + S_{gf} + \int_{x,\tau} \frac{m^2}{2}(A^a_\mu - \bar{A}^a_\mu)^2$$

Several motivations:

- Perturbative gauge-fixed Yang-Mills breaks down at low energies (Landau pole, Gribov copies...), there is no analytical model for this region.
- A gluon mass term seems to dominate the (unknown) gauge-fixed action in the IR; decoupling behaviour on the lattice. CF could be an effective model.
- The CF model is renormalizable, avoids the Landau pole and lifts the degeneracy between Gribov copies. Perturbative window into non-perturbative region.

RESULTS - $T_c(MeV)$

	Lattice	FRG-BG[2]	CF-BG, 1-lp[3]	CF-BG, 2-lp[4]	CF-CS, 1-lp[5]
SU(2)	295	230	238	284	**265**
SU(3)	270	275	185	254	**267**

BG: Background effective action
CS: Centersymmetric Landau gauge
.

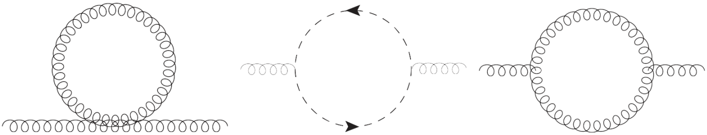
[2] L. Fister and J. M. Pawlowski, Phys.Rev. D88 (2013) 045010

[3] U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys.Lett. B742 (2015) 61-68.

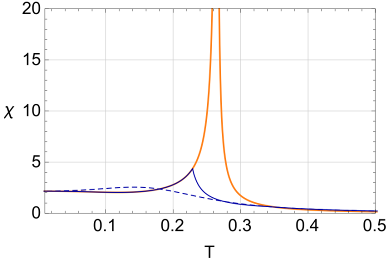
[4] U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys.Rev. D93 (2016) 105002.

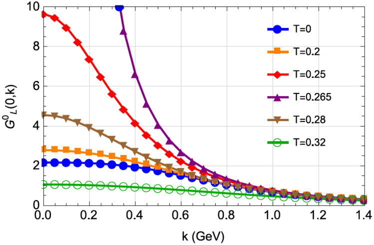
[5] **DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022)**

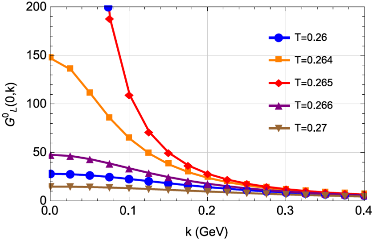
FEYNMAN DIAGRAMS GLUON PROPAGATOR



RESULTS: SU(2) GLUON PROPAGATOR





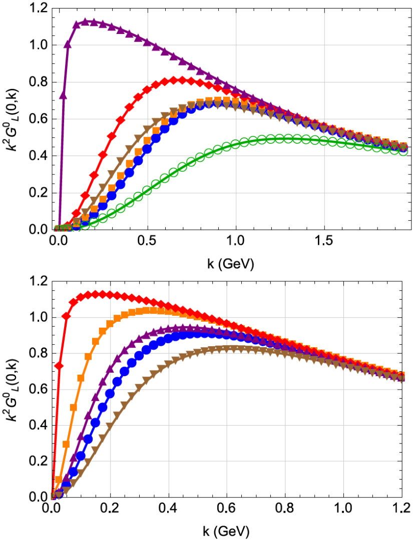


Landau gauge
Background field effective action
Centersymmetric Landau gauge

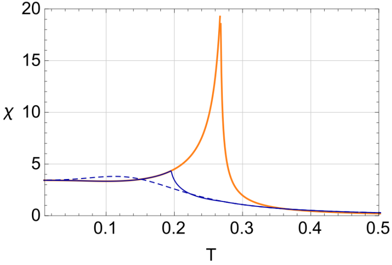
m=0.68 GeV, μ=1 GeV, g=7.5

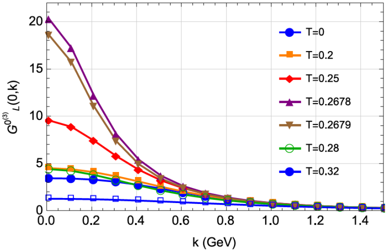
*DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022).

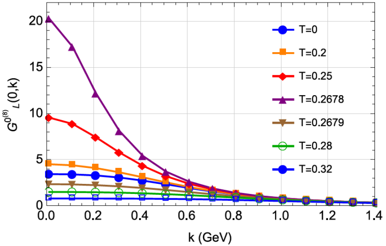
SU(2) DRESSING FUNCTION



RESULTS: SU(3) GLUON PROPAGATOR





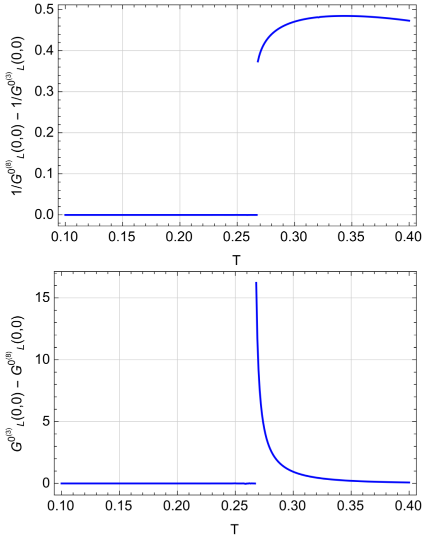


Landau gauge
Background field effective action
Centersymmetric Landau gauge

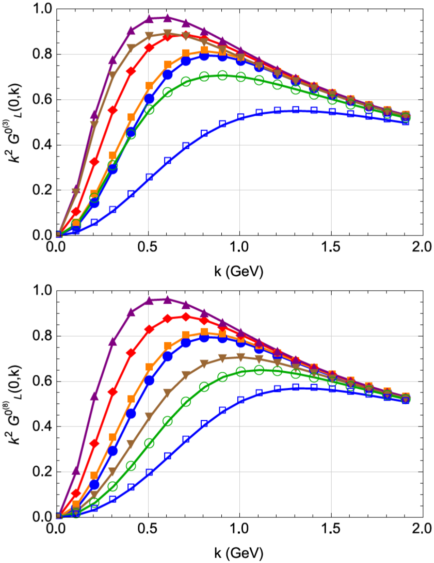
m=0.54 GeV, μ=1 GeV, g=4.9

*DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022).

SU(3) PROPAGATOR DIFFERENCE



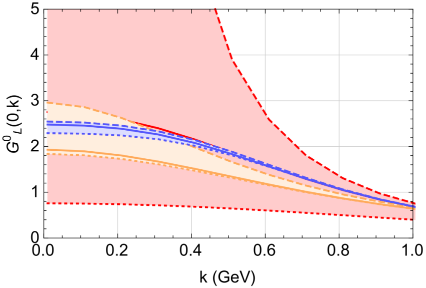
SU(3) DRESSING FUNCTION



CONCLUSION AND OUTLOOK

- We have performed, for the first time, calculations of the gluon one-and two-point correlator in the centersymmetric Landau gauge.
- We find a good agreement with lattice data for T_c.
- We find that for $SU(2)$, the deconfinement transition is signaled by a divergence of the longitudinal gluon propagator for $k \to 0$.
- For $SU(3)$, the difference between the propagators in the neutral color mode is an order parameter for the transition.
- This model can be tested on the lattice by changing the boundary conditions in the Landau gauge [with O. Oliveira and P. Silva].
- Ideas for future works: RG improvement, transversal propagator and dynamically generated mass [with D. Dudal and D. Vercauteren].

DISCUSSION

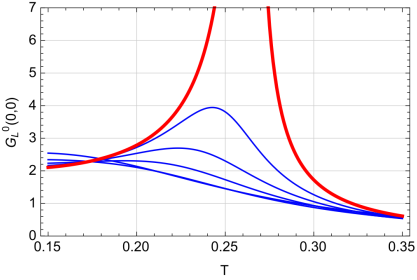


Landau gauge

Self-consistent Landau gauge

Center-symmetric Landau Gauge

DISCUSSION



DISCUSSION