From the Quark and Gluon Structure of Nuclei to the Search for Dark Matter



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Outline

- I. Nuclei from Quarks
 - start from a QCD-inspired model of *hadron* structure
 - develop a quantitative theory of nuclear structure

II. Search for observable effects of the change in hadron structure in-medium

III. New: Indication of a BSM particle: Dark Photon





I. Insights into nuclear structure

- what is the atomic nucleus?

There are two very different views....





Quark Structure matters/doesn't matter

 Nuclear femtography: the science of mapping the quark and gluon structure of *atomic nuclei* is just beginning (EIC motivation)

 "Considering quarks is in contrast to our modern understanding of nuclear physics... the basic degrees of freedom of QCD (quarks and gluons) have to be considered only at higher energies. The energies relevant for nuclear physics are only a few MeV"





What do we know?

- Since 1970s: Dispersion relations → intermediate range NN attraction is a strong Lorentz scalar
- In relativistic treatments (RHF, RBHF, QHD...) this leads to mean scalar field on a nucleon ~300 to 500 MeV!!
- This is not small up to half the nucleon mass
 death of "wrong energy scale" arguments
- Largely cancelled by large vector mean field BUT these have totally different dynamics: ω⁰ just shifts energies, σ seriously modifies internal hadron dynamics
- Latter not naturally described by EFT with N and π alone



Self-consistent solution for confined quarks in a hadron in nuclear matter

$$[i\gamma^{\mu}\partial_{\mu} - (m_q - g_{\sigma}{}^q\bar{\sigma}) - \gamma^0 g_{\omega}{}^q\bar{\omega}]\psi = 0$$

 $\int_{Bag} d\vec{r} \overline{\psi}(\vec{r}) \psi(\vec{r})$

Source of σ changes:

and hence mean scalar field changes...

and hence quark wave function changes....

THIS PROVIDES A NATURAL SATURATION MECHANISM (VERY EFFICIENT BECAUSE QUARKS ARE LIGHT)

source is suppressed as mean scalar field increases (i.e. as density increases)







SELF-CONSISTENCY

Quark-Meson Coupling Model (QMC): Role of the Scalar Polarizability of the Nucleon

The response of the nucleon internal structure to the scalar field is of great interest... and importance

$$M * (\mathbf{r}) = M - g_{\sigma} \sigma(\mathbf{r}) + \frac{d}{2} (g_{\sigma} \sigma(\mathbf{r}))^{2}$$

Non-linear dependence through the scalar polarizability d ~ 0.22 R in original QMC (MIT bag)

Indeed, in nuclear matter at mean-field level, this is the ONLY place the response of the internal structure of the nucleon enters.





Summary : Scalar Polarizability

 Consequence of polarizability in atomic physics is many-body forces:



$$V = V_{12} + V_{23} + V_{13} + V_{123}$$

- same is true in nuclear physics
- Three-body forces (NNN, HNN, HHN...) generated with NO new parameters

 critical in neutron stars







Application to nuclear structure





Derivation of Density Dependent Effective Force

Physical origin of density dependent forces of Skyrme type within the quark meson coupling model

P.A.M. Guichon^{a,*}, H.H. Matevosyan^{b,c}, N. Sandulescu^{a,d,e}, A.W. Thomas^b

Nuclear Physics A 772 (2006) 1–19

- Start with classical theory of MIT-bag nucleons with structure modified in medium to give M_{eff} (σ).
- Quantise nucleon motion (non-relativistic), expand in powers of derivatives
- Derive equivalent, local energy density functional:

$$\langle H(\vec{r}) \rangle = \rho M + \frac{\tau}{2M} + \mathcal{H}_0 + \mathcal{H}_3 + \mathcal{H}_{\text{eff}} + \mathcal{H}_{\text{fin}} + \mathcal{H}_{\text{so}}$$

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Derivation of EDF (cont.)

$$\begin{aligned} \mathcal{H}_{0} + \mathcal{H}_{3} &= \rho^{2} \bigg[\frac{-3G_{\rho}}{32} + \frac{G_{\sigma}}{8(1 + d\rho G_{\sigma})^{3}} - \frac{G_{\sigma}}{2(1 + d\rho G_{\sigma})} + \frac{3G_{\omega}}{8} \bigg] \\ &+ (\rho_{n} - \rho_{p})^{2} \bigg[\frac{5G_{\rho}}{32} + \frac{G_{\sigma}}{8(1 + d\rho G_{\sigma})^{3}} - \frac{G_{\omega}}{8} \bigg], \end{aligned}$$

$$\begin{aligned} \mathcal{H}_{\text{eff}} = \left[\left(\frac{G_{\rho}}{8m_{\rho}^{2}} - \frac{G_{\sigma}}{2m_{\sigma}^{2}} + \frac{G_{\omega}}{2m_{\omega}^{2}} + \frac{G_{\sigma}}{4M_{N}^{2}} \right) \rho_{n} + \left(\frac{G_{\rho}}{4m_{\rho}^{2}} + \frac{G_{\sigma}}{2M_{N}^{2}} \right) \rho_{p} \right] \tau_{n} \\ + p \leftrightarrow n, \end{aligned}$$

$$\begin{aligned} \mathcal{H}_{\text{fin}} &= \left[\left(\frac{3G_{\rho}}{32m_{\rho}^{2}} - \frac{3G_{\sigma}}{8m_{\sigma}^{2}} + \frac{3G_{\omega}}{8m_{\omega}^{2}} - \frac{G_{\sigma}}{8M_{N}^{2}} \right) \rho_{n} \\ &+ \left(\frac{-3G_{\rho}}{16m_{\rho}^{2}} - \frac{G_{\sigma}}{2m_{\sigma}^{2}} + \frac{G_{\omega}}{2m_{\omega}^{2}} - \frac{G_{\sigma}}{4M_{N}^{2}} \right) \rho_{p} \right] \nabla^{2}(\rho_{n}) + p \leftrightarrow n, \\ \mathcal{H}_{\text{so}} &= \nabla \cdot J_{n} \left[\left(\frac{-3G_{\sigma}}{8M_{N}^{2}} - \frac{3G_{\omega}(-1+2\mu_{s})}{8M_{N}^{2}} - \frac{3G_{\rho}(-1+2\mu_{v})}{32M_{N}^{2}} \right) \rho_{n} \right] \\ &+ \left(\frac{-G_{\sigma}}{4M_{N}^{2}} + \frac{G_{\omega}(1-2\mu_{s})}{4M_{N}^{2}} \right) \rho_{p} \right] + p \leftrightarrow n. \end{aligned}$$

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Note the totally new, subtle density dependence

Systematic approach to finite nuclei

J.R. Stone, P.A.M. Guichon, P. G. Reinhard & A.W. Thomas: (Phys Rev Lett, 116 (2016) 092501)

• Constrain 3 basic quark-meson couplings (g_{σ}^{q} , g_{ω}^{q} , g_{ρ}^{q}) so that nuclear matter properties are reproduced within errors

 $\begin{array}{l} -17 < \text{E/A} < -15 \ \text{MeV} \\ 0.14 < \rho_0 < 0.18 \ \text{fm}^{-3} \\ 28 < \text{S}_0 < 34 \ \text{MeV} \\ \text{L} > 20 \ \text{MeV} \\ 250 < \text{K}_0 < 350 \ \text{MeV} \end{array}$

- Fix at overall best description of finite nuclei with 5 parameters (3 for the EDF +2 pairing pars)
- Benchmark comparison: SV-min 16 parameters (11+5 pairing)





Overview of 106 Nuclei Studied – Across Periodic Table

Element	Z	N	Element	Z	N
С	6	6 -16	Pb	82	116 - 132
0	8	4 - 20	Pu	94	134 - 154
Са	20	16 - 32	Fm	100	148 - 156
Ni	28	24 - 50	No	102	152 - 154
Sr	38	36 - 64	Rf	104	152 - 154
Zr	40	44 -64	Sg	106	154 - 156
Sn	50	50 - 86	Hs	108	156 - 158
Sm	62	74 - 98	Ds	110	160
Gd	64	74 -100			

Ν	Z	Ν	Z
20	10 - 24	64	36 - 58
28	12 - 32	82	46 - 72
40	22 - 40	126	76 - 92
50	28 - 50		



i.e. We look at most challenging cases of p- or n-rich nuclei

Not fit

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Superheavy Binding : 0.1% accuracy



Stone et al., PRL 116 (2016) 092501 For detailed study of SHE see: arXiv:1901.06064







Includes better approximation to σ field, π Fock terms and term in σ^3





Giant Monopole Resonances



FIG. 13. GMR energies for ²⁰⁸Pb, ¹⁴⁴Sm, ¹¹⁶Sn, and ⁹⁰Zr from experiment and for the QMC π -II and SVmin models. Experimental data are taken from Table 1 of Ref. [24].



Kay Martinez et al., Phys Rev C100 (2019) 024333



Neutron distributions





Deformation of Gd isotopes





More Nuclear Structure

All results for QMC-III from PhD thesis of Kay Martinez - now at Silliman University (Philippines) (publications in preparation)





QMC π3

- Just 5 parameters^{*}: m_{σ} , quark couplings to σ , ω and ρ and λ_3 the strength of σ^3 term
- Tensor term included: $H^{J}_{\sigma,\omega,\rho} = \left(\frac{G_{\sigma}(1-dv_{0})^{2}}{4m_{\sigma}^{2}} - \frac{G_{\omega}}{4m_{\omega}^{2}}\right) \sum_{m} \vec{J}_{m}^{2}$ $- \frac{G_{\rho}}{4m_{\rho}^{2}} \sum_{m,m'} S_{m,m'} \vec{J}_{m} \cdot \vec{J}_{m'},$ and $H^{J}_{S} = -\frac{G_{\sigma} - G_{\omega}}{16M^{2}} \sum_{m} \vec{J}_{m}^{2} + \frac{G_{\rho}}{16M^{2}} \sum_{mm'} S_{m,m'} \vec{J}_{m} \cdot \vec{J}_{m'}.$ with $\vec{J}_{m} = i \sum_{i \in F_{m}} \sum_{\sigma\sigma'} \vec{\sigma}_{\sigma'\sigma} \times [\vec{\nabla}\phi^{i}(\vec{r},\sigma,m)]\phi^{i*}(\vec{r},\sigma',m), \quad \vec{J} = \vec{J}_{p} + \vec{J}_{n},$
- Pairing interaction (simple BCS) derived in the model

$$V_{\text{pair}}^{\text{QMC}} = -\left(\frac{G_{\sigma}}{1 + d'G_{\sigma}\rho(\vec{r})} - G_{\omega} - \frac{G_{\rho}}{4}\right)\delta(\vec{r} - \vec{r}')$$
$$d' = d + \frac{1}{3}G_{\sigma}\lambda_3,$$





Binding Energies – All Known Even-Even Nuclei





Charge Radii





Model	<i>rms</i> residual (fm)	<i>rms</i> % deviation	
QMCπ-III	0.024	0.50	
$QMC\pi$ -II	0.029	0.66	
$QMC\pi$ -I	0.028	0.65	
QMC-I	0.030	0.66	
SV-min	0.024	0.61	
UNEDF1	0.029	0.65	
$DD-ME\delta$	0.035	0.78	



Separation energies: Drip Lines



TABLE 7.3: Comparison of *rms* residuals for separation energies (in MeV) from QMC and from other nuclear models.

Model	S_{2n}	S_{2p}	δ_{2n}	δ_{2p}	Qα
QMCπ-III	0.97	0.95	1.24	1.28	1.07
$QMC\pi$ -II	1.03	1.08	1.20	1.25	1.19
SV-min	0.77	0.82	0.87	1.00	0.79
UNEDF1	0.74	0.82	0.85	0.90	0.80
DD-ME δ	1.01	1.05	1.12	1.11	1.30
FRDM	0.50	0.55	0.61	0.75	0.61







Deformation



TABLE 7.4: Comparison of β_2 *rms* residuals and *rms* % deviations from QMC π -III and from other nuclear models. There are a total of 324 even-even nuclei with available data for β_2 included for comparison.

Model	rms residual	rms % deviation
QMCπ-III	0.11	28
SV-min	0.16	59
UNEDF1	0.15	53
$DD-ME\delta$	0.14	40
FRDM	0.11	30





Shape Co-Existence for Z=N







The Superheavy Region First study:

PHYSICAL REVIEW C 100, 044302 (2019)

Physics of even-even superheavy nuclei with 96 < Z < 110 in the quark-meson-coupling model

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> > P. A. M. Guichon[‡] CEA/IRFU/SPhN Saclay, F91191, France

A. W. Thomas[§] CSSM and CoEPP, Department of Physics, University of Adelaide, SA 5005, Australia

Updated and expanded here (Martinez thesis)





Binding Energies



TABLE 6.1: Comparison of *rms* percent deviations and *rms* residuals from QMC and from other nuclear models for SHE with available data.

	<i>rms</i> residual (MeV)	<i>rms</i> % deviation	
anding agreement	0.52	0.03	$QMC\pi$ -III
anding agreement	2.04	0.11	QMC <i>π</i> -II [54]
	2.42	0.12	QMC <i>π</i> -I [53]
SPECIAL RESEARCH	1.50	0.08	QMC-I [8]
	2.25	0.11	FRDM [23]
JUDAI	6.99	0.36	SV-min [24]
	1.31	0.07	UNEDF1 [28]
UTUR OF	2.28	0.12	DD-MEδ [66]
STRUC			



Many Almost Degenerate Minima in Superheavy Region









Overview of Initial Work on Finite Nuclei

- The effective force was derived at the quark level based upon the changing structure of a bound nucleon
- Has many less parameters but reproduces nuclear properties at a level comparable with the best phenomenological Skyrme forces
- Looks like standard nuclear force
- BUT underlying theory also predicts modified internal structure and hence modified
 - DIS structure functions
 - elastic form factors.....





Nuclear DIS Structure Functions: The EMC Effect

To address questions like this one MUST start with a theory that quantitatively describes nuclear structure and allows calculation of structure functions

- very, very few examples.....





The EMC Effect: Nuclear PDFs

- Observation stunned and electrified the HEP and Nuclear communities 40 years ago
- What is it that alters the quark momentum in the nucleus?





Theoretical Understanding

- Still numerous proposals but few consistent theories
- Initial studies used MIT bag¹ to estimate effect of self-consistent change of structure in-medium
 but better to use a covariant theory
- For that Bentz and Thomas² re-derived change of nucleon structure in-medium in the NJL model
- This set the framework for sophisticated studies by Bentz, Cloët and collaborators over the last decade

¹ Thomas, Michels, Schreiber and Guichon, Phys. Lett. B233 (1989) 43 ² Bentz and Thomas, Nucl. Phys. A696 (2001) 138







EMC Effect for Finite Nuclei

(There is also a spin dependent EMC effect - as large as unpolarized)



FIG. 7: The EMC and polarized EMC effect in ¹¹B. The empirical data is from Ref. [31].

FIG. 9: The EMC and polarized EMC effect in $^{27}\mathrm{Al.}\,$ The empirical data is from Ref. [31].



Cloët, Bentz &Thomas, Phys. Lett. B642 (2006) 210 Related work by Miller and co-workers



Approved JLab Experiment

- Effect in ⁷Li is slightly suppressed because it is a light nucleus and proton does not carry all the spin (simple WF: $P_p = 13/15$ & $P_n = 2/15$)
- Experiment now approved at JLab [E12-14-001] to measure spin structure functions of ⁷Li (GFMC: $P_p = 0.86$ & $P_n = 0.04$)



Spin EMC measurement is critical as the proposed explanation in terms of <u>SRC</u> through the stensor force gives <u>NO spin EMC effect (arXiv:1809.06622</u>)

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For the EIC: Polarized Gluon EMC Effect

Xuangong Wang et al.

J. Phys. G: Nucl. Part. Phys. 49 (2022) 03LT01



Figure 3. (Left panel) Unpolarized EMC ratios for the structure functions $F_{2A}(x)/F_{2N}(x)$ (solid) and the unpolarized gluon distributions $g_A(x)/g_p(x)$ (dashed). (Right panel) Polarized EMC ratios for the structure functions $g_{1A}(x)/g_{1p}(x)$ (solid) and polarized gluon distributions $\Delta g_A(x)/\Delta g_p(x)$ (dashed). The empirical data points are the unpolarized nuclear matter results for the EMC ratio from reference [53].





Letters

Hypothesis that SRC gives the EMC Effect

 $F_2^{A} = (Z - n_{SRC}^{A})F_2^{p} + (N - n_{SRC}^{A})F_2^{n} + n_{SRC}^{A}(F_2^{p*} + F_2^{n*})$



New Observation on SRC and EMC Effect



Wanli Xing et al., to be published

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Recent Results Concerning a Dark Photon





Dark Matter and Dark Energy

Over the past two decades we have learnt that, in spite of the successes of the Standard Model, most of the matter in the Universe is something else





It interacts very weakly but has major gravitational effects





New Underground Laboratory under construction

- Approximately midway between Adelaide and Melbourne
- 1km underground in an active gold mine









SUPL construction images









SABRE Experiment





PE

to







Evidence for a New BSM Particle? A Dark Photon





Dark Photon

It is not the solution to DM problem but may be a portal connecting the sectors

- There are a number of formulations of the dark photon
- Initially purely vector couplings, then more like a Z´
- For us it is a U_Y(1) boson interacting with Standard Model particles through kinetic mixing

$$\mathcal{L} \supset -\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \frac{\bar{m}_{A'}^2}{2} A'_{\mu} A'^{\mu} + \frac{\epsilon}{2\cos\theta_W} F'_{\mu\nu} B^{\mu\nu}$$

• So that the couplings are:

$$C_Z^v = (\cos \alpha - \epsilon_W \sin \alpha) \bar{C}_Z^v + \epsilon_W \sin \alpha \cot \theta_W C_{\gamma}^v,$$

$$C_Z^a = (\cos \alpha - \epsilon_W \sin \alpha) \bar{C}_Z^a,$$

and

$$C_{A_D}^v = -(\sin\alpha + \epsilon_W \cos\alpha) \bar{C}_Z^v + \epsilon_W \cos\alpha \cot\theta_W C_{\gamma}^v$$





Dark Photon (cont.)

Where the Standard Model Z couplings are

 $\bar{C}_Z \sin 2\theta_W = T_3^f - 2q_f \sin^2 \theta_W, \qquad \bar{C}_Z^a \sin 2\theta_W = T_3^f$

and the mixing parameters are

$$\tan \alpha = \frac{1}{2\epsilon_W} \left[1 - \epsilon_W^2 - \rho^2 - \sin(1 - \rho^2) \sqrt{4\epsilon_W^2 + (1 - \epsilon_W^2 - \rho^2)^2} \right]$$

and

$$\epsilon_W = \frac{\epsilon \tan \theta_W}{\sqrt{1 - \epsilon^2 / \cos^2 \theta_W}}$$
$$\rho = \frac{\bar{m}_{A'} / \bar{m}_{\bar{Z}}}{\sqrt{1 - \epsilon^2 / \cos^2 \theta_W}}$$

with $\boldsymbol{\epsilon}$ the mixing parameter in the Lagrangian





New JAM Analysis of World Data

Following earlier work of Wang and Thomas

J. Phys. G: Nucl. Part. Phys. 47 (2020) 015102







Use JAM Framework to Analyse World Data to determine Exclusion Limits

Hunt-Smith et al., arXiv2302.11126

- NLO analysis with $Q^2 > 1.69 \text{ GeV}^2$ and $W^2 > 10 \text{ GeV}^2$
- Over 3,000 fixed target and HERA data







BUT if we accept that a dark photon may exist

X² analysis strongly prefers it!

reaction	$\chi^2_{\rm dof}({\rm dark})$	$\chi^2_{\rm dof}({\rm base})$	$N_{\rm dof}$
fixed target DIS	1.01	1.05	1495
HERA NC	1.02	1.03	1104
HERA CC	1.13	1.18	81
Drell-Yan	1.18	1.16	205
Z rapidity	1.08	1.05	56
W asymmetry	1.04	1.07	97
jets	1.16	1.15	200
total	1.03	1.05	3283







Hypothesis that SM is correct EXCLUDED at 6.5 σ



Stable against omission of HT corrections and increase of lower Q² cut to 10 GeV²



Hunt-Smith et al., arXiv2302.11126



Improvement over full range of Q² and x







PDFs consistent with NO DP Analysis





Uncertainties based on analysis with 200 replicas



Best fit parameters





Hunt-Smith et al., arXiv2302.11126



I. Summary



- Intermediate range NN attraction is STRONG Lorentz scalar
- This modifies the intrinsic structure of the bound nucleon

 profound change in shell model :
 what occupies shell model states are NOT free nucleons
- Scalar polarizability is a natural source of three-body forces (NNN, HNN, HHN...)

 clear physical interpretation
- Naturally generates effective HN and HNN forces with no new parameters and predicts heavy neutron stars





II. Summary

• Need empirical confirmation of changing baryon structure:

- Response Functions & Coulomb sum rule
- EMC effect; spin EMC (not too long...)
- Change in Λ decay rate in nuclei?
- Initial systematic study of finite nuclei very promising With just 5 parameters:
 - Binding energies typically within 0.29% across periodic table
 - Super-heavies (Z > 100) especially good: 0.03%
 - Systematics of charge radii, deformations, shell and subshell closures pretty good
- Finally, we showed an exciting hint of the existence of a new BSM particle in the mass region 1-8 GeV





Special Mentions.....



Guichon



Tsushima



Saito



Stone



Krein



Matevosyan

ADELAIDE UNIVERSITY TRALIA



Cloët



Whittenbury



Simenel



Bentz





Motta



Antic



Kalaitzis



Martinez



Latest papers

- QMC π3;
 Martinez et al., Phys Rev C102 (2020) 034304
- Review: Guichon *et al.*, PPNP 100 (2018) 262
- SHE:

Stone et al., arXiv: 1901.06064

 Systematic application to finite nuclei: Stone et al., Phys Rev Lett 116 (2016) 092501





Key papers on QMC

- Many-body forces:
 - 1. Guichon, Matevosyan, Sandulescu, Thomas, Nucl. Phys. A772 (2006) 1.
 - 2. Guichon and Thomas, Phys. Rev. Lett. 93 (2004) 132502
- Built on earlier work on QMC: e.g.
 - 3. Guichon, Phys. Lett. B200 (1988) 235
 - 4. Guichon, Saito, Rodionov, Thomas, Nucl. Phys. A601 (1996) 349
- Major review of applications of QMC to many nuclear systems:
 - 5. Saito, Tsushima, Thomas,
 - Prog. Part. Nucl. Phys. 58 (2007) 1-167 (hep-ph/0506314)
 - 6. Guichon et al., Prog. Part. Nucl. Phys. 100 (2018) 262



References to: Covariant Version of QMC

- Basic Model: (Covariant, chiral, confining version of NJL)
- •Bentz & Thomas, Nucl. Phys. A696 (2001) 138
- Bentz, Horikawa, Ishii, Thomas, Nucl. Phys. A720 (2003) 95
- Applications to DIS:
- Cloet, Bentz, Thomas, Phys. Rev. Lett. 95 (2005) 052302
- Cloet, Bentz, Thomas, Phys. Lett. B642 (2006) 210
- Applications to neutron stars including SQM:
- Lawley, Bentz, Thomas, Phys. Lett. B632 (2006) 495
- Lawley, Bentz, Thomas, J. Phys. G32 (2006) 667







Experimental Constraints





From Graham et al., Annu. Rev. Nucl. Part. Sci. 71 (2021) 37



Very large scalar mean-fields are a fact

1970

R. BROCKMANN AND R. MACHLEIDT

TABLE II. Results of a relativistic Dirac-Brueckner calculation in comparison to the tential *B*. As a function of the Fermi momentum k_F , it is listed: the energy per nucleon vector potentials U_S and U_V , and the wound integral κ .

		Relativistic				
$\frac{k_F}{(\mathrm{fm}^{-1})}$	€ / A (MeV)	Μ̃/Μ	U_S (MeV)	U_V (MeV)	к (%)	
0.8	-7.02	0.855	-136.2	104.0	23.1	
0.9	-8.58	0.814	-174.2	134.1	18.8	
1.0	- 10.06	0.774	-212.2	164.2	16.1	
1.1	-11.18	0.732	-251.3	195.5	12.7	
1.2	-12.35	0.691	-290.4	225.8	11.9	
1.3	-13.35	0.646	-332.7	259.3	12.5	
1.35	-13.55	0.621	- 355.9	278.4	13.0	
1.4	-13.53	0.601	-374.3	293.4	13.8	
1.5	-12.15	0.559	-413.6	328.4	14.4	
1.6	-8.46	0.515	-455.2	371.0	15.8	





Brockmann and Machleidt, Phys Rev C52 (1990) 1965

Suggests a different approach : QMC Model

(Guichon, Saito, Tsushima et al., Rodionov et al., Stone - see Saito et al., Prog. Part. Nucl .Phys. 58 (2007) 1 and Guichon et al., Prog. Part. Nucl. Phys. 100 (2018) 262-297 for reviews)

- Start with quark model (MIT bag/NJL...) for all hadrons
- Introduce a relativistic Lagrangian with σ, ω and ρ mesons coupling to non-strange quarks
- Hence, initially only 4 parameters

 $(m_{\sigma}\,,\,g^{\sigma,\omega,\rho}\,_{q})$

- determine by fitting to:
 - $\rho_{0\,,}\,$ E/A and symmetry energy
- same in dense matter & finite nuclei
- Must solve <u>self-consistently</u> for the internal structure of baryons in-medium







Proton and Neutron Drip Lines







α Decay Half-Lives



FIGURE 6.10: Comparison of $\log_{10}(T_{1/2})$ predictions from FRDM, SV-min and QMC π -III along with values obtained from available data.



SPEcial RESEARCH CENTER FOR THE SUBATION THE STRUCTURED

Not as good as UNDEF1 and SV-min comparable to FRDM

Λ- and Ξ-Hypernuclei in QMC

	$^{89}_{\Lambda} \mathrm{Yb} \ (\mathrm{Expt.})$	$^{91}_{\Lambda}{\rm Zr}$	${}^{91}_{\Xi^0}\mathrm{Zr}$	$^{208}_{\Lambda} \mathrm{Pb} \ (\mathrm{Expt.})$	$^{209}_{\Lambda}{ m Pb}$	$^{209}_{\Xi^0}\mathrm{Pb}$
$1s_{1/2}$	-22.5	-24.0	-9.9	-27.0	-26.9	-15.0
$1p_{3/2}$		-19.4	-7.0		-24.0	-12.6
$1p_{1/2}$	-16.0(1p)	-19.4	-7.2	-22.0 (1p)	-24.0	-12.7
$1d_{5/2}$		-13.4	-3.1		-20.1	-9.6
$2s_{1/2}$		-9.1			-17.1	-8.2
$1d_{3/2}$	-9.0~(1d)	-13.4	-3.4	-17.0~(1d)	-20.1	-9.8
$1f_{7/2}$		-6.5			-15.4	-6.2
$2p_{3/2}$		-1.7			-11.4	-4.2
$1f_{5/2}$	-2.0~(1f)	-6.4		-12.0~(1f)	-15.4	-6.5
$2p_{1/2}$		-1.6			-11.4	-4.3

Also predicts **E** – hypernuclei bound by 5-15 MeV – being tested at J-PARC

"The first evidence of a bound state of Ξ^{-14} N system", K. Nakazawa et al., Prog. Theor. Exp. Phys. (2015) Guichon *et al*., Nucl.Phys. A814 (2008) 66; see also 1998



Change in Structure of Bound Λ

- Effect of the σ mean field is to modify the wave functions of the light quarks in the N
- Hence, the rates of vector and axial vector strangeness changing weak decays change in-medium





