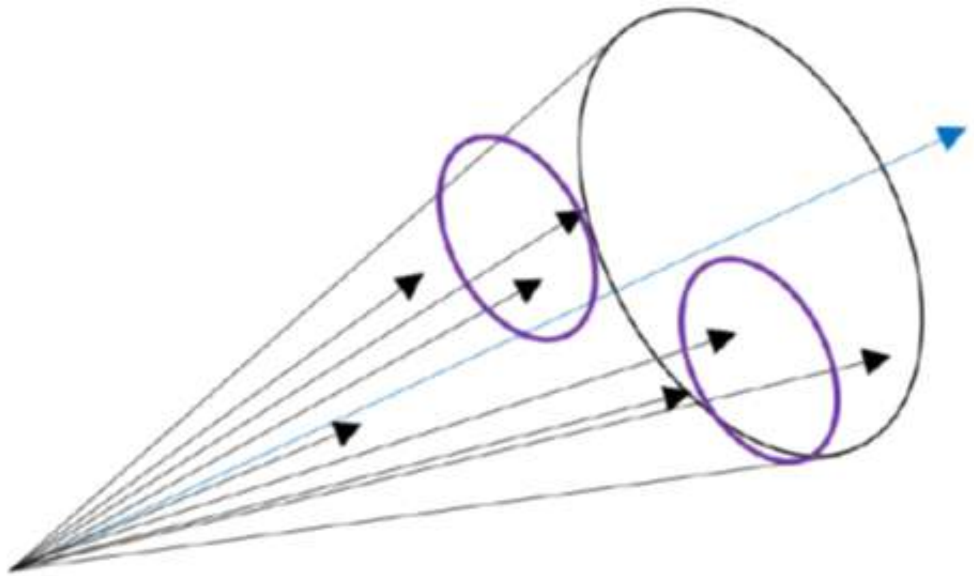
An Overview of Jets at the EIC

Brian Page

POETIC 2023

May 4th, 2023



Brookhaven National Laboratory

Outline

- ❑ Introduction: EIC Science, Jet History and Advantages

- ❑ Select Jet Measurements

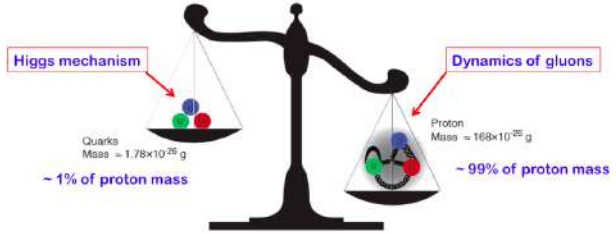
- ❑ Practical Aspects of Jet Measurements at EIC
  - ➢ Partonic -> Particle -> Jet Kinematics
  - ➢ Detector Considerations
  - ➢ Crossing Angle

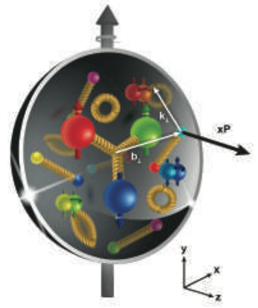
- ❑ Conclusions

The EIC Physics Pillars

How are the sea quarks and gluons, and their spins, distributed in space and momentum inside the nucleon?

How do the nucleon properties emerge from them and their interactions?





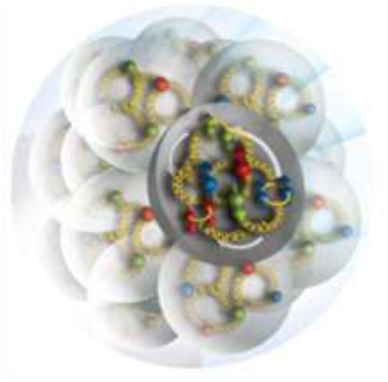
How do color-charged quarks and gluons, and colorless jets, interact with a nuclear medium?

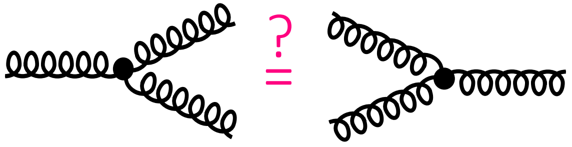
How do the confined hadronic states emerge from these quarks and gluons?

How do the quark-gluon interactions create nuclear binding?

How does a dense nuclear environment affect the quarks and gluons, their correlations, and their interactions?

What happens to the gluon density in nuclei? Does it saturate at high energy, giving rise to a gluonic matter with universal properties in all nuclei, even the proton?





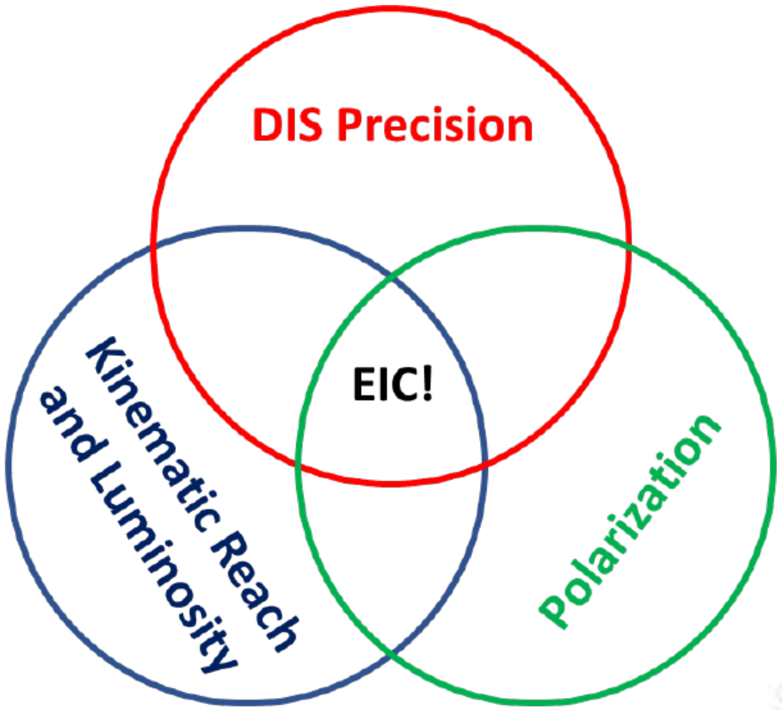
The Ultimate QCD Microscope

- Understanding internal dynamics of QCD bound states will require precision measurements over a wide kinematic regime – the EIC fits the bill

- Equally important will be high resolution, hermetic detectors with good PID and far-forward detection capabilities to fully characterize the final states of collisions

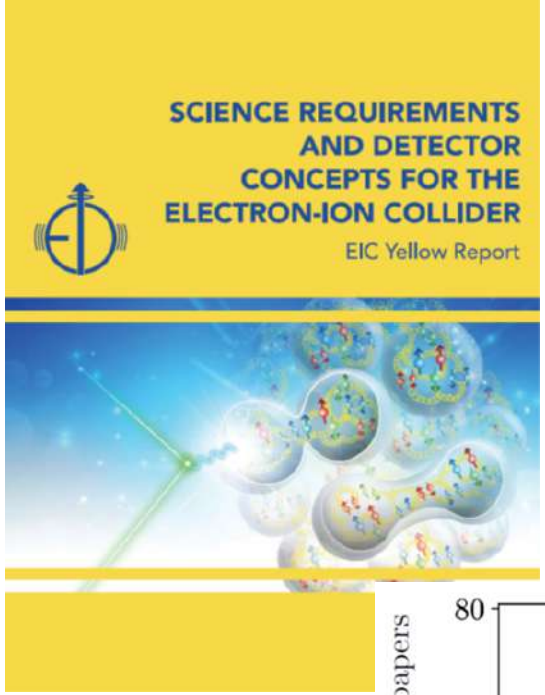
- Finally, need observables which are well understood theoretically, and which connect to the properties we want to explore

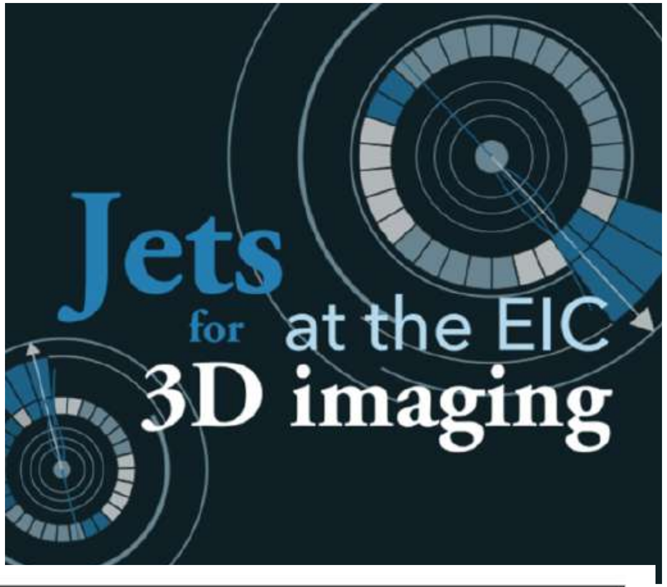
- This talk will focus on the unique capabilities jet observables will bring to the EIC science program

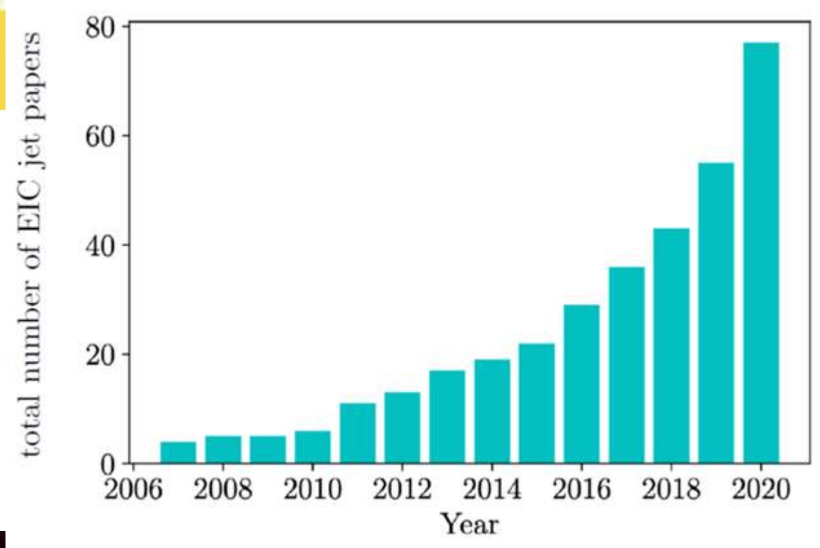


Jets at the EIC: A Brief History

- Initially, jets did not receive much attention and were only mentioned in passing in the initial EIC Whitepaper
  - EIC a relatively low energy machine – will jets be a useful probe?
  - SIDIS is well established and understood – will jets be redundant?
- Last ~10+ years have seen a significant uptick in interest in jet physics at the EIC as evidenced by:
  - Increasing number of papers
  - Dedicated treatment in EICUG Yellow Report and recent Snowmass process
  - Dedicated WGs in EIC proto-collaborations and current project detector collaboration - ePIC
- All this activity is driven by the realization that jets will have a robust and unique role to play in the EIC science program







Jet Physics at the EIC

Jets have several properties which will make them important tools for realizing the EIC physics program

- Well understood theoretically and experimentally

- Excellent proxies for the underlying parton kinematics

- Showers probe QCD from hard interaction to hadronization scale within the same event – can explore dynamics at different time (angular) scales

- Precision tools exist to probe these shower properties - substructure

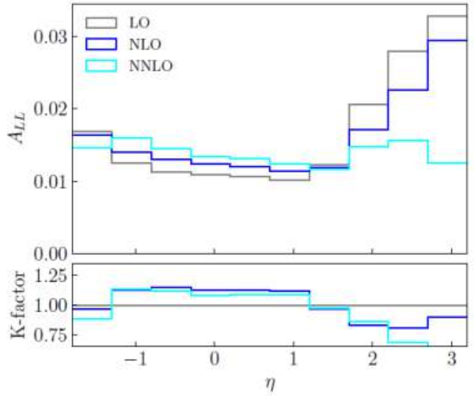
Jet Physics at the EIC

The importance of jet probes was reflected in the EIC Yellow Report where they touched on nearly every major physics topic (Nucl. Phys. A, Vol 1026, 122447)

Jets have several properties which will make them important tools for realizing the EIC physics program

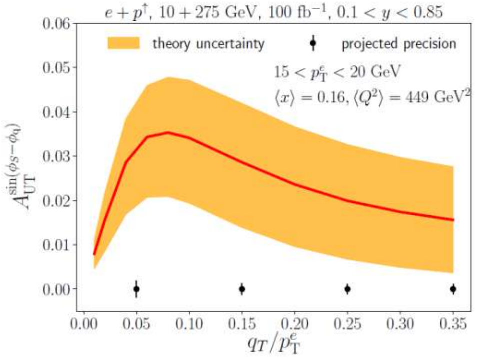
- Well understood theoretically and experimentally
- Excellent proxies for the underlying parton kinematics
- Showers probe QCD from hard interaction to hadronization scale within the same event – can explore dynamics at different time (angular) scales
- Precision tools exist to probe these shower properties - substructure

<u>Global properties and parton structure of hadrons</u>



Borsa, de Florian, Pedron `20

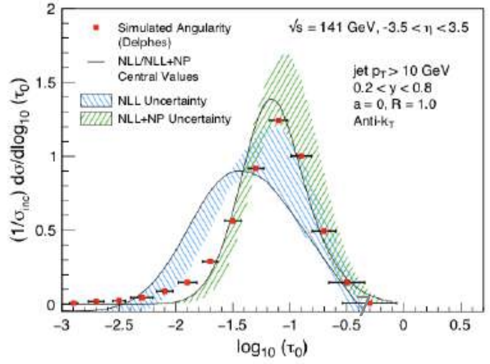
<u>Multi-dimensional imaging of nucleons, nuclei and mesons</u>



Arratia, Kang, Prokudin, Ringer `20

<u>The nucleus: a laboratory for QCD</u>

<u>Understanding hadronization</u>



J. Adam *et al* 2022 *JINST* 17 P10019

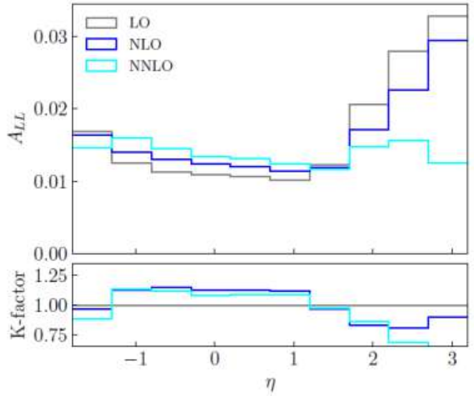
Jet Physics at the EIC

The importance of jet probes was reflected in the EIC Yellow Report where they touched on nearly every major physics topic (Nucl. Phys. A, Vol 1026, 122447)

Jets have several properties which will make them important tools for realizing the EIC physics program

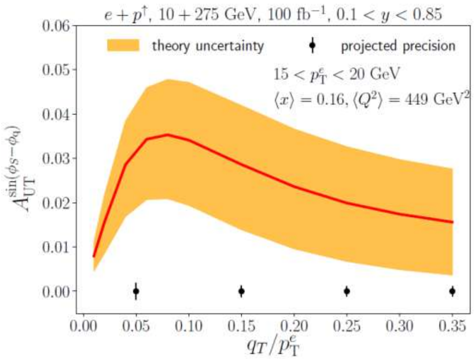
- Well understood theoretically and experimentally
- Excellent proxies for the underlying parton kinematics
- Showers probe QCD from hard interaction to hadronization scale within the same event – can explore dynamics at different time (angular) scales
- Precision tools exist to probe these shower properties - substructure

<u>Global properties and parton structure of hadrons</u>



Borsa, de Florian, Pedron `20

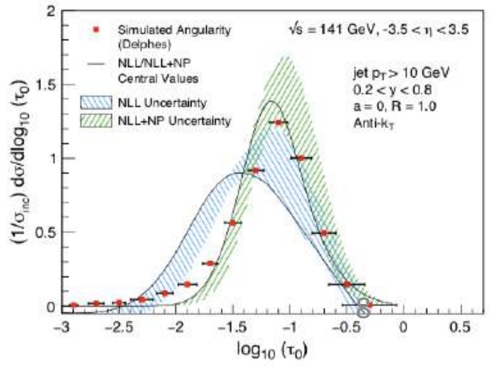
<u>Multi-dimensional imaging of nucleons, nuclei and mesons</u>



Arratia, Kang, Prokudin, Ringer `20

<u>The nucleus: a laboratory for QCD</u>

<u>Understanding hadronization</u>



J. Adam *et al* 2022 *JINST* 17 P10019

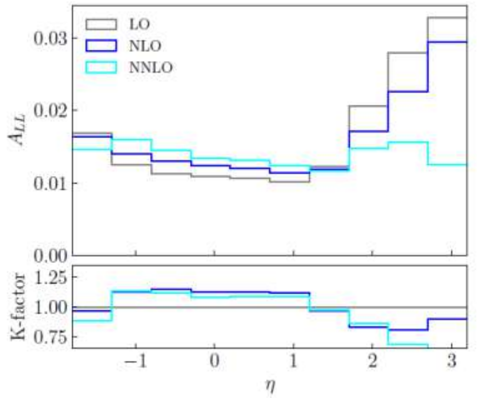
Jet Physics at the EIC

The importance of jet probes was reflected in the EIC Yellow Report where they touched on nearly every major physics topic (Nucl. Phys. A, Vol 1026, 122447)

Jets have several properties which will make them important tools for realizing the EIC physics program

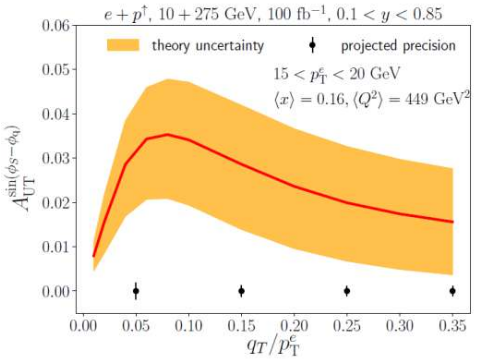
- Well understood theoretically and experimentally
- Excellent proxies for the underlying parton kinematics
- Showers probe QCD from hard interaction to hadronization scale within the same event – can explore dynamics at different time (angular) scales
- Precision tools exist to probe these shower properties - substructure

Global properties and parton structure of hadrons



Borsa, de Florian, Pedron `20

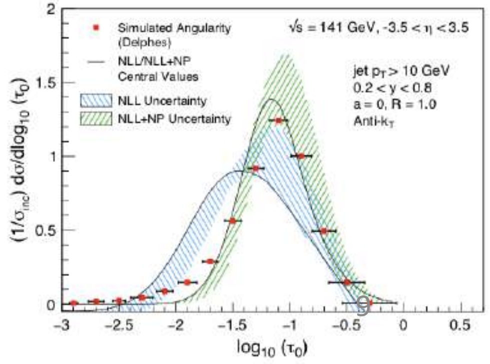
Multi-dimensional imaging of nucleons, nuclei and mesons



Arratia, Kang, Prokudin, Ringer `20

The nucleus: a laboratory for QCD

Understanding hadronization



J. Adam *et al* 2022 *JINST* 17 P10019

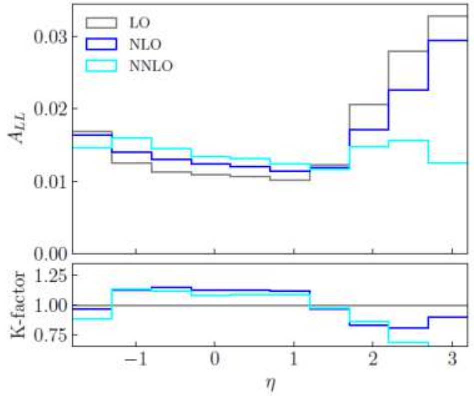
Jet Physics at the EIC

Jets have several properties which will make them important tools for realizing the EIC physics program

> The importance of jet probes was reflected in the EIC Yellow Report where they touched on nearly every major physics topic (Nucl. Phys. A, Vol 1026, 122447)

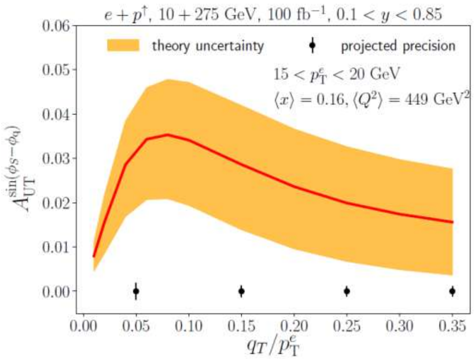
- Well understood theoretically and experimentally
- Excellent proxies for the underlying parton kinematics
- Showers probe QCD from hard interaction to hadronization scale within the same event – can explore dynamics at different time (angular) scales
- Precision tools exist to probe these shower properties - substructure

<u>Global properties and parton structure of hadrons</u>



Borsa, de Florian, Pedron \`20

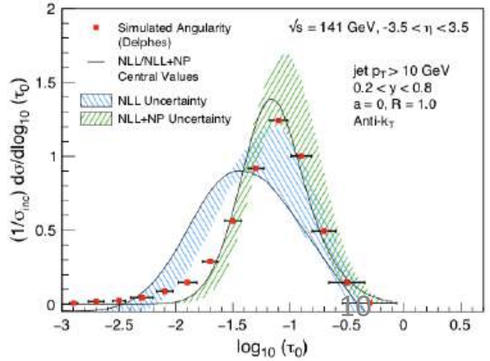
<u>Multi-dimensional imaging of nucleons, nuclei and mesons</u>



Arratia, Kang, Prokudin, Ringer \`20

<u>The nucleus: a laboratory for QCD</u>

<u>Understanding hadronization</u>



J. Adam *et al* 2022 *JINST* 17 P10019

Jet Physics at the EIC

Jets have several properties which will make them important tools for realizing the EIC physics program

➢ Well understood theoretically and experimentally

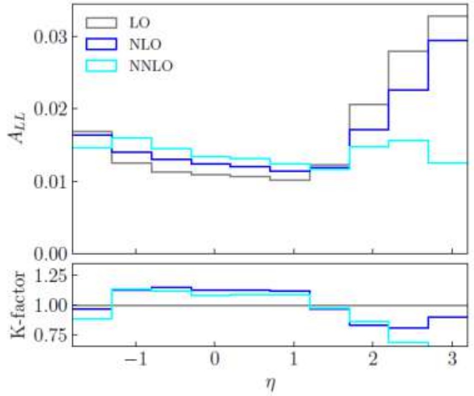
➢ Excellent proxies for the underlying parton kinematics

➢ Showers probe QCD from hard interaction to hadronization scale within the same event – can explore dynamics at different time (angular) scales

➢ Precision tools exist to probe these shower properties - substructure

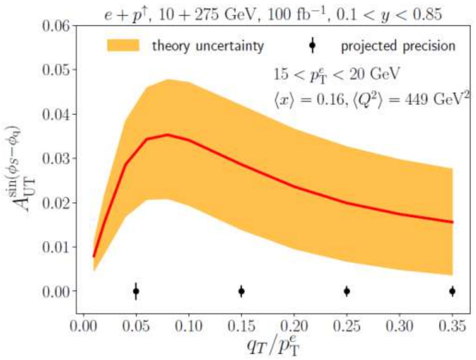
The importance of jet probes was reflected in the EIC Yellow Report where they touched on nearly every major physics topic (Nucl. Phys. A, Vol 1026, 122447)

Global properties and parton structure of hadrons



Borsa, de Florian, Pedron `20

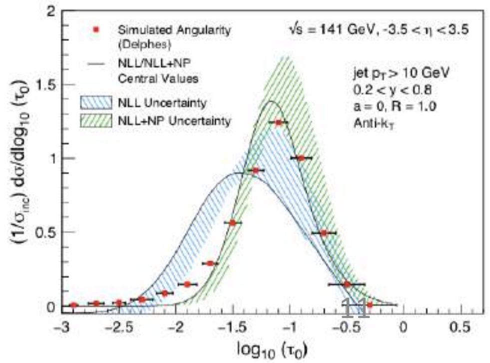
Multi-dimensional imaging of nucleons, nuclei and mesons



Arratia, Kang, Prokudin, Ringer `20

The nucleus: a laboratory for QCD

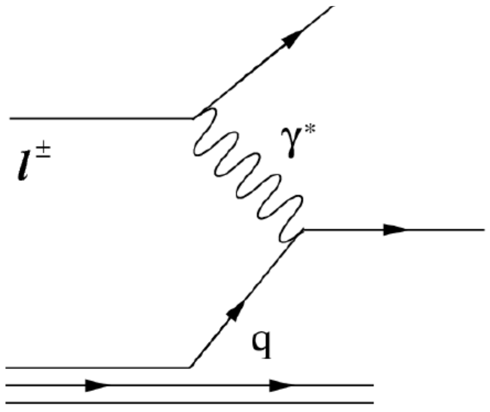
Understanding hadronization



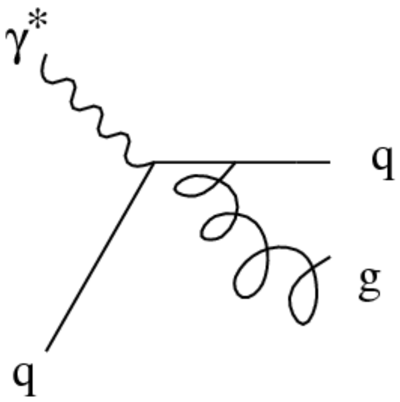
J. Adam *et al* 2022 *JINST 17* P10019

Relevant Subprocesses

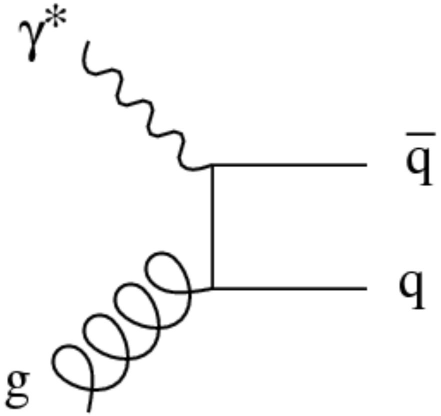
DIS



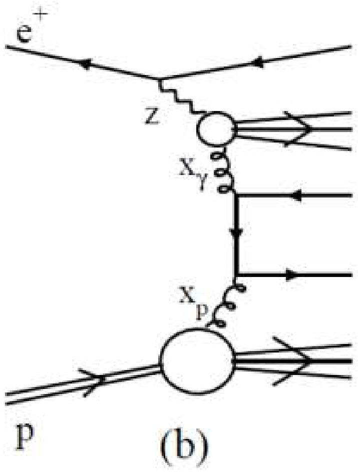
QCD-Compton (QCDC)



Photon-Gluon Fusion (PGF)



Resolved



- ❑ Leading order process gives rise to a single jet (not counting target remnant) whose kinematics are largely determined by the underlying event kinematics

- ❑ Higher-order corrections to this process can give rise to back-to-back jet configurations (dijets) which break the dependencies on event kinematics

- ❑ At low Q2, the hadronic (resolved) nature of the virtual photon becomes important and parton – parton (2 -> 2) scattering can give rise to dijet states

- ❑ Jets can also arise from diffractive events for example

Jets in the Yellow Report

Global properties and parton structure of hadrons

- Unpolarized parton structure of the proton and neutron
- Spin structure of the proton and neutron
- Inclusive and hard diffraction
- Global event shapes and the strong coupling constant

Multi-dimensional imaging of nucleons, nuclei and mesons

- Imaging of quarks and gluons in momentum space
- Wigner functions

The nucleus: a laboratory for QCD

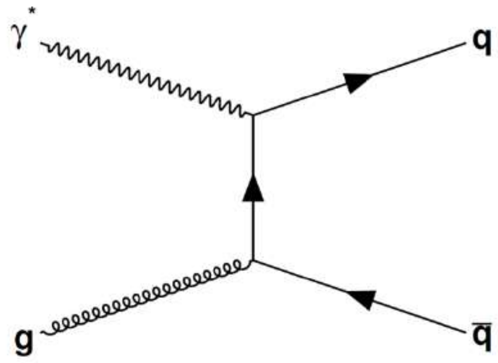
- High parton densities and saturation
- Particle propagation in matter and transport properties
- Special opportunities with jets and heavy quarks

Understanding hadronization

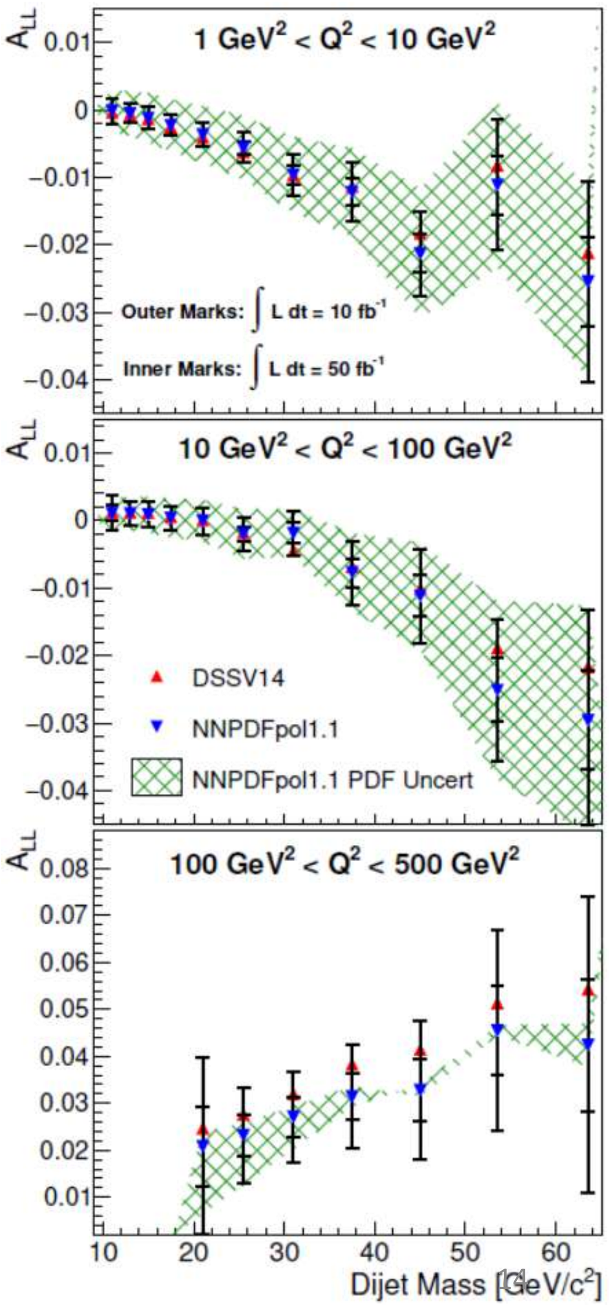
- Hadronization in the vacuum
- Hadronization in the nuclear environment

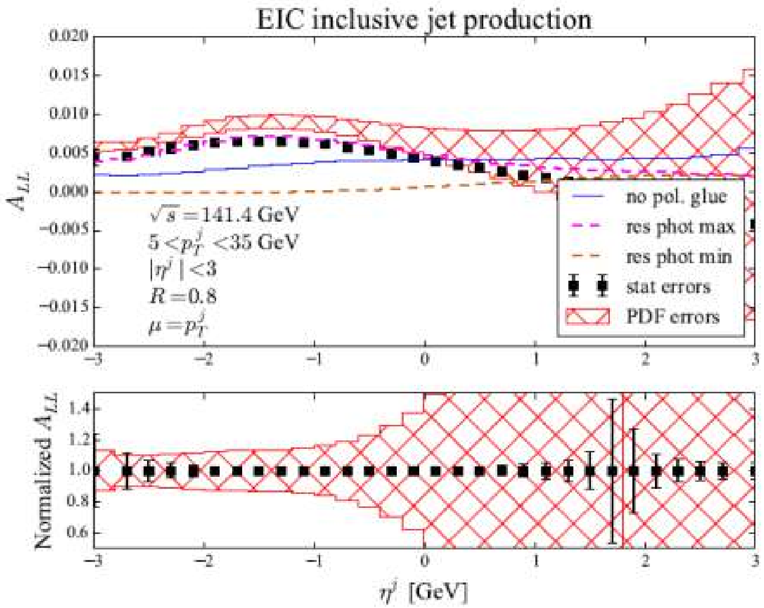
Longitudinal Spin Structure

- Recent results on inclusive jet A_{LL} at NLO and NNLO both with and without tagged lepton
- Will place strong constraints on helicity distributions
- Feasibility study for dijet A_{LL} in the Breit frame also performed – access to gluon via PGF process

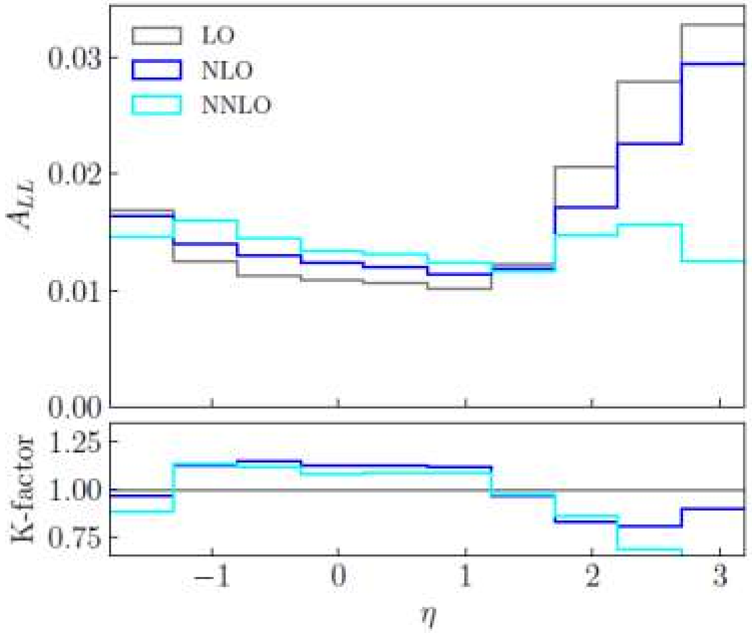


Page, Chu, Aschenauer `20





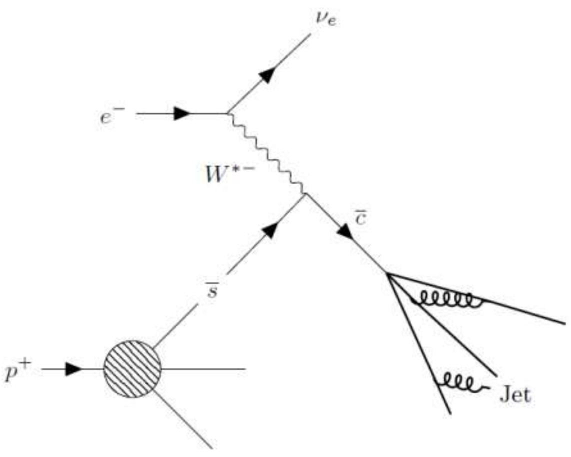
Boughezal, Petriello, Xing `18

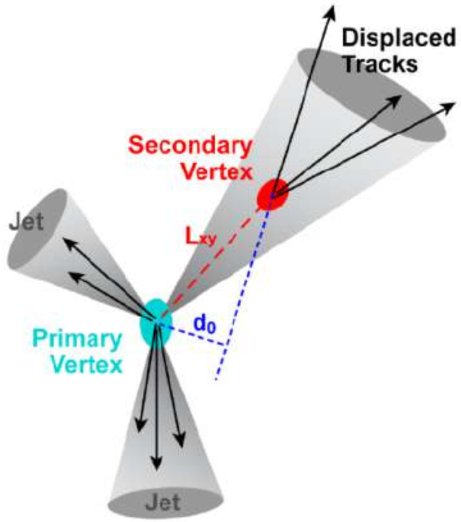


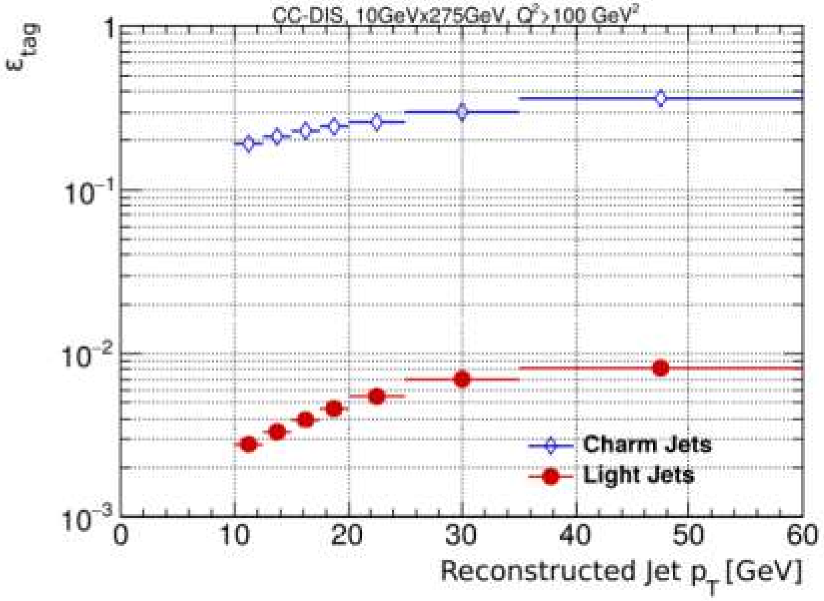
Borsa, de Florian, Pedron `20

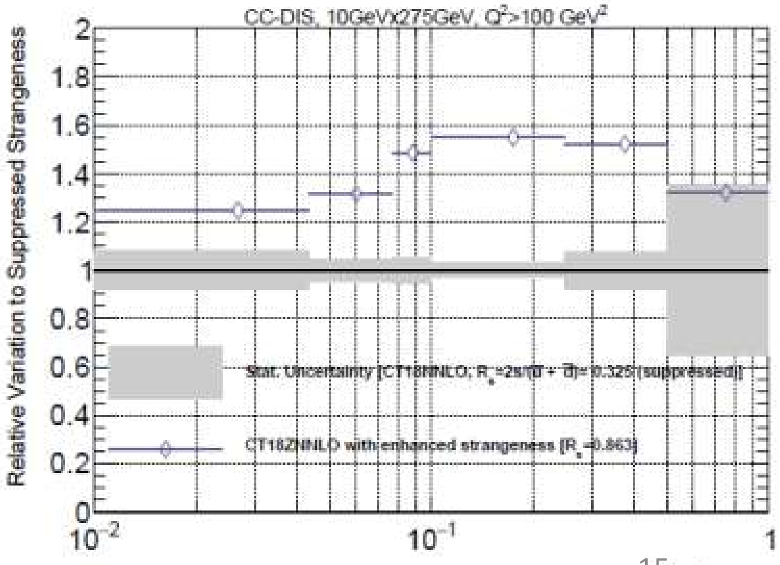
Strangeness PDF: Charm Jets

Arratia, Furletova, Hobbs, Olness, Sekula `20





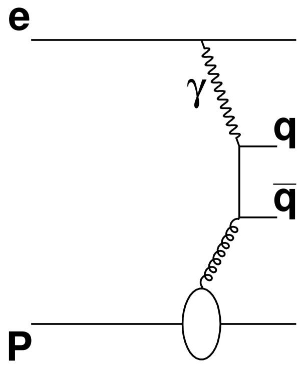


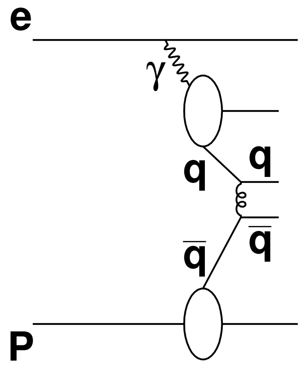


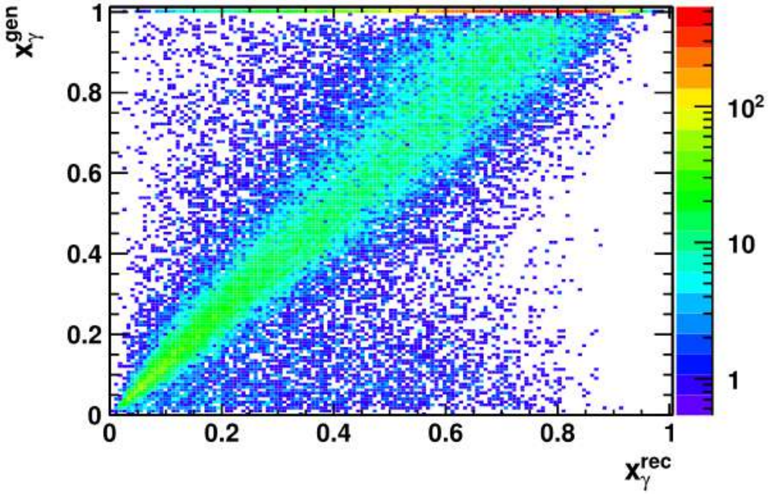
- Tension exists between neutrino DIS and SIDIS measurements of strange content and LHC extractions
- EIC is sensitive to strange content via charm production in charged-current DIS
- Charm is tagged within a jet via the presence of displaced tracks – good charm efficiency is seen, and methods are being refined
- Charm jet measurements at EIC should be able to discriminate between low and high strangeness scenarios

(Polarized) Photon Structure

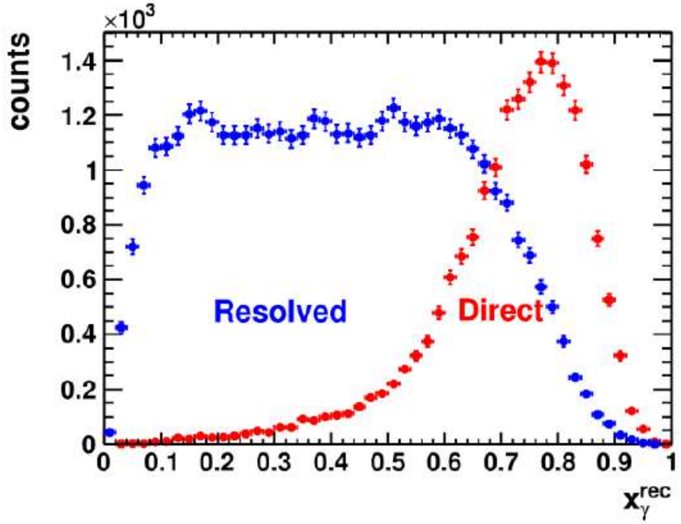
- At low Q2, virtual photon can behave hadronically and initiate 2->2 type scattering events
- Results in a quark/anti-quark final state with high transverse momentum
- Dijet allows to reconstruct event characteristics to separate signal and background and characterize the structure of the photon

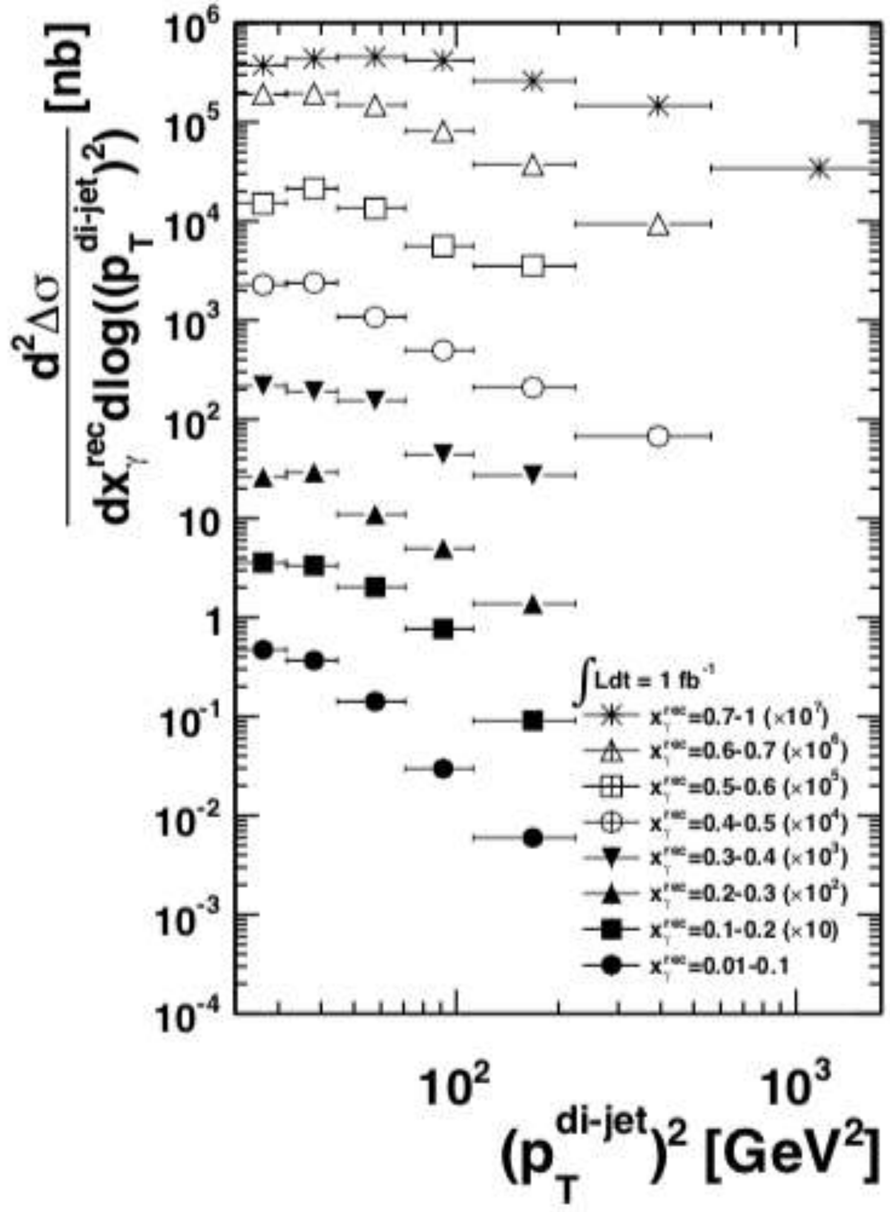






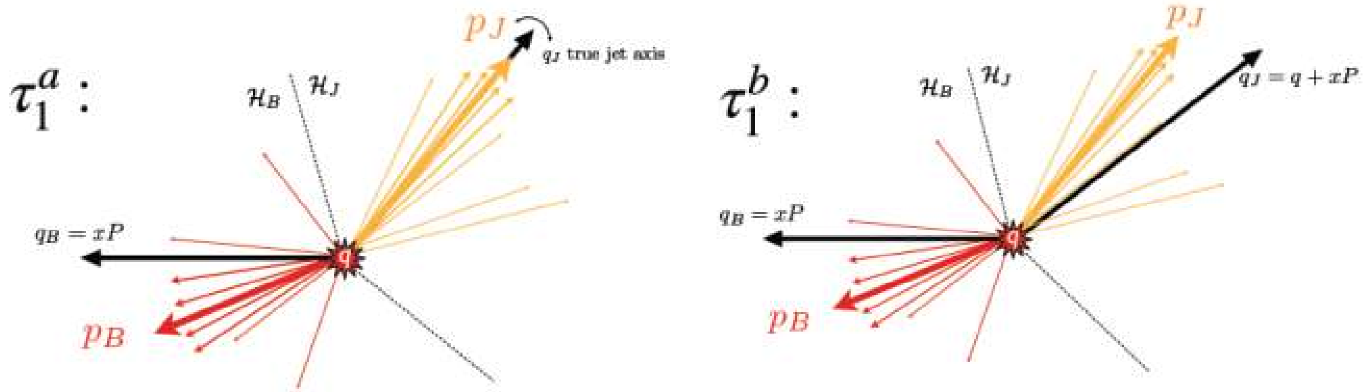
$$x_\gamma^{rec} = \frac{1}{2E_e y}\left(p_{T1}e^{-\eta_1} + p_{T2}e^{-\eta_2}\right)$$

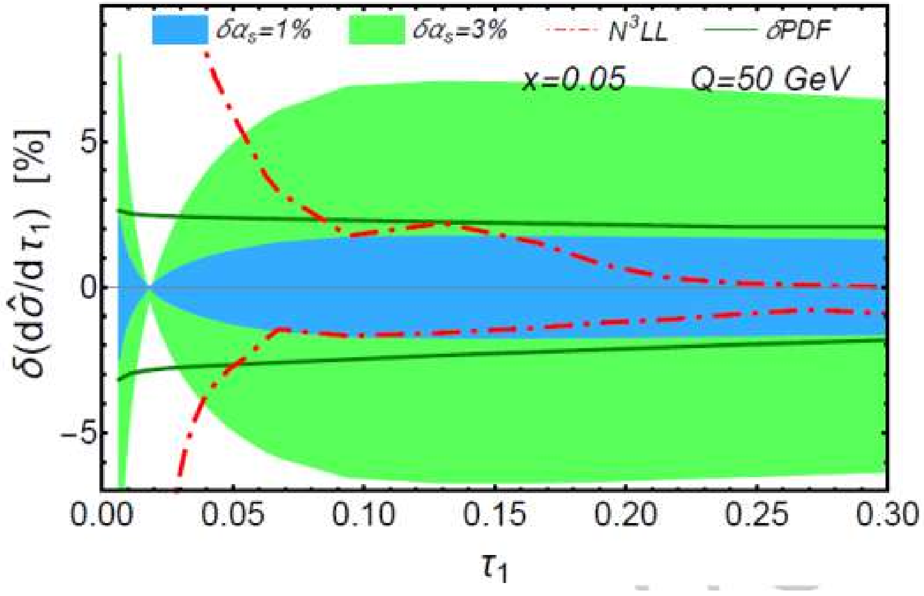


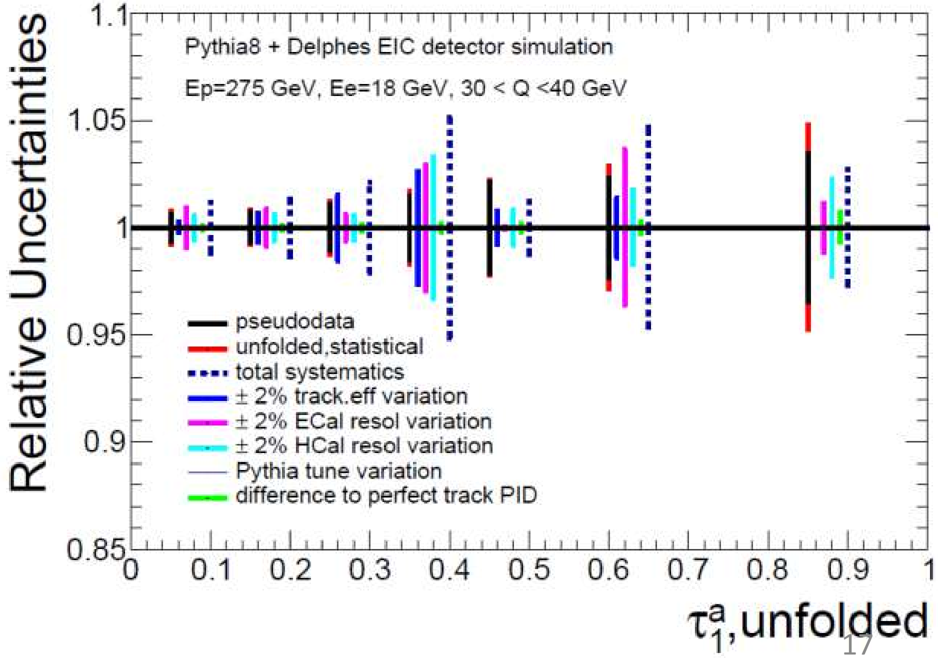


Global Event Shapes

$$\tau_1 = \frac{2}{Q^2} \sum_{i \in X} \min\{q_B \cdot p_i, q_J \cdot p_i\}$$







- Global event shapes offer possibility of very high precision measurements for extractions of non-perturbative parameters such as the strong coupling constant

- Detailed feasibility studies of 1-jettiness observable were carried out as part of the Yellow Report effort taking into account uncertainties on tracking efficiency and calorimeter resolution

- At N^{3}LL, roughly 1% precision is possible, challenging experimental problem, but recent studies show promise

Jets in the Yellow Report

Global properties and parton structure of hadrons

- Unpolarized parton structure of the proton and neutron
- Spin structure of the proton and neutron
- Inclusive and hard diffraction
- Global event shapes and the strong coupling constant

Multi-dimensional imaging of nucleons, nuclei and mesons

- Imaging of quarks and gluons in momentum space
- Wigner functions

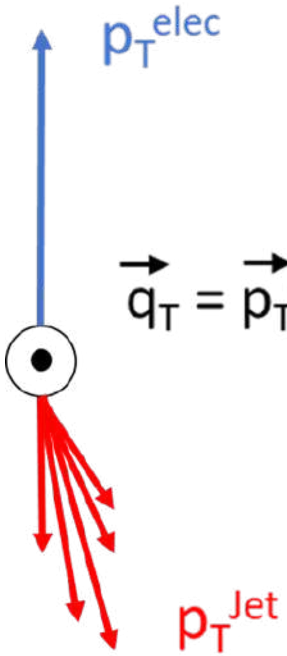
The nucleus: a laboratory for QCD

- High parton densities and saturation
- Particle propagation in matter and transport properties
- Special opportunities with jets and heavy quarks

Understanding hadronization

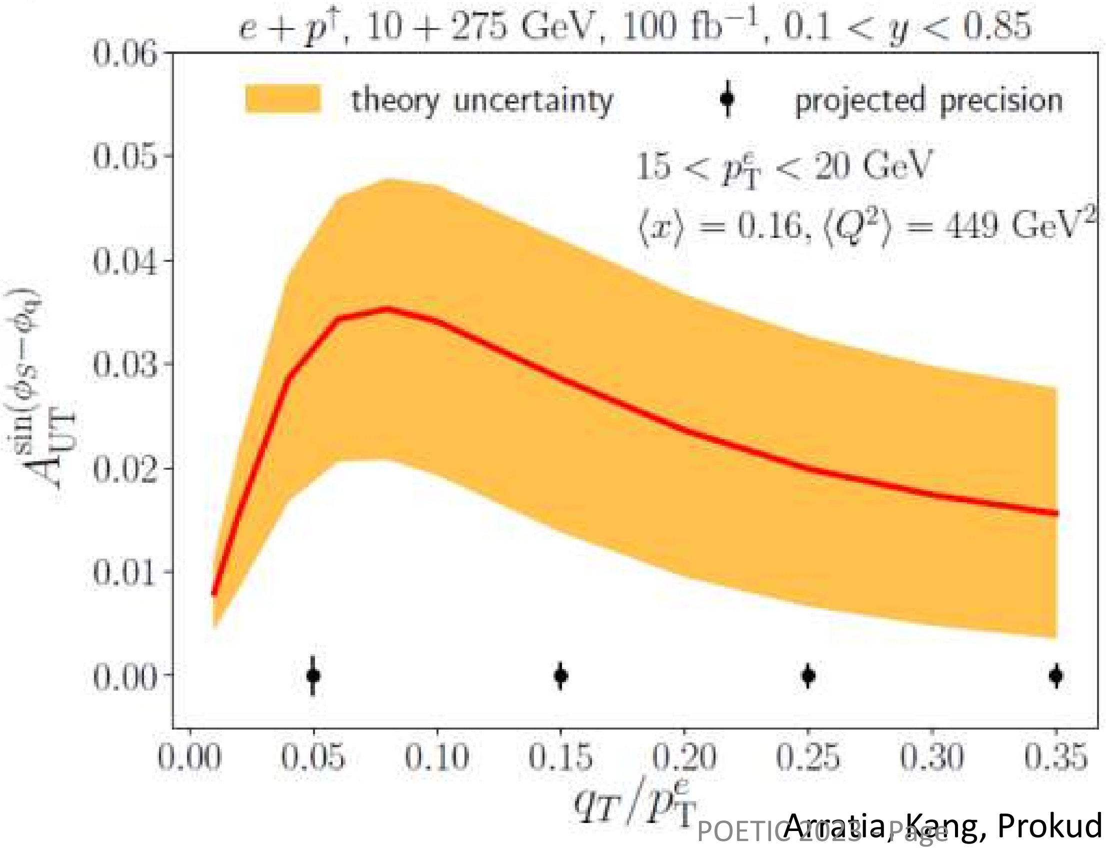
- Hadronization in the vacuum
- Hadronization in the nuclear environment

Lepton-Jet Correlations: Sivers TMD



$\vec{q}_T = \vec{p}_T^{\,elec} + \vec{p}_T^{\,Jet}$

$$e + p(\vec{s}_T) \to e + (\text{jet}(\vec{q}_T)\, h(z_h, \vec{j}_T)) + X$$

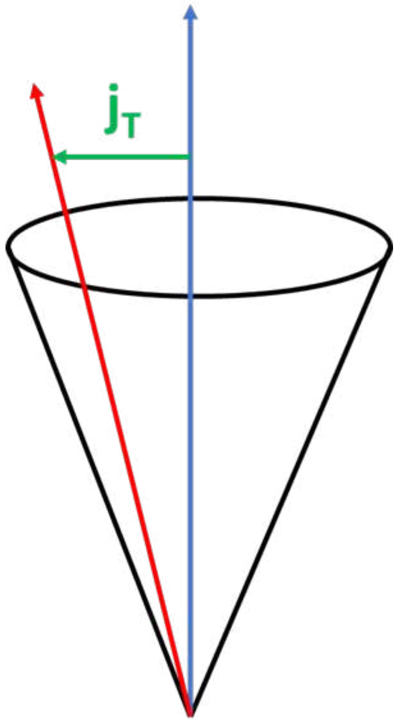


Arratia, Kang, Prokudin, Ringer `20

- ❑ Jet measurements for 3D imaging of nucleons at the EIC is emerging as a fruitful field
- ❑ Jets are complementary to standard SIDIS extractions of TMDs and provide better surrogates for parton kinematics while allowing cleaner separation of target and current fragmentation regions
- ❑ Jet measurements allow independent constraints on TMD PDFs and FFs from a single measurement
- ❑ Azimuthal correlation between jet and lepton sensitive to TMD PDFs (Sivers)

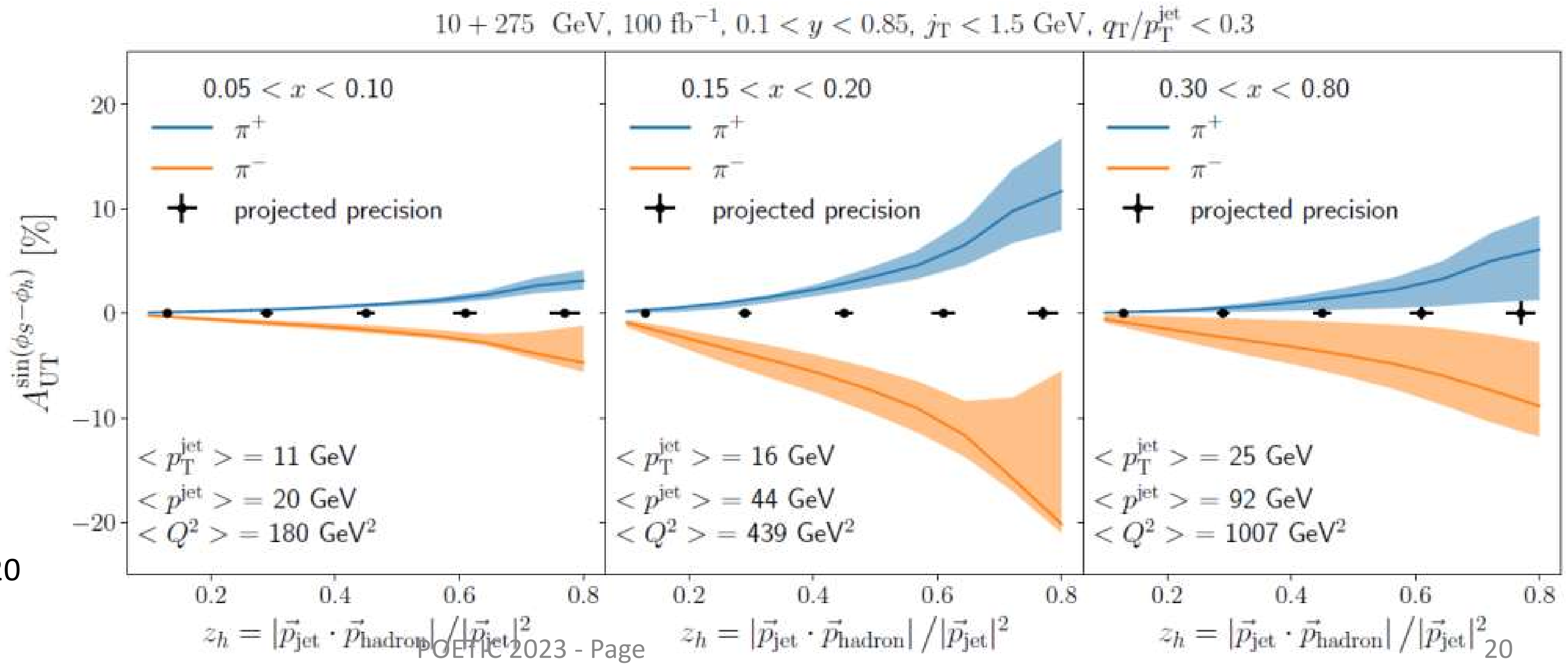
Hadron-in-Jet: Collins TMD

$$e + p(\vec{s}_T) \rightarrow e + (\text{jet}(\vec{q}_T)\, h(z_h, \vec{j}_T)) + X$$

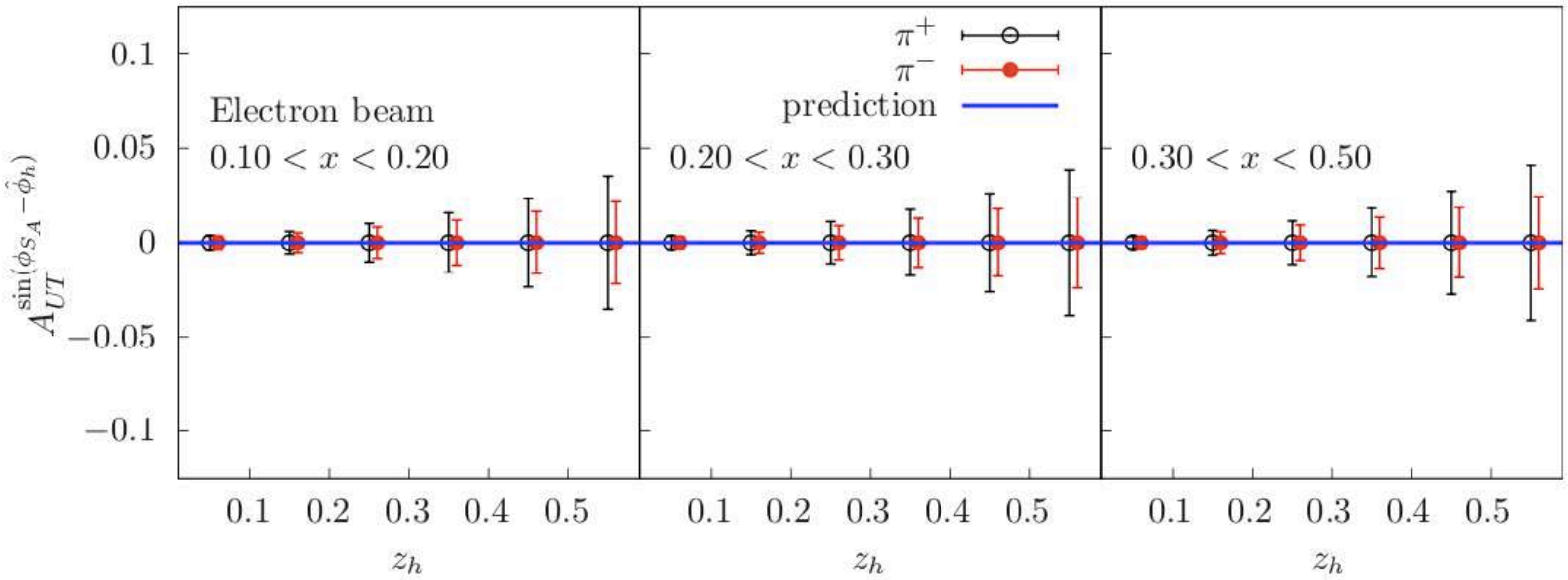


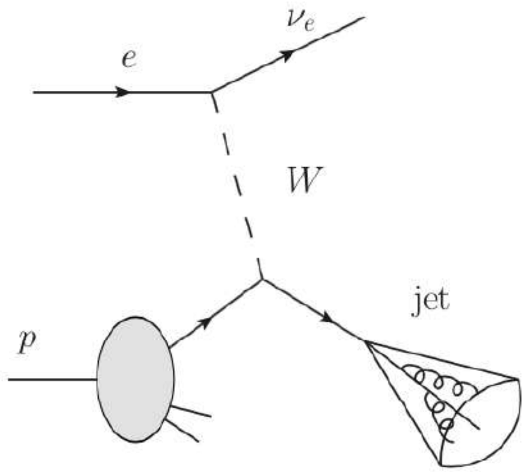
Arratia, Kang, Prokudin, Ringer `20

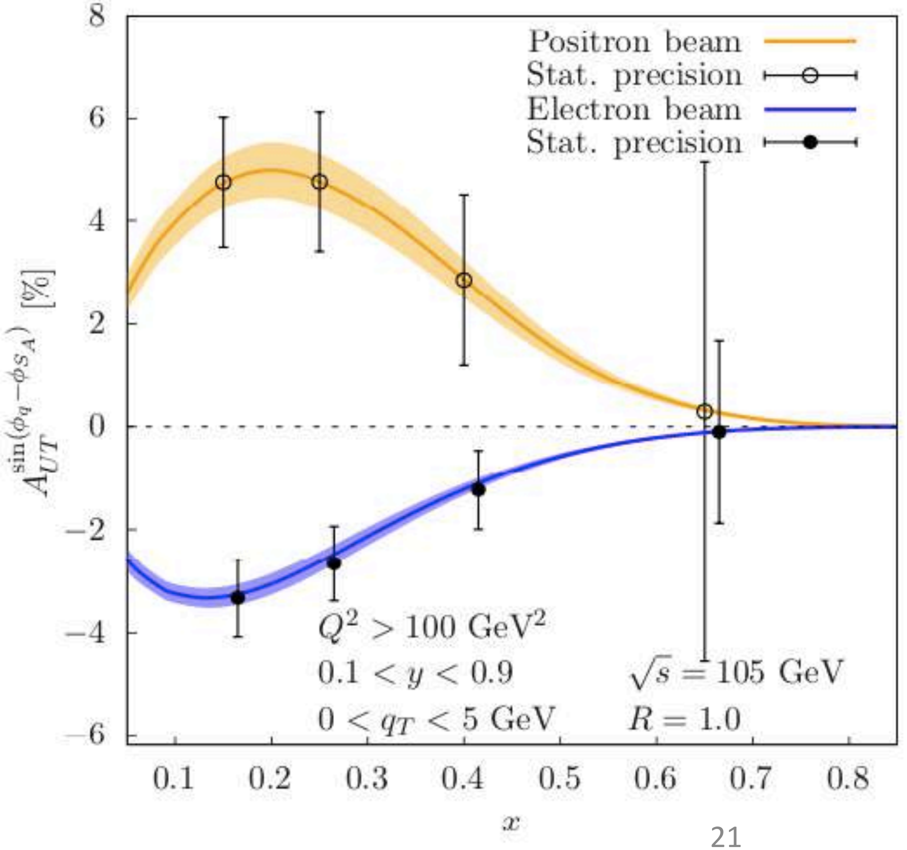
- Measurement of hadrons within jet give access to TMD FFs
- Relevant variables are j_T – transverse momentum of the hadron with respect to the jet and z – fraction of jet momentum carried by hadron
- Collins asymmetry correlates proton spin vector with j_T
- Identified hadrons allow for flavor separation of Collins FF



Charged Current Lepton-Jet Correlations



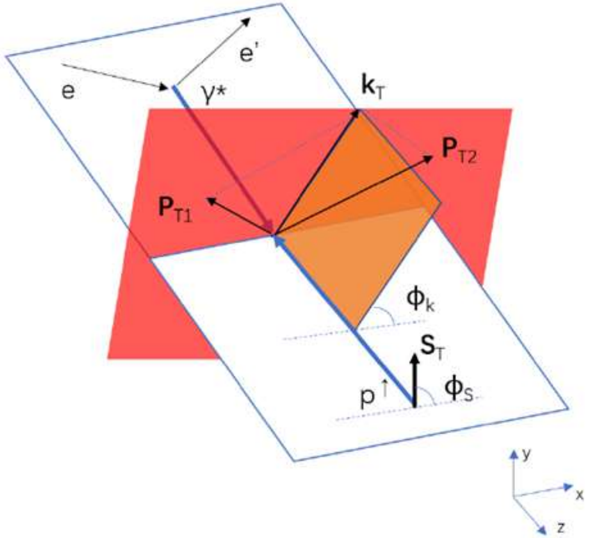




- Previous lepton-jet and hadron-in-jet measurements can also be carried out in charged current DIS where the outgoing neutrino is deduced via missing transverse energy
- Charge conservation leads to flavor separation
- Chiral odd nature of interaction means Collins asymmetries vanish at leading order

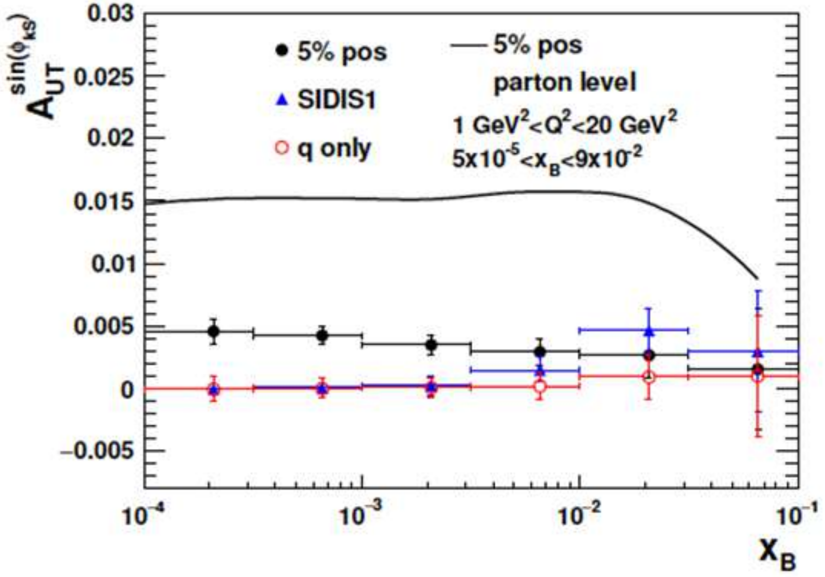
Dijet Correlations: Gluon Sivers TMD

Phys. Rev. D 98, 034011 (2018)

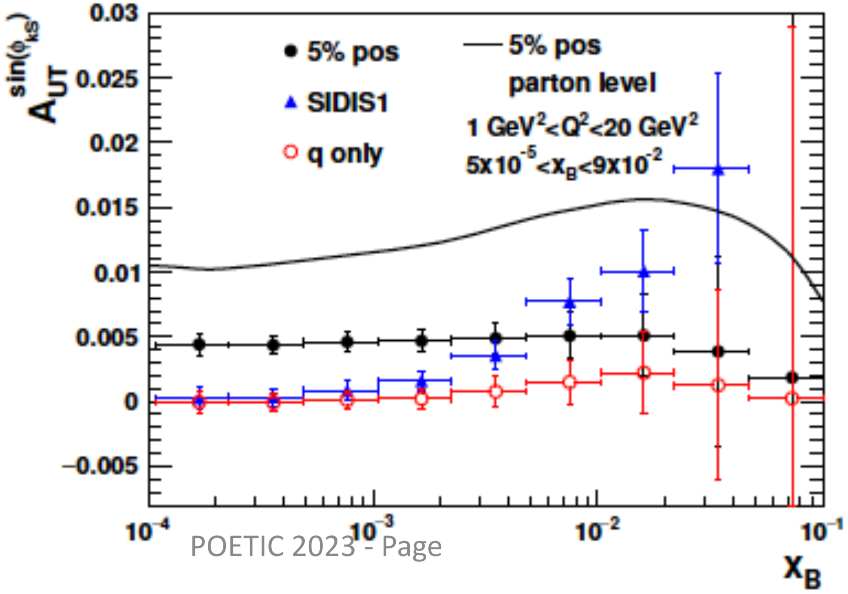


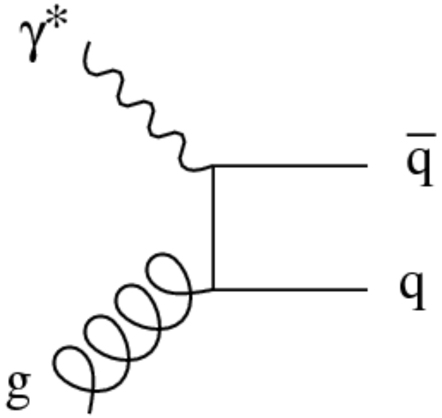
- Modulations of the angle between the proton spin vector and the sum of the di-parton system provide access to gluon sivers function
- Use of dijets has several advantages over di-hadrons including lower dilution of asymmetry and better separation between models of gluon sivers effect
- Jets don't suffer from uncertainties arising due to fragmentation (although hadronization still a concern)

Di-Hadrons



Dijets





Jets in the Yellow Report

Global properties and parton structure of hadrons

- Unpolarized parton structure of the proton and neutron
- Spin structure of the proton and neutron
- Inclusive and hard diffraction
- Global event shapes and the strong coupling constant

Multi-dimensional imaging of nucleons, nuclei and mesons

- Imaging of quarks and gluons in momentum space
- Wigner functions

The nucleus: a laboratory for QCD

- High parton densities and saturation
- Particle propagation in matter and transport properties
- Special opportunities with jets and heavy quarks

Understanding hadronization

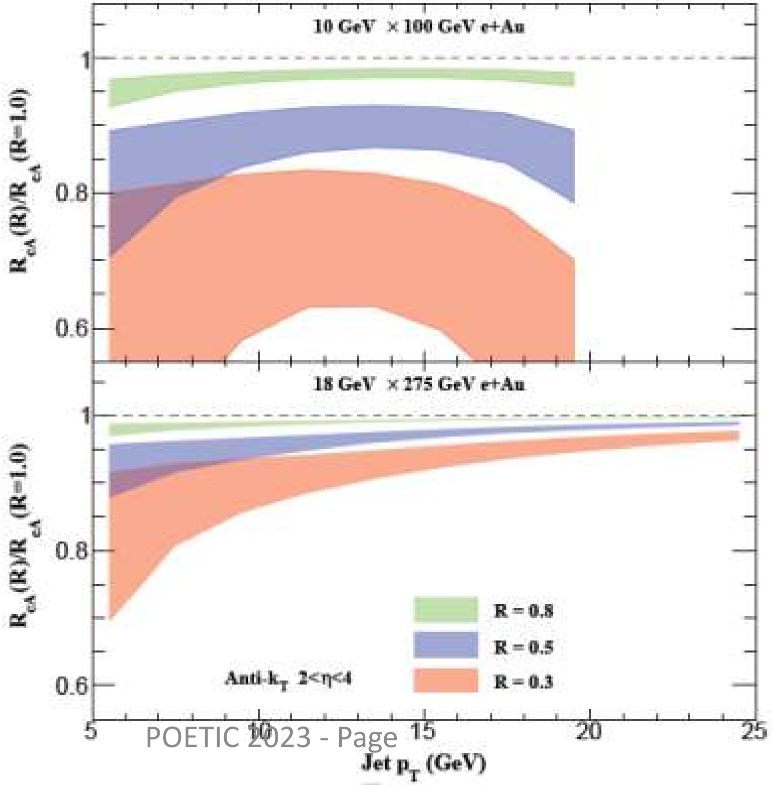
- Hadronization in the vacuum
- Hadronization in the nuclear environment

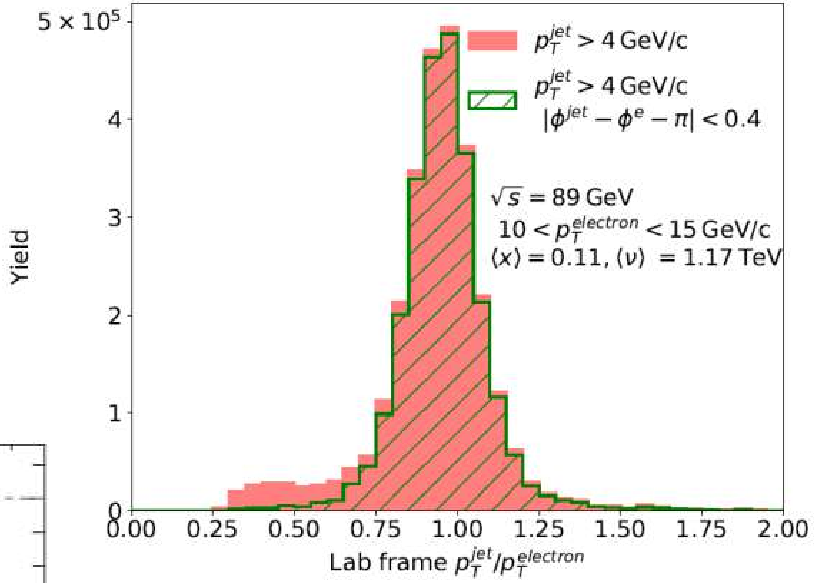
Jets in the Medium: CNM Properties

- Many opportunities to study the properties of cold nuclear matter with jets
- Simple comparisons of jet yields in ep vs eA will be informative – double ratio $R_{eA}(R)/R_{eA}(R = 1.0)$ will reduce impact from nPDFs and enhance final state effects
- Lepton – Jet correlations in Born level DIS can be thought of as analogous to boson – Jet measurements with the lepton as the tag and the jet as the probe of the medium
- Dijets and gamma-dijet correlations also expected to be powerful probes of saturation / small-x dynamics

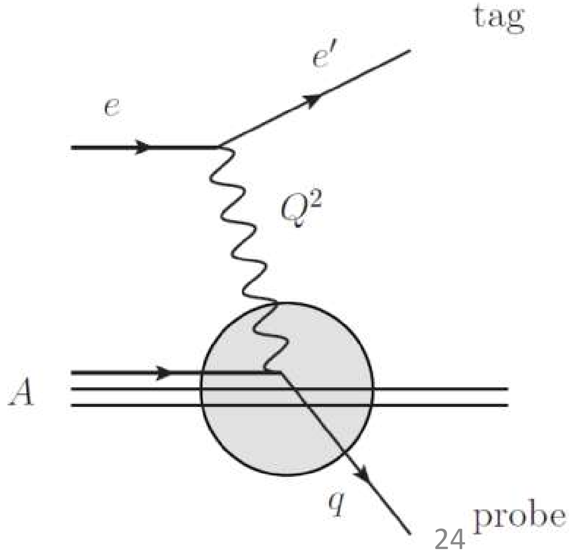
$$R_{eA}(R) = \frac{1}{A}\frac{\int_{\eta 1}^{\eta 2} d\sigma/d\eta dp_T|_{e+A}}{\int_{\eta 1}^{\eta 2} d\sigma/d\eta dp_T|_{e+p}}$$

Li & Vitev '20





Arratia, Song, Ringer, Jacak '20



Jets in the Yellow Report

Global properties and parton structure of hadrons

- Unpolarized parton structure of the proton and neutron
- Spin structure of the proton and neutron
- Inclusive and hard diffraction
- Global event shapes and the strong coupling constant

Multi-dimensional imaging of nucleons, nuclei and mesons

- Imaging of quarks and gluons in momentum space
- Wigner functions

The nucleus: a laboratory for QCD

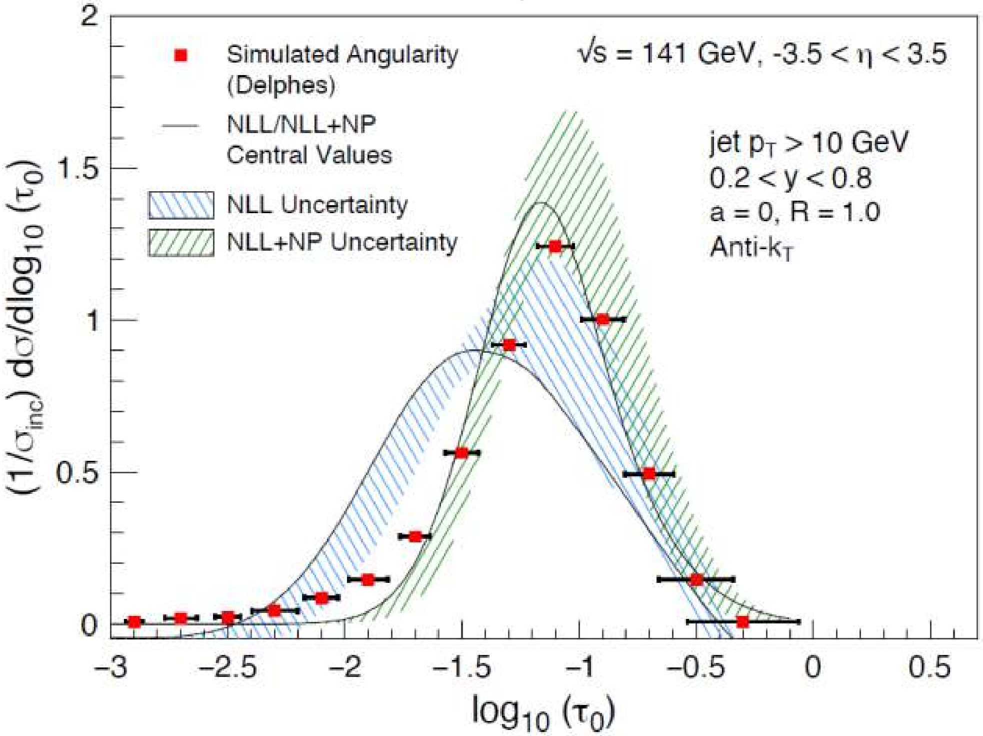
- High parton densities and saturation
- Particle propagation in matter and transport properties
- Special opportunities with jets and heavy quarks

Understanding hadronization

- Hadronization in the vacuum
- Hadronization in the nuclear environment

Jet Substructure: Angularity

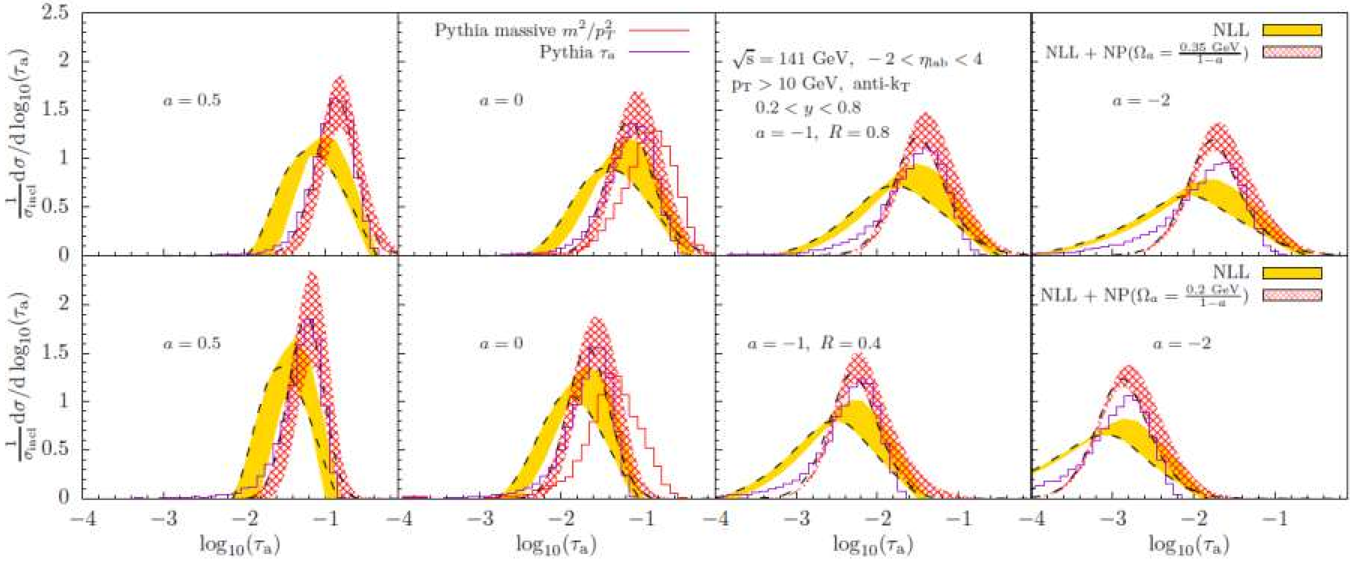
$$\tau_a \equiv \frac{1}{p_T} \sum_{i \in J} p_T^i \left(\Delta R_{iJ}\right)^{2-a}$$



Aschenauer, Lee, Page, Ringer `20 & ATHENA Proposal

$$F_\kappa(k) = \left(\frac{4k}{\Omega_\kappa^2}\right) \exp\left(-\frac{2k}{\Omega_\kappa}\right)$$ NP Effects

- Jet angularity are a family of one-parameter substructure observables correlating momentum and radial distance of particles in a jet
- Different choices of ‘a’ parameter interpolate between familiar substructure observables such as mass and broadening
- Sensitive to hadronization effects via convolution with the non-perturbative shape function Ω_1

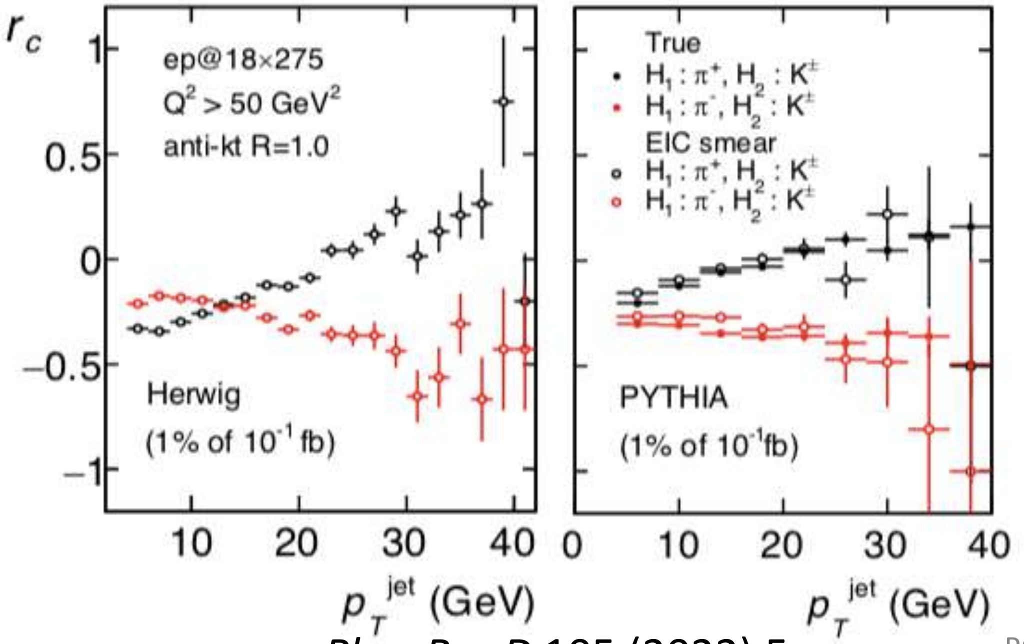


Leading Di-hadron Correlations

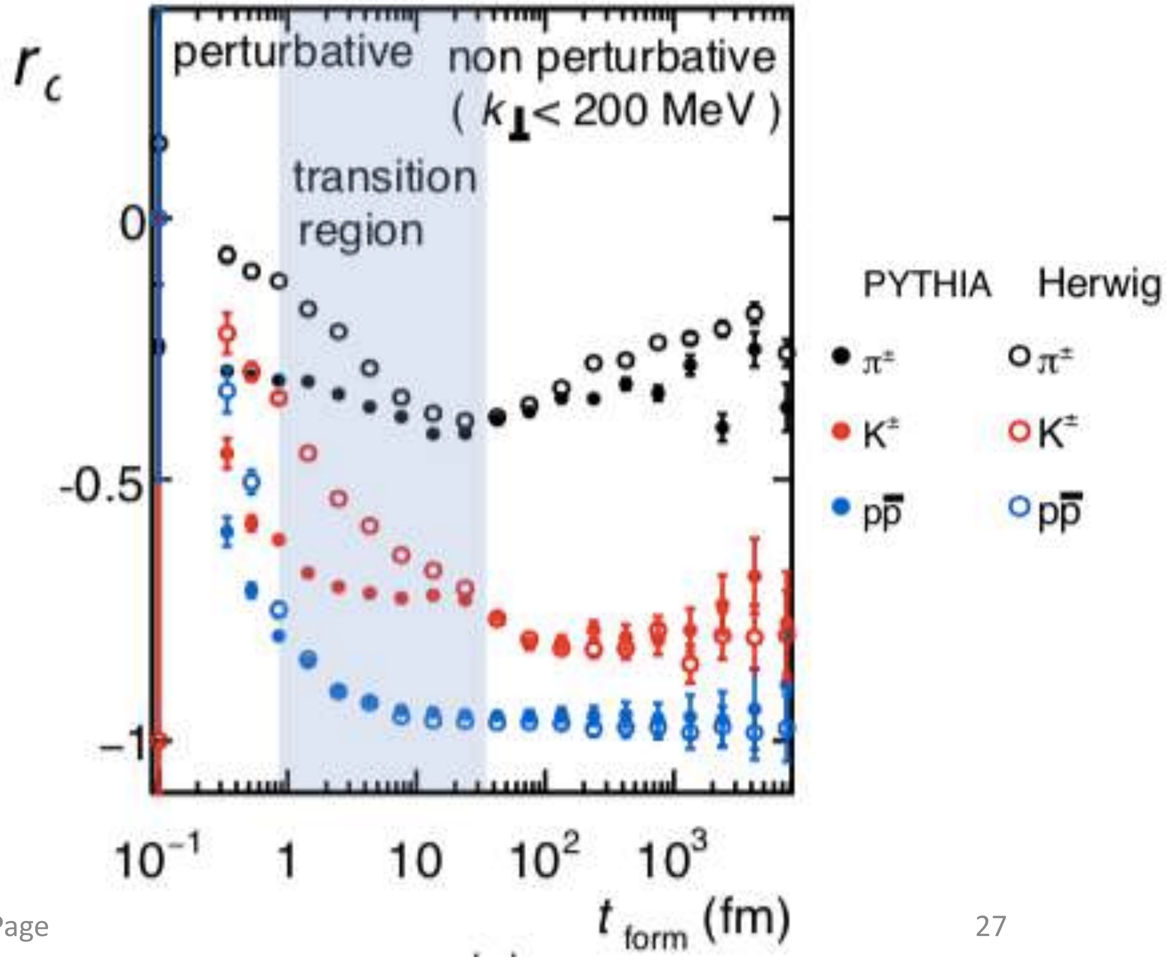
$$r_c(X) = \frac{d\sigma_{h_1 h_2}/dX - d\sigma_{h_1 \overline{h_2}}/dX}{d\sigma_{h_1 h_2}/dX + d\sigma_{h_1 \overline{h_2}}/dX}$$

- Quantify distributions of charge, flavor, spin, etc within a jet -> study hadronization process
- Define the ratio r_c to explore the charge and flavor correlations between leading hadrons within a jet
- See differences in how charge and flavor are distributed vs jet p_T and formation time between different hadronization models

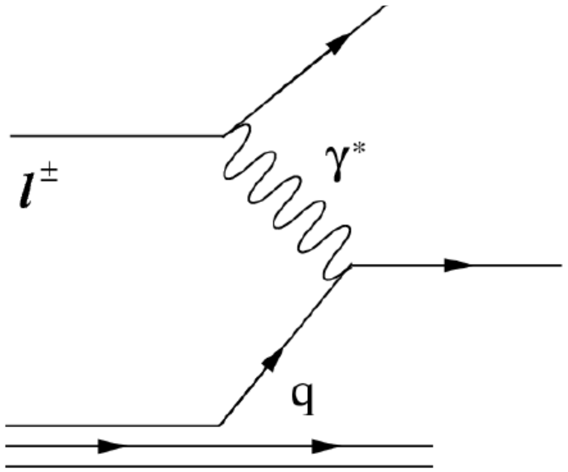
$t_{form} = \frac{z(1-z)p}{k_\perp^2}; p = p_1 + p_2; z = \frac{p_2}{p}$; $k_\perp$= rel trans mom

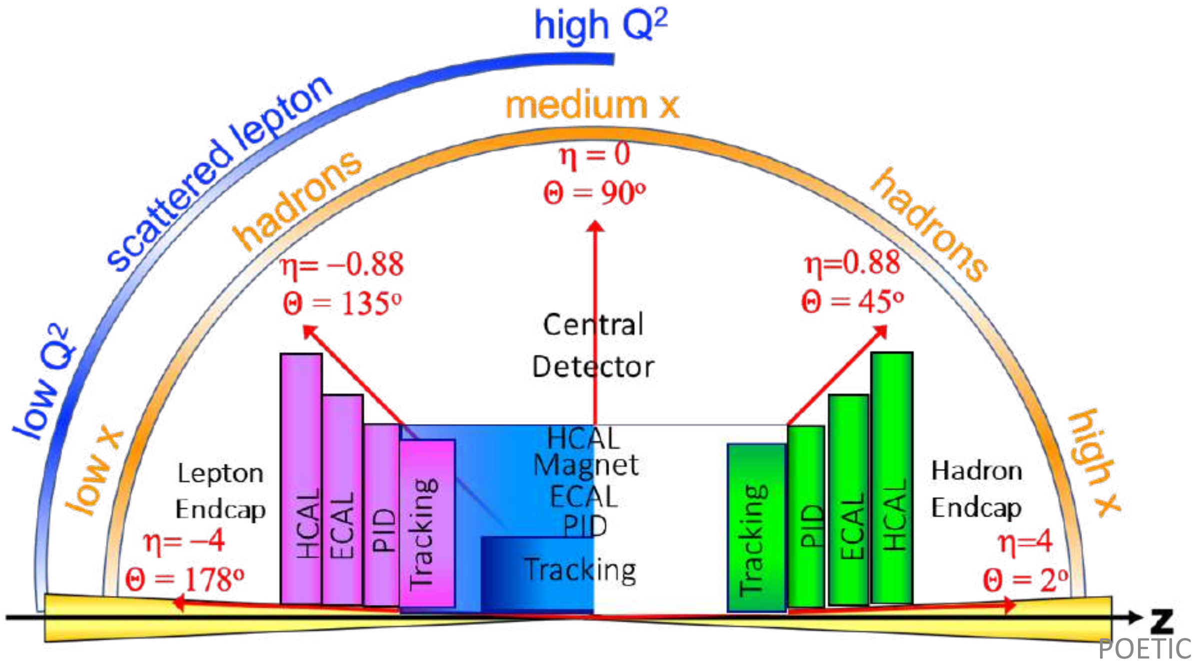


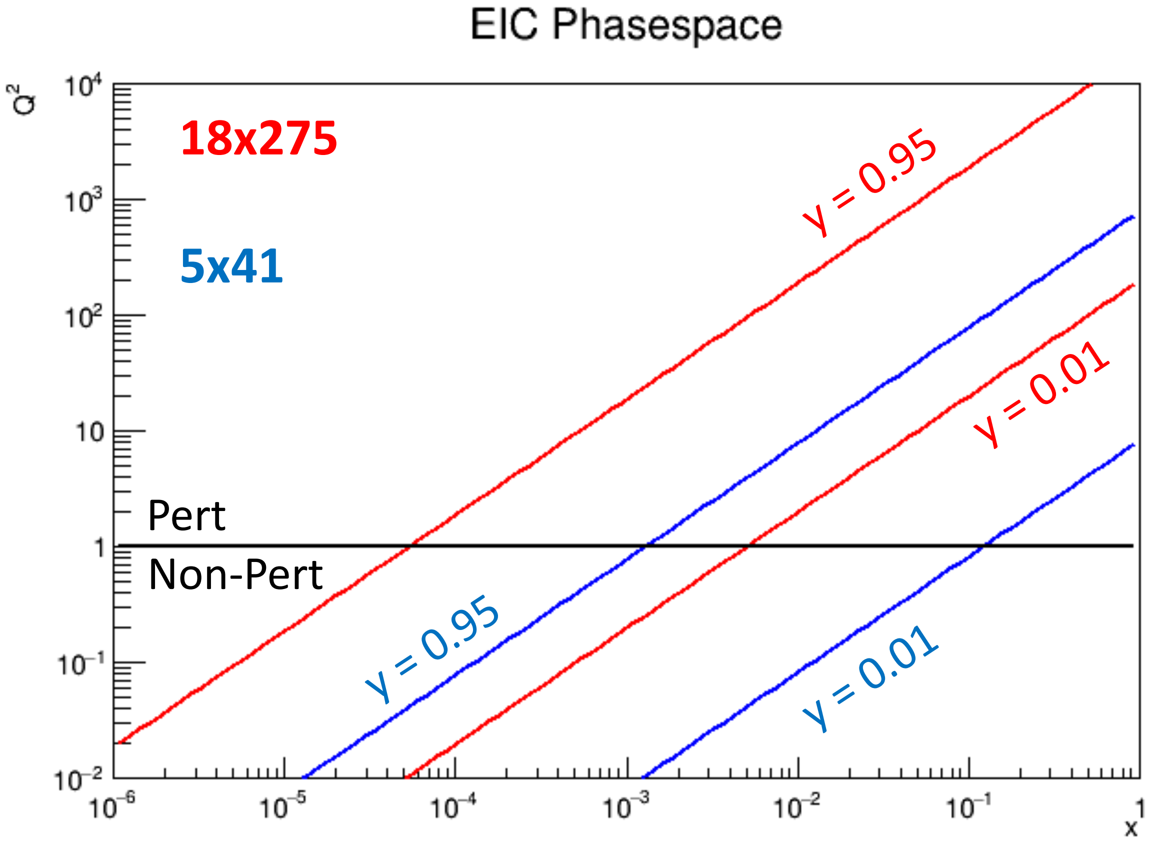
Phys.Rev.D 105 (2022) 5



DIS Event Kinematics





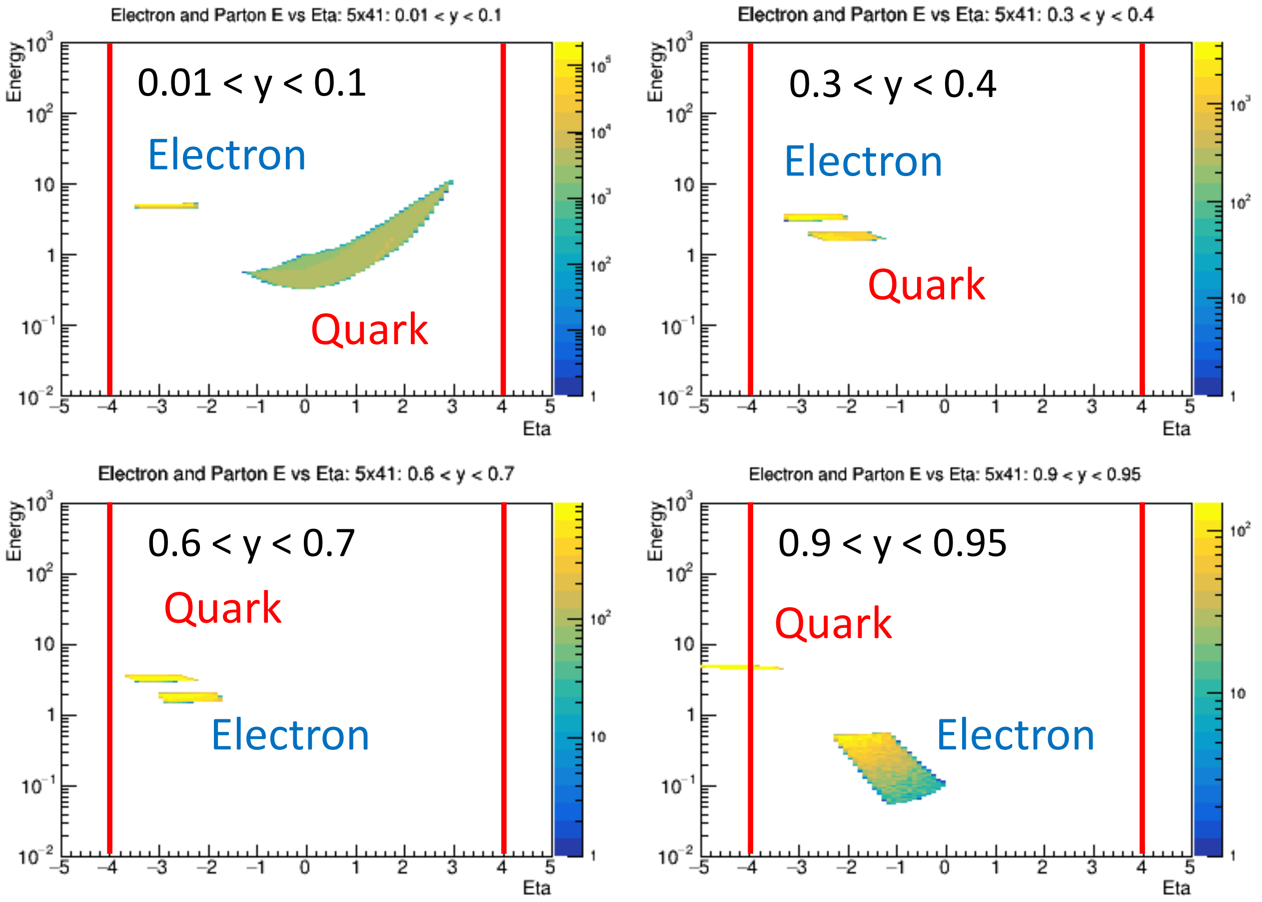


- ❑ For the leading order process, jet location and energy are dictated by the event kinematics (x, Q^2, y)
- ❑ For a given Q^2, inelasticity determines x value probed and pseudorapidity of the jet
  - ➢ Low y -> high x, jet at positive pseudorapidity
  - ➢ High y -> low x, jet at negative pseudorapidity

Electron and Struck Quark (5x41)

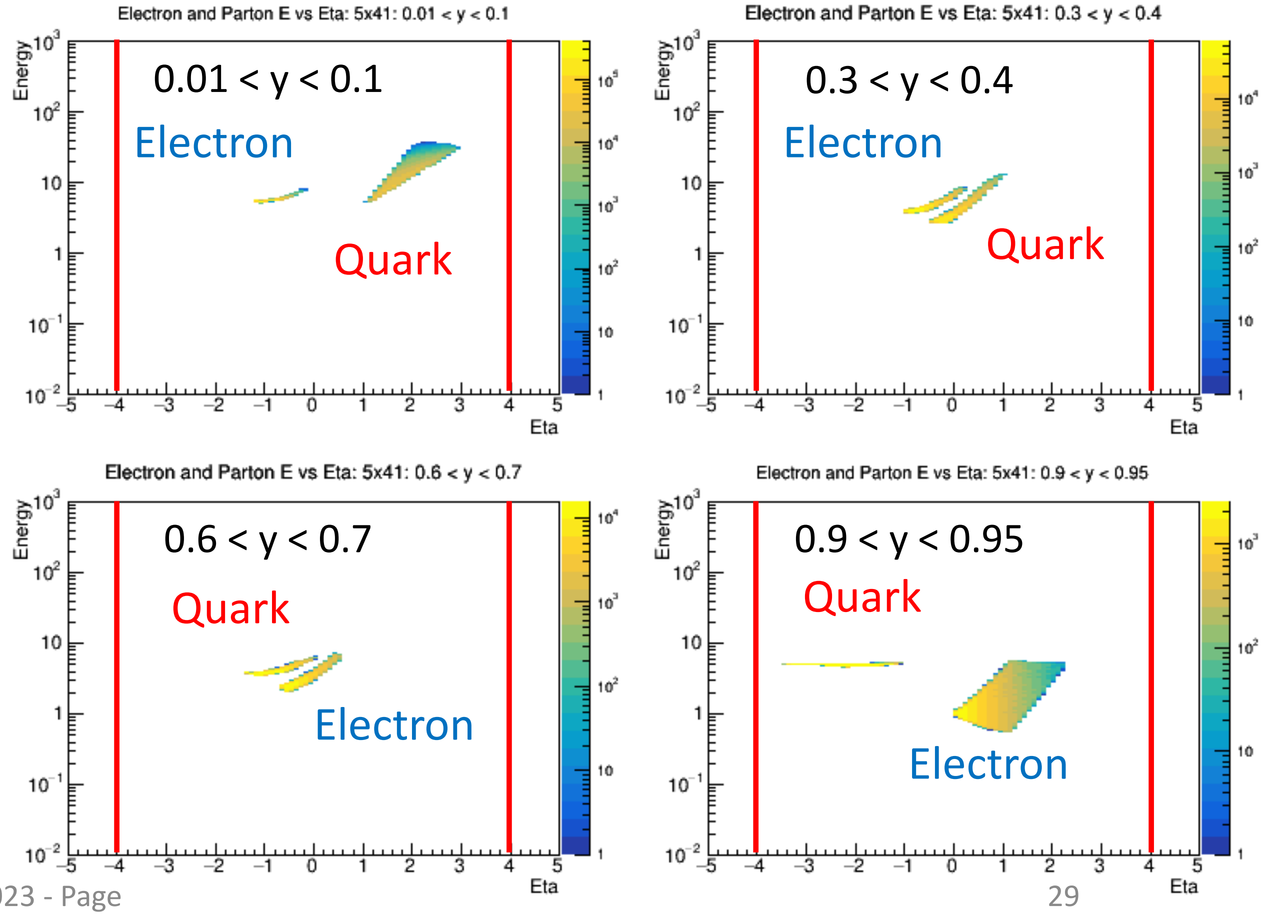
- Look at energy vs pseudorapidity of the scattered electron and struck quark as a function of y and Q^2

- For fixed Q^2, as y increases, electron eta increases while parton eta decreases



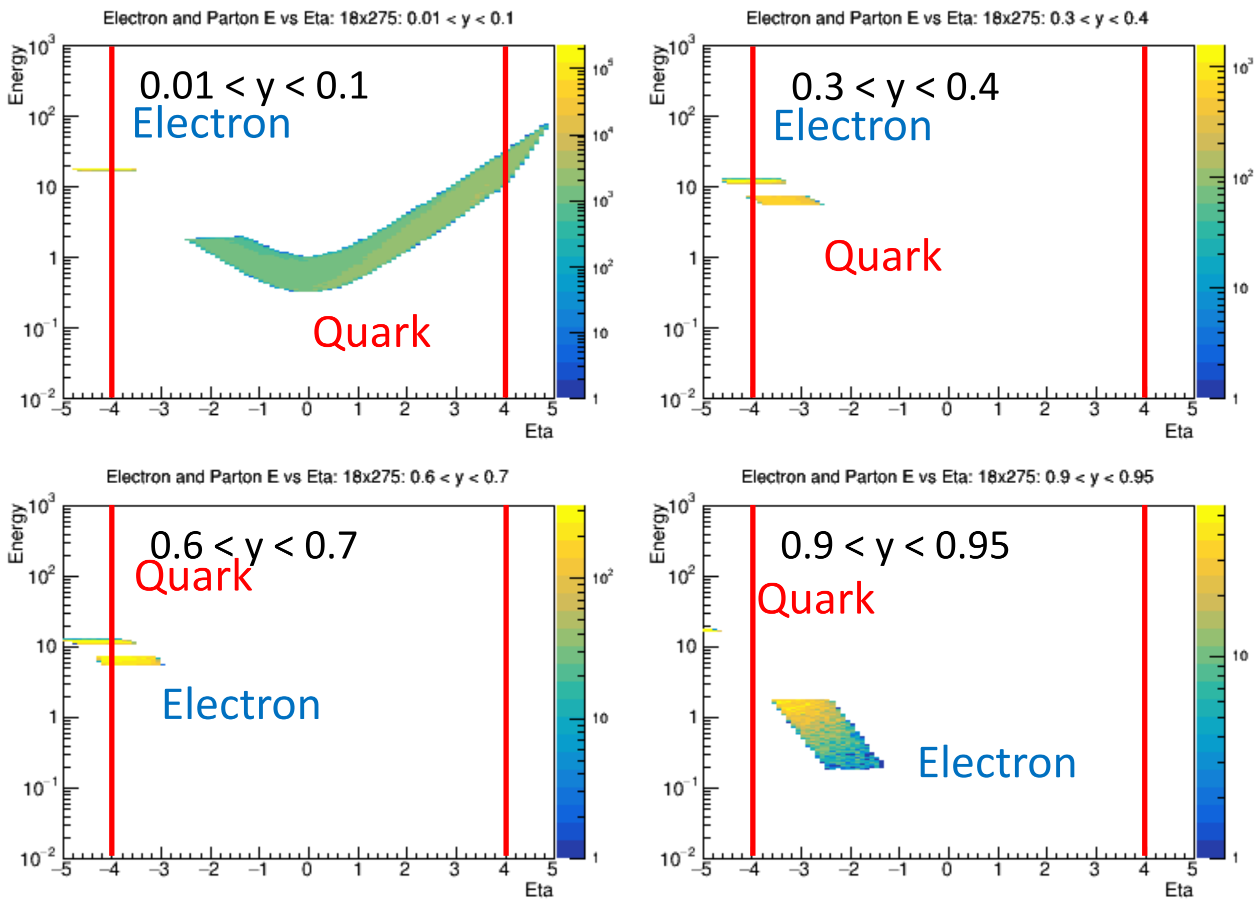
$0.1 < Q^2 < 1.0$

$10 < Q^2 < 100$



- As Q^2 increases, both the scattered electron and struck quark move to larger eta for all values of y

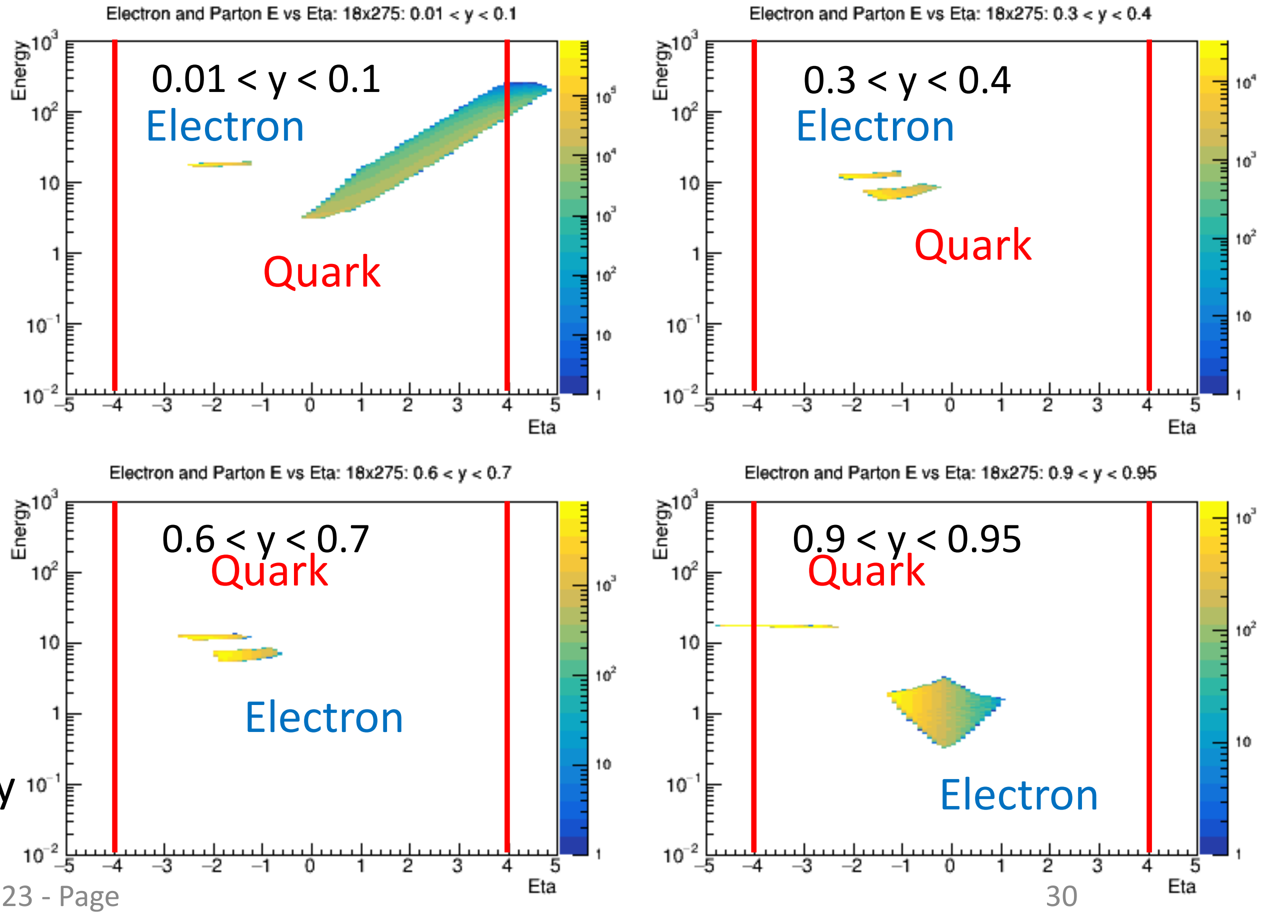
Electron and Struck Quark (18x275)



0.1 < Q² < 1.0

- ❑ At fixed Q^2 and y, larger beam energies push the electron and parton to lower pseudorapidity values

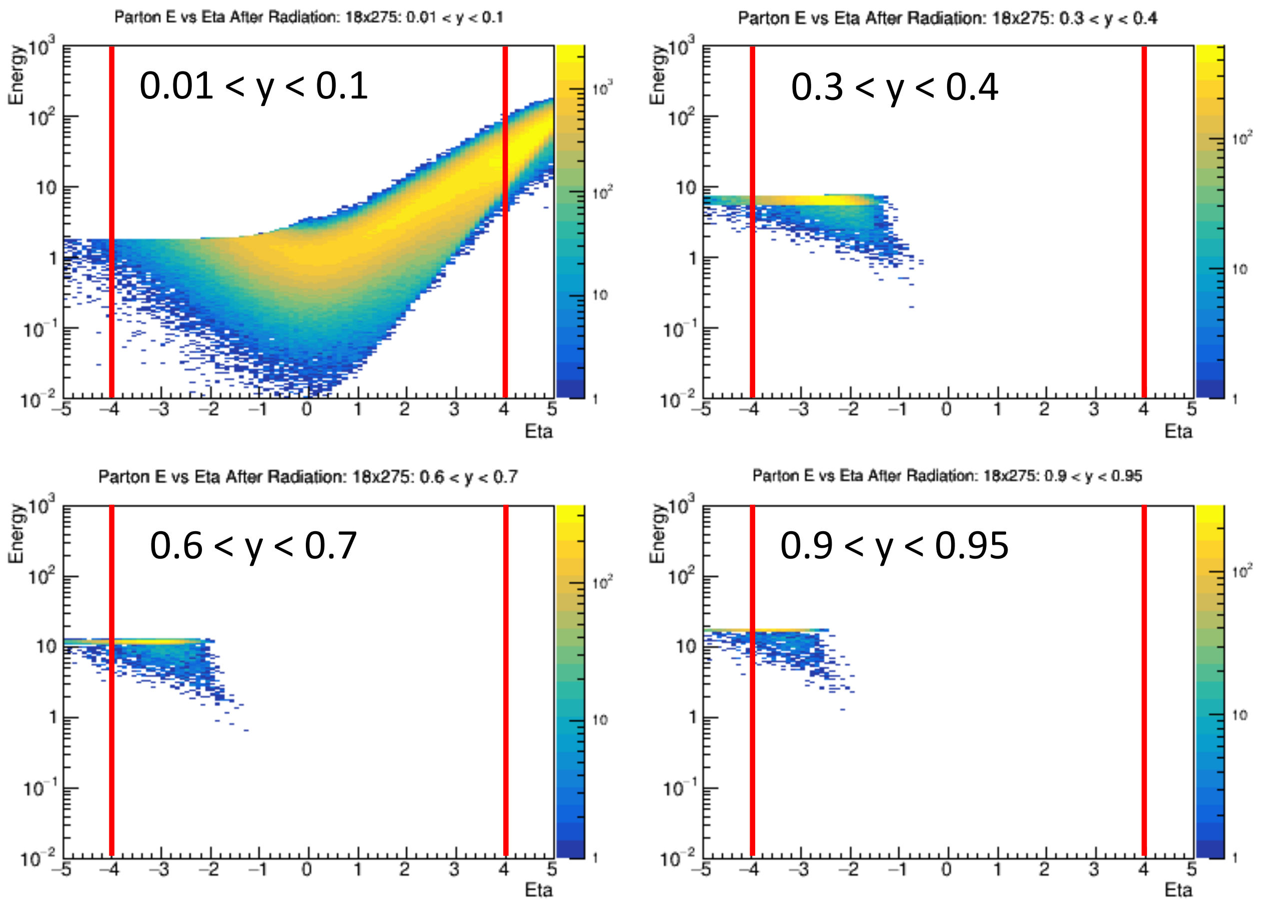
10 < Q² < 100



- ❑ As y -> 0, the struck quark can take the full ion beam energy
- ❑ As y -> 1, the struck quark takes the full electron beam energy
- ❑ Different detector considerations in forward and backward regions

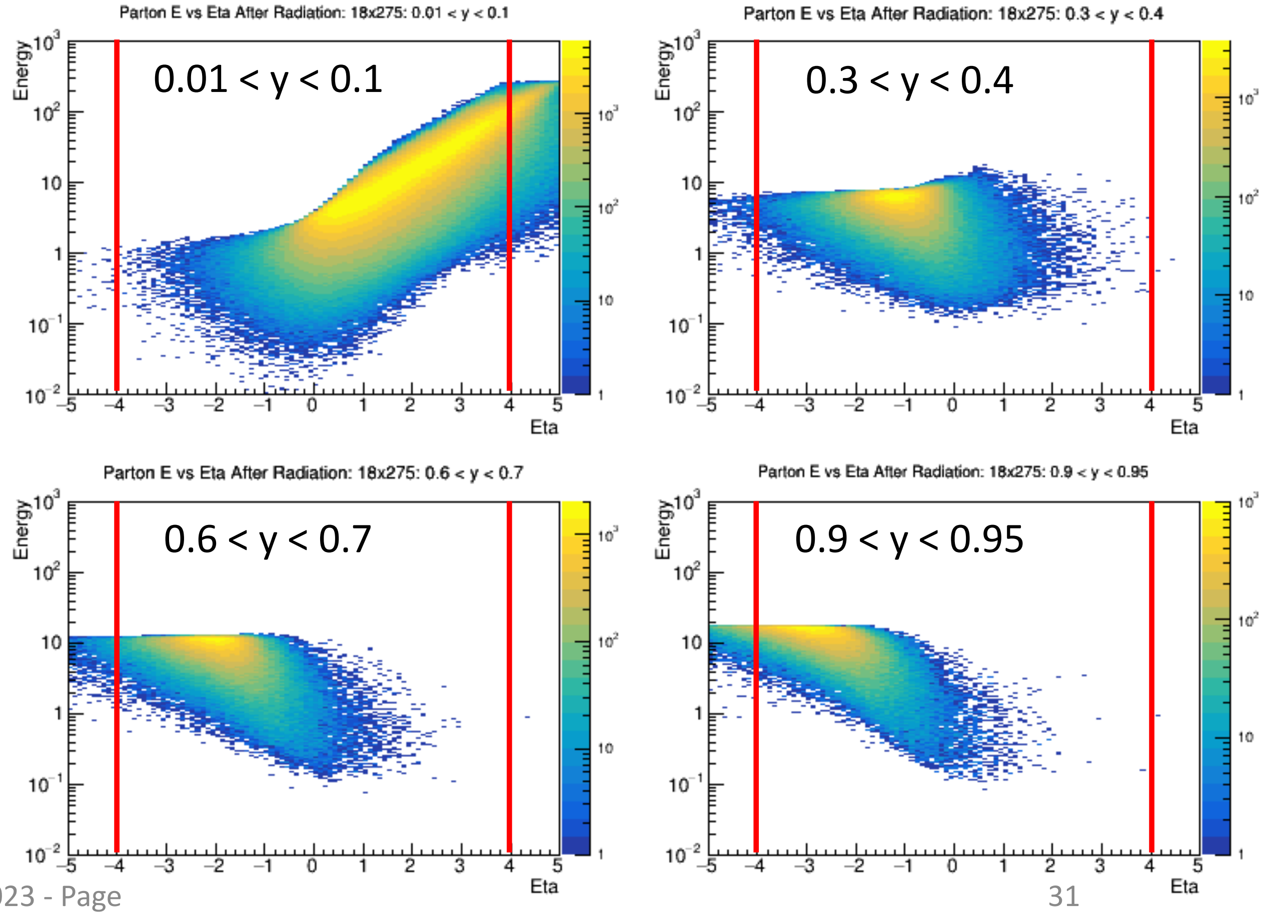
Struck Quark + FSR (18x275)

- For Born-level process, struck quark kinematics are correlated with event kinematics
- Final state radiation can alter quark kinematics significantly



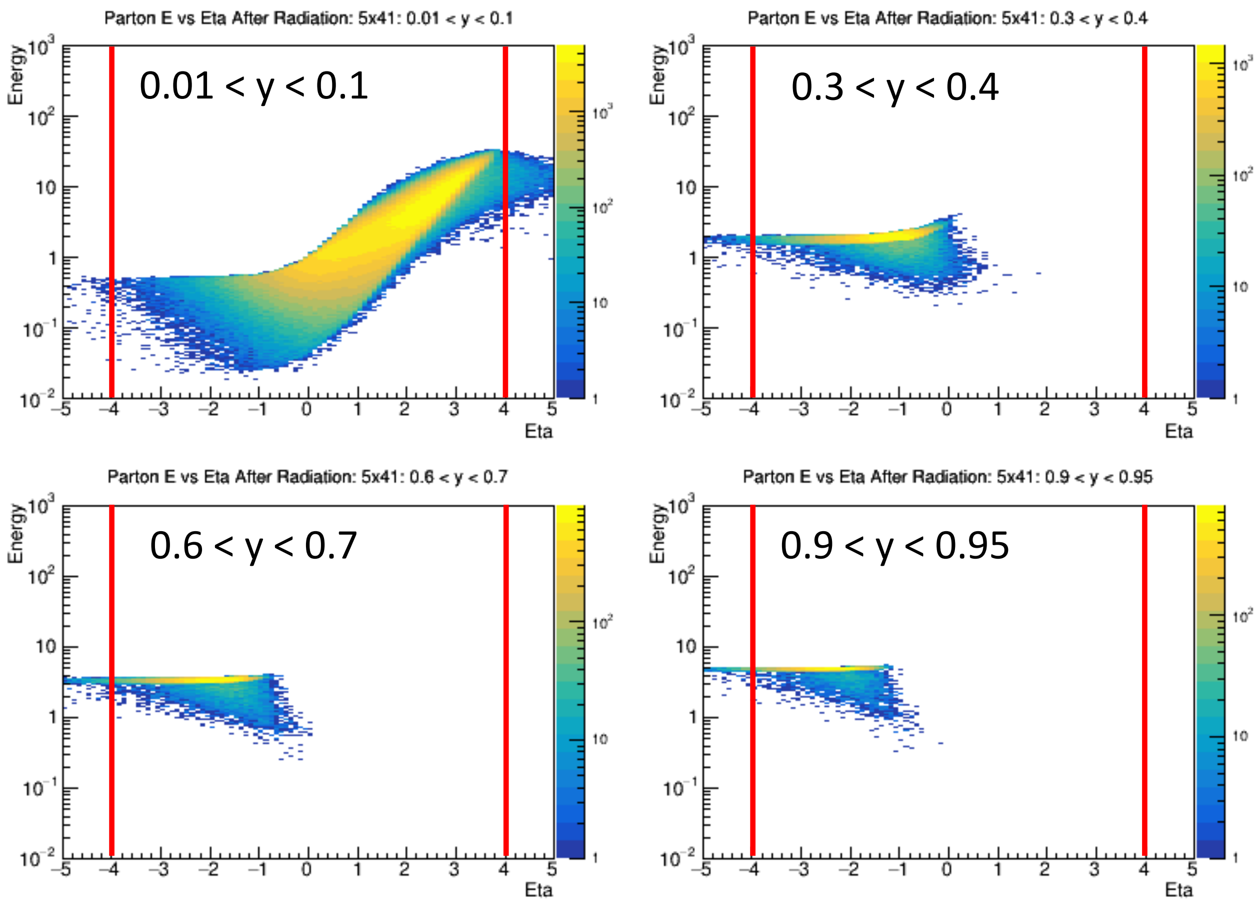
$0.1 < Q^2 < 1.0$

$10 < Q^2 < 100$



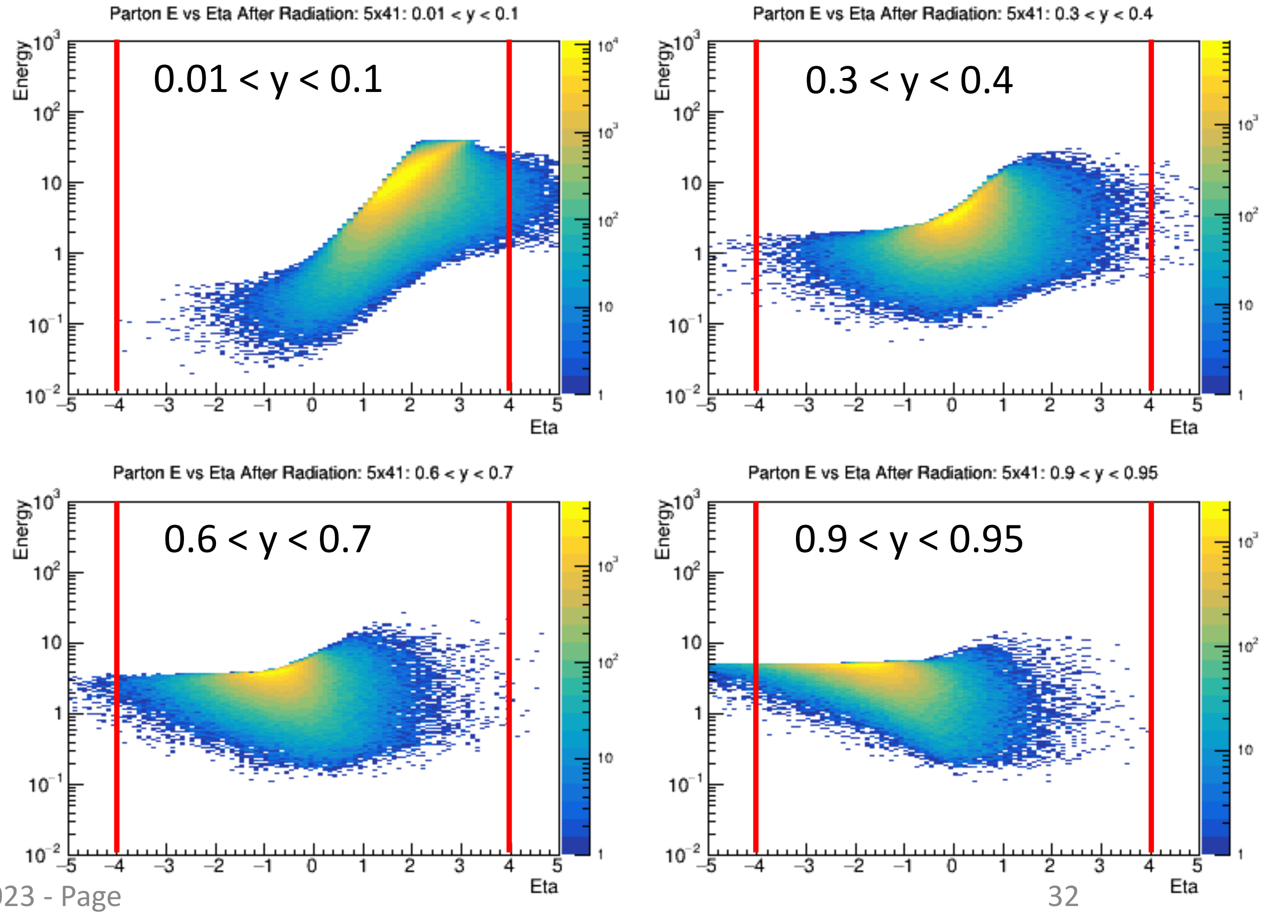
Struck Quark + FSR (5x41)

- Behavior as a function of Q^2, y, and beam energy is still seen

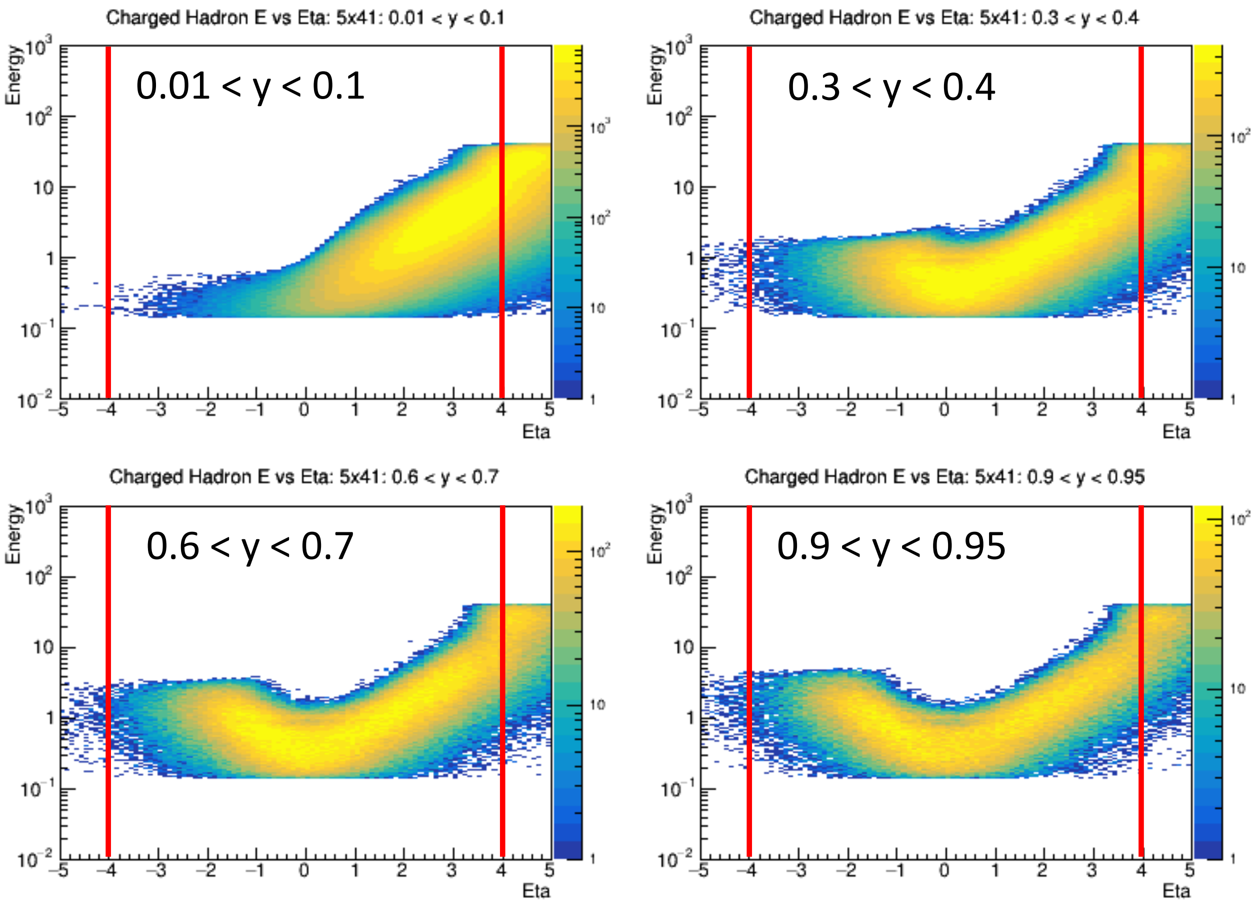


$0.1 < Q^2 < 1.0$

$10 < Q^2 < 100$



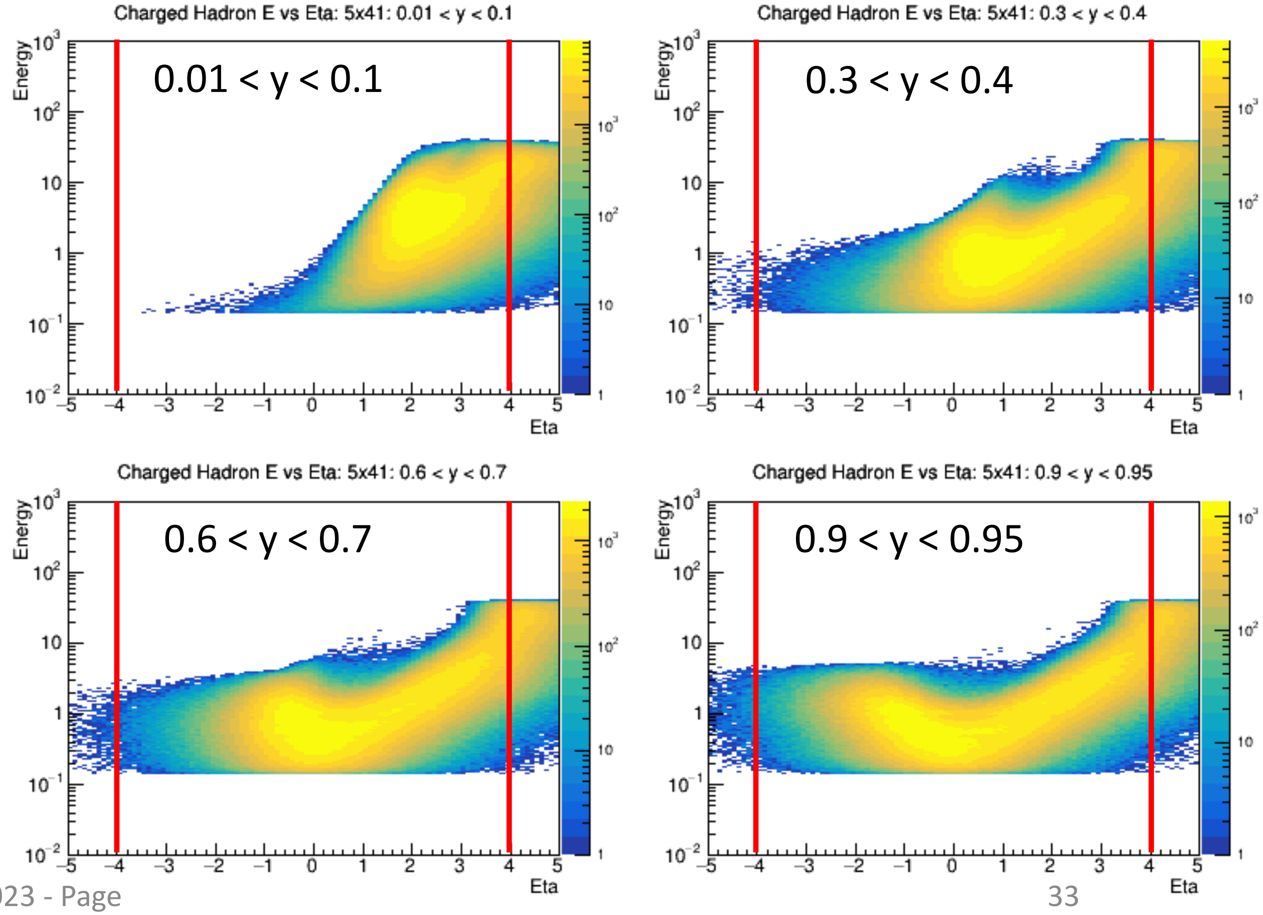
(Charged) Particle Distributions (5x41)



$0.1 < Q^2 < 1.0$

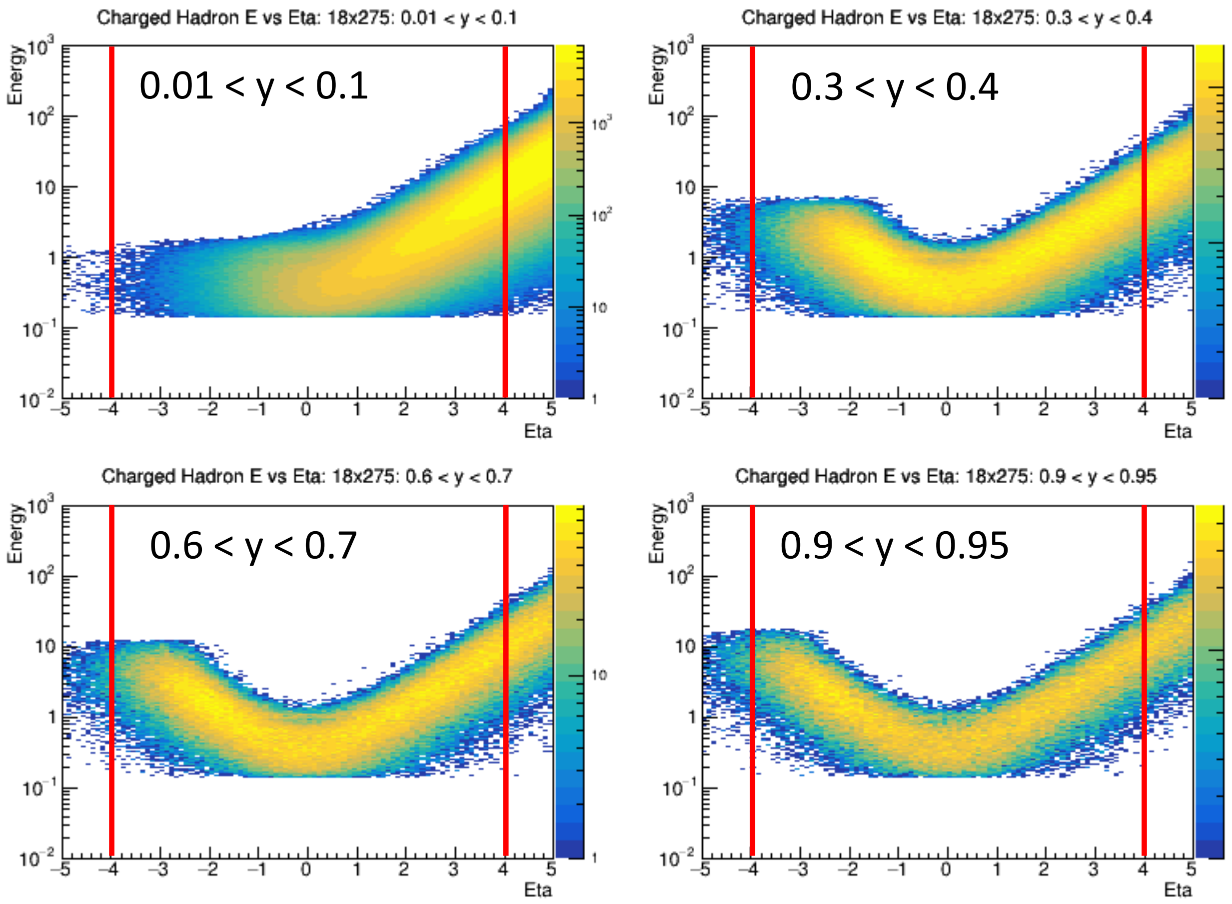
- ❑ Of course, it is final state hadrons which are measured
- ❑ Differences with y and Q2 are now somewhat less pronounced

$10 < Q^2 < 100$



- ❑ Particle production not associated only with struck parton
- ❑ Gammas and neutrals follow same pattern

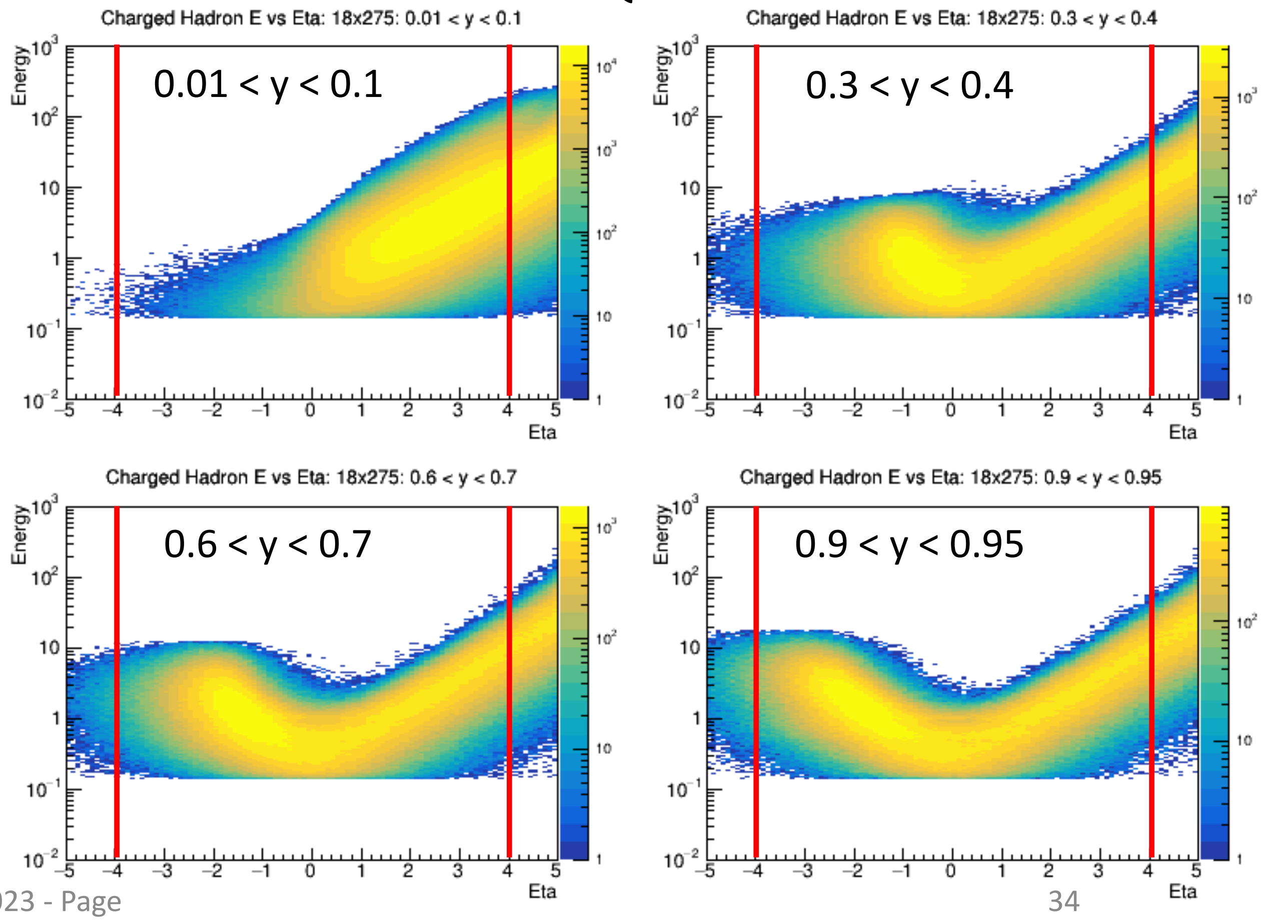
(Charged) Particle Distributions (18x275)



0.1 < Q² < 1.0

- ❑ Higher energy extends energy range somewhat, but basic distributions are similar

10 < Q² < 100

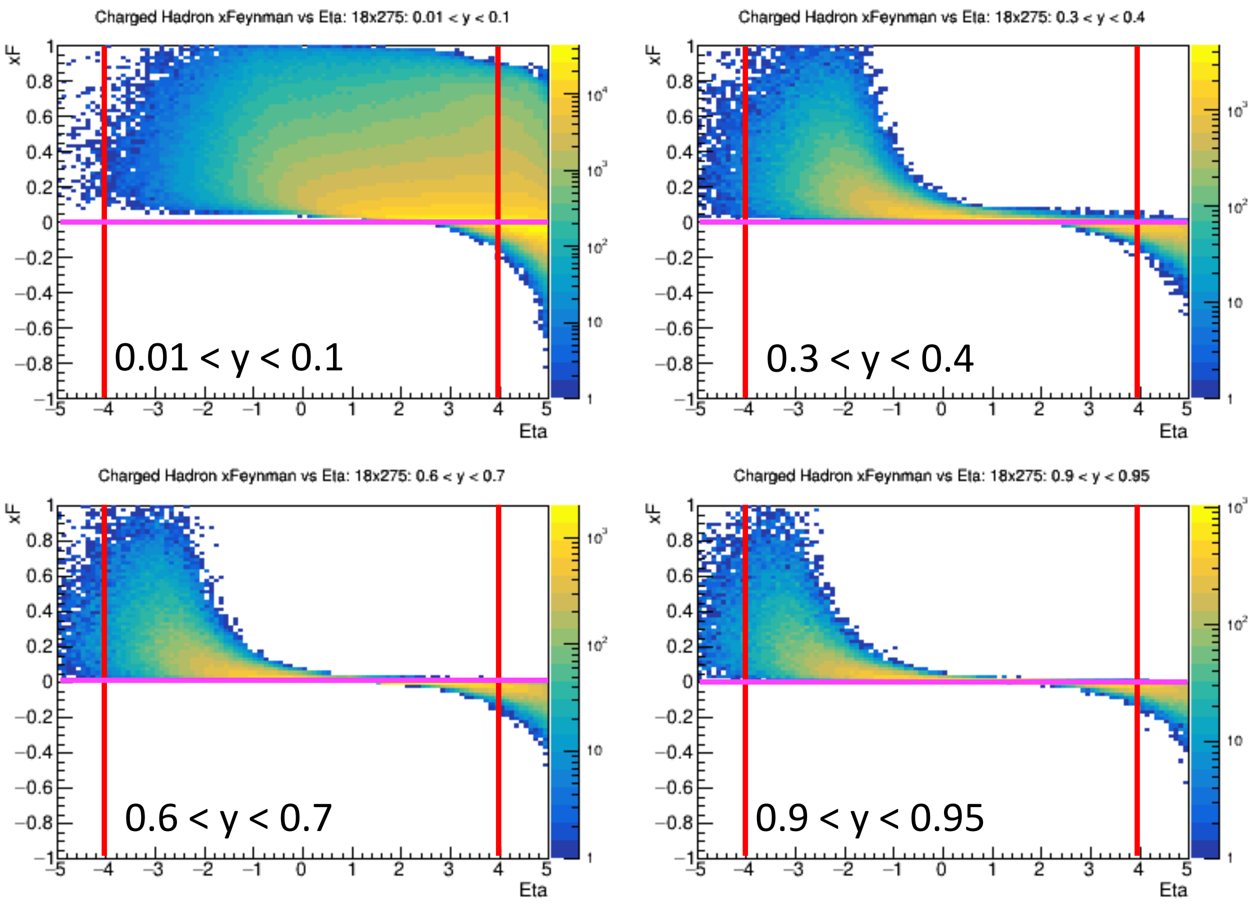


- ❑ How well can we reconstruct the parton kinematics from these particles?
- ❑ Can we form jets away from the struck parton?

(Charged) Particle x-Feynman (18x275)

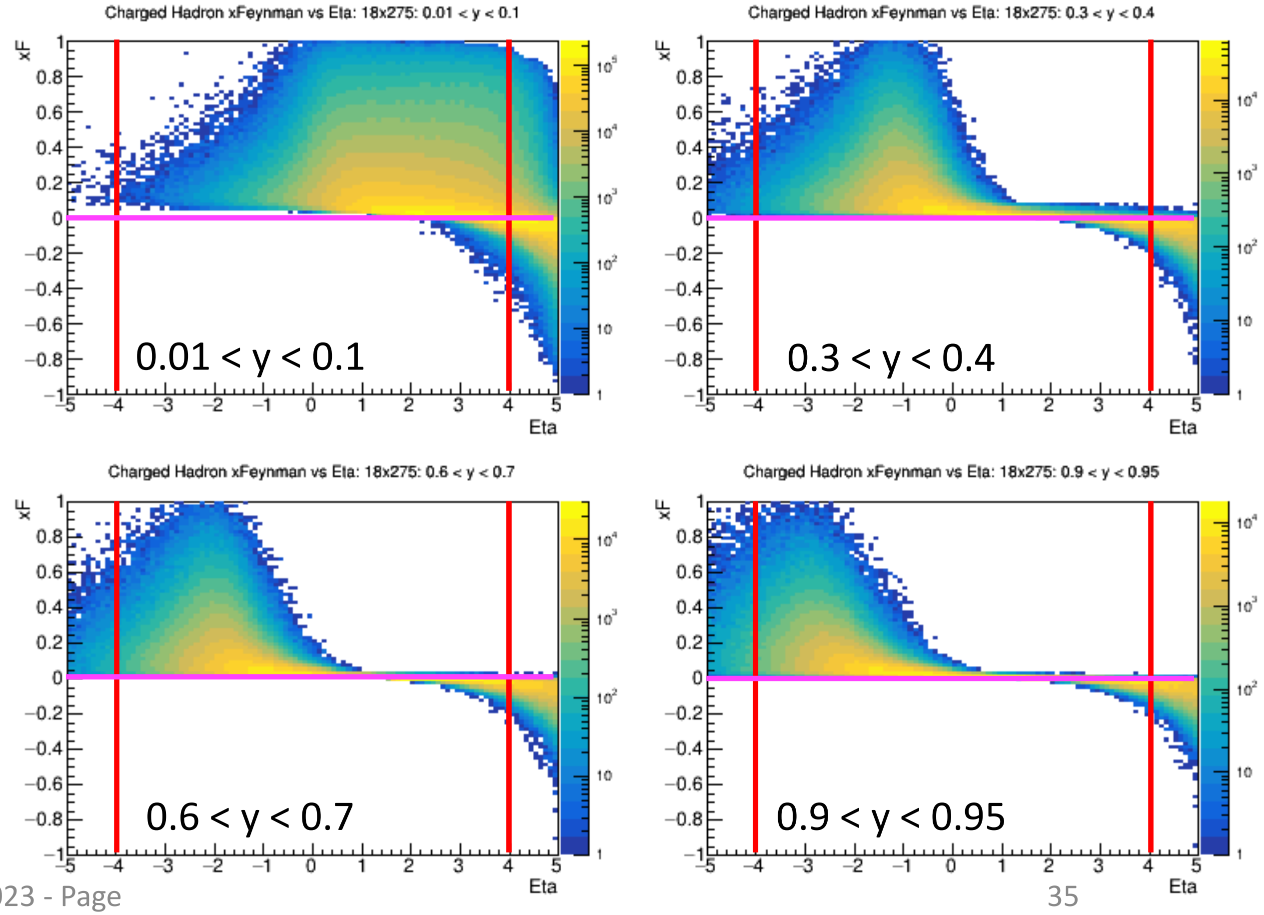
- ❑ Define xF as 2*particle_pZ/W in the hadron-boson center of mass frame

- ❑ Z-axis defined w.r.t. the virtual photon



$0.1 < Q^2 < 1.0$

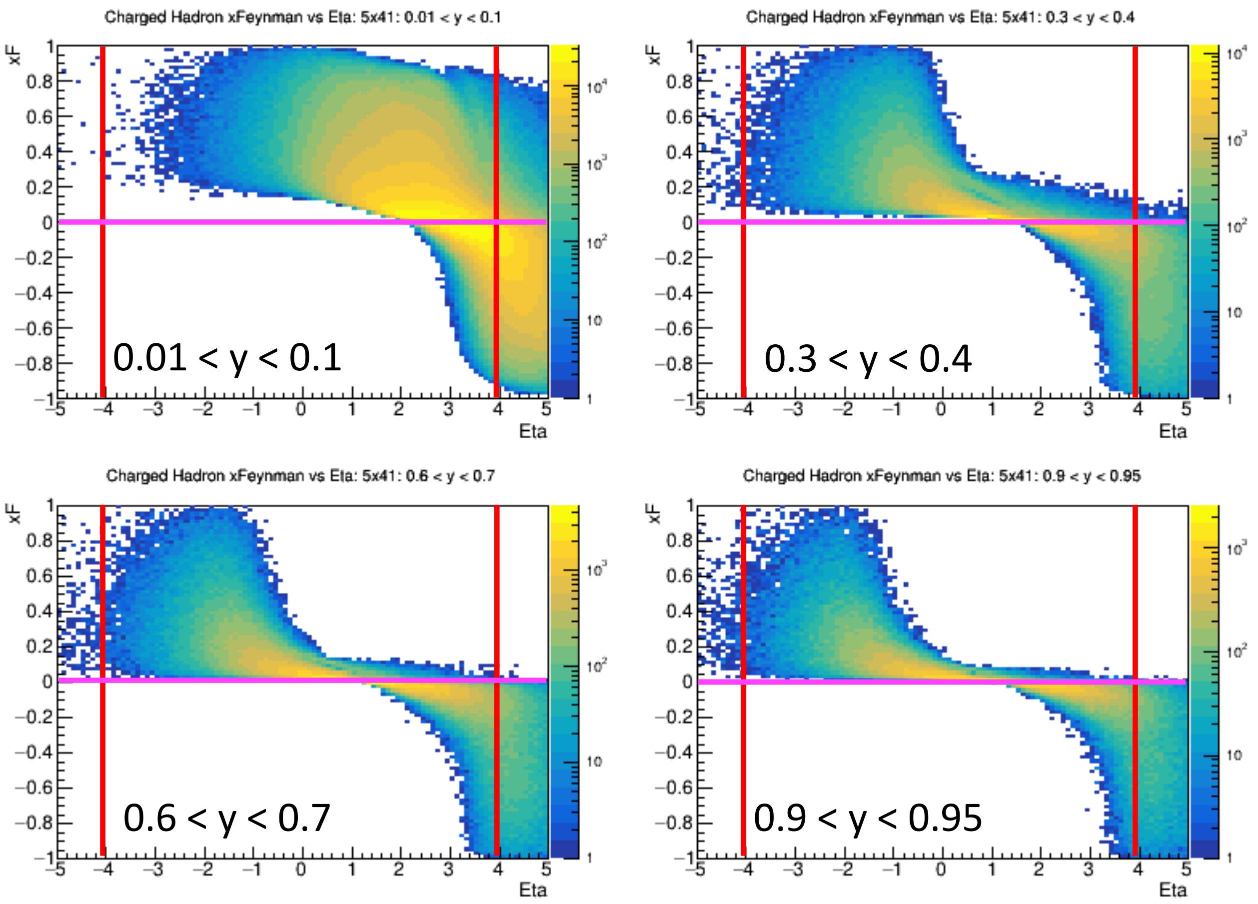
$10 < Q^2 < 100$



- ❑ Positive xF indicates particles more associated with the struck quark

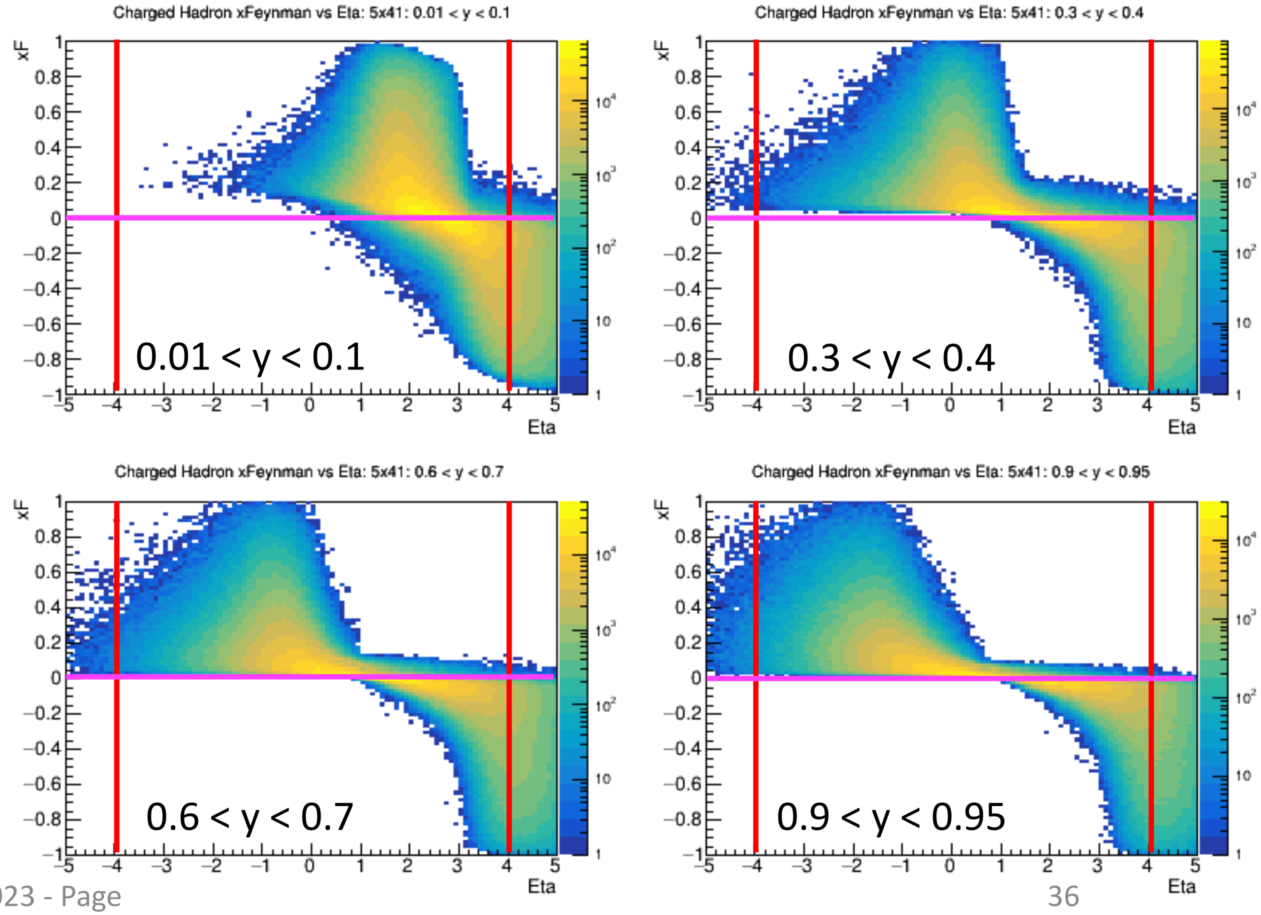
(Charged) Particle x-Feynman (5x41)

- ❑ Larger negative xF particles come into the detector acceptance for lower beam energies



$0.1 < Q^2 < 1.0$

$10 < Q^2 < 100$



- ❑ Study differences in fragmentation in these regions
- ❑ Look at correlations between these regions

Jet Algorithms

Anti_k_T

$$d_{ij} = \min[p_{ti}^{-2}, p_{tj}^{-2}]\Delta R_{ij}/R$$

EE_kT (Spherically Invariant)

$$d_{ij} = 2 * \min[E_i^2, E_j^2](1 - \cos\Delta_{ij})$$

Centauro

$$d_{ij} = \left[(\Delta f_{ij})^2 + 2 f_i f_j (1 - \cos\Delta\phi_{ij})\right]/R^2$$

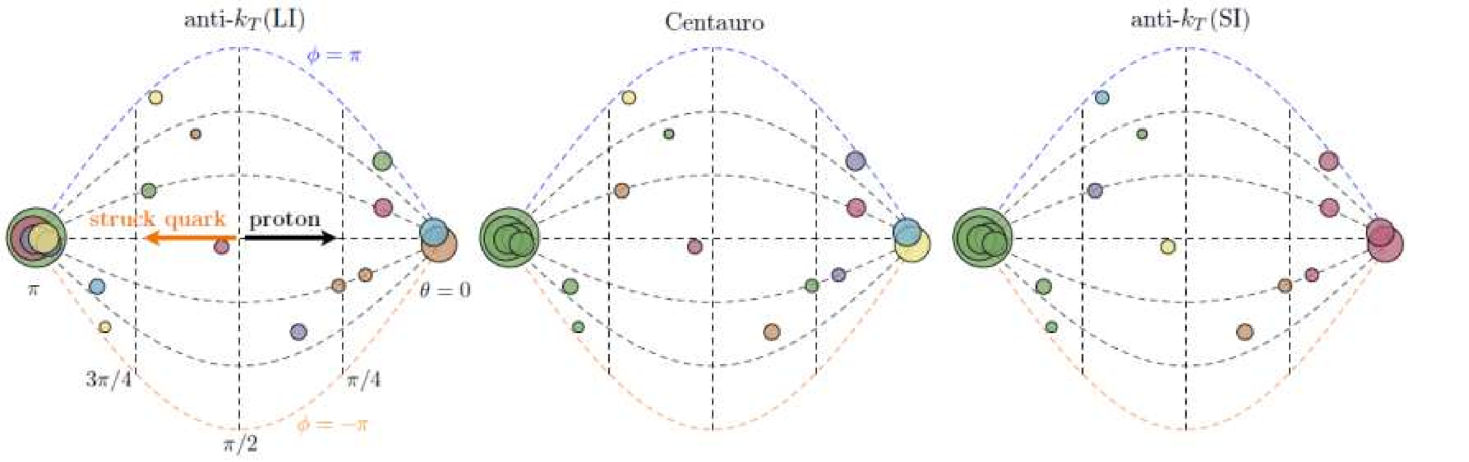
Asymmetric measure is necessary

$$f(x) = x + \mathcal{O}(x^2)$$

$$\bar{\eta}_i = -\frac{2Q}{\bar{n}\cdot q}\frac{p_i^{\perp}}{n\cdot p_i}$$

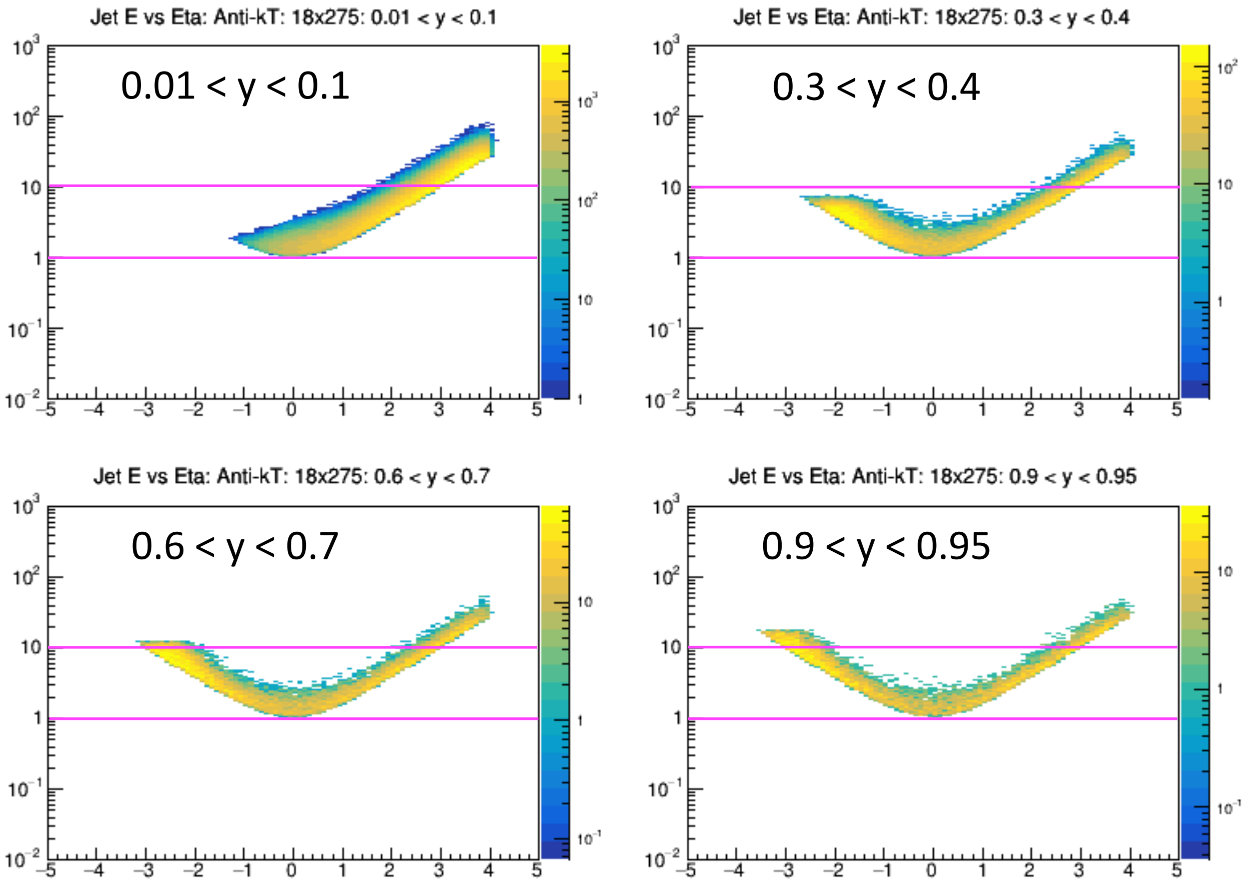
$$\bar{\eta}_i(\mathrm{BF}) = 2p_i^{\perp}/p_i^{+}$$

- Sequential recombination algorithms, especially Anti_k_T, have been the "industry standard" at hadron colliders for a number of years
- Is this appropriate for very forward jets or Born-level jets in the Breit frame where transverse momenta are by definition small?
- Look at alternative distance measures such as spherically invariant and symmetric EE_k_T or longitudinally invariant and anti-symmetric centauro algorithms



Longitudinally invariant Symmetric | Longitudinally invariant Asymmetric | Spherically invariant Symmetric

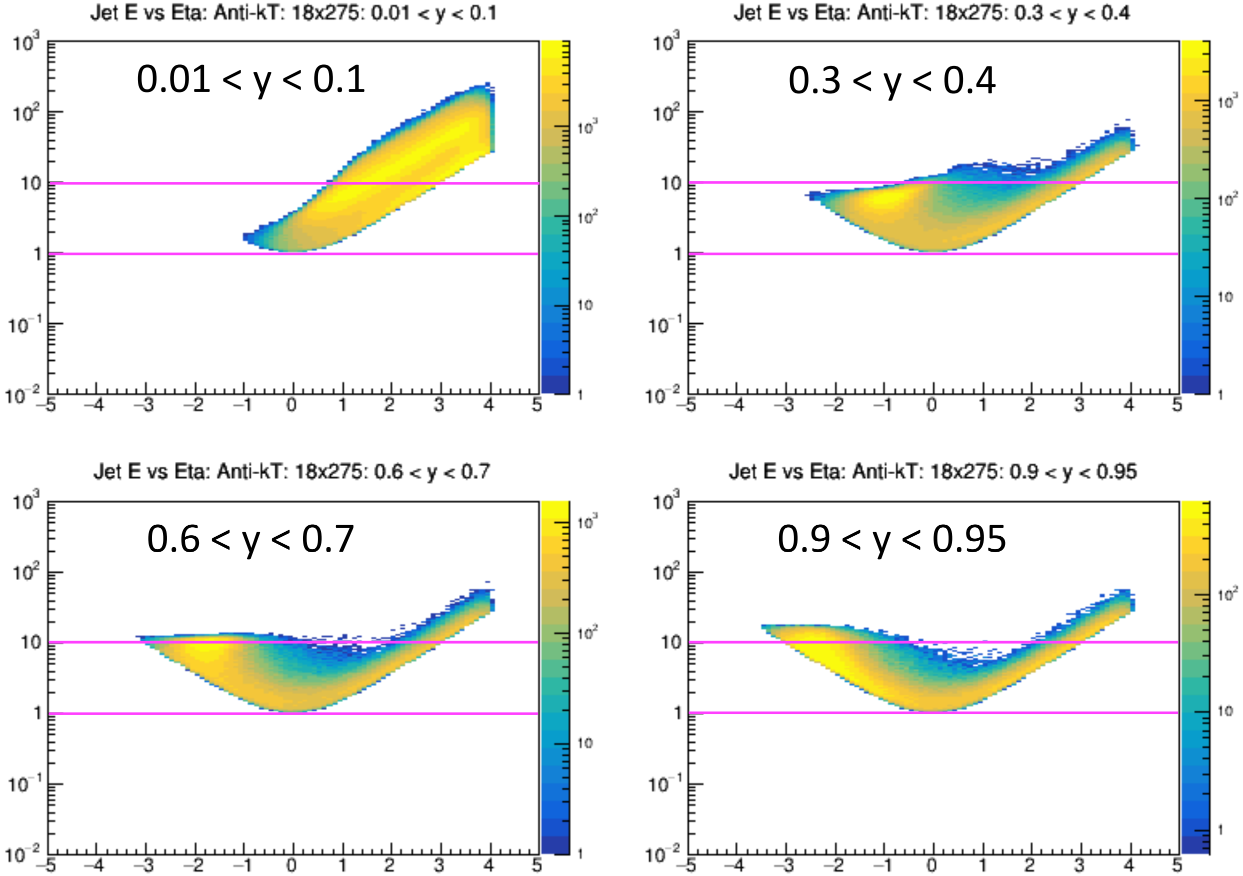
Jet Distributions: Anti_k_T (18x275)



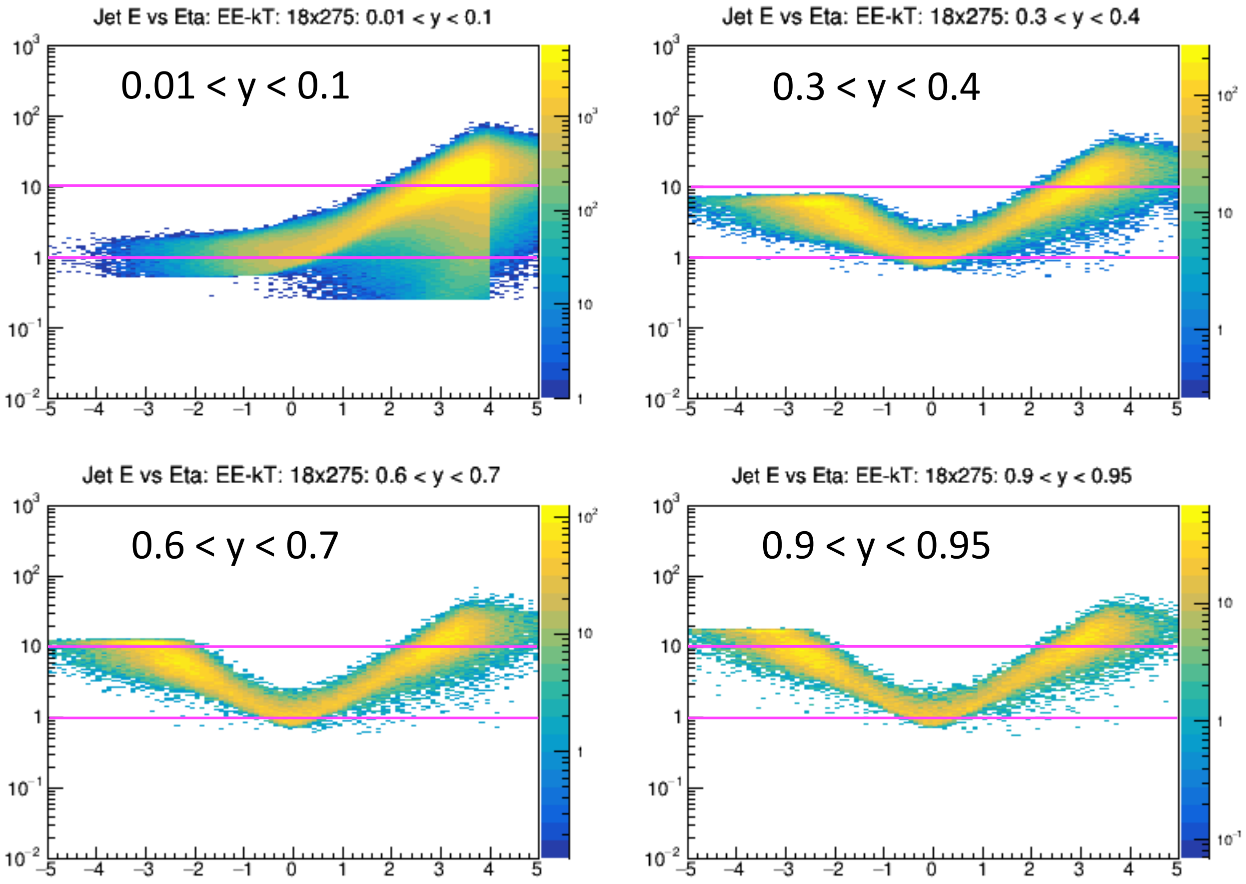
$0.1 < Q^2 < 1.0$

- Run inclusive Anti_k_T on all stable particles (|eta| < 4) with 1 GeV minimum p_T cut
- Jets roughly follow particle distributions

$10 < Q^2 < 100$



Jet Distributions: EE_k_T (18x275)



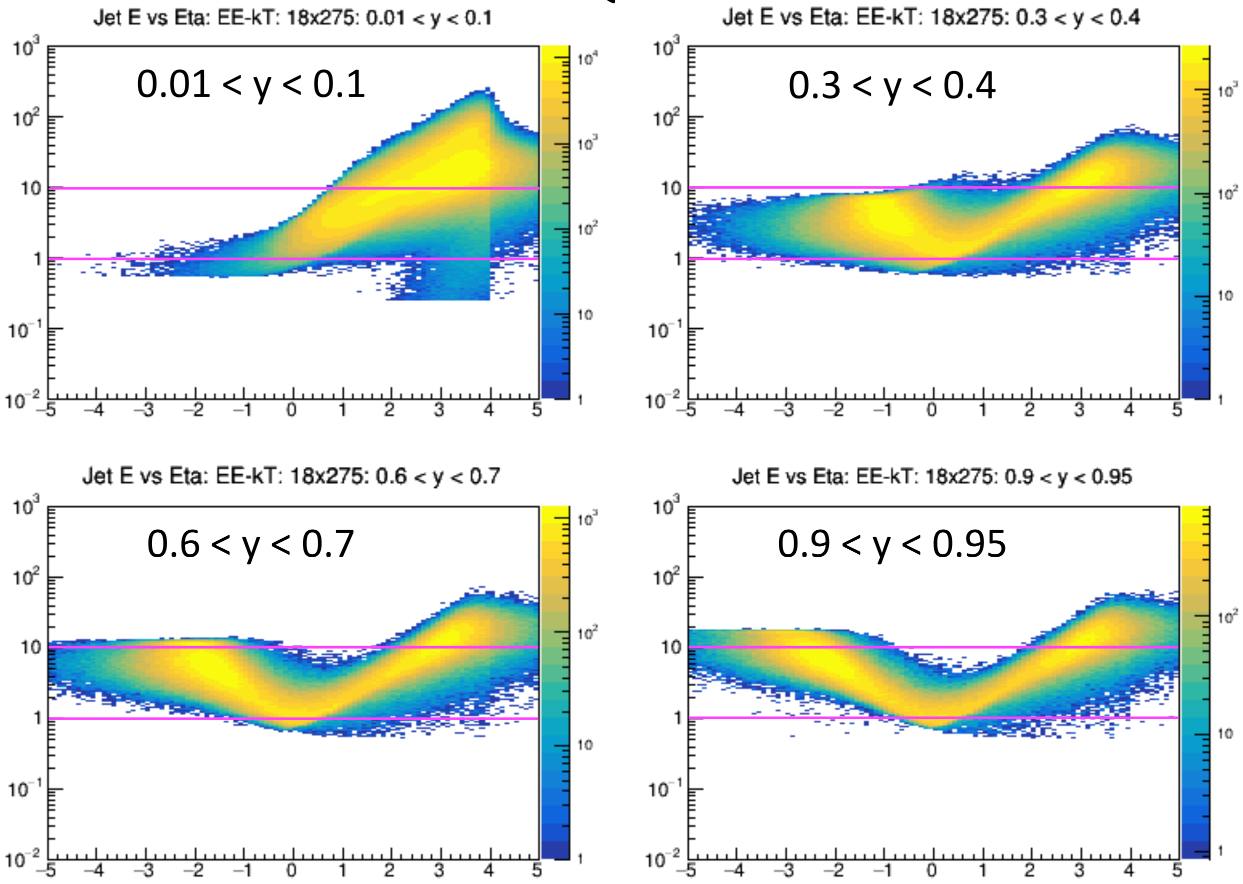
0.1 < Q^2 < 1.0

- ❑ Need to understand the artifact around eta = 4, must be related to particle eta cut

- ❑ Overall distributions are similar for EE_k_T algorithm

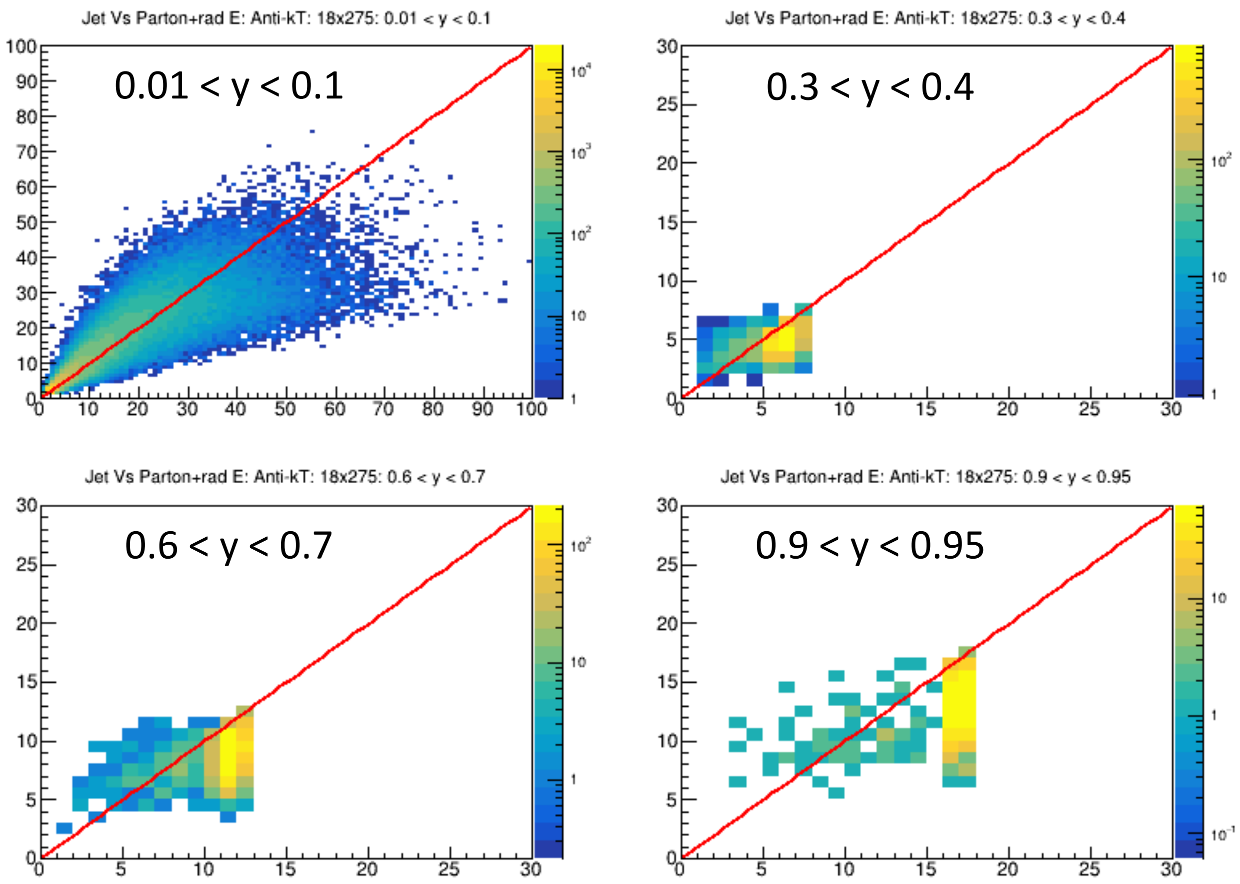
- ❑ In general, see larger number of jets, more jets at higher eta, and more jets away from struck quark

10 < Q^2 < 100



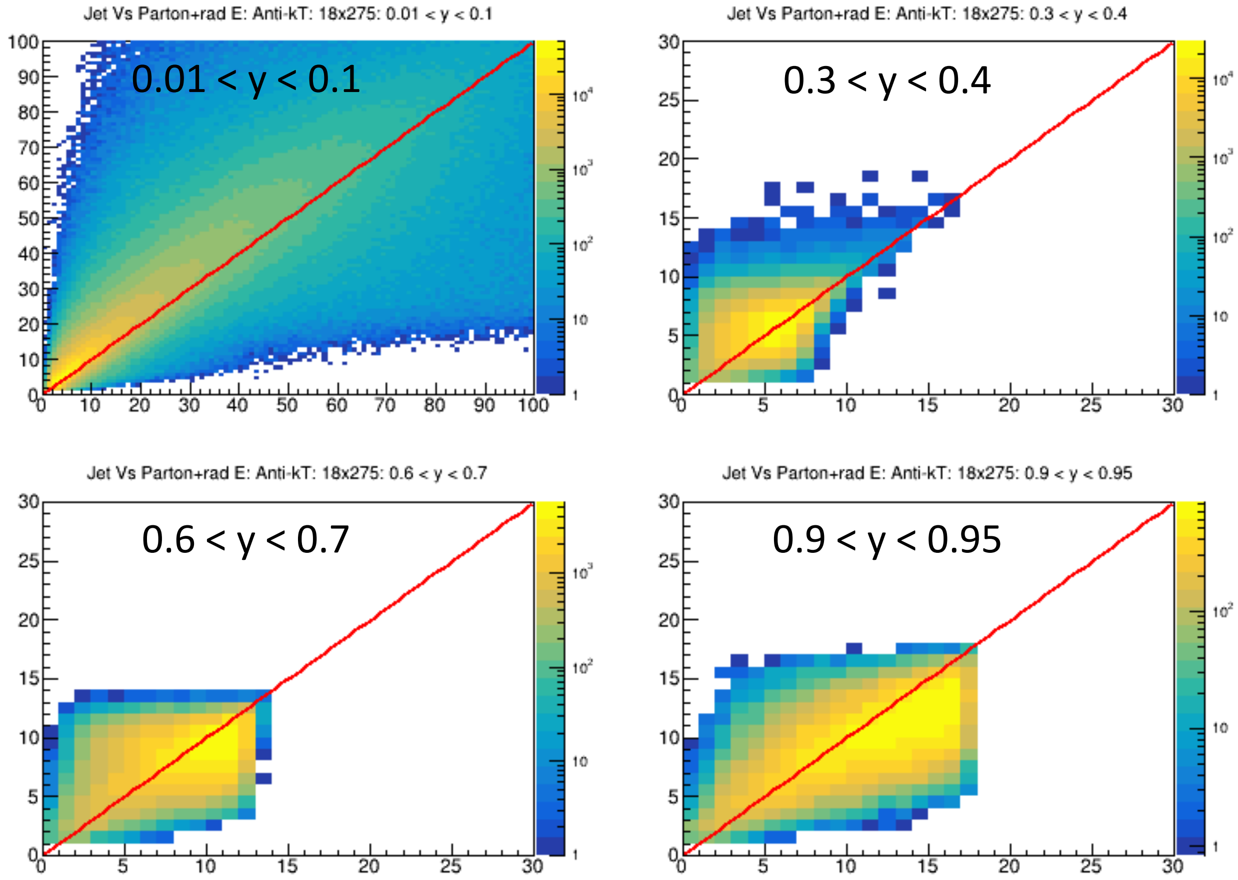
Jet – Parton Energy Comparison: Anti_k_T (18x275)

- ❑ How well do jets represent the parton?
- ❑ Plot jet energy vs parton+FSR energy for different Q^2 and inelasticity



0.1 < Q^2 < 1.0

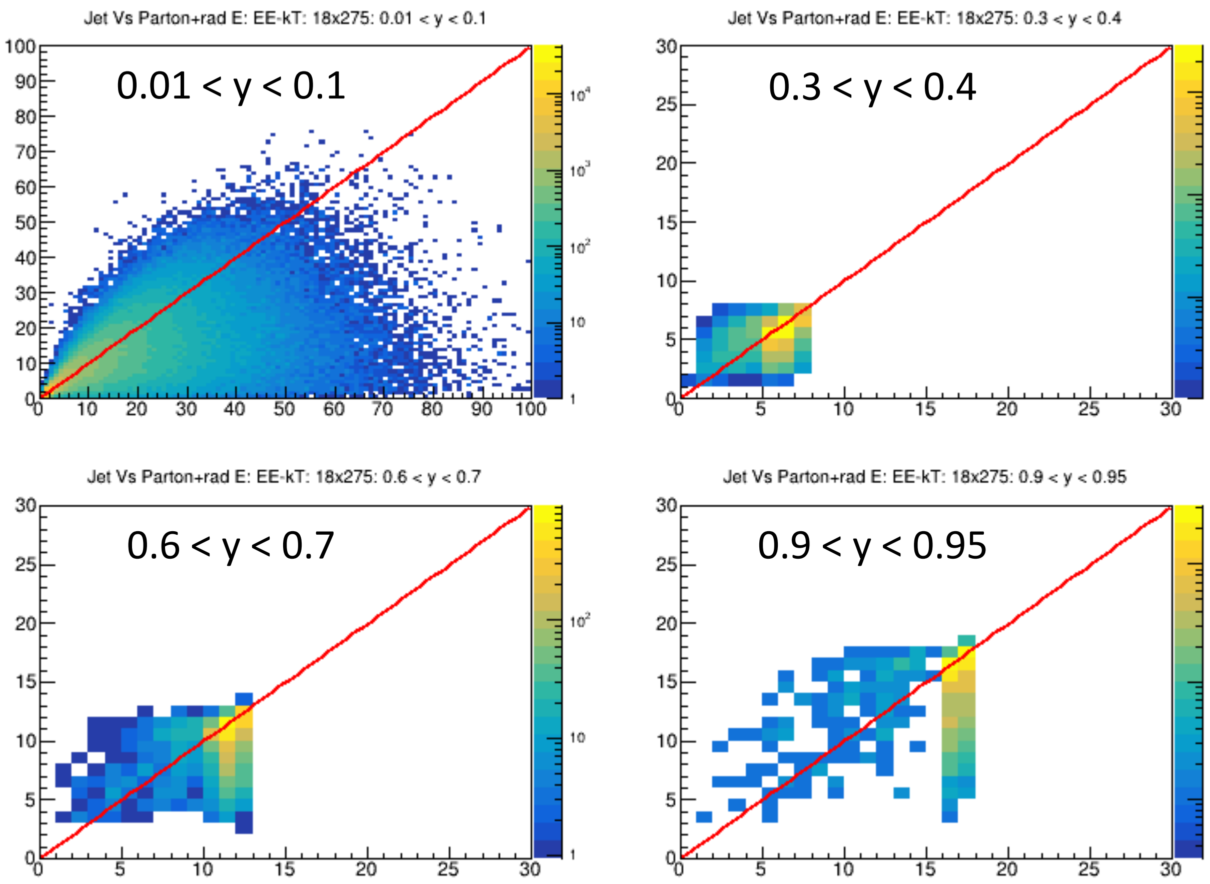
10 < Q^2 < 100



- ❑ Performance degrades somewhat at larger y (backward jets)

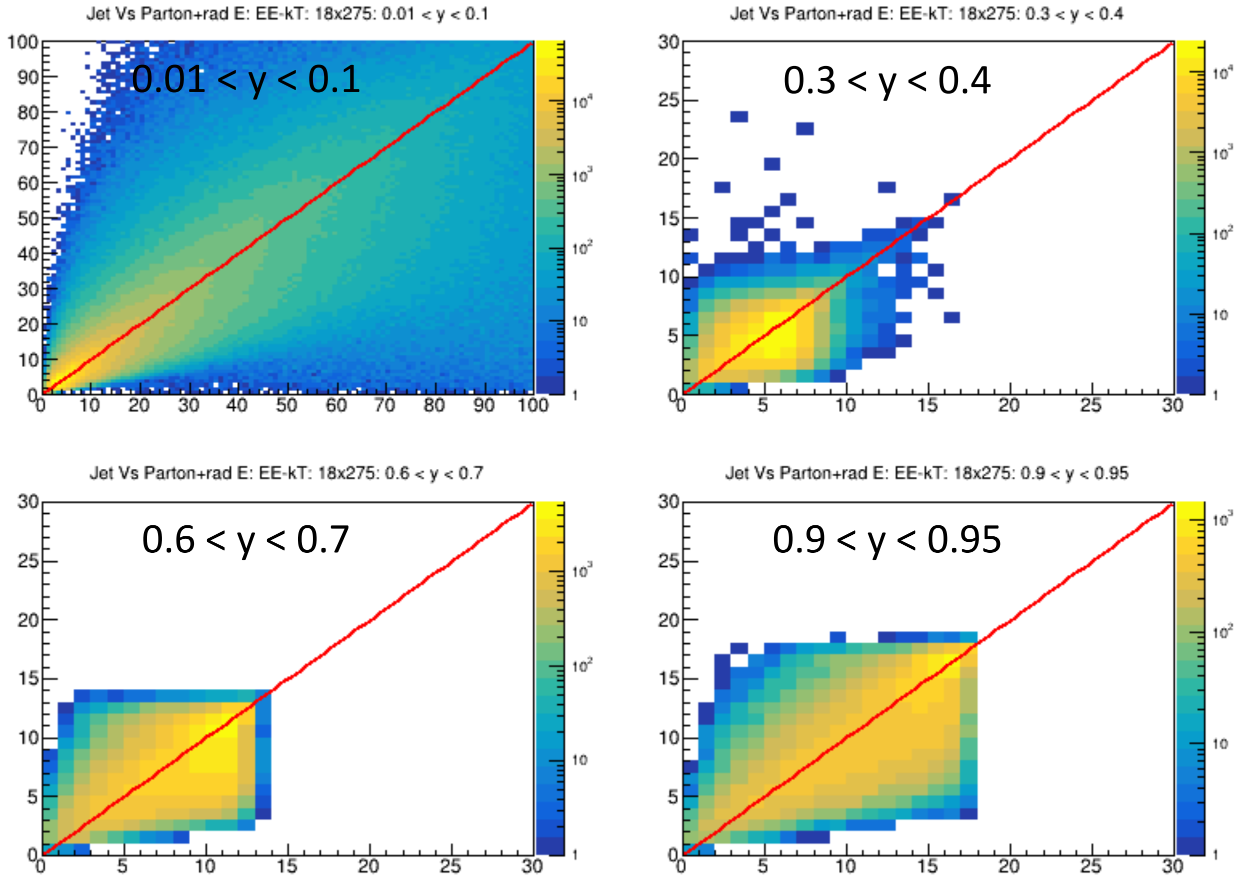
Jet – Parton Energy Comparison: EE_k_T (18x275)

- Better agreement between parton and jet seen with EE_kT algorithm, especially at high y

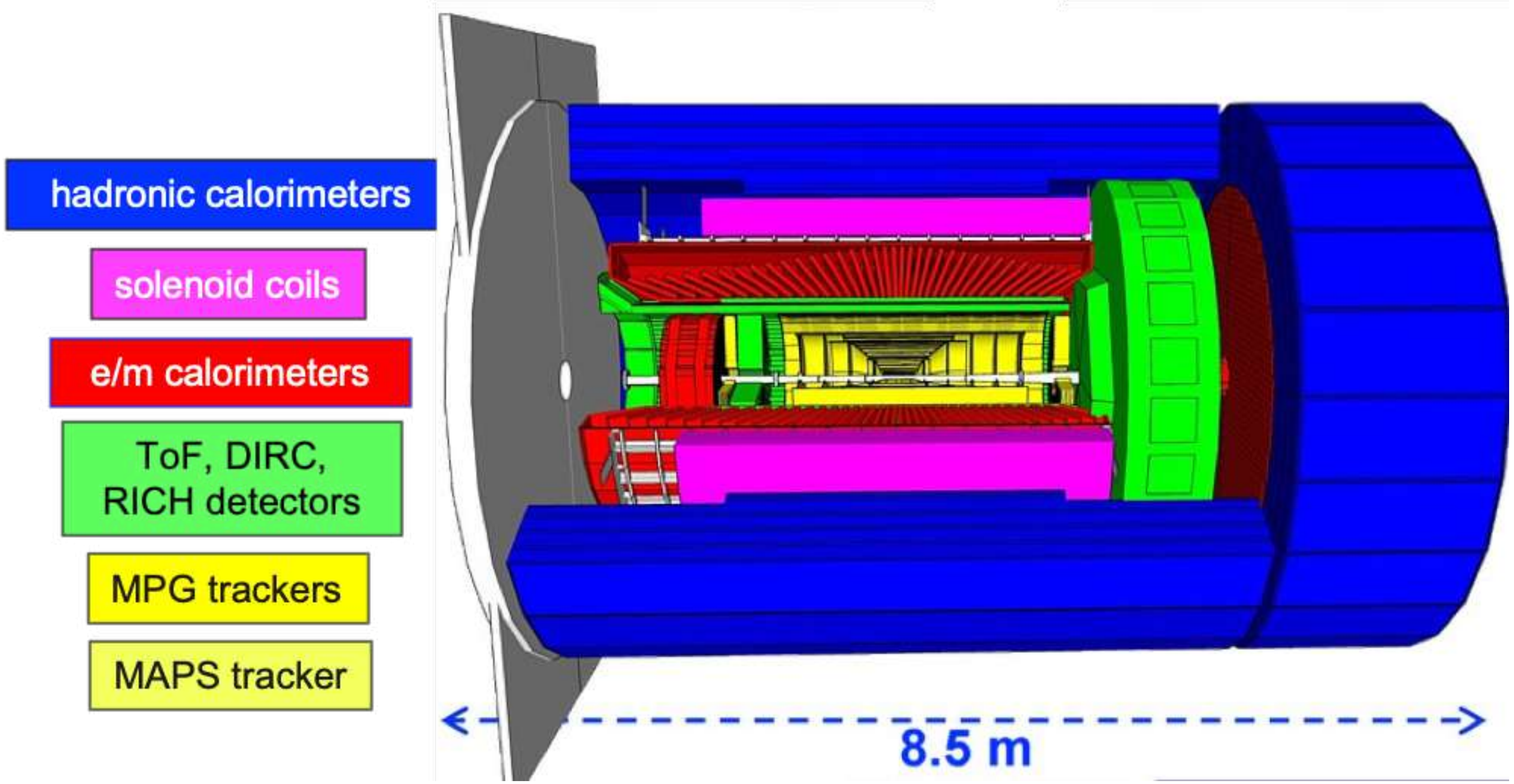


0.1 < Q^2 < 1.0

10 < Q^2 < 100



Detector Configuration



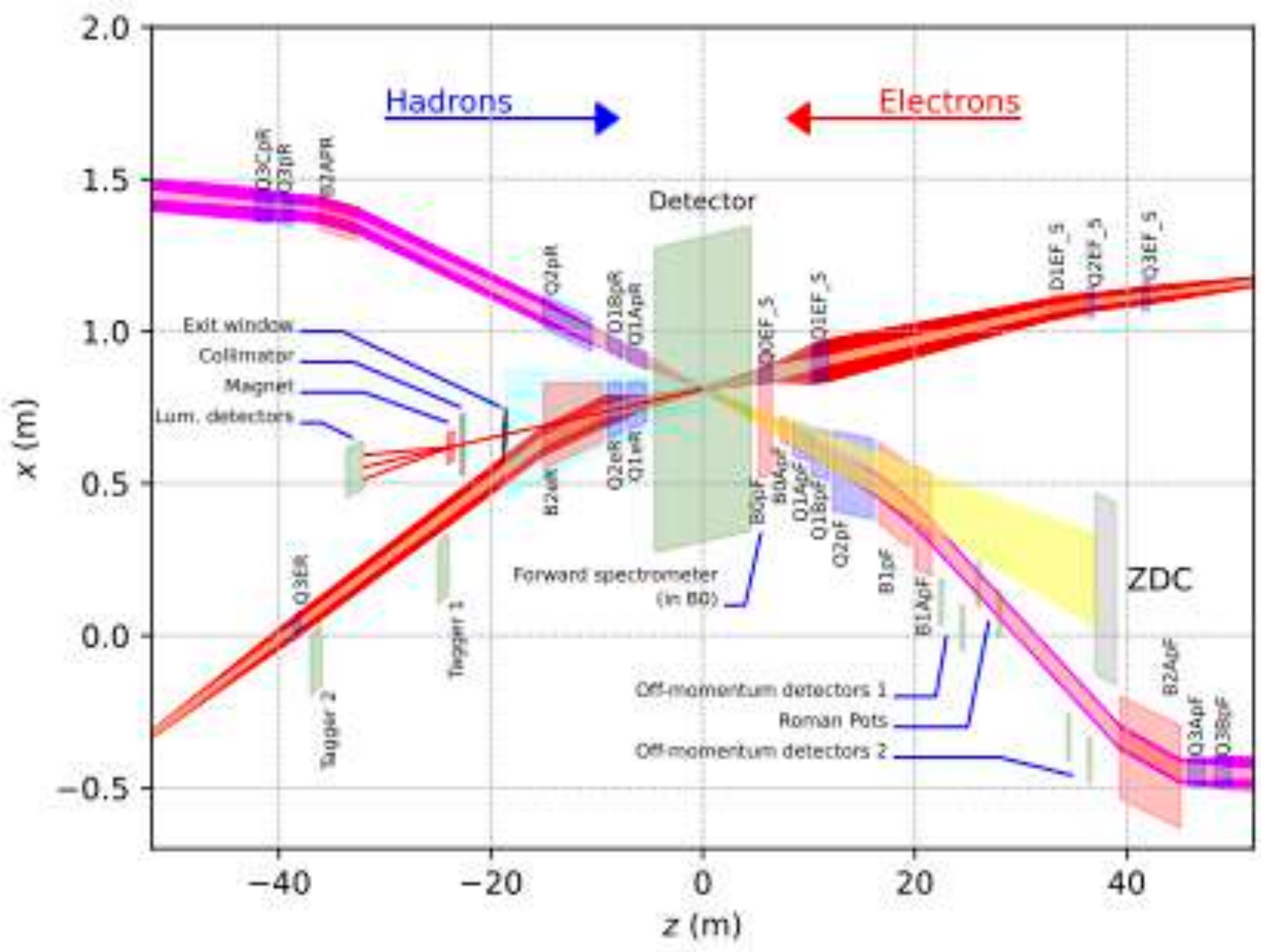
- Tracking
  - New 1.7 T solenoid
  - Si MAPS Tracker
  - MPGDs (uRWELL/uMegas)

- PID
  - hpDIRC
  - pfRICH
  - dRICH
  - AC-LGAD (~30 ps TOF)

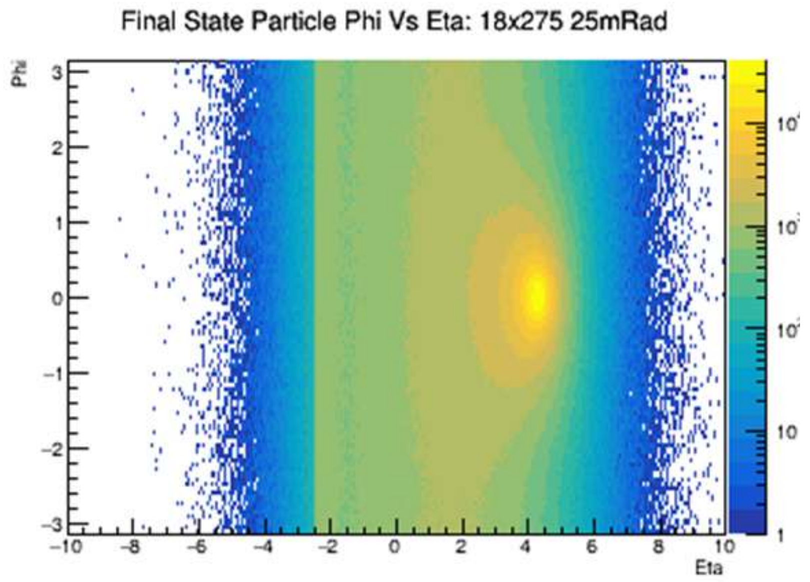
- Calorimetry
  - Imaging Barrel EMCal
  - PbW04 EMCal (backward)
  - Finely segmented EMCal+Hcal (forward)
  - Outer Barrel Hcal (sPHENIX)
  - Backward Hcal (tail-catcher)

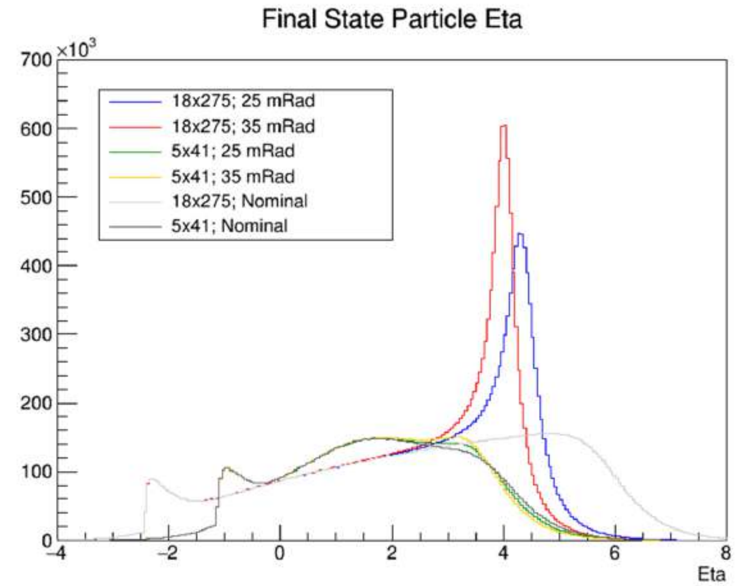
Beam Crossing Angle

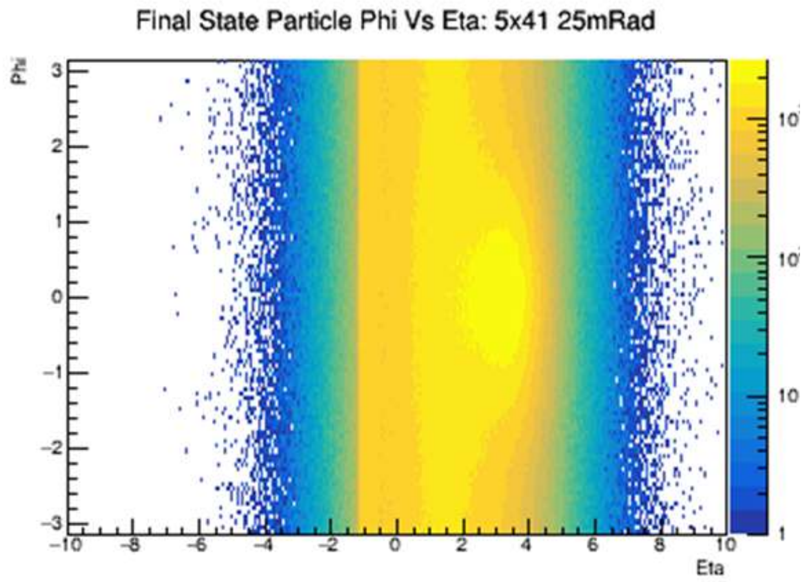
- ❑ Both interaction regions at the EIC will feature significant beam crossing angles (25 mRad for IP6 and 35 mRad for IP8)
- ❑ Crossing angles needed to avoid parasitic collisions which would degrade beams
- ❑ Presence of crossing angle will affect acceptance and detector design in the proton-going endcap – a region of great importance for many jet measurements
- ❑ A summary of beam effects, including crossing angle, bunch crabbing, divergence, etc as well as methods for simulating these can be found in the technical note here: https://zenodo.org/record/6514605#.ZETiIOzMJAY

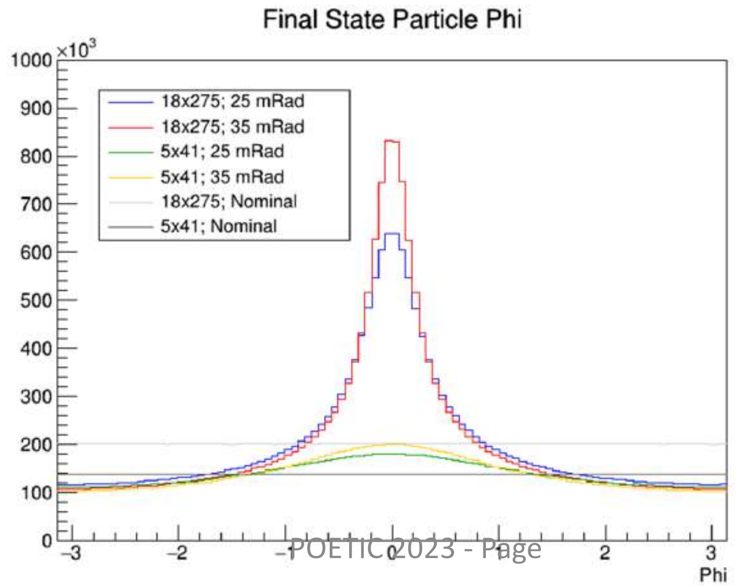


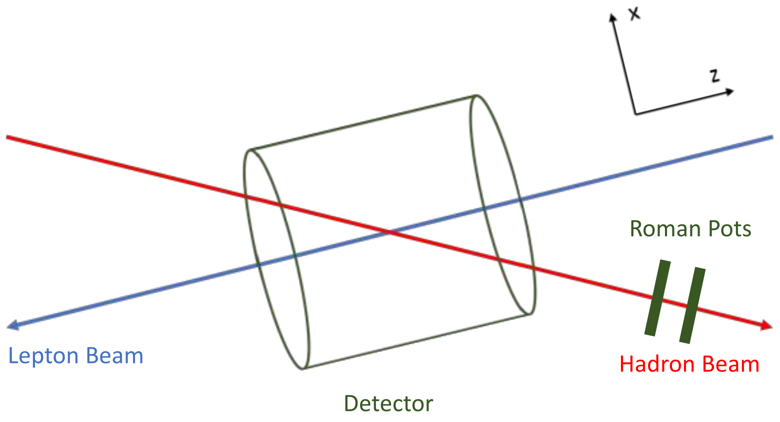
Final State Particle Distributions









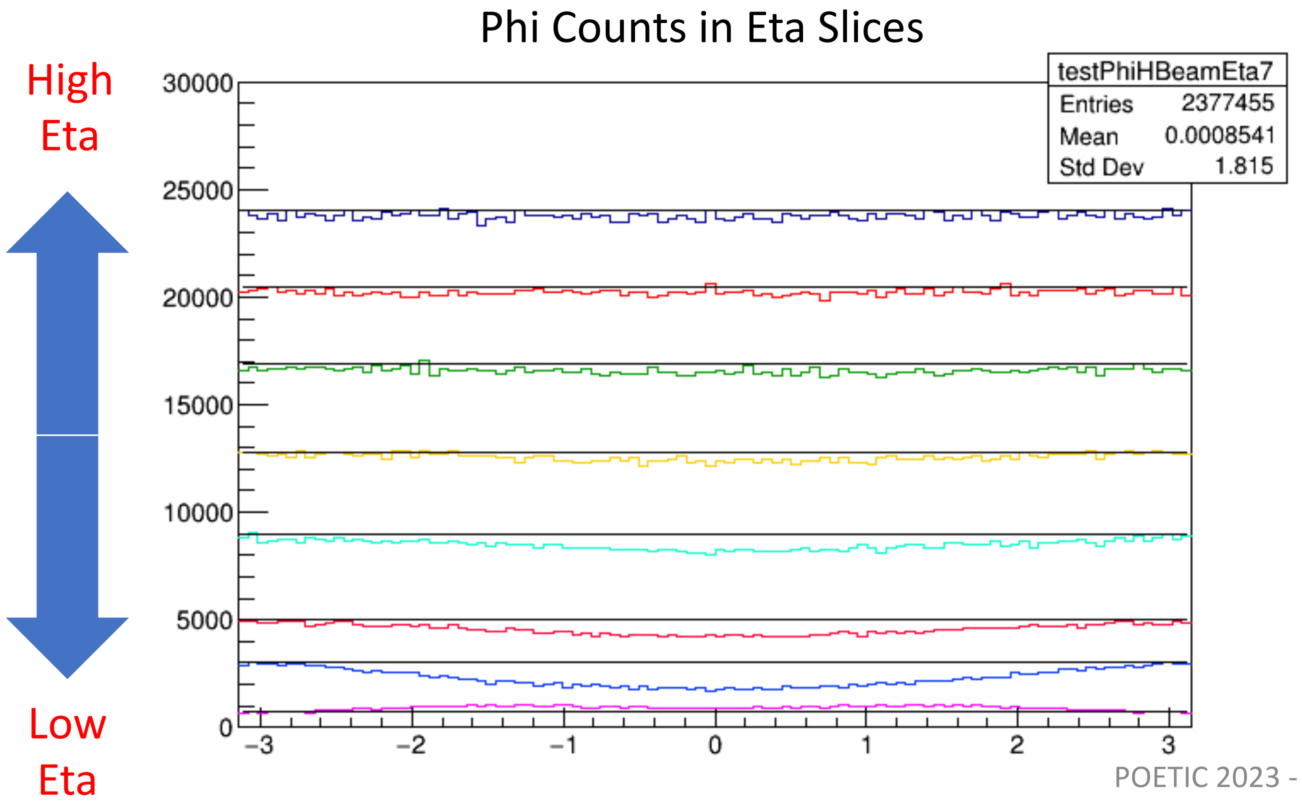


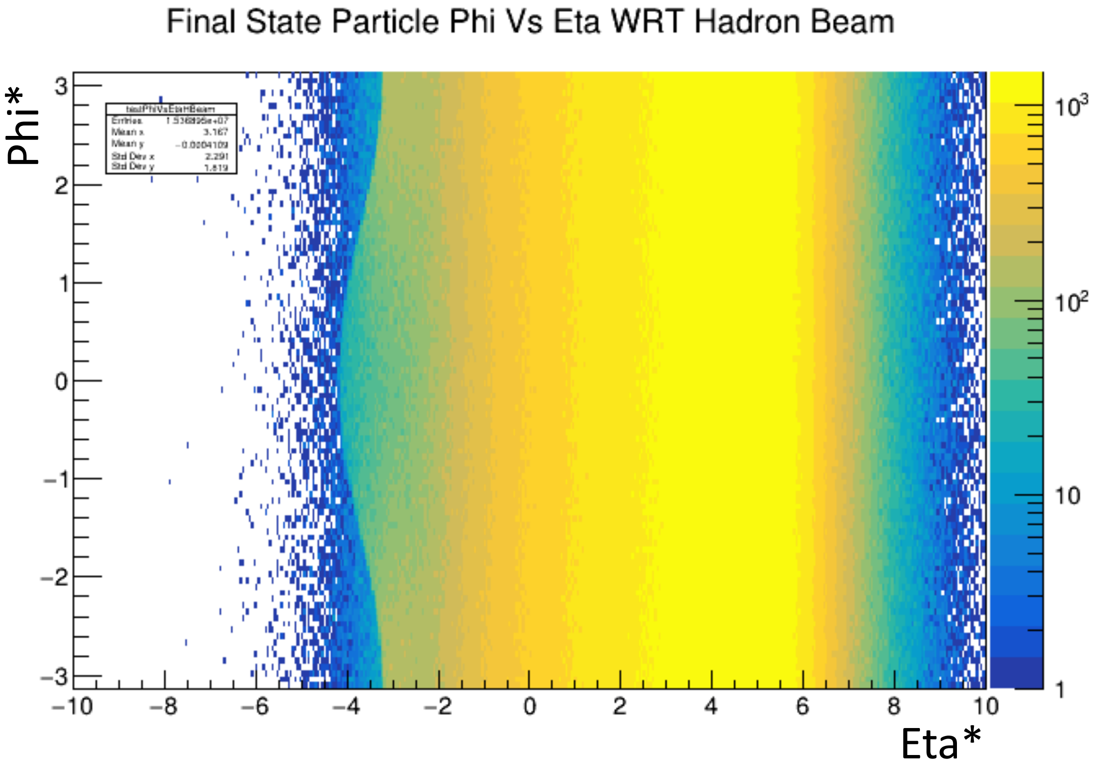
- Detector solenoid must align with electron beam to minimize synchrotron radiation: “lab frame” -> electron beam = z-axis
- When measuring in lab frame coordinates – see a hot spot in eta/phi corresponding to the beam direction
- More pronounced for more relativistic beams
- How do we mitigate these features?

Coordinates W.R.T. Hadron Beam

- ”Physics” in the forward region should be consistent around the hadron beam regardless of where the beam is pointing

- In some sense, the features seen above are simply artifacts of measuring about the ”wrong” axis -> instead, define eta and phi with respect to the hadron beam direction (Eta*, Phi*)





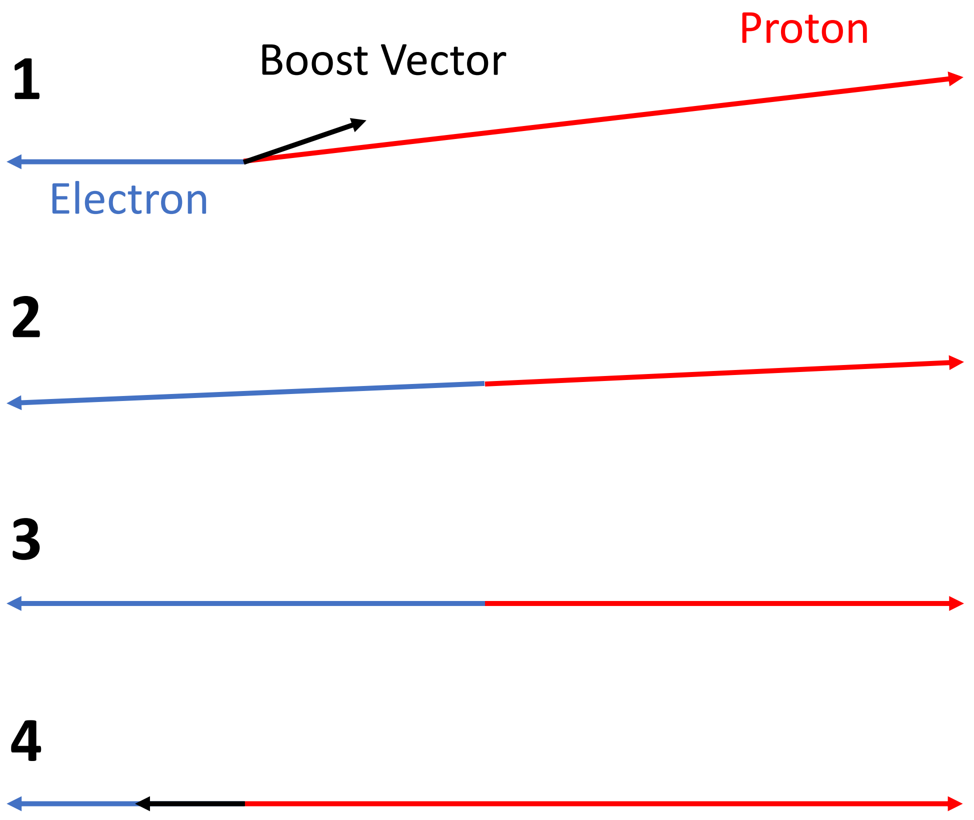
- When defined w.r.t. the hadron beam, the concentrations in eta and phi disappear

- However, because there is no common beam axis, the particle distribution along the electron-going direction becomes distorted

- Can avoid these distortions by boosting to a frame in which the beams are collinear

Head-On (Minimum Boost) Frame

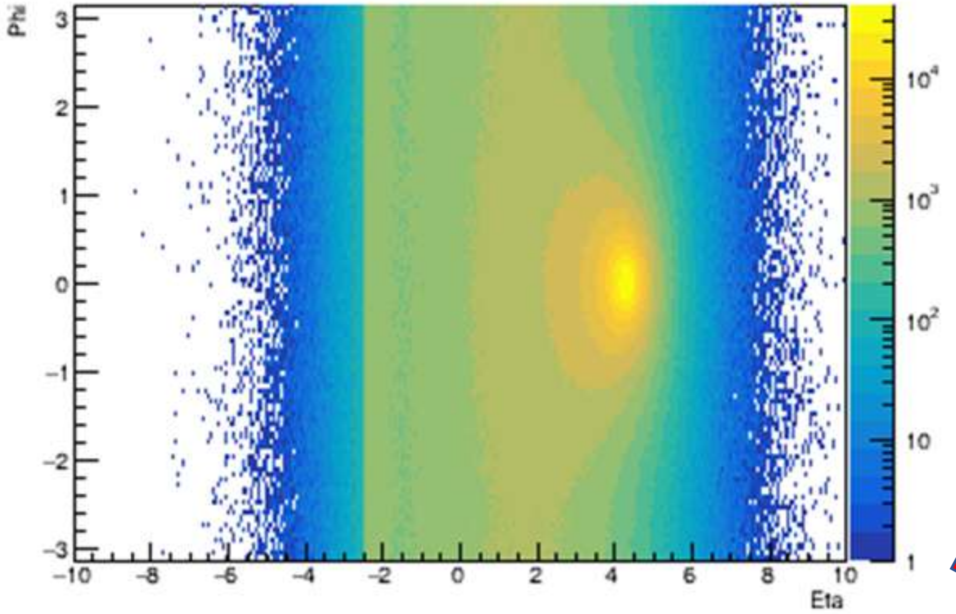
- ❑ Can boost and rotate into a frame in which the beams are collinear (no crossing angle) and energies are very close to the original (minimum boost)
- ❑ This should give an undistorted distribution of particles at high and low eta simultaneously



1. Initial Configuration in the Lab Frame includes a relative angle between the beams
2. Boost by sum of beam 4-momenta to get to CM Frame
3. Rotate about y-axis to eliminate x-component of momentum
4. Boost back along z to (nearly) restore original beam energies

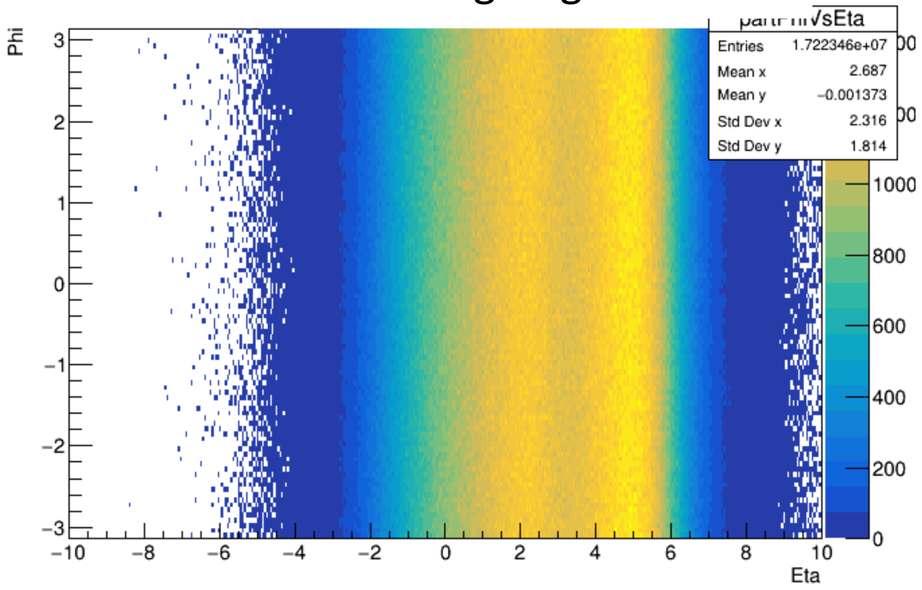
Head-On Frame Particle Distributions

Lab Frame Distribution

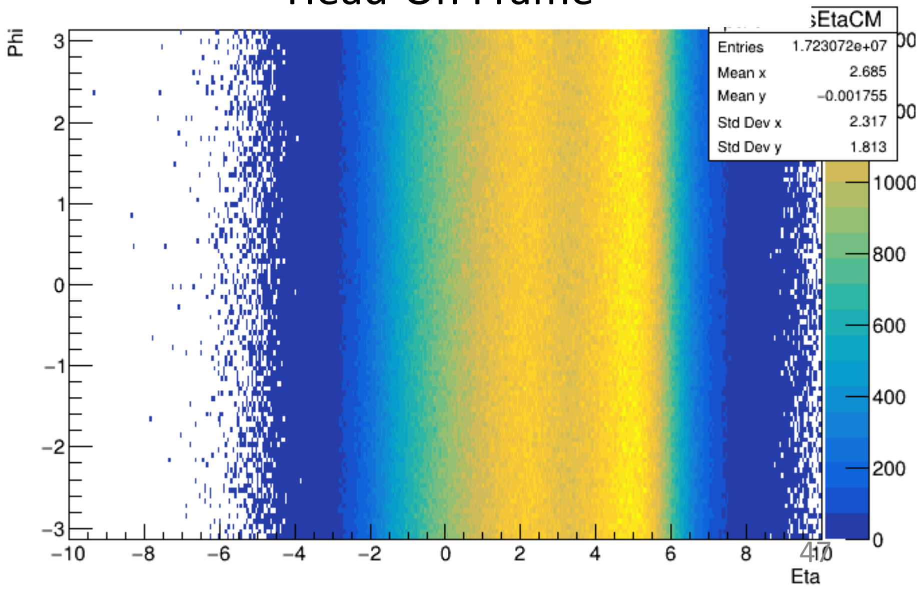


Transform to Head-On

No Crossing Angle



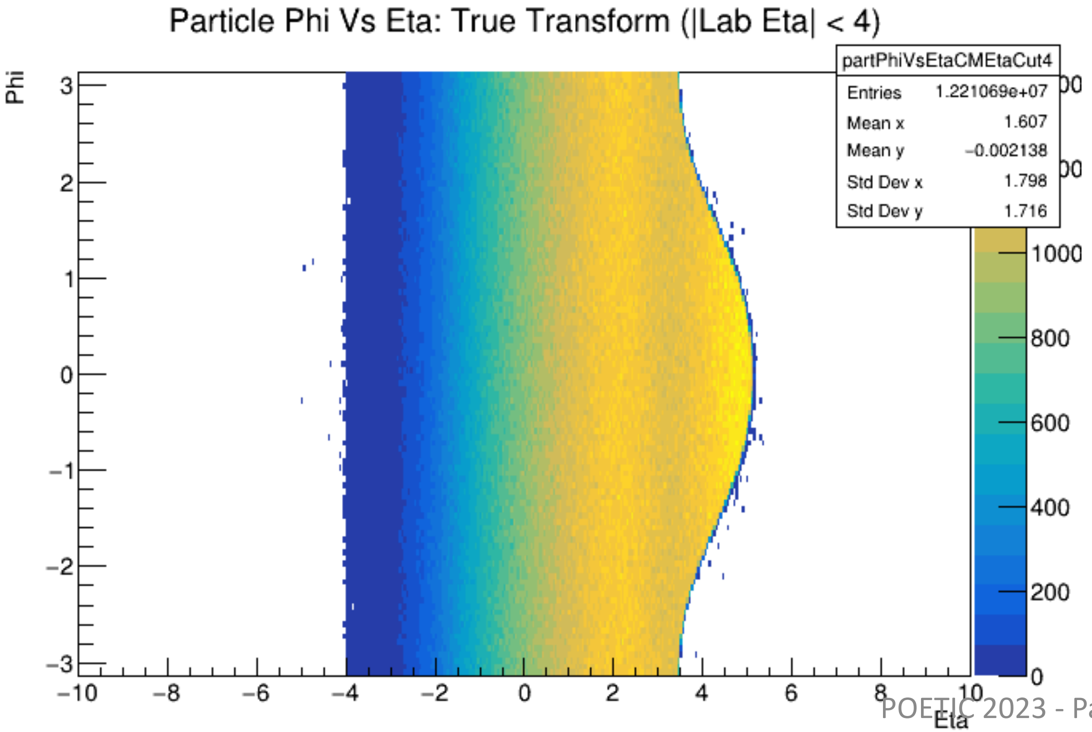
Head-On Frame



- Transformation to the head-on frame removes all features in the final state particle distribution for forward and backward regions simultaneously
- Resulting distribution matches that from default simulation with no crossing angle introduced

Detector Acceptance Considerations

- The head-on frame distributions shown previously assumed infinite acceptance – what effect will finite detector acceptance have?
- Displacement between beams means that acceptance cuts in the lab frame (w.r.t. the electron beam) will introduce phi-dependent acceptance features in head-on frame
- Try defining acceptance cuts w.r.t. the hadron beam instead



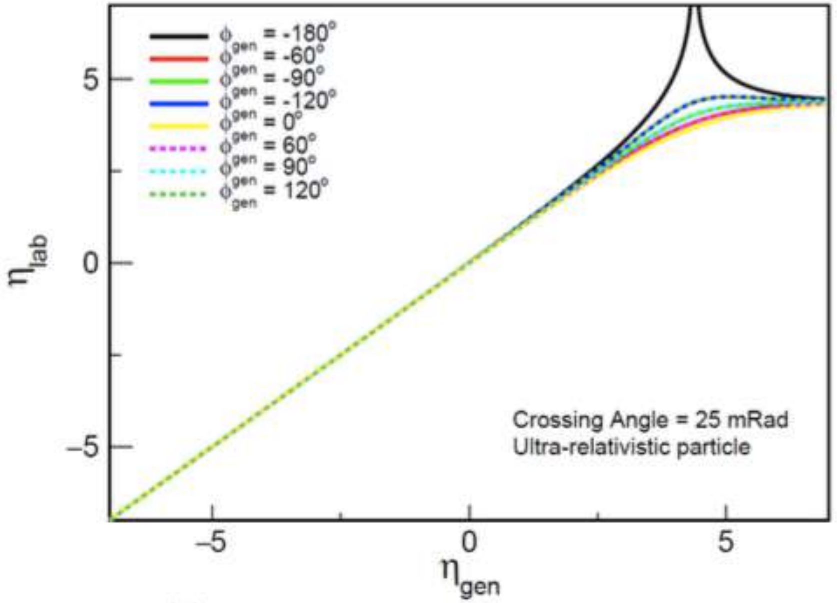
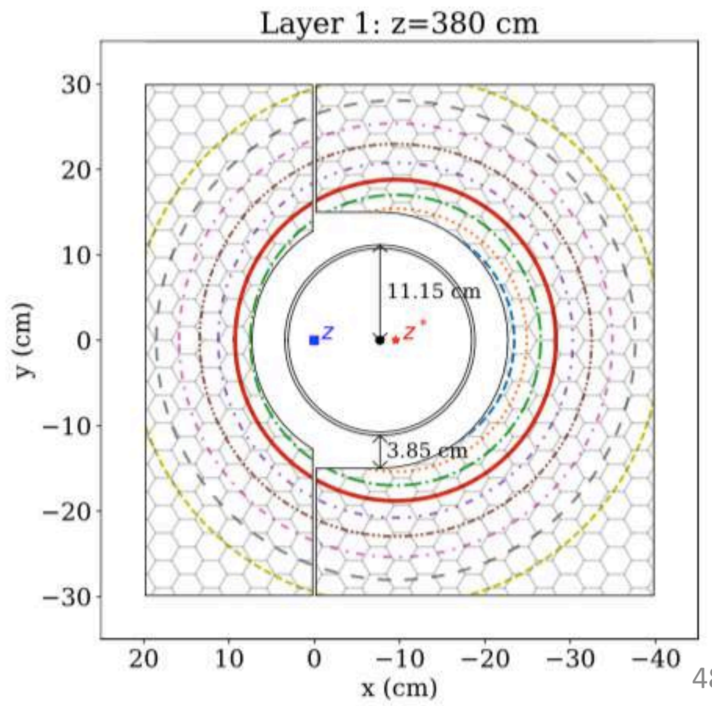


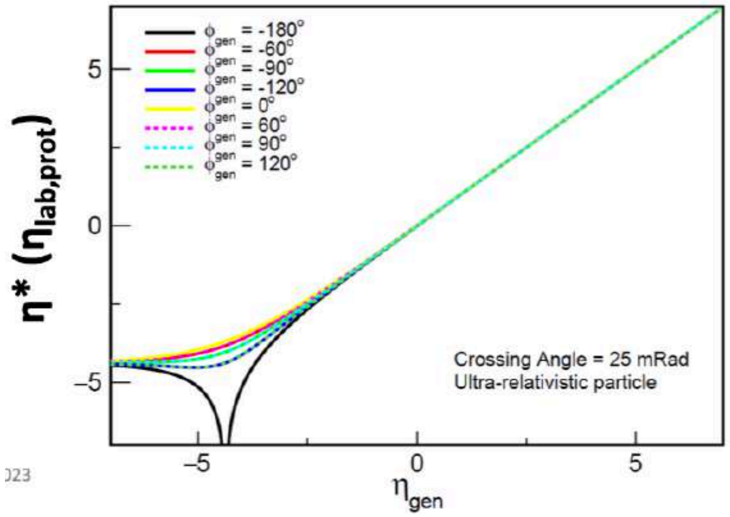
Figure by Barak Schmookler

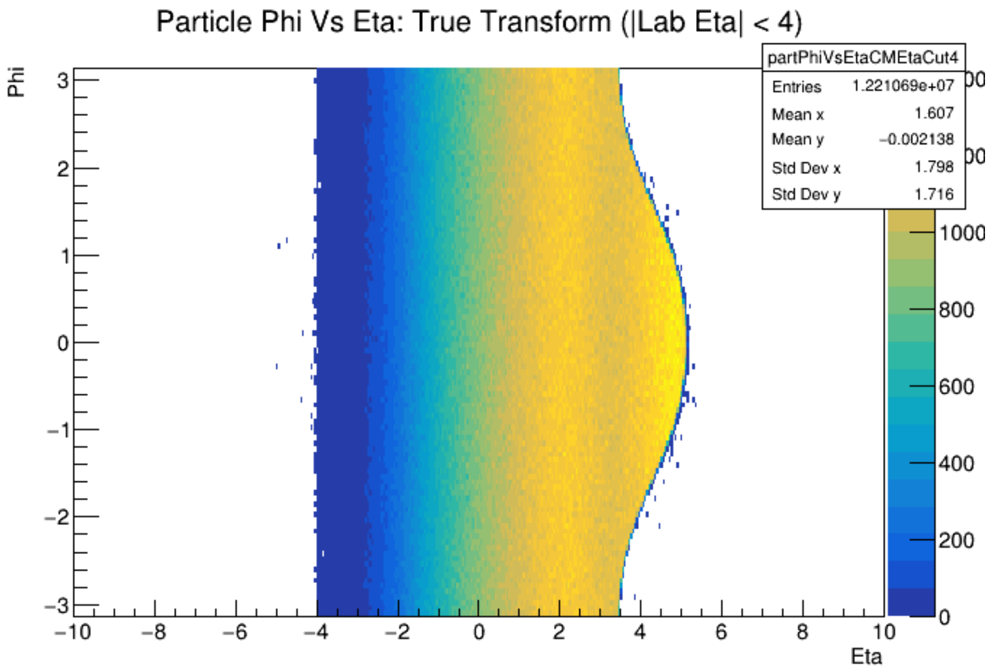


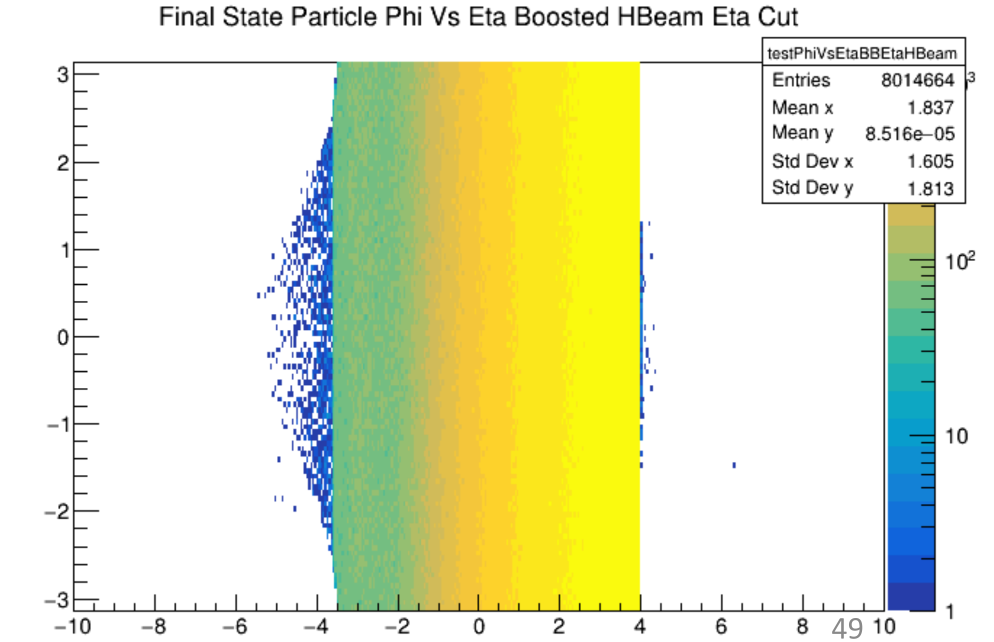
arXiv:2208.05472

Defining Acceptance Cuts

- ❑ The beam line shape in the endcap region is complicated, but mostly follows the hadron beam direction
- ❑ The z-axis in the head-on frame corresponds to the direction of the lab frame proton beam -> defining detector acceptance w.r.t. the hadron beam should eliminate the phi-dependent artifact
- ❑ Both plots on the right show the phi vs eta distribution where these quantities are defined in the head-on frame
  - ➢ Top plot applies a cut for |eta| < 4 where eta is defined relative to the electron beam
  - ➢ Bottom plot applies a cut for |eta| < 4 where eta is defined relative to the hadron beam



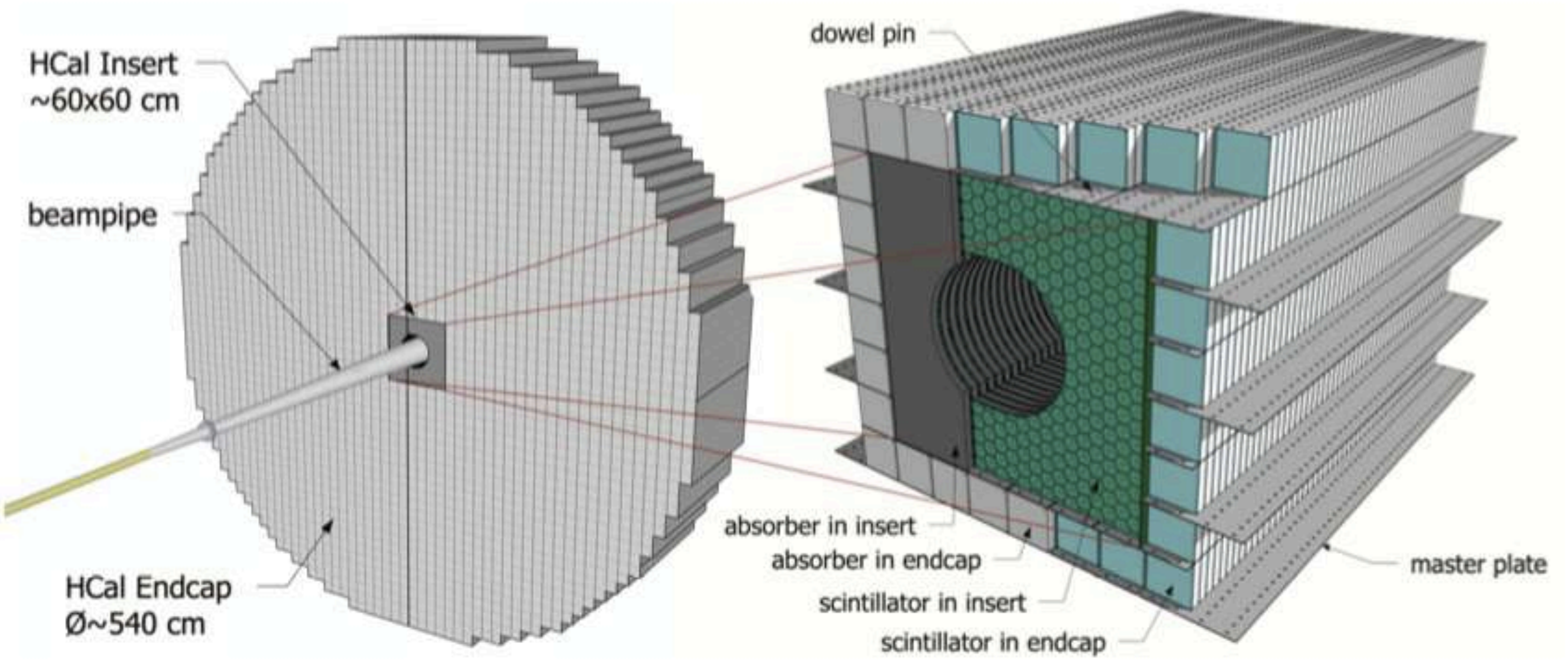


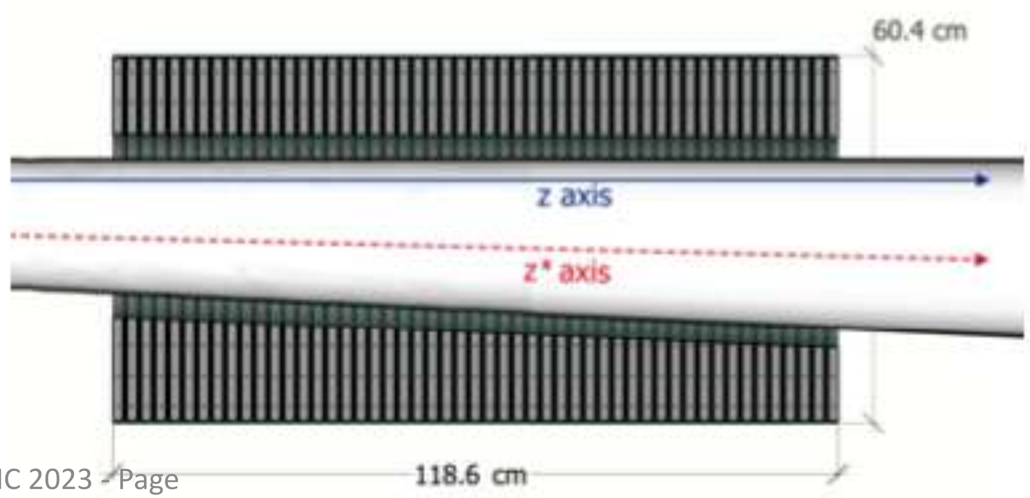


Detector Considerations: Calorimeter Insert

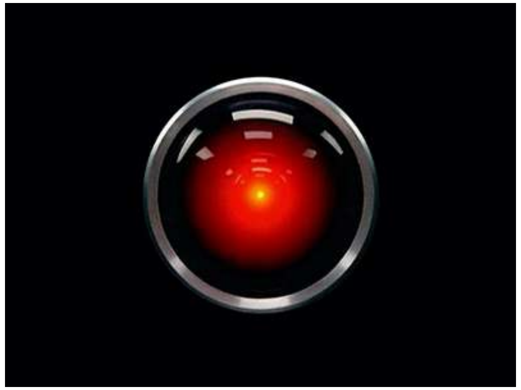
arXiv:2208.05472

- ❑ Precision calorimetry in the most forward regions of the central detector ($3 < \eta < 4$) will be important as tracking performance degrades rapidly while particle energy is the largest
- ❑ Coverage in this region important for classification of the hadronic-final-state, detection of highest energy jets, and tagging beam-induced backgrounds
- ❑ Increasing calorimeter acceptance will also provide more uniform acceptance in the head-on frame as discussed above
- ❑ Propose a steel/tungsten – scintillator sampling calorimeter insert in this region
- ❑ Cutouts in each layer will vary along z to follow the shape of the beampipe and maximize coverage





Future Directions



- AI and Machine Learning are quickly becoming common tools in many areas of physics – what about EIC jets?
  - AI/ML assist for design / calibration / reconstruction for relevant detector components
  - 'Physics Aware' AI which can extract underlying physics directly from event / jet observables
  - See arXiv:2012.06582; https://iris-hep.org/projects/ml4jets.html; YR Sec 11.12
- Events at EIC will be relatively sparse – can we adapt tools & techniques developed for jets to characterize the full event?
  - See for example arXiv:2101.02708 for event level grooming
- Major aspect of EIC physics program will be exploration of low-x / saturation phenomena
  - Implies low energy / p_T jets – develop tools needed to understand these jets theoretically and experimentally
- Expand use of jets to understand non-perturbative aspects of QCD

Summary

- The last ~10 years have seen a rapid increasing in the amount of theoretical and experimental work exploring the use of jets at the EIC

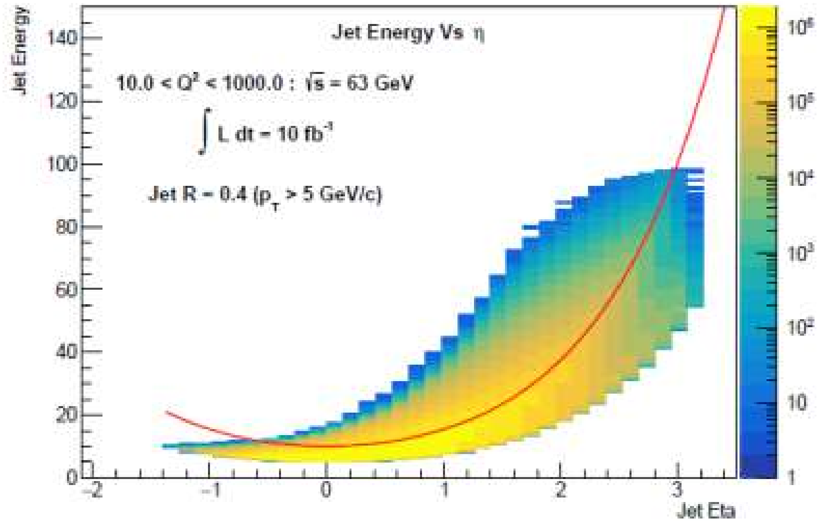
- Jet observables will contribute to nearly every area of the EIC science program, both complementing more traditional inclusive and semi-inclusive measurements and providing unique capabilities of their own

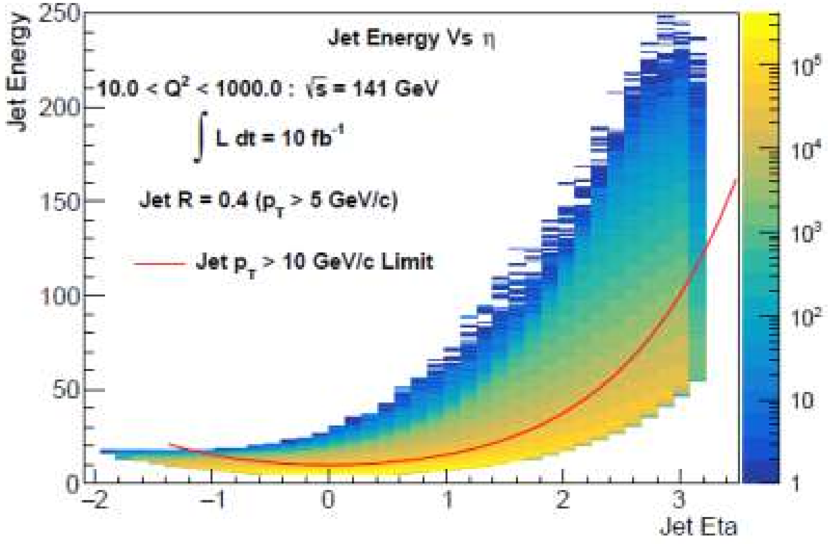
- The ability to accurately measure jets from low to high energy over the full detector span will be critical to probing both high and low-x physics and exploring interesting correlations between jets

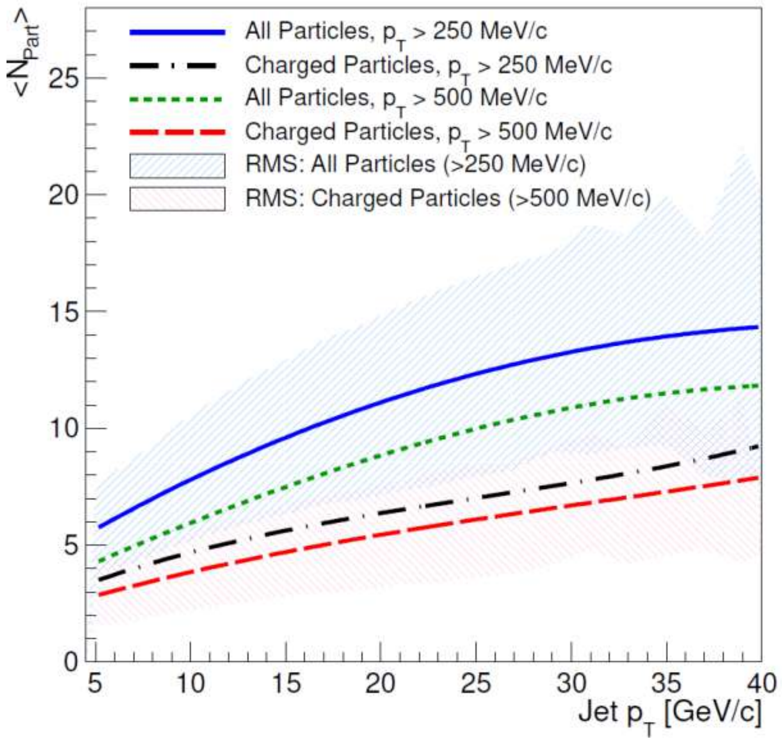
- Further development of jet clustering and substructure techniques, AI/ML applications, extensions to wider event classifications, etc will surely lead to new applications for jets in the future – data-taking is 10 years from now, the most exciting jet measurements may be something we have not even thought of yet!

Backup

Jet Kinematics

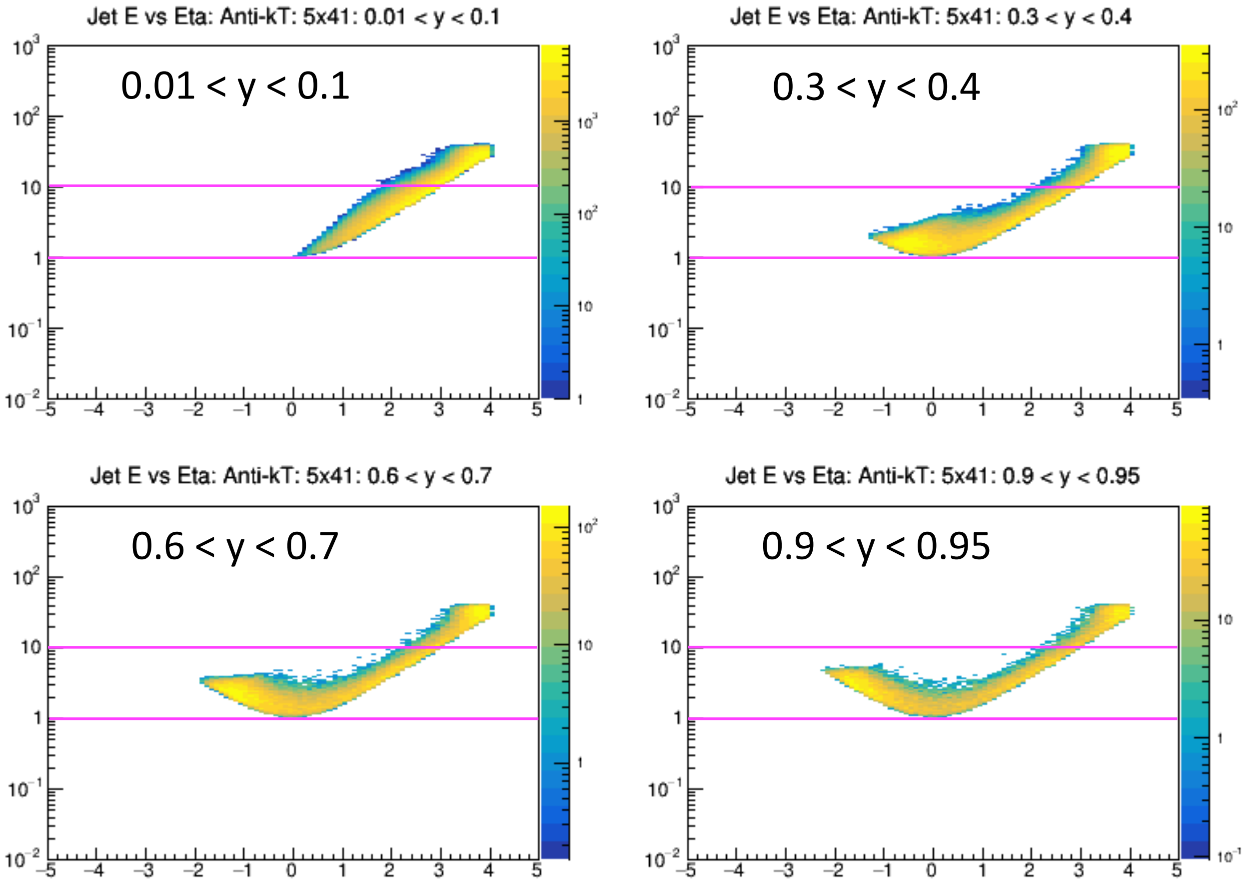






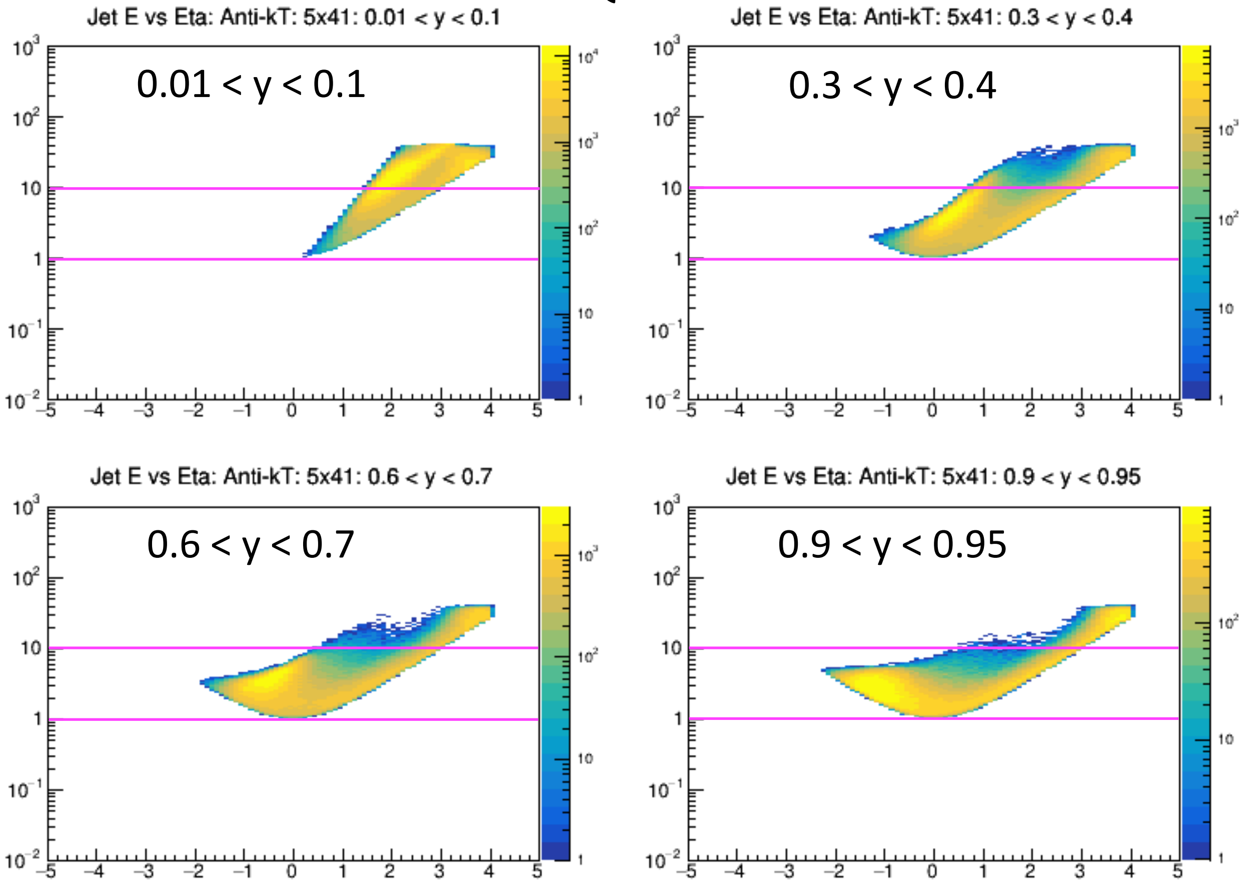
- ❑ Bulk of jets produced at the EIC will be low energy / low pT
- ❑ Pushing to analyze the lowest energy jets will provide access to the lowest x values, which will be important for spin structure and saturation studies
- ❑ In addition to being relatively low energy, jets will be quite broad and have few particles
- ❑ Must ensure theory and MC can make robust predictions for low energy jets and hadronization models can handle low multiplicity jets

Jet Distributions: Anti_kT (5x41)

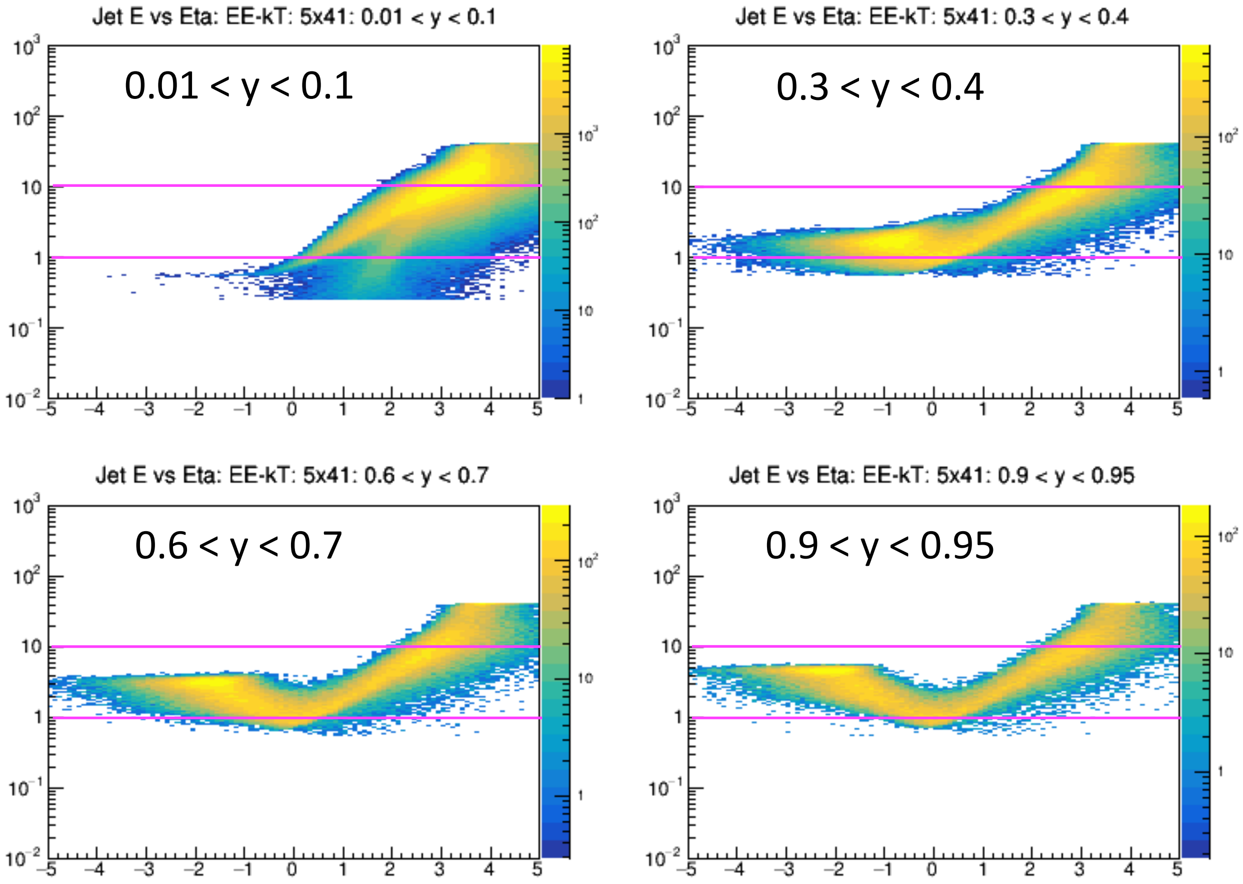


$0.1 < Q^2 < 1.0$

$10 < Q^2 < 100$

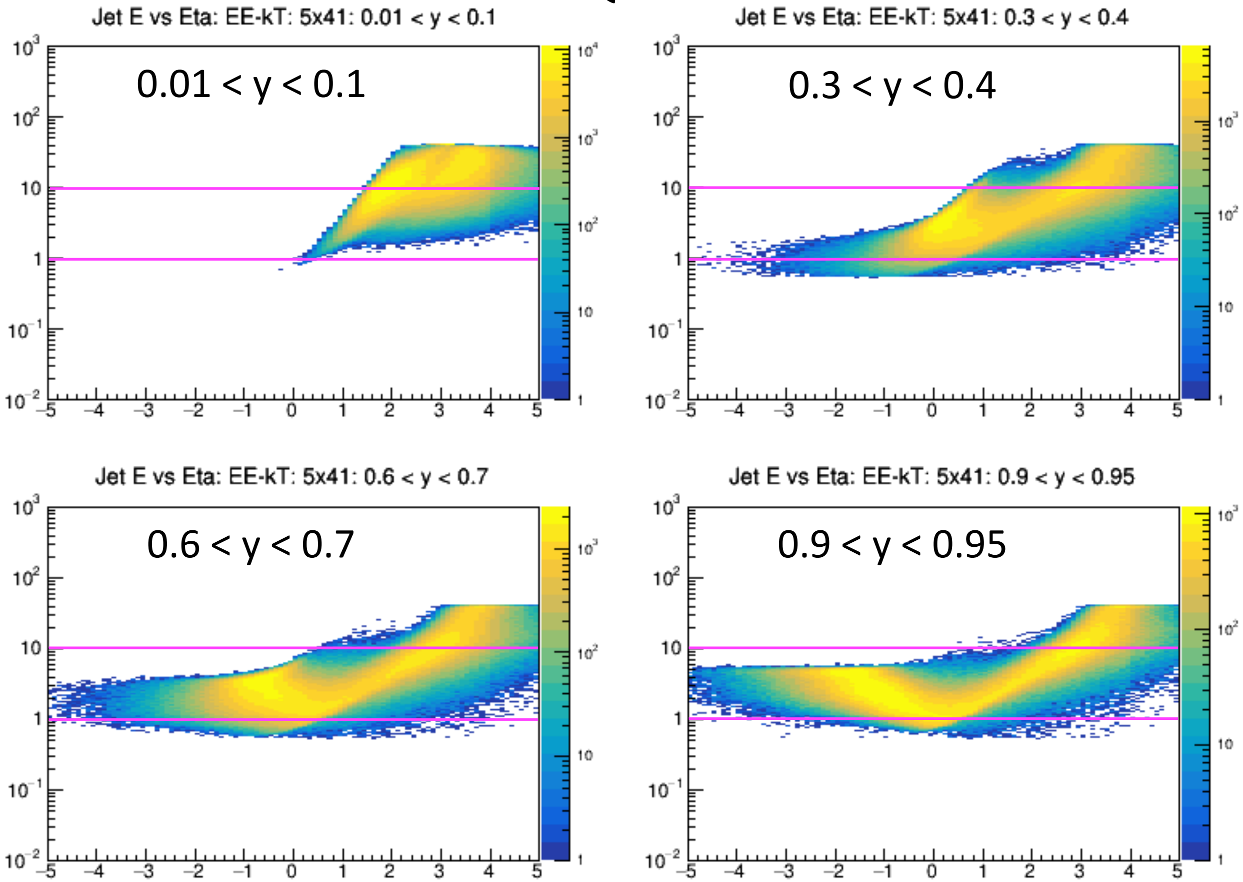


Jet Distributions: EE_kT (5x41)

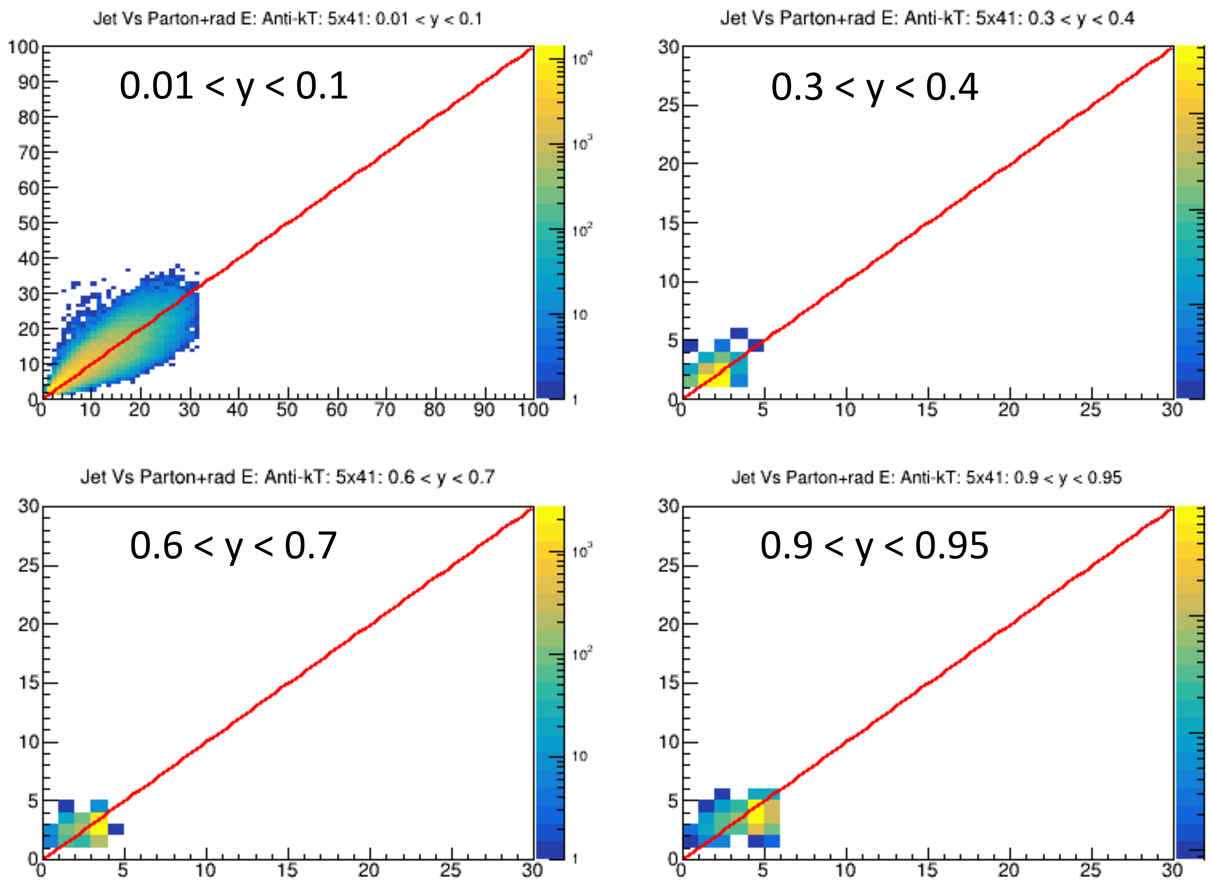


$0.1 < Q^2 < 1.0$

$10 < Q^2 < 100$

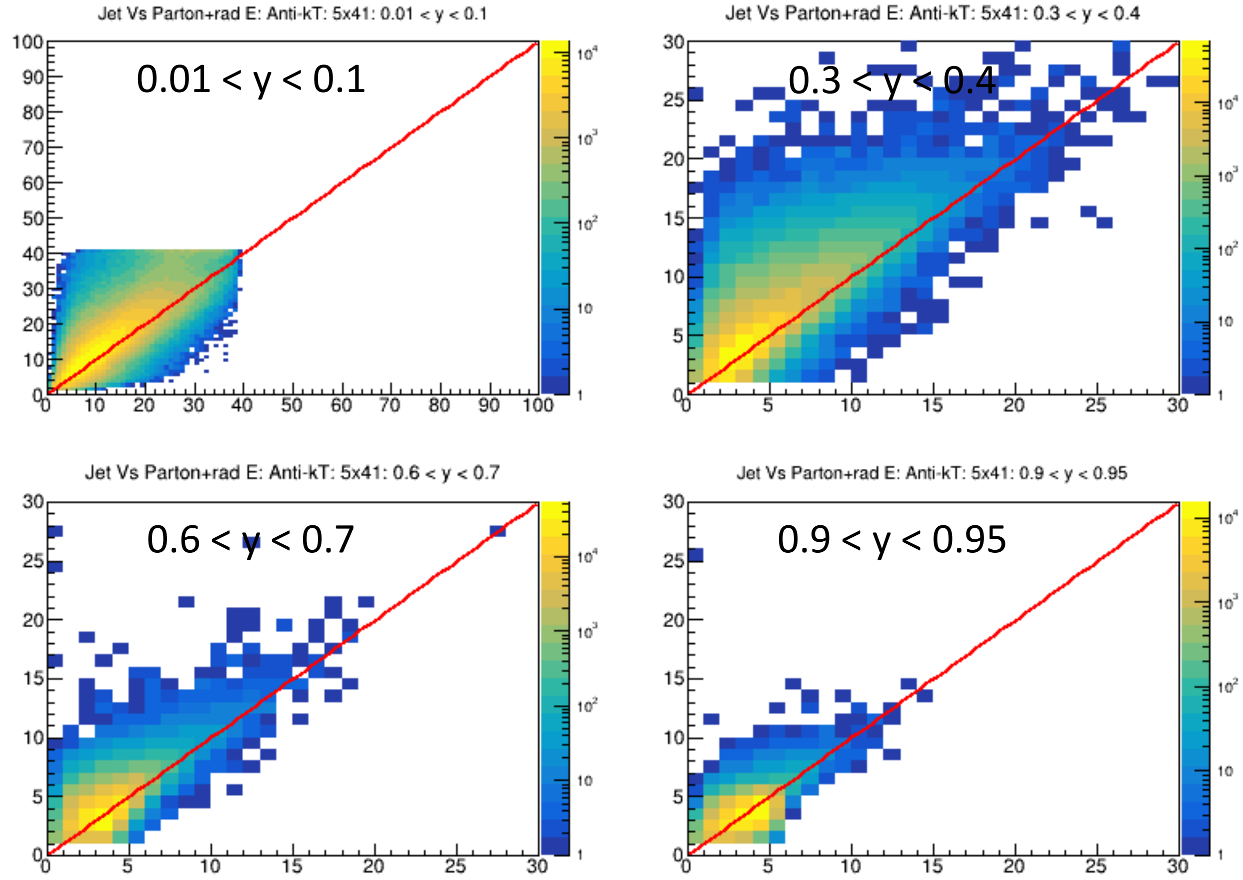


Jet – Parton Energy Comparison: Anti_kT (5x41)

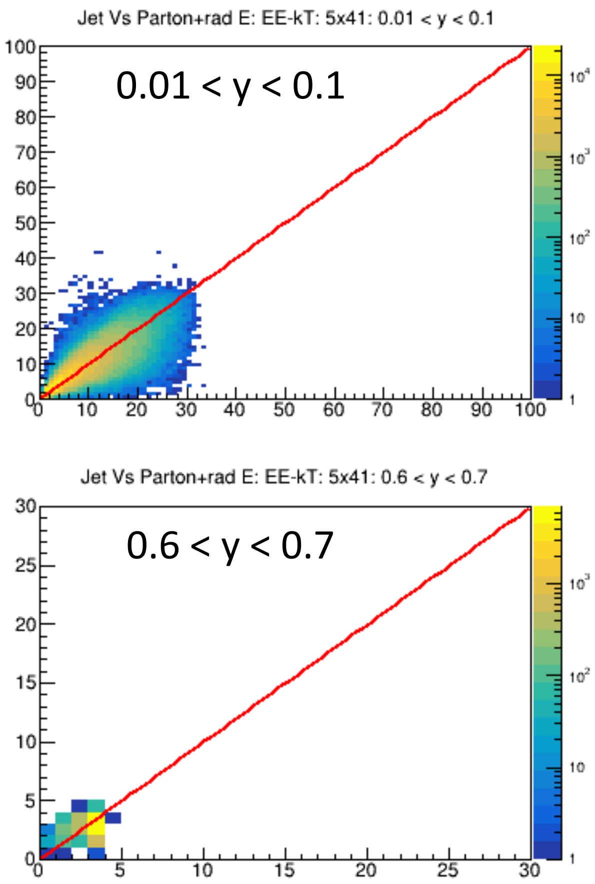


0.1 < Q^2 < 1.0

10 < Q^2 < 100

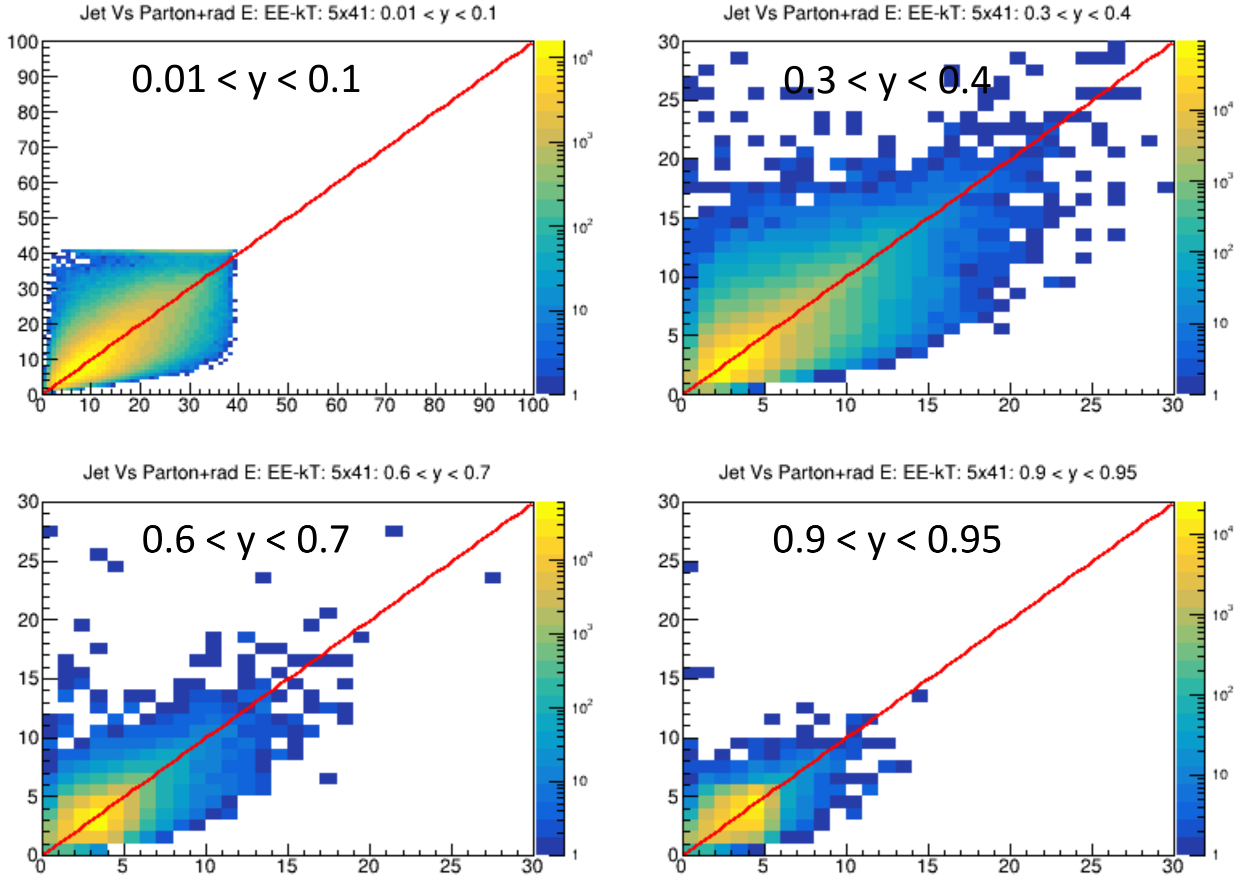


Jet – Parton Energy Comparison: EE_kT (5x41)



0.1 < Q^2 < 1.0

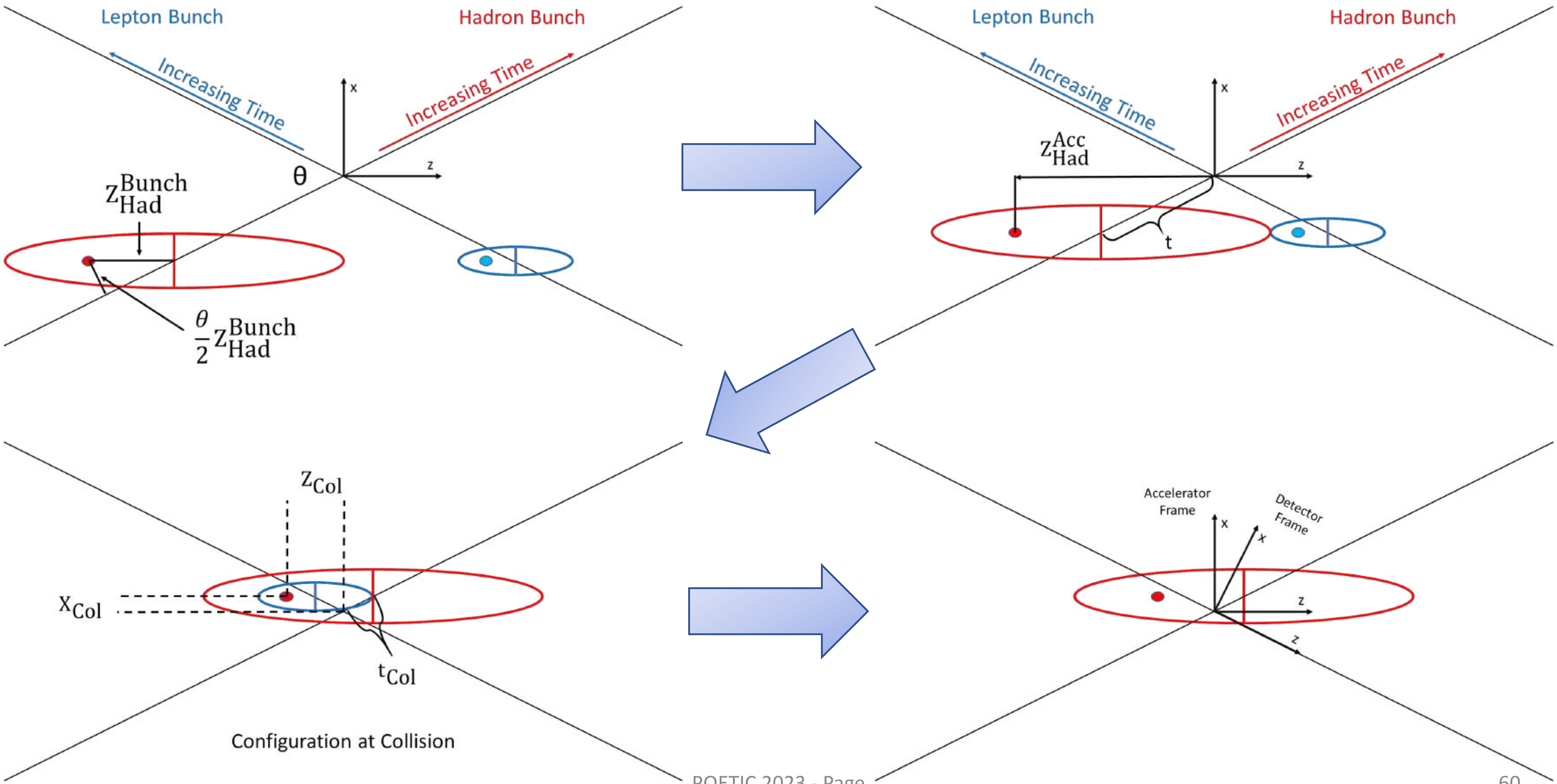
10 < Q^2 < 100



PYTHIA-8 Vertex Model

- ❑ Determine the x, y, and z vertex of the collision along with the time of collision
- ❑ Assume each bunch is rotated through half the beam crossing angle and assume it stays in a fixed orientation throughout the colliding region
- ❑ Assume particles in bunch are distributed along z as a gaussian with a sigma of the RMS bunch length cited in CDR table 3.3, 3.4. Correct collision distribution should follow automatically
- ❑ Assume particles in bunch are distributed in x,y as a gaussian with sigma given by Sqrt(RMS emittance x beta*) as given in CDR table 3.3, 3.4
- ❑ Procedure:
  1. Chose z (in in-bunch coordinates) of colliding particle in hadron and lepton bunch
  2. Propagate bunches until colliding particles overlap – this sets collision z, t, and a central x offset
  3. Randomly sample an x value according to beam widths and add to central x offset. Randomly sample a y value
  4. Rotate system from ‘accelerator frame’ into ‘detector frame’

PYTHIA-8 Vertex Model



PYTHIA-8 Vertex Model

$$z_{\mathrm{Had}}^{\mathrm{Acc}} = \mathrm{Cos}\left(\frac{\theta}{2}\right) \times t + z_{\mathrm{Had}}^{\mathrm{Bunch}}$$

$$z_{\mathrm{Lep}}^{\mathrm{Acc}} = -\mathrm{Cos}\left(\frac{\theta}{2}\right) \times t + z_{\mathrm{Lep}}^{\mathrm{Bunch}}$$

Z-position of interacting bunch from each beam as a function of time given by this set of equations

Collisioin occurs when Z_Had and Z_Lep are equal – can then solve the system to get time, z-position, and x-position of collision

$$t_{\mathrm{Col}} = \frac{\left(z_{\mathrm{Lep}}^{\mathrm{Bunch}} - z_{\mathrm{Had}}^{\mathrm{Bunch}}\right)}{2 \times \mathrm{Cos}\left(\frac{\theta}{2}\right)}$$

$$z_{\mathrm{Col}} = \frac{\left(z_{\mathrm{Lep}}^{\mathrm{Bunch}} + z_{\mathrm{Had}}^{\mathrm{Bunch}}\right)}{2}$$

$$x_{\mathrm{Col}} = t_{\mathrm{Col}} \times \mathrm{Sin}\left(\frac{\theta}{2}\right).$$

PYTHIA-8 Vertex Distributions