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Thomas Guhr:

Random Matrix Theory: from Single- to Many-Body Quantum Chaos

1. Phenomenological Model for the Spacing Distribution

The nearest-neighbor-spacing distribution $p(s)$ is one of the most important spectral observables. It is easily accessible in experiments and numerical simulations, in contrast to the energy correlation functions $X_k(\xi_1 \dots, \xi_k)$ of higher order k .

- Compare $p(s)$ und $X_2(r)$, $r = \xi_1 - \xi_2$, we recall that the unfolded energy correlation functions only depend on the differences $\xi_i - \xi_j$. Which $X_k(\xi_1 \dots, \xi_k)$ are needed to calculate $p(s)$, i.e. to which order k do they enter $p(s)$?
- In a phenomenological model, the spacing distribution satisfies the integral equation

$$p(s) = \mu(s) \int_s^\infty p(s') ds' .$$

Solve this equation for arbitrary $\mu(s)$, use $p(0) = \mu(0)$. Show that the normalization $\int_0^\infty p(s) ds = 1$ holds, if $\int_0^\infty \mu(s) ds \rightarrow \infty$.

- Which are the nearest-neighbor-spacing distributions that you get for the choices $\mu(s) = 1$ and $\mu(s) = \pi s/2$? Make qualitative drawings. What is the phenomenological meaning of the function $\mu(s)$?

2. Heuristic Derivation of Wigner's Surmise

Wigner put forward the surmise that the nearest-neighbor-spacing distribution $p(s)$ for Gaussian random matrices is well approximated by the formula

$$p(s) = p^{(W)}(s) = \frac{\pi}{2} s \exp\left(-\frac{\pi}{4} s^2\right) .$$

Show that this is the spacing distribution for real symmetric 2×2 matrices

$$H = \begin{bmatrix} H_{11} & H_{12} \\ H_{12} & H_{22} \end{bmatrix} .$$

The matrix elements are Gaussian distributed,

$$P(H) = \frac{\sqrt{2}}{\sqrt{2\pi v^2}^3} \exp\left(-\frac{1}{2v^2} \text{tr } H^2\right)$$

with variance v^2 . The eigenvalues of H are E_1 and E_2 . The dimensionless distance to the next neighbor is in this case just $s = |E_1 - E_2|/D$, where D formally represents the mean level spacing. The nearest-neighbor-spacing distribution is then given as

$$p(s) = \int_{-\infty}^{+\infty} dH_{11} \int_{-\infty}^{+\infty} dH_{22} \int_{-\infty}^{+\infty} dH_{12} P(H) \delta(s - |E_1 - E_2|/D) .$$

Compute $p(s)$ by making the following steps:

- (a) The matrix H is diagonalized by an orthogonal 2×2 matrix $U(\varphi)$ which depends on an angle φ . We have $H = U^T E U$ with $E = \text{diag}(E_1, E_2)$. Carry out a change of variables into eigenvalue-angle coordinates. Hint: the volume element transforms as

$$dH_{11} dH_{22} dH_{12} = \frac{1}{4} |E_1 - E_2| dE_1 dE_2 d\varphi ,$$

where E_1 and E_2 take values in \mathbb{R} . Do the angular integration.

- (b) Calculate the eigenvalue integrals, use $\tilde{E} = (E_1 + E_2)/2$ and $\varepsilon = E_2 - E_1$ as new coordinates.
- (c) Determine v^2 and D from the two normalization conditions

$$\int_0^{\infty} p(s) ds = 1 \quad \text{and} \quad \int_0^{\infty} s p(s) ds = 1 .$$

What is the meaning of the second condition?

- (d) Why is this derivation heuristic? Why is this not a full derivation of the spacing distribution in Random Matrix Theory?

3. Level Number Variance for Regular Systems

The level number variance $\Sigma^2(L)$ was introduced in the lecture by dividing the unfolded spectrum into M windows. The window $m = 1, \dots, M$ contains $\nu_m(L)$ levels. The level number variance is then

$$\Sigma^2(L) = \langle \nu^2(L) \rangle - \langle \nu(L) \rangle^2 ,$$

where $\langle \dots \rangle$ is the average over all windows. On the unfolded scale, there are $L \pm \sqrt{\Sigma^2(L)}$ levels in an interval of length L . Thus, in contrast to the nearest-neighbor-spacing distribution, the level number variance gives information on long-range properties of the level statistics. For chaotic systems, $\Sigma^2(L)$ behaves logarithmically for larger L . For regular systems, one often finds $\Sigma^2(L) \sim L$. The case of the harmonic oscillator in one dimension is very special.

- (a) Calculate $\Sigma^2(L)$ for the unfolded spectrum of the harmonic oscillator in one dimension. Begin with considering integer $L \in \mathbb{N}$. What is then $\Sigma^2(L)$? Show $\Sigma^2(L) = \Sigma^2(L+1)$. In the general case, decompose $L = [L] + x$ where $0 \leq x < 1$ with $[L]$ being the largest integer smaller L .
- (b) Show $\Sigma^2(L) = L$ for a system with Poissonian spectrum.

4. Spectral Rigidity

The spectral rigidity $\Delta_3(L)$ is closely related to the level number variance. In the lecture, it was introduced for window m as the minimum

$$\Delta_{3m}(L) = \frac{1}{L} \min_{A,B} \int_0^L d\xi (\nu_m(\xi) - A\xi - B)^2 ,$$

where $\nu_m(\xi)$ is the number of levels in window m .

- (a) Express A and B in terms of $\nu_m(\xi)$ by minimizing the above integral.
- (b) Re-insert A and B in the above expression to show that $\Delta_{3m}(L)$ can be written as

$$\Delta_{3m}(L) = \frac{1}{L} \int_0^L \nu_m^2(\xi) d\xi - \frac{1}{L} \vec{I}^T C \vec{I} ,$$

where the vector \vec{I} is defined as

$$\vec{I} = \begin{bmatrix} \int_0^L \xi \nu_m(\xi) d\xi \\ \int_0^L \nu_m(\xi) d\xi \end{bmatrix} .$$

What is the meaning of these two integrals? Determine the matrix C .

5. Spherical Functions — Itzykson–Zuber Integral

To compute spectral transitions, group integrals are needed. Unfortunately, only the case of integrals over the unitary group is explicitly doable for arbitrary matrix dimension N . Consider the simplest orthogonal case, the integral over $\text{SO}(2)$,

$$\Phi_2^{(1)}(x, k) = \int_{\text{SO}(2)} \exp(i \text{tr } x U^T k U) d\mu(U) ,$$

here we have $x = \text{diag}(x_1, x_2)$, $k = \text{diag}(k_1, k_2)$. To calculate it, proceed as follows:

- (a) Write $U \in \text{SO}(2)$ as a rotation matrix with angle φ in two real dimensions. Hint: $d\mu(U) = d\varphi$.
- (b) Calculate $\text{tr } xU^T kU$, write it in terms of only $\cos 2\varphi$. Hint: differences and sums $x_1 \pm x_2$, $k_1 \pm k_2$ appear.
- (c) Which function results from the remaining integral over φ ?

6. Spectrum of Chiral (Random) Matrices

In chiral Random Matrix Theory, one starts from the Dirac operator in the chiral basis and replaces it by matrices of the form

$$D = \begin{bmatrix} 0 & M \\ M^\dagger & 0 \end{bmatrix},$$

such that $D = D^\dagger$. Here, D should not be confused with the mean level spacing introduced previously. We notice that the matrix M has no further symmetries. In particular, M can be a rectangular matrix of dimension $N \times L$. Without loss of generality, we may assume that $L \geq N$. Moreover, we restrict ourselves to the case that M and hence D are real. Chiral symmetry reflects itself in a pointwise symmetry of the spectrum. Show the following:

- (a) In the case $L > N$, there are $L - N$ eigenvalues which are zero. They are called zero modes.
- (b) All non-zero eigenvalues λ of D come in pairs. What does that imply for the level density?
- (c) The spectrum of chiral random matrices is closely related to that of Wishart correlation matrices.