# Gravitationally Mediated Entanglement: Newtonian Fields, Gravitons and Black Holes

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Princeton University

ICTP: Witnessing Quantum Aspects of Gravity in a Lab

D. Danielson, G.S., & R.M. Wald Phys. Rev. D 105, 086001 (2022) [arXiv:2112.10798]

D. Danielson, G.S. & R.M. Wald [arXiv:2205.06279],[arXiv:2301.00026],[2407.02567]

see also: A. Belenchia, R. M. Wald, F. Giacomini, E. Castro-Ruiz, C. Brukner & M. Aspelmeyer [arXiv:1807.070105]

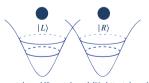
J. Wilson-Gerow, A. Dugad & Y. Chen [arXiv:2405.00804]

and A. Biggs & J. Maldacena [arXiv:2405.02227]

September 25, 2024

## (Gedanken) experiments and Quantum Gravity





[Carney et al. 2019],[Bose et al. 2017], [Marletto et al. 2017], ...

"One should think about designing a gedankenexperiment which uses a gravitational link and at the same time shows quantum interference."

[ Feynman, 1957 Chapel Hill Conference]

- ➤ To probe properties of quantum gravity, it's useful to consider situations where both quantum theory and gravity play an essential role.
- ► What aspects of quantum gravity can we learn from such (gedanken)experiments?

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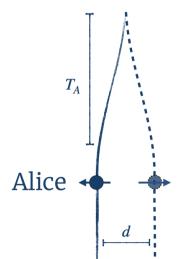
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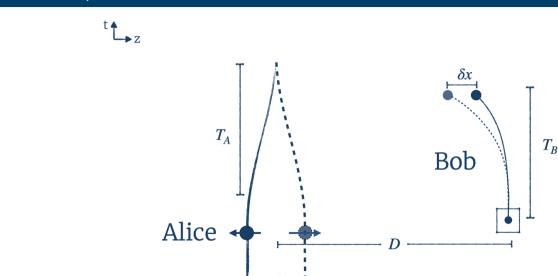




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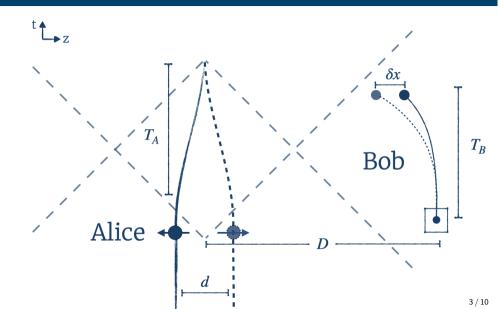
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- ▶ Bob performs this experiment at spacelike separation from Alice's recombination.



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$$\mathscr{D} = 1 - |\langle \Psi_1 | \Psi_2 \rangle| = 1 - e^{-\frac{1}{2} \langle N \rangle_{\Psi_1 - \Psi_2}} \text{ where } \langle N \rangle_{\Psi_1 - \Psi_2} \sim (q_A d / T_A)^2$$

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▶ Both inequalities cannot be satisfied for  $T_A$ ,  $T_B < D$ . If Alice can maintain coherence, then Bob does not have enough time to get "which path" information. If Bob has enough time, then Alice decoheres herself be entangling radiation.

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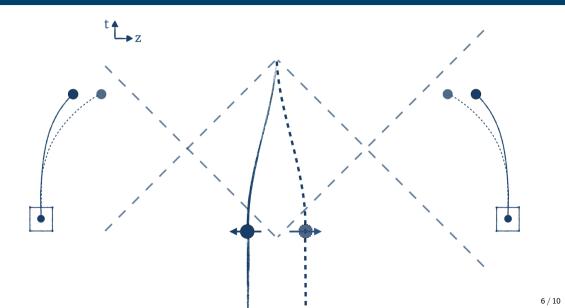
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► The analysis and conclusions parallel the EM case with "dipole" ↔ "quadrupole". Quantized radiation and vacuum fluctuations of the gravitational field are essential to avoid contradictions with causality and complementarity

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To prove this we will obtain a precise relationship between causality and decoherence

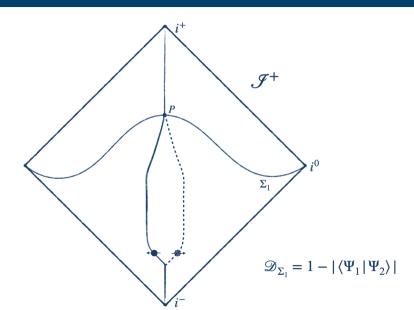
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- ▶ However, after the recombination, the components of her particle now share a common Coulomb field. On any Cauchy surface  $\Sigma_1$  after the recombination, the total quantum state with the Coulomb field subtracted is

$$rac{1}{\sqrt{2}}(|A_1;\uparrow\rangle\otimes|\Psi_1\rangle_{\Sigma_1}+|A_2;\downarrow\rangle\otimes|\Psi_2\rangle_{\Sigma_1})$$

where  $|\Psi_1\rangle$  and  $|\Psi_2\rangle$  are genuine radiation states which contain *all* of the "which path" information on  $\Sigma_1$ .



#### Decoherence due to Bob

▶ Consider the case where Alice recombines her particle slowly and emits negligible radiation. Let  $|B_0\rangle$  be the initial state of Bob's apparatus (with any number of assistants) which is initially unentangled with Alice's particle

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▶ Due to the Coulomb/Newtonian interaction, Bob's apparatus will evolve to  $|B_1\rangle$  if Alice's particle followed path 1 and  $|B_2\rangle$  if Alice's particle followed path 2

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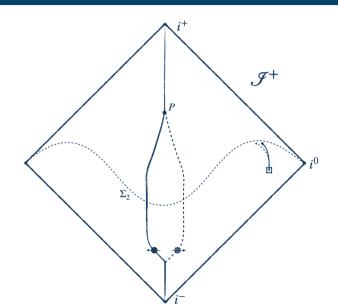
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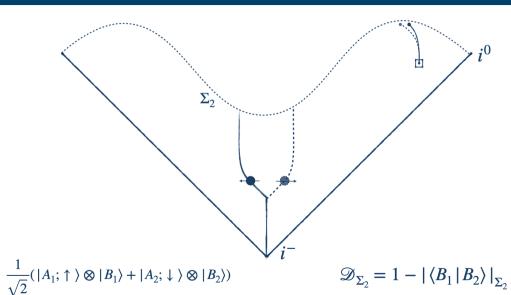
▶ The decoherence at any time  $\Sigma_2$  where Bob's apparatus is interacting with the Coulomb/Newtonian field of Alice' particle

$$\mathscr{D}_{\Sigma_2} = 1 - |\langle B_1 | B_2 \rangle|_{\Sigma_2}$$

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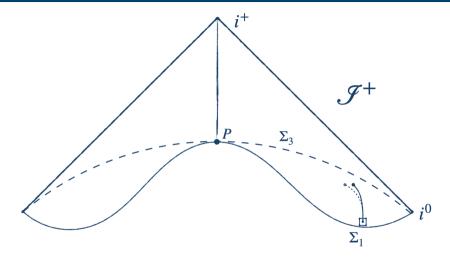
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Viewpoints (1) and (2) are equivalent!

 $\blacktriangleright$  Consider the joint evolution of the Alice + Bob + radiation quantum state from  $\Sigma_1$ 

$$\frac{1}{\sqrt{2}}(|A_1;\uparrow\rangle\otimes|\Psi_1\rangle_{\Sigma_1}+|A_2;\uparrow\rangle\otimes|\Psi_2\rangle_{\Sigma_1})\otimes|B_0\rangle$$

to a Cauchy surface  $\Sigma_3$  which is to *future* of both Alice's recombination and Bob's experiment.



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► Alice's particle does not change under this evolution and Bob's apparatus simply becomes entangled with the radiation states emitted by Alice

$$\frac{1}{\sqrt{2}}(|A_1;\uparrow\rangle\otimes|\Psi_1'\rangle_{\Sigma_3}\otimes|B_1\rangle+|A_2;\uparrow\rangle\otimes|\Psi_2'\rangle_{\Sigma_3}\otimes|B_2\rangle)$$

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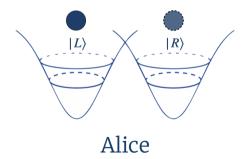
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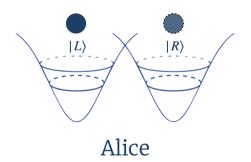
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Unitarity of the evolution implies

$$\langle \Psi_1' | \Psi_2' \rangle_{\Sigma_3} \, \langle \mathcal{B}_1 | \mathcal{B}_2 \rangle = \langle \Psi_1 | \Psi_2 \rangle_{\Sigma_1} \implies | \, \langle \mathcal{B}_1 | \mathcal{B}_2 \rangle \, | \geq | \, \langle \Psi_1 | \Psi_2 \rangle \, |_{\Sigma_1} \implies \mathscr{D}_{\mathrm{Alice}} \geq \mathscr{D}_{\mathrm{Bob}}$$

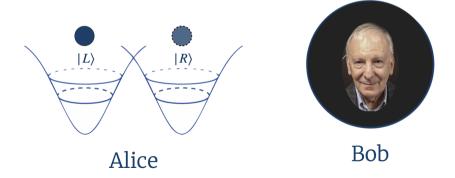
- ▶ If both Alice and Bob follow their protocols, then both (1) Newtonian-mediated entanglement and (2) on-shell graviton entanglement are equivalent veiwpoints
- ▶ It is essential that both (1) and (2) or, alternatively, neither (1) nor (2) are valid descriptions of the entanglement in order to provide a consistent description of a quantum spatial superposition and avoid contradictions with causality and complementarity.
- Additionally, we have shown that  $\mathscr{D}_{Alice} \geq \mathscr{D}_{Bob}$ , generalizing the analysis of [Belenchia et al, 2018]. However, read in the other direction, this implies that any quantum spatial superposition must be at least as decohered as any Bob(s) at spacelike separation.

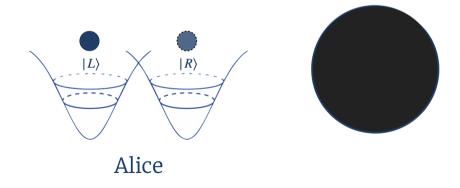






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- ▶ Any physical body would decohere the superposition by gravitationally interacting with its internal degrees of freedom. As Daine will explain tomorrow, a black hole decoheres quantum superpositions as if it contains internal degrees of freedom.