

Spin polarization in off-equilibrium medium at weak and strong coupling



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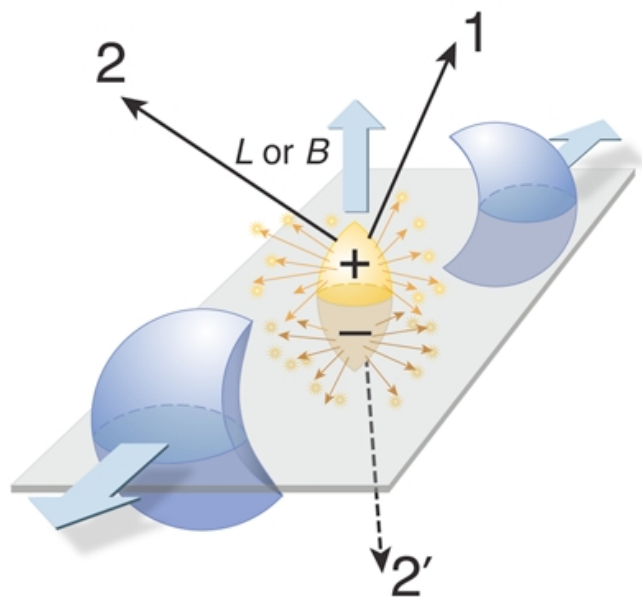
The 9th International Conference on Chirality, Vorticity and
Magnetic Field in Quantum Matter, Sao Paulo, July 7-11, 2025

Li, SL, 2505.21883
SL, Tian, 2410.22935

Outline

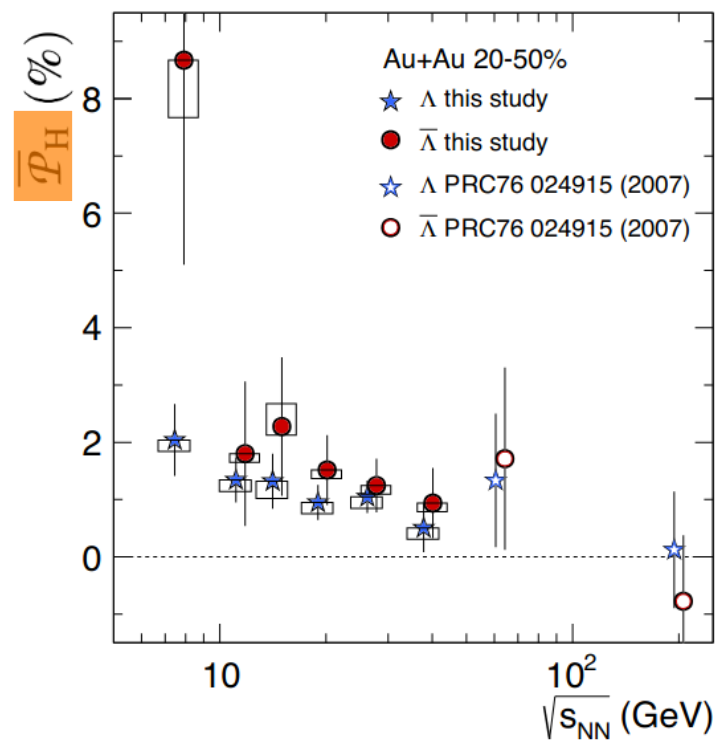
- ♦ Spin polarization in heavy ion collisions: sensitivity on structure/interaction of particles
- ♦ Spectral function of particle in off-equilibrium state
- ♦ Polarized spectral function at weak coupling
- ♦ Polarized spectral function at strong coupling
- ♦ Summary and outlook

Global polarization of Λ



$$L_{ini} \sim 10^5 \hbar \rightarrow S_{final}$$

Liang, Wang, PRL 2005, PLB 2005



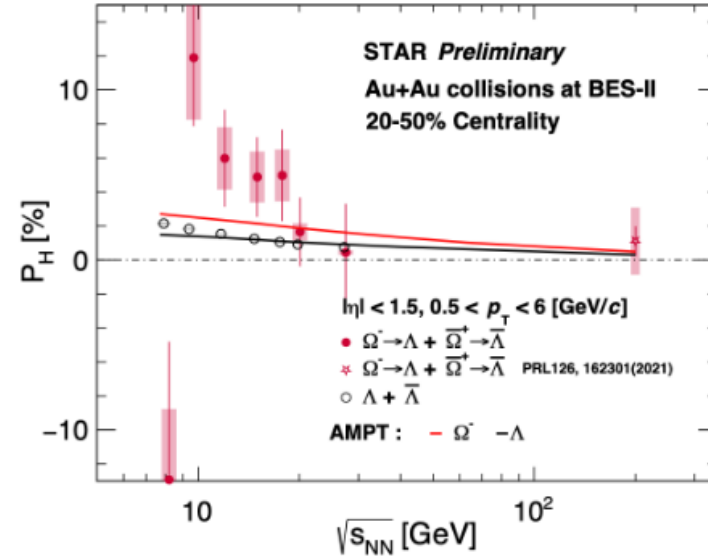
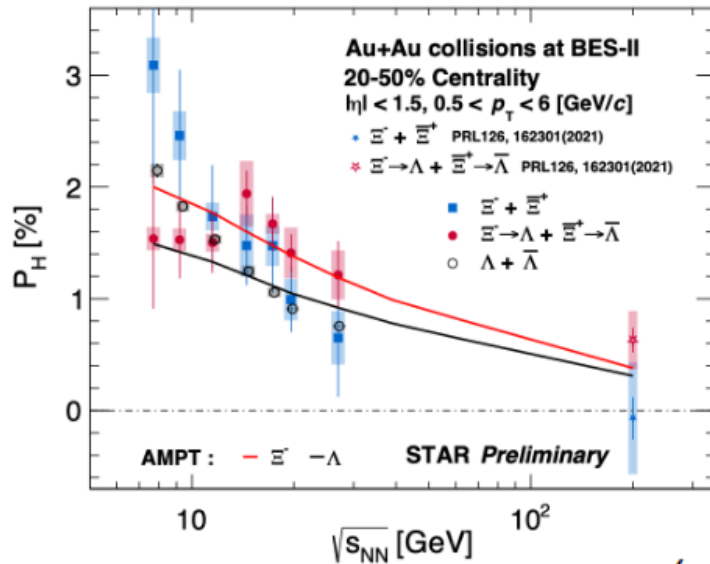
STAR collaboration,
Nature 2017

$$e^{-\beta(H_0 - \mathbf{S} \cdot \boldsymbol{\omega})}$$

Voloshin's talk

Becattini et al, PRC 2017

Global polarization of Ξ , Ω



point particle

$$P_H \approx \frac{(s+1)}{3} \frac{\omega}{T}$$

Becattini et al, PRC 2017

quark model

$$P_\Lambda = P_s \quad P_\Xi \approx \frac{4P_s - P_d}{3} \quad P_\Omega \approx \frac{5P_s}{3}$$

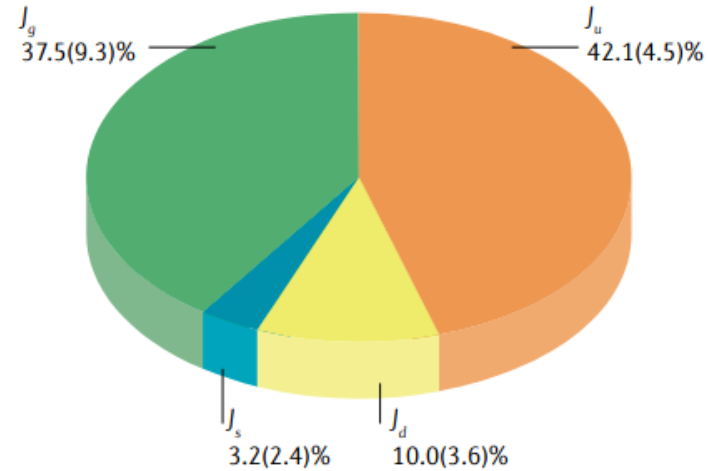
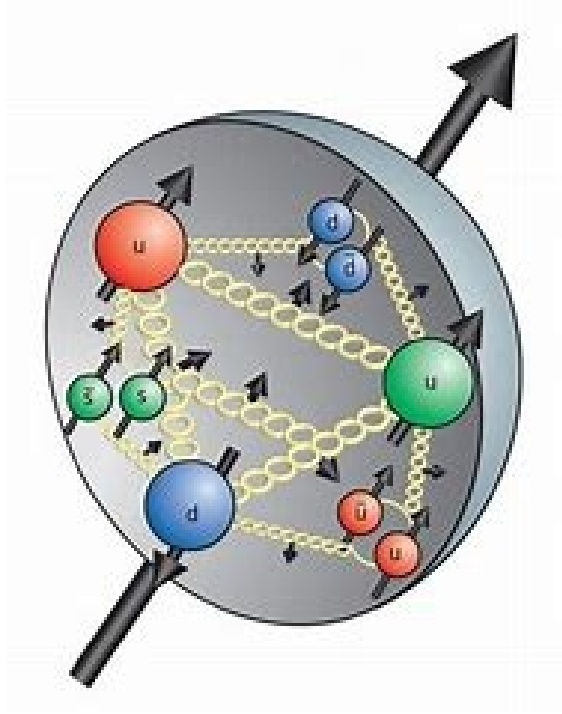
Liang, Wang, PRL 2005

→ $P_\Omega \approx \frac{5P_\Lambda}{3} \quad P_\Xi \approx P_\Lambda$

same prediction from non-interacting nature of models

disfavored by data?

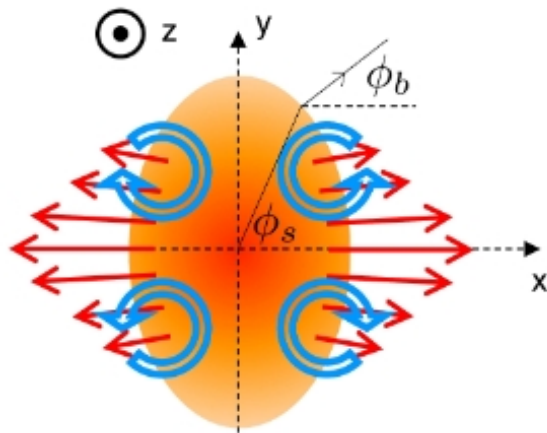
Interaction effect significant



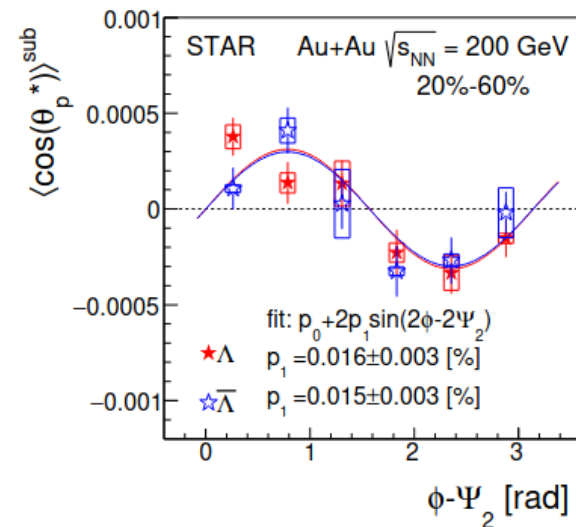
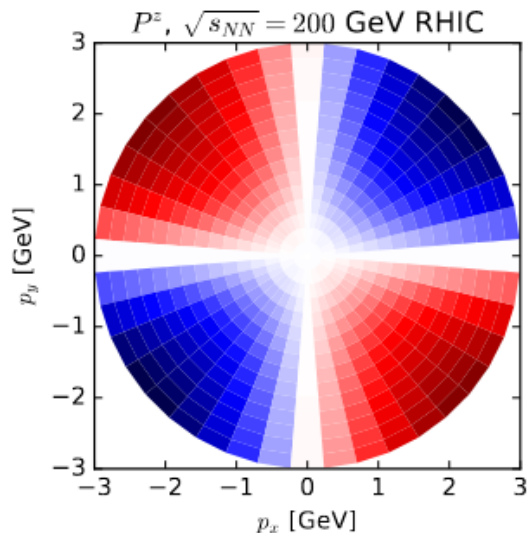
$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + L_q + \Delta G + L_g$$

~40%

Local polarization of Λ : vorticity only



$$S^i \sim \omega^i$$

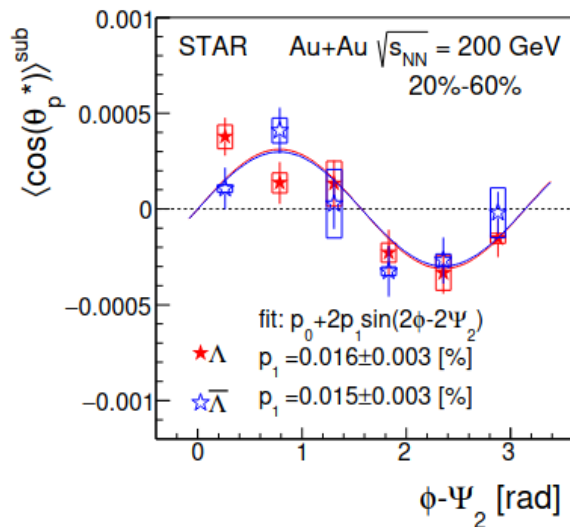


Becattini, Karpenko, PRL 2018
 Wei, Deng, Huang, PRC 2019
 Wu, Pang, Huang, Wang, PRR 2019
 Fu, Xu, Huang, Song, PRC 2021

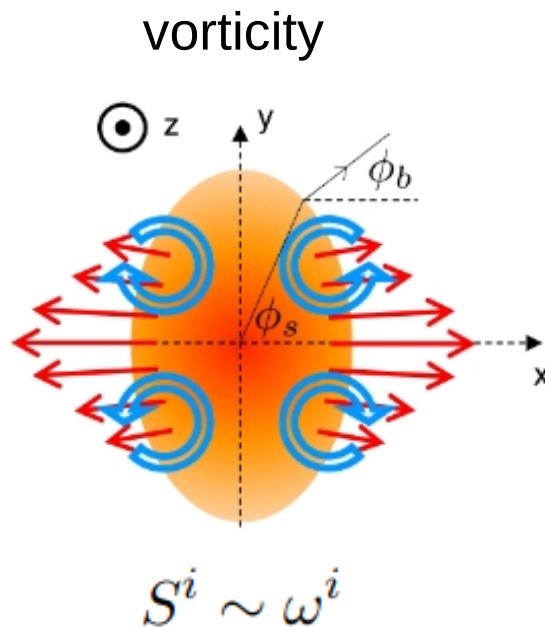
STAR collaboration, PRL
 2019

wrong sign

Local polarization of Λ : vorticity + shear



STAR collaboration, PRL
2019



wrong sign

shear

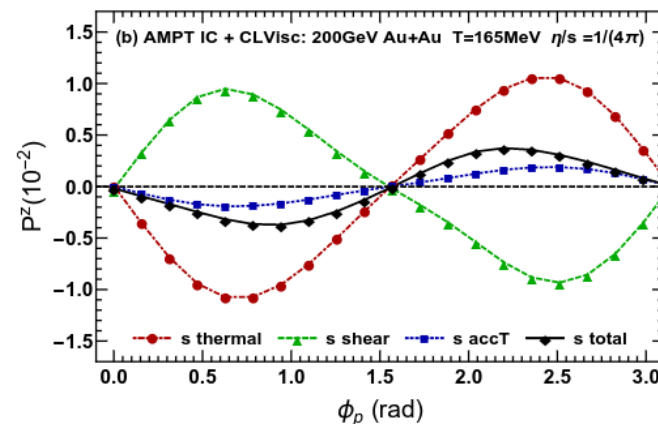
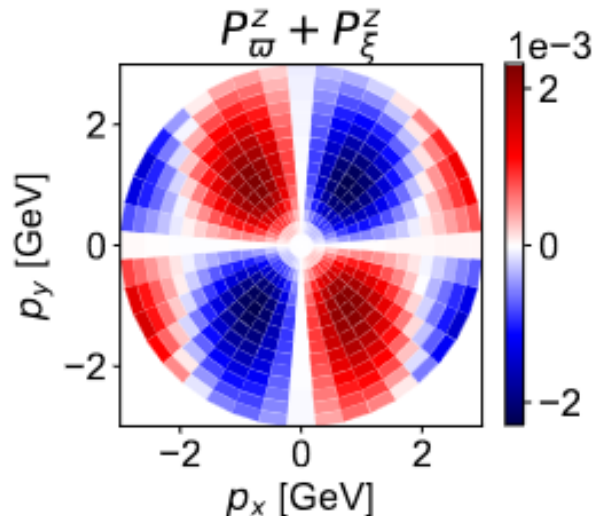
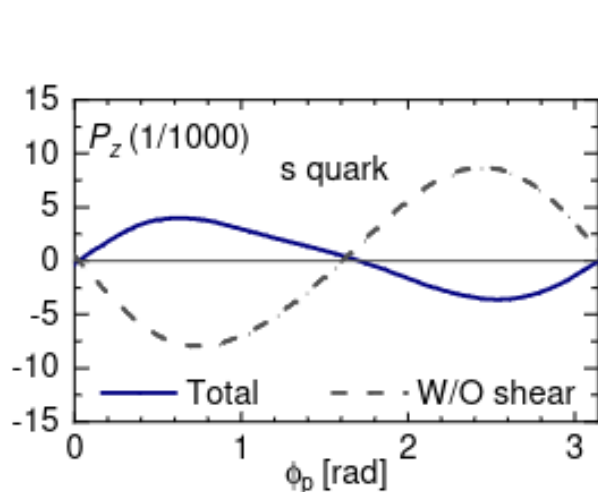
$$S^i \sim \epsilon^{ijk} \hat{p}_k \hat{p}_l \sigma_{jl}$$

$$S^z \sim (\langle p_y^2 \rangle - \langle p_x^2 \rangle) \partial_y u_x$$

right sign

Hidaka, Pu, Yang, PRD 2018
Liu, Yin, JHEP 2021
Becattini, et al, PLB 2021

Sensitivity on interaction: shear



two scenarios

Λ : point particle

Λ : quark model

$$S^\mu(p^\alpha) = \frac{1}{4m} \frac{\int d\Sigma \cdot p n_0 (1 - n_0) \mathcal{A}^\mu}{\int d\Sigma \cdot p n_0}$$



Fu, Liu, Pang, Song, Yin, PRL 2021

Becattini, et al, PRL 2021

Yi, Pu, Yang, PRC 2021

C. Yi's talk

Interaction effect?

Interaction effect in medium: weak coupling picture

free theory

$$S^i \sim \left(1 \omega^i + 1 \epsilon^{ijk} \hat{p}_j \hat{p}_l \sigma_{kl} + 1 \epsilon^{ijk} \hat{p}_j \frac{\partial_k \beta}{\beta} \right) \delta(p^2 - m^2)$$

- distribution function correction

$$(\partial_t + \hat{p} \cdot \nabla_x) f(t, x, p) = -C[f] \quad \delta f \propto \frac{\nabla_x f}{g^4}$$

deviation of steady state
from local equilibrium

$$S^i \sim \left((1 + O(1)) \epsilon^{ijk} \hat{p}_j \hat{p}_l \sigma_{kl} \right) \delta(P^2 - m^2)$$

SL, Wang, JHEP 2022, PRD 2025
Fang, Pu, Yang, PRD 2024
Fang, Pu, PRD 2025

- spectral function correction
in off-equilibrium state

focus of this talk

SL, Tian, 2410.22935
Fang, Pu, Yang, 2503.13320

Spectral function of probe “particle” in medium

$$\rho_{\alpha\beta}(\omega, \vec{p}) = \int d^4x e^{ip \cdot x} \left\langle \Psi_{\alpha}(x) \Psi_{\beta}^{\dagger}(0) + \Psi_{\beta}^{\dagger}(0) \Psi_{\alpha}(x) \right\rangle$$

$\omega, p \gg \partial_X$ particle localized in fluid element

$\langle \dots \rangle = \text{Tr}[D \dots]$ D: density matrix

probe particle in off-equilibrium fluid $D = D_{\text{probe}} \otimes D_{\text{fluid}}$.

$$\begin{array}{ccc} D^{(0)}_{\text{probe}} \otimes D^{(1)}_{\text{fluid}} & & D^{(1)}_{\text{probe}} \otimes D^{(0)}_{\text{fluid}} \\ \rho = 2 \text{Im } G^R & \longrightarrow & S^i \sim \text{tr}[\rho \sigma^i] \end{array}$$

Off-equilibrium spectral function at weak coupling

- Probe quark

Spectral function from self-energy

$$\frac{i}{2}\not{\partial}S_R(X,P) + \not{P}S_R(X,P) - \left(\Sigma_R(X,P)S_R(X,P) + \frac{i}{2}\{\Sigma_R(X,P), S_R(X,P)\}_{\text{PB}} \right) = -1.$$

$$S_R = S_R^{(0)} + S_R^{(1)} + \cdots,$$

$$\{A, B\}_{\text{PB}} = \partial_P A \cdot \partial_X B - \partial_X A \cdot \partial_P B.$$

$$S_R^{(0)} = -\frac{1}{\not{P}} - \frac{1}{\not{P}}\Sigma_R\frac{1}{\not{P}}.$$

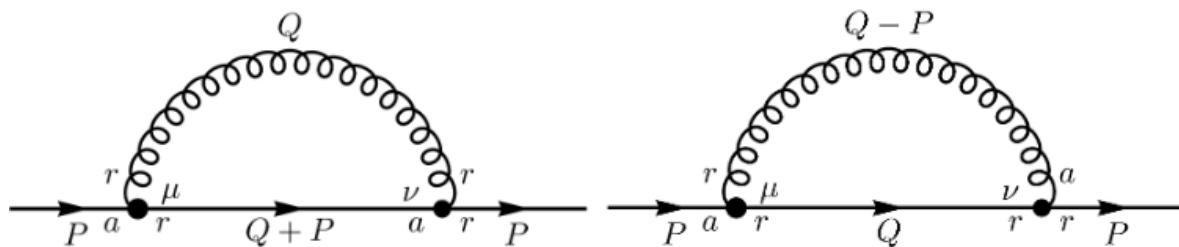
$$S_R^{(1)} = -\frac{1}{\not{P}}\delta\Sigma_R\frac{1}{\not{P}} + \gamma^5\gamma^\beta P^\nu T^{\mu\lambda}\epsilon_{\beta\lambda\mu\nu}\frac{-1}{(P^2)^2}$$

$$T_{\mu\lambda} = \partial_{[\mu}\Sigma_{\lambda]}^R$$

off-equilibrium self-energy

gradient of equilibrium self-energy

Equilibrium self-energy



$$P \gg T, \quad P^2 \ll p/\beta,$$

energetic quark
nearly on-shell

$$\frac{\Sigma_{ar}}{g^2 C_F} = 2i\cancel{P}(A+B) + 4ip_0\gamma^0 A.$$

→ $p_0 = p(1 + 8A)$

modified dispersion:

$$A = \frac{1}{2(2\pi)^2} \frac{-i\pi}{2p\beta}.$$

finite damping

$$S_R^{(1)} = \gamma^5 \gamma_\beta P_\nu \partial_\mu \Sigma_\lambda^R \frac{-1}{(P^2)^2} \epsilon^{\beta\lambda\mu\nu} \sim O(T\partial) \quad D^{(0)}_{\text{probe}} \otimes D^{(1)}_{\text{fluid}}$$

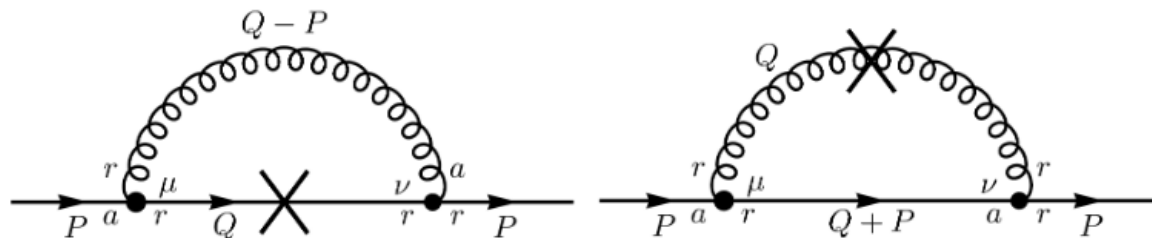
Off-equilibrium self-energy

off-equilibrium propagators

Hidaka, Pu, Yang 2017

Huang et al 2020

Hattori et al 2020



$$\frac{\delta \Sigma_{ar}}{g^2 C_F} = \gamma^5 \gamma^\mu \mathcal{A}_\mu,$$

$$\mathcal{A}^0 = i\omega^i p_i \beta (-4\delta A - 2\delta B),$$

$$\mathcal{A}^k = i\omega_{\parallel}^k \beta p (-4\delta A - 2\delta B) + \omega_{\perp}^k \beta p (-4\delta A - \delta B) + \epsilon^{ijk} \hat{p}_i \hat{p}_l \sigma_{jl} \beta p (-\delta B) + \epsilon^{ijk} p_i \partial_j \beta (-\delta C)$$

$$\delta S_R^{(0)} = -\frac{1}{\not{P}} \delta \Sigma_R \frac{1}{\not{P}} \sim O(p\partial)$$

dominant contribution

$P \gg T$

$$D^{(1)}_{\text{probe}} \otimes D^{(0)}_{\text{fluid}}$$

SL, Tian, 2410.22935

Weak coupling summary

free $\gamma^5 \gamma_i \frac{2\pi}{2} \left(\omega^i + \epsilon^{ijk} \hat{p}_k \frac{\partial_j \beta}{\beta} + \epsilon^{ijk} \hat{p}_k \hat{p}_l \sigma_{lj} \right) \beta \tilde{f}(p),$

interaction $\int dp_0 \delta\rho(P) \tilde{f}(p_0)$

$$O(\lambda) = \frac{g^2 C_F}{2(2\pi)^2} \frac{\pi}{2p} \gamma^5 \gamma_i \left[0.95 \omega_{\parallel}^i + 1.48 \omega_{\perp}^i - 0.52 \epsilon^{ijk} \hat{p}_j \hat{p}_l \sigma_{kl} - 0.02 \epsilon^{ijk} \hat{p}_j \frac{\partial_k \beta}{\beta} \right] \tilde{f}(p)$$

Polarized spectral function

Off-equilibrium spectral function at strong coupling

- Probe baryon

Holographic model for baryon

$$S = i \int d^{D+1}x \sqrt{-g} \bar{\psi} (\Gamma^M \nabla_M - m) \psi,$$

Iqbal, Liu 2009

5D Dirac fermion



4D Weyl fermion

$$\psi = \begin{pmatrix} \psi_R \\ \psi_L \end{pmatrix}$$

“lump of quark/gluon”

One baryon = R + L Weyl fermions

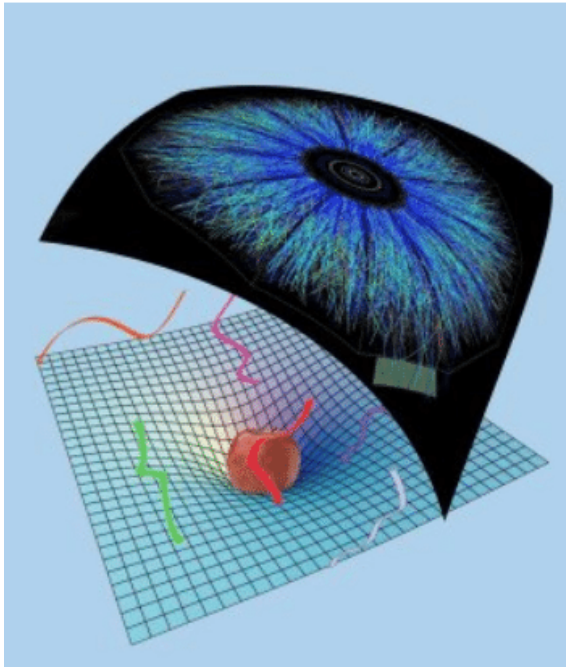
focus on spectral function in off-equilibrium QGP

Holographic model for QGP

$$ds^2 = -2u_\mu(x)dx^\mu dr - r^2 f(b(x)r)u_\mu u_\nu dx^\mu dx^\nu + r^2 P_{\mu\nu} dx^\mu dx^\nu \quad \text{local equilibrium}$$

$$+ 2r^2 b F(br) \sigma_{\mu\nu} dx^\mu dx^\nu + \frac{2}{3} r u_\mu u_\nu \partial_\lambda u^\lambda dx^\mu dx^\nu - r u^\lambda \partial_\lambda (u_\mu u_\nu) dx^\mu dx^\nu \quad \text{steady state}$$

Bhattacharyya et al,
JHEP 2008



$$T^{\mu\nu} = \overset{O(\partial^0)}{(\pi T)^4 (\eta^{\mu\nu} + 4 u^\mu u^\nu)} - 2 \overset{O(\partial)}{(\pi T)^3} \sigma^{\mu\nu}$$

local equilibrium η
steady state

$$\sigma_{ij} \quad \cancel{\partial_0 b} = \frac{1}{3} \cancel{\partial_i u_i} \quad \partial_i b = \partial_0 u_i$$

T-grad =
acceleration

Gradient corrections

$$(\Gamma^M \nabla_M - m) \psi = 0$$

gradient correction to Dirac field $\longrightarrow$ gradient correction to baryon

$$D_{\text{fluid}}^{(0)} \otimes D_{\text{baryon}}^{(1)}$$

$$(\Gamma^M \nabla_M - m) \psi = 0$$

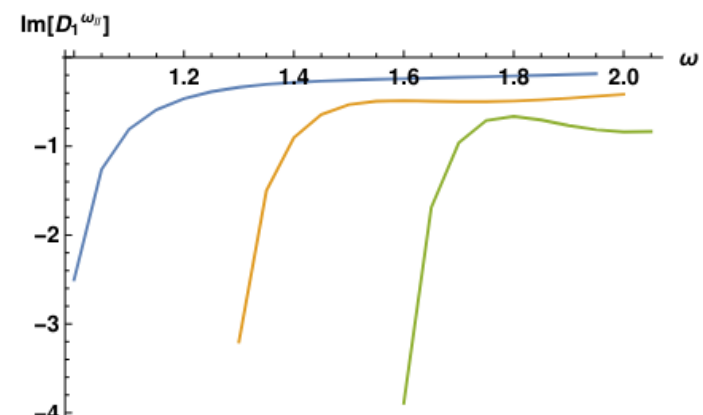
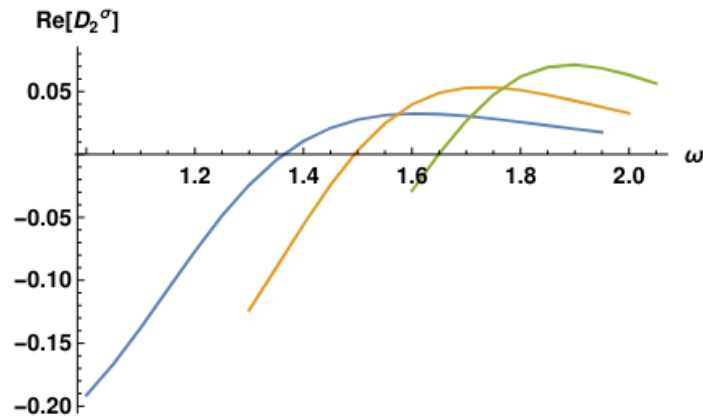
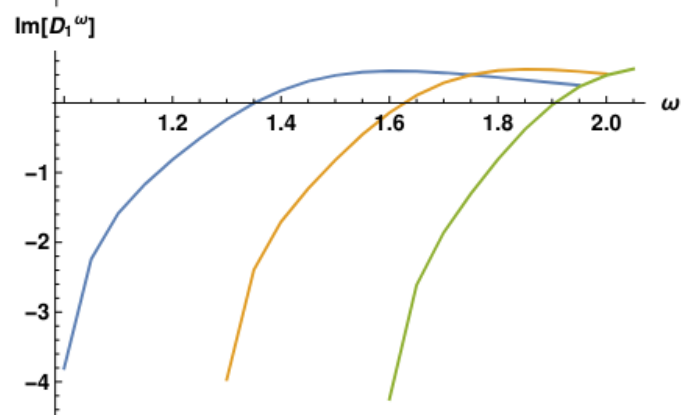
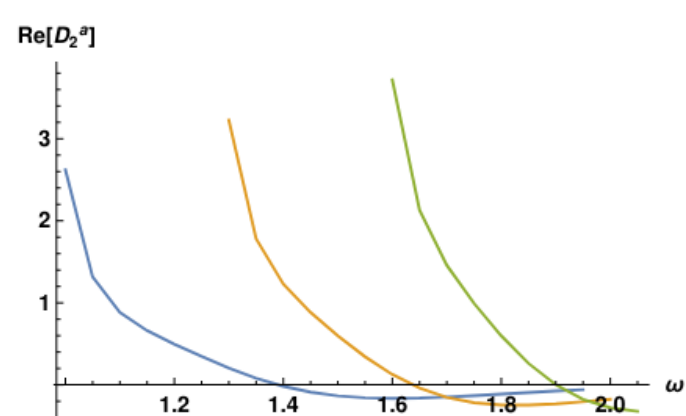
gradient correction to Dirac operator $\longrightarrow$ gradient correction to medium

$$D_{\text{fluid}}^{(1)} \otimes D_{\text{baryon}}^{(0)}$$

$$D_{\text{fluid}}^{(1)\text{le}} \quad D_{\text{fluid}}^{(1)\text{ss}}$$

$$\frac{1}{2} \text{tr} [(\delta\rho^L + \delta\rho^R)\sigma_k] = 4 \left(-\text{Re}[D_2^a]\epsilon^{ijk}\hat{p}_j\partial_0 u_i - \text{Re}[D_2^\sigma]\epsilon^{ijk}\hat{p}_j\hat{p}_l\sigma_{il} + \text{Im}[D_1^\omega]\omega_k + \text{Im}[D_1^{\omega_\parallel}]\omega_k^\parallel \right)$$

$$D_{\text{fluid}}^{(0)} \otimes D_{\text{baryon}}^{(1)}$$



— $p=0.9$
— $p=1.2$
— $p=1.5$

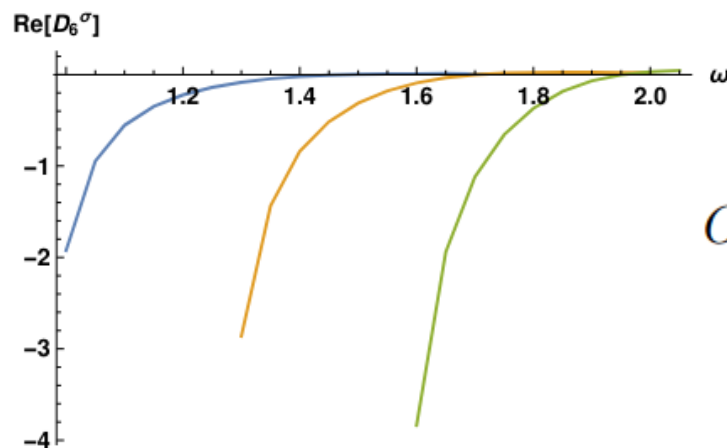
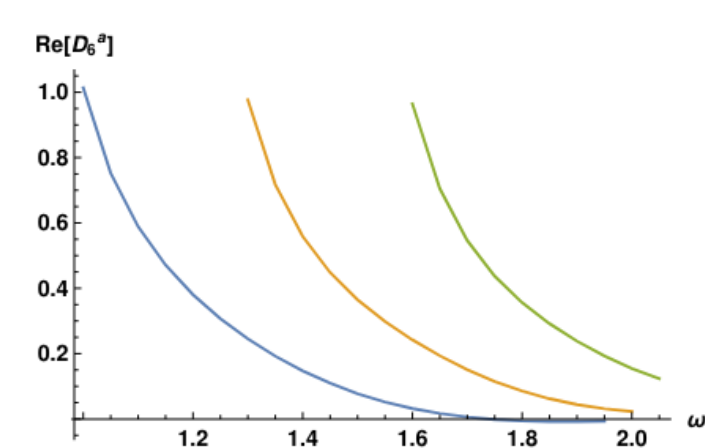
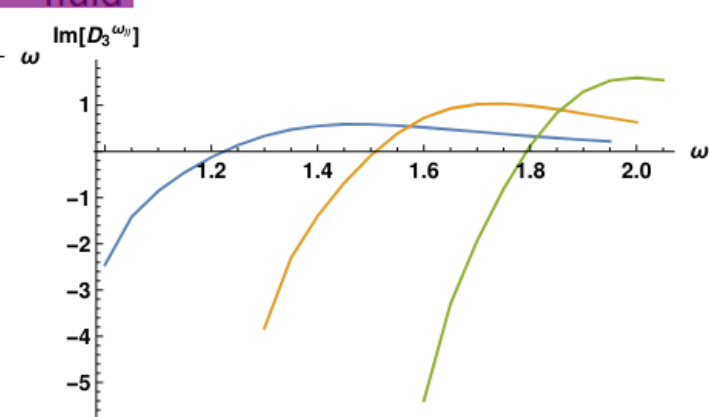
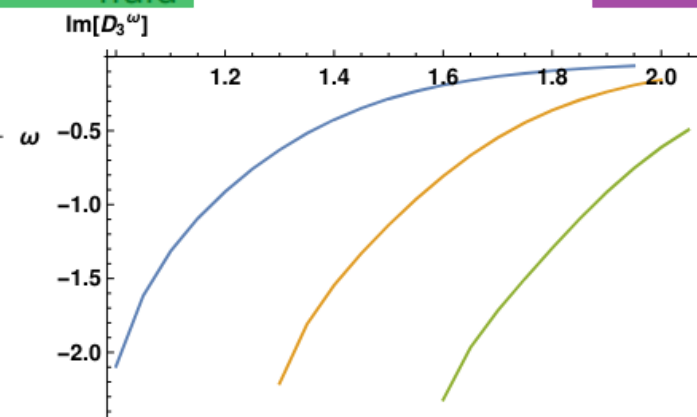
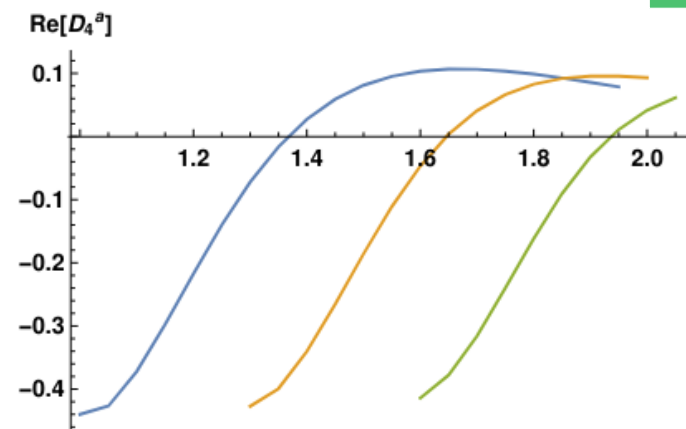
$$\pi T = 1.$$

$$O(\lambda^0)$$

$$\frac{1}{2}\text{tr}[(\delta\rho^L + \delta\rho^R)\sigma_k] = 4\left[-\text{Re}[D_4^a + D_6^a]\epsilon^{ijk}\hat{p}_j\partial_0 u_i - \text{Re}[D_6^\sigma]\epsilon^{ijk}\hat{p}_j\hat{p}_l\sigma_{il} + \text{Im}[D_3^\omega]\omega_k + \text{Im}[D_3^\omega]\omega_k^\parallel\right]$$

$D^{(1)\text{le}}_{\text{fluid}}$

$D^{(1)\text{ss}}_{\text{fluid}}$



$p=0.9$
 $p=1.2$
 $p=1.5$

$O(\lambda^0)$ $\pi T = 1.$

Li, SL, 2505.21883

Strong coupling summary

$\text{tr}[(\delta\rho^L + \delta\rho^R)\sigma_k]$	$\epsilon^{ijk}\hat{p}_j\partial_0 u_i$	$\epsilon^{ijk}\hat{p}_j\hat{p}_l\sigma_{il}$	$\omega_k^\perp$	$\omega_k^\parallel$
$D_{\text{probe}}^{(1)}$	$-4\text{Re}[D_2^a]$	$-4\text{Re}[D_2^\sigma]$	$4\text{Im}[D_1^\omega]$	$4\text{Im}[D_1^\omega + D_1^{\omega^\parallel}]$
$D_{\text{fluid}}^{(1)\text{le}}$	$-4\text{Re}[D_4^a]$		$4\text{Im}[D_3^\omega]$	$4\text{Im}[D_3^\omega + D_3^{\omega^\parallel}]$
$D_{\text{fluid}}^{(1)\text{ss}}$	$-4\text{Re}[D_6^a]$	$-4\text{Re}[D_6^\sigma]$		

Polarized spectral function

Weak (quark) vs Strong (baryon)

Common

- ♦ Polarized spectral function in off-equilibrium state

$$D^{(1)}_{\text{probe}} \quad D^{(1)\text{le}}_{\text{fluid}} \quad D^{(1)\text{ss}}_{\text{fluid}}$$

- ♦ Correction suppressed for energetic particles

Difference

- ♦ Weak $O(\lambda)$ vs strong $O(\lambda^0)$

Summary

- ♦ Sensitivity of polarization to baryon structure
- ♦ Spectral function in off-equilibrium state
- ♦ vorticity, shear, T-grad generically lead to (interaction dependent) polarized spectral function

Outlook

- ♦ Mass dependence of polarized spectral function

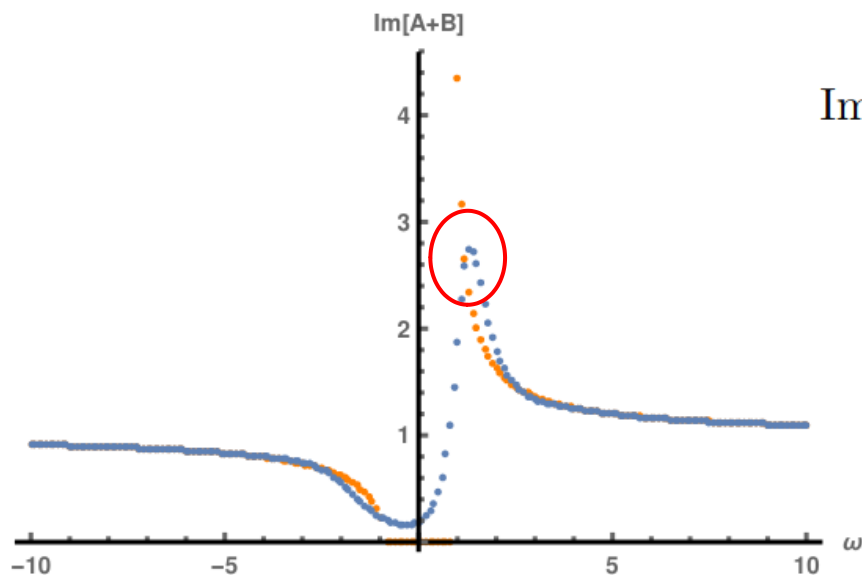
The 10th International Conference on Chirality, Vorticity and Magnetic Field in Quantum Matter, Zhuhai, Nov 16-20, 2026



Holography: equilibrium spectral function

$$G_R = A + B \hat{p} \cdot \vec{\sigma}$$

$$\rho = 2\text{Im}G_R = \text{Im}[A + B] (1 - \hat{p} \cdot \vec{\sigma}) + \text{Im}[A - B] (1 + \hat{p} \cdot \vec{\sigma})$$



orange: vacuum

Iqbal, Liu 2009

$$\text{Im}[A + B] = \text{sgn}(\omega) \theta(\omega^2 - p^2) \left(\frac{\omega + p}{\omega - p} \right)^{1/2}$$

no spacelike spectral

blue: equilibrium QGP

soften the singularity

develop spacelike spectral

focus on **timelike spectral for baryon**